

Evolution of structure functions at NLO without PDFs

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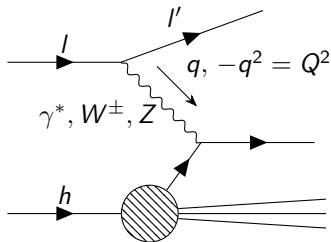
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Center of Excellence in Quark Matter

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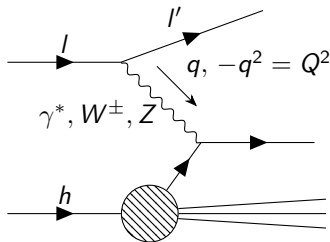
Background

- Deep Inelastic Scattering (DIS):



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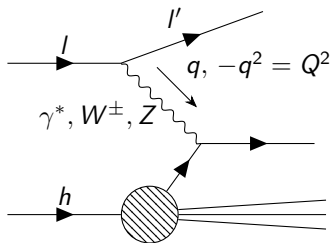
- Deep Inelastic Scattering (DIS):



- Measured cross sections expressed in terms of structure functions
- **Structure functions** expressed in terms of **parton distribution functions (PDFs)**
 $F_i(x, Q^2) = \sum_j C_{ij}(Q^2, \mu^2) \otimes f_j(\mu^2) \quad j = q, \bar{q}, g \quad \mu = \text{factorization scale}$

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- The conventional procedure in collinear factorization:
 - ▶ PDFs are fitted to DIS data (to structure functions)
 - ▶ Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) evolution: PDFs to higher Q^2

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 - ▶ Parametrize non-observable quantities
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- Physical basis \equiv set of linearly independent DIS observables
- DGLAP evolution of observables in a physical basis
 - ▶ Avoiding the problems with PDFs
 - ▶ More straightforward to compare to experimental data

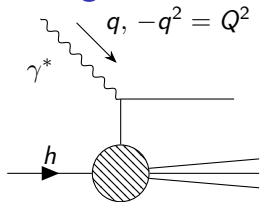
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- Previously discussed e.g. in Harland-Lang and Thorne [1811.08434](#), Hentschinski and Stratmann [1311.2825](#), W.L. van Neerven and A. Vogt [hep-ph/9907472](#)

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- The novelty of our work:
 - ▶ Momentum space
 - ▶ Full three-flavor basis at NLO
- NLO physical basis [2412.09589](#) continuation for LO work [2304.06998](#)

Straightforward example with only two observables



$$F_i(x, Q^2) = \sum_j C_{F_i f_j}(Q^2, \mu^2) \otimes f_j(\mu^2),$$

where $F_i = F_2, F_L / \frac{\alpha_s}{2\pi}$, and $f_j = \Sigma, g$

Quark singlet:

$$\Sigma(x, \mu^2) = \sum_q^{n_f} [q(x, \mu^2) + \bar{q}(x, \mu^2)], \quad n_f = 3$$

Gluon PDF: $g(x, \mu^2)$

First step: invert the linear mapping (difficult because $f \otimes g = \int_x^1 \frac{dz}{z} f(z) g(\frac{x}{z})$)

$$f_j(\mu^2) = \sum_i C_{F_i f_j}^{-1}(Q^2, \mu^2) \otimes F_i(Q^2) + \mathcal{O}(\alpha_s^2)$$

DGLAP evolution in physical basis

$$\begin{aligned} \frac{dF_i(x, Q^2)}{d \log(Q^2)} &= \sum_j \frac{dC_{F_i f_j}(Q^2, \mu^2)}{d \log(Q^2)} \otimes f_j(\mu^2) \\ &= \sum_j \frac{dC_{F_i f_j}(Q^2, \mu^2)}{d \log(Q^2)} \otimes \sum_k C_{F_k f_j}^{-1}(Q^2, \mu^2) \otimes F_k(Q^2) + \mathcal{O}(\alpha_s^3) \end{aligned}$$

Scheme and scale dependence at NLO

DGLAP evolution in physical basis:

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Kernels \mathcal{P}_{ik} are independent of the factorization scheme and scale

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Kernels \mathcal{P}_{ik} are independent of the factorization scheme and scale

\mathcal{P}_{ij} 's determined by:

