

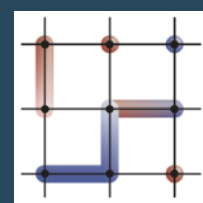
Tensor Networks: from quantum information to quantum many-body systems and QFT

Mari-Carmen Bañuls

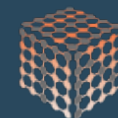
collaborators J.I. Cirac (MPQ), K. Cichy (Poznan), E. Zohar (HUJI), K. Jansen (DESY),
Y.-P. Lin, Y.-J. Kao (NTU), D.T.-L. Tan, C.-J. D. Lin (NCTU), M. Heller (Gent), J.
Knaute (HUJI), V. Svensson (Lund), P. Emonts, I. Papaestathiou, D. Robaina, (MPQ),



MAX PLANCK INSTITUTE
OF QUANTUM OPTICS



DFG FOR 5522



T-NiSQ

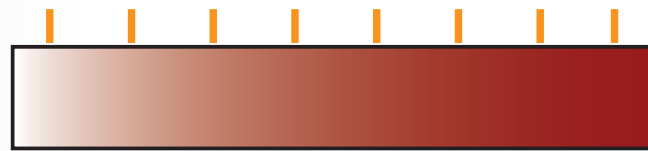
Tensor Networks in Simulation of Quantum Matter

3.4.2025

what are Tensor Networks?