- Splitting functions
- Coefficient functions

→ The scheme and scale dependence exactly cancels between these two

Inverting the gluon PDF at NLO

Simple example without quarks

Invert $g(x)$ from $\tilde{F}_L = C_{F_L g}^{(1)} \otimes g + \frac{\alpha_s}{2\pi} C_{F_L g}^{(2)} \otimes g$ where $\tilde{F}_L(x, Q^2) \equiv \frac{2\pi}{\alpha_s} \frac{F_L(x, Q^2)}{x}$

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Define inverse of $C_{F_L g}^{(1)}$ as: $g(x) = \hat{P}(x) \left[C_{F_L g}^{(1)} \otimes g \right]$ with $\hat{P}(x) \equiv \frac{1}{8T_R n_f \bar{e}_q^2} \left[x^2 \frac{d^2}{dx^2} - 2x \frac{d}{dx} + 2 \right]$

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Get $C_{F_L g}^{(1)} \otimes g$ from \tilde{F}_L : $C_{F_L g}^{(1)} \otimes g = \tilde{F}_L - \frac{\alpha_s}{2\pi} C_{F_L g}^{(2)} \otimes g$

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$$g(x) = \hat{P}(x) \left[\tilde{F}_L(x) - \frac{\alpha_s}{2\pi} C_{F_L g}^{(2)} \otimes g \right]$$

Plug in $g(x) = \hat{P}(x) \tilde{F}_L(x) + \mathcal{O}(\alpha_s)$ to the right hand side

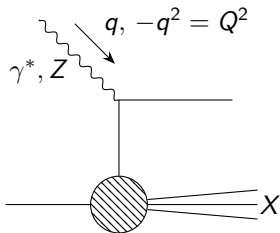
$$g(x) = \hat{P}(x) \tilde{F}_L(x) - \frac{\alpha_s(Q^2)}{2\pi} \hat{P}(x) \left[C_{F_L g}^{(2)} \otimes \hat{P} \tilde{F}_L \right] + \mathcal{O}(\alpha_s^2)$$

Six observable basis

- Full three-flavor basis: $u, \bar{u}, d, \bar{d}, s = \bar{s}$, and g
→ Need six linearly independent DIS structure functions

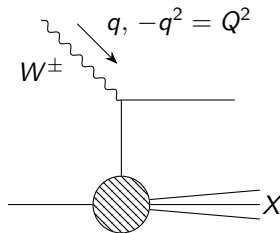
Six observable basis

- Full three-flavor basis: $u, \bar{u}, d, \bar{d}, s = \bar{s}$, and g
→ Need six linearly independent DIS structure functions
- We choose the NLO structure functions:



Neutral current γ^*, Z

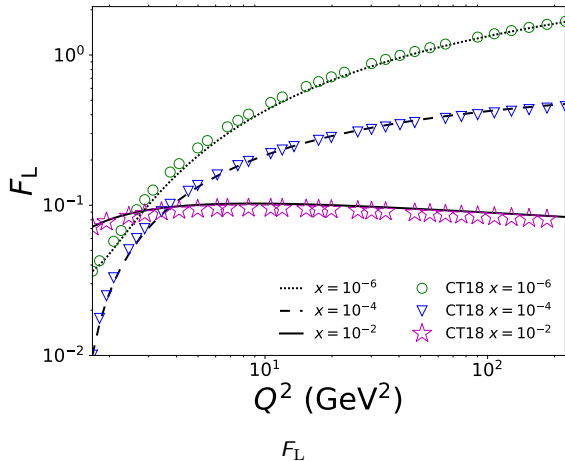
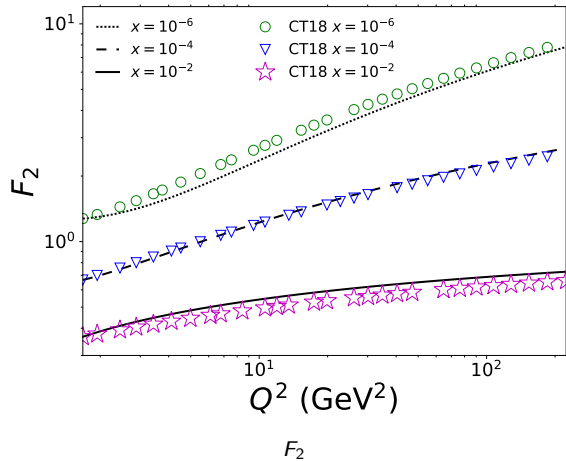
- γ^* exchange → F_2 and F_L
- Z boson exchange → F_3