TNS: entanglement-based ansatzes for
quantum many-body states

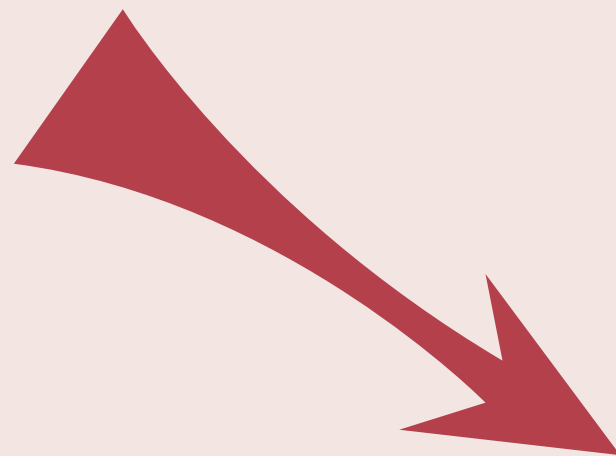
arbitrary many-
body state



d^N

$$|\Psi\rangle = \sum_{i_j} c_{i_1 \dots i_N} |i_1 \dots i_N\rangle$$

exponential

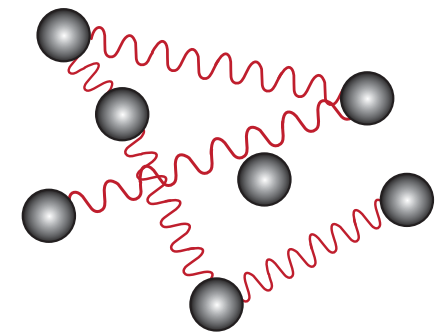


polynomial

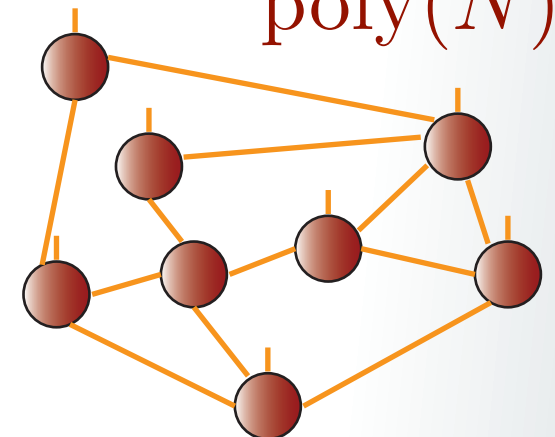
TNS: restricted
family

$\{|i\rangle\}_{i=0}^{d-1}$

N

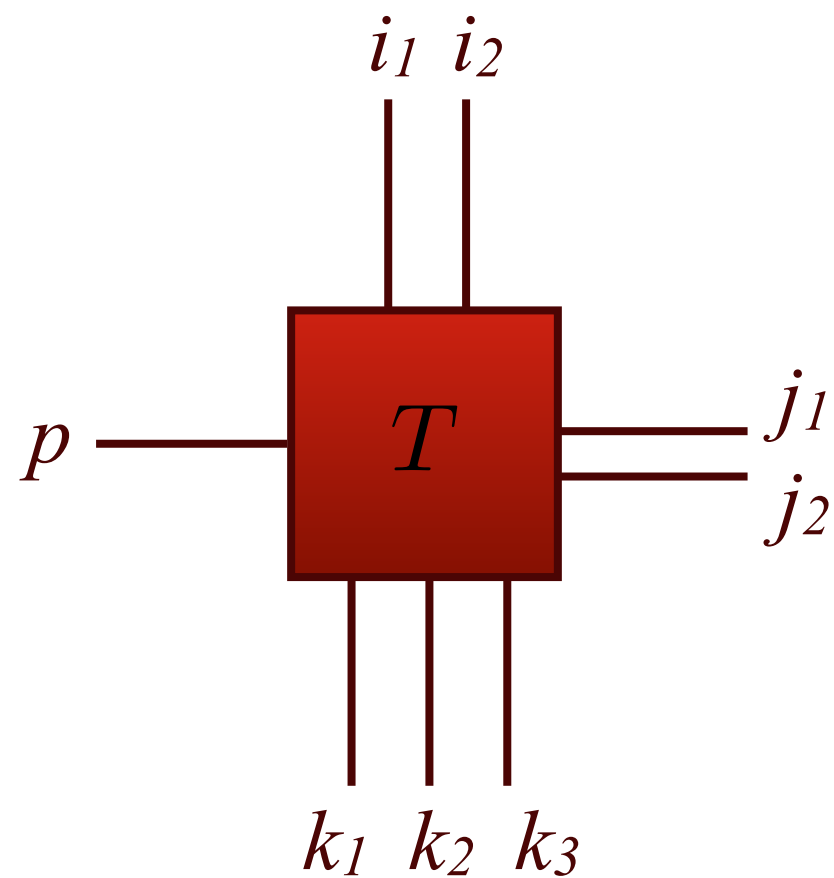


$\text{poly}(N)$



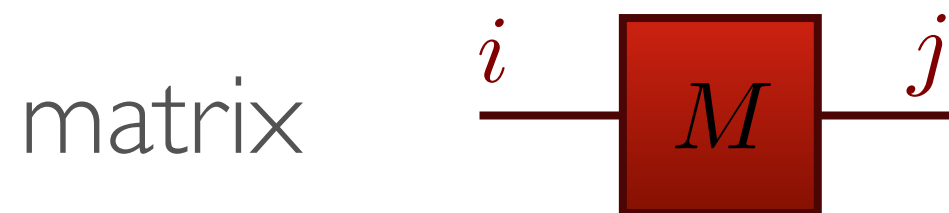
pictorial representation

tensor = multidimensional array