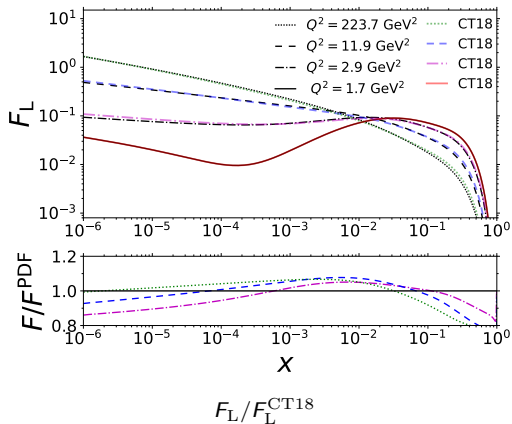
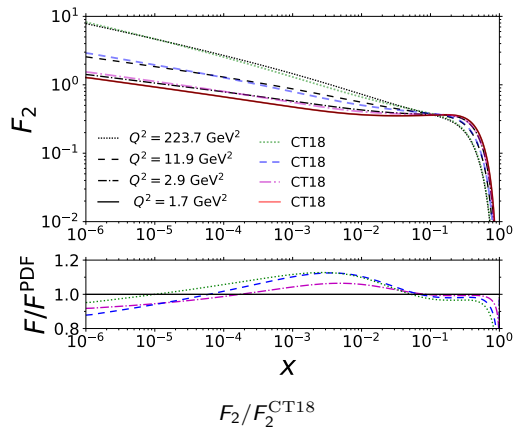
Charged current W^\pm

- W^- exchange → $F_3^{W^-}$ and $F_{2c}^{W^-}$
- $\Delta F_2^W = F_2^{W^-} - F_2^{W^+}$

Comparison with conventional DGLAP evolution



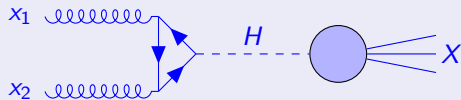
Comparison with conventional DGLAP evolution



- Similar Q^2 evolution
- Differences in values from:
 - ▶ uncertainty in PDFs from scheme and scale (error band not shown)
 - ▶ perturbative truncation

Cross sections in terms of physical basis

Example of Higgs production by gluon fusion

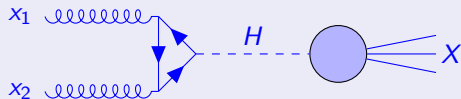


$$\sigma(p + p \rightarrow H + X) = \int dx_1 dx_2 g(x_1, \mu) g(x_2, \mu) \hat{\sigma}_{gg \rightarrow H+X}(x_1, x_2, \frac{m_H^2}{\mu^2}),$$

where m_H is the Higgs mass, $g(x_1, \mu)$ and $g(x_2, \mu)$ are the **gluon PDFs**

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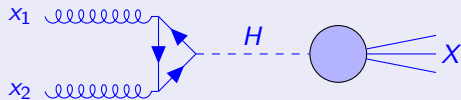
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Plug in the gluon PDF in physical basis: $g(x, \mu^2) = \sum_j C_{jg}^{-1}(Q^2, \mu^2) \otimes F_j(Q^2)$

where $F_j = F_2, F_L/\frac{\alpha_s}{2\pi}, F_3, \Delta F_2^W, F_3^{W-}, F_{2c}^{W-}$

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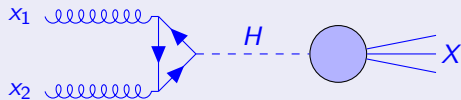
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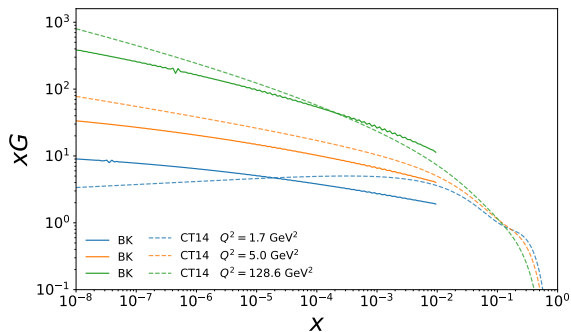
Harland-Lang and Thorne [1811.08434](#):

explicit μ dependence vanishes and terms $\log(Q^2/m_H^2)$ are left behind

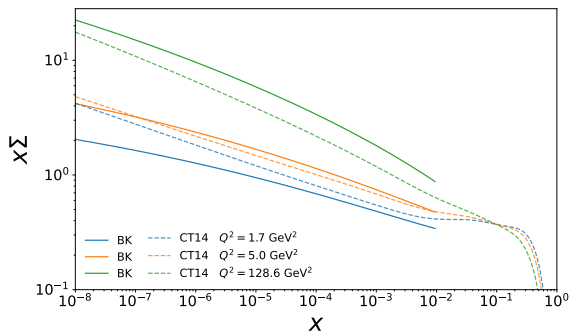
→ no need to choose relation between μ and Q or m_H

PDFs from BK-evolved structure functions

Now we have analytical form to calculate gluon PDF and quark singlet from F_2 and F_L in dipole picture



LO gluon



LO quark singlet

- Weaker x -evolution with BK-evolved $F_{2,L}$
- Bigger difference in gluon than in quark singlet

Comparison to BK-evolved $F_{2,L}$ (work in preparation)

Goal

Set BK-evolved F_2 and F_L as initial condition for (2-observable) physical basis DGLAP evolution
→ compare BK vs. DGLAP dynamics

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• *However..*

- ▶ LO DGLAP evolution (and NLO PDFs) in physical basis includes convolutions e.g.

$$P_{qq} \otimes F_2 = \int_x^1 \frac{dz}{z} P_{qq}(z) F_2\left(\frac{x}{z}\right)$$

→ need $F_{2,L}$ initial values up to $x = 1$

- ▶ Validity region for BK-evolved $F_{2,L}$ only up to $x \sim 10^{-2}$

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Quick fix: BK-improved initial condition

- Initial values for $F_{2,L}$:

- ▶ at $x \leq 10^{-2}$ from BK/dipole picture
- ▶ at $x > 10^{-2}$ from DGLAP/collinear factorization

- Scale collinear factorization $F_{2,L}$ so that they match dipole picture values

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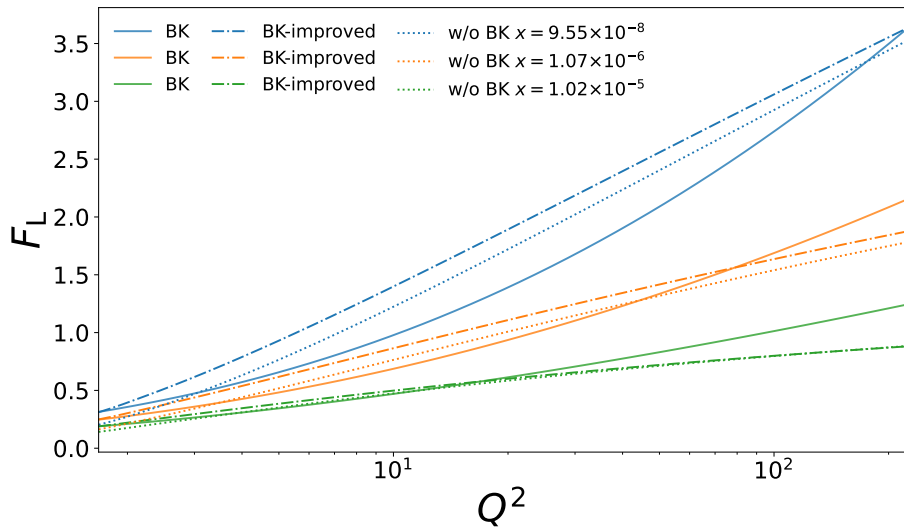
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- Scheme dependence of PDFs play a role at NLO in α_s
→ Scheme and scale dependence avoided in the physical basis
- What next:
 - ▶ BK vs. DGLAP comparison
 - ▶ Express LHC cross sections, e.g. Drell-Yan, in physical basis
 - ▶ Include heavy quarks

Backup: Comparison to BK-evolved F_L (work in preparation)



Backup: Comparison to BK-evolved F_2 (work in preparation)