pictorial representation

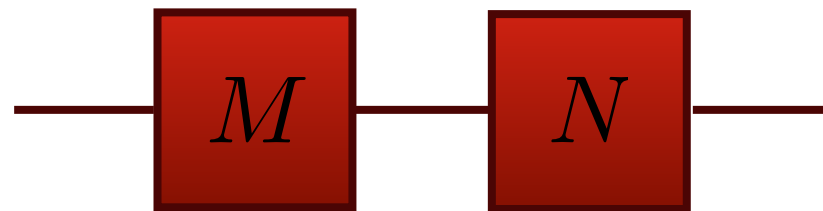
tensor = multidimensional array



M_{ij}

pictorial representation

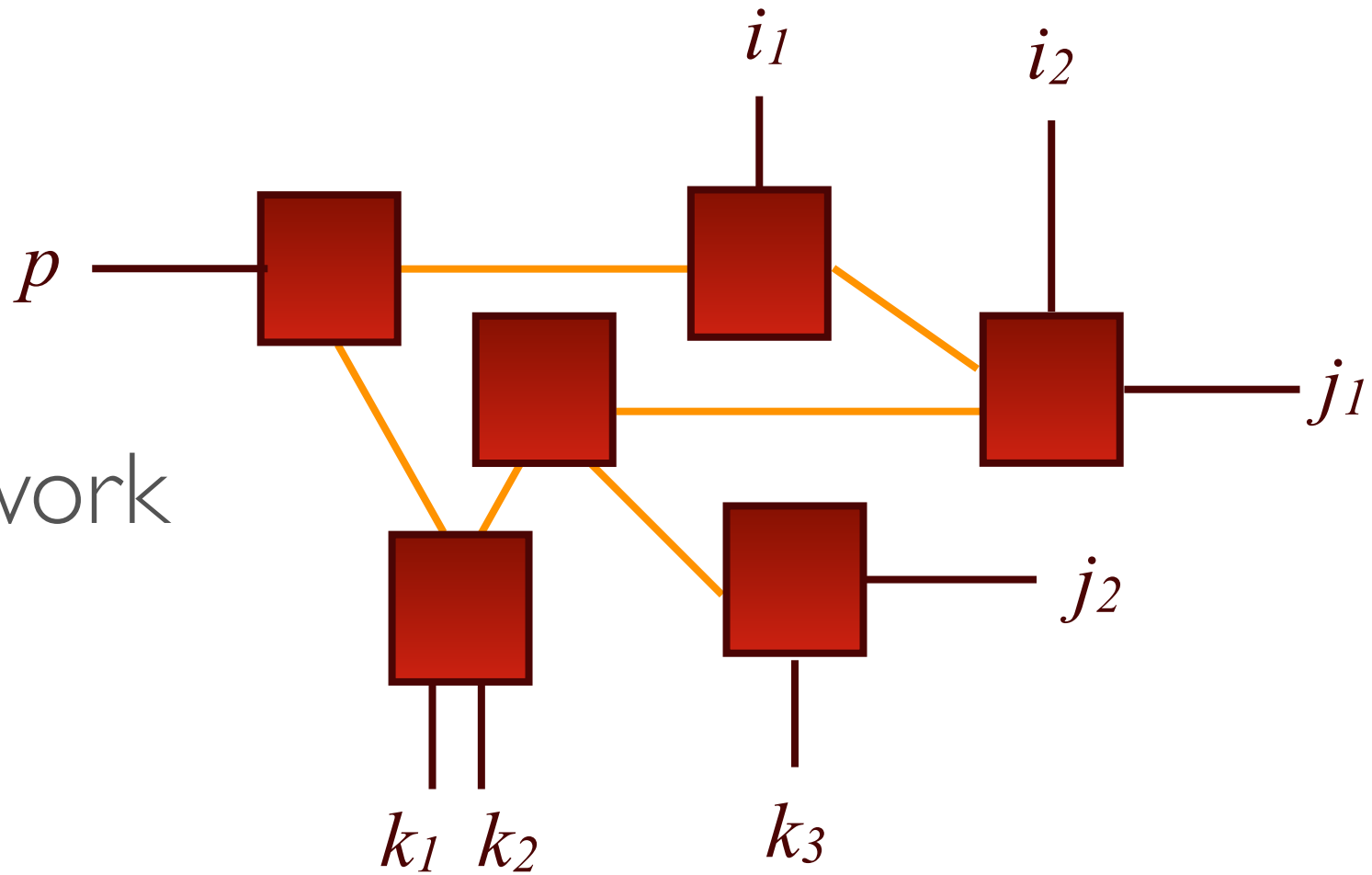
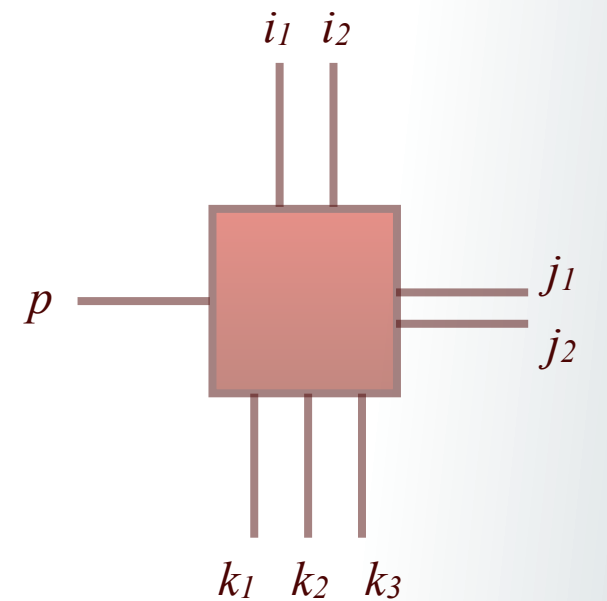
tensor = multidimensional array



$$M \cdot N = \sum_j M_{ij} N_{jk}$$

tensor network

tensor = multidimensional array

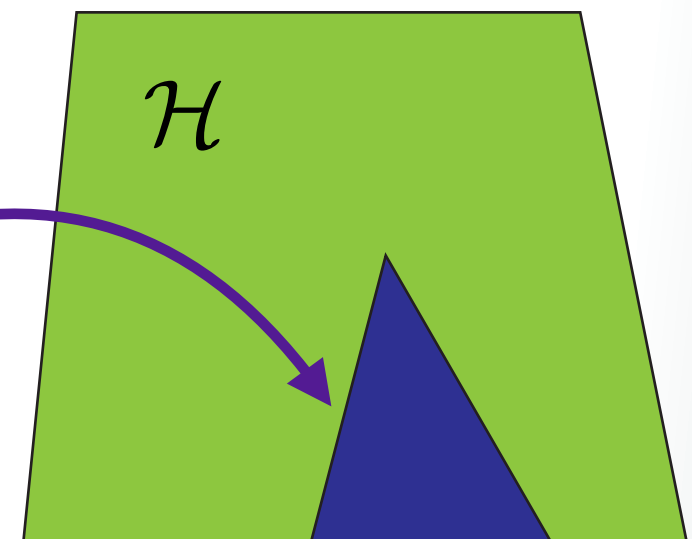

$$\{T_{i_1 i_2, j_1 j_2, k_1 k_2 k_3, p}\}_{\{i, j, k, p\}}$$


tensor network (TN)

why are TNS useful?

we are looking for states with small
entanglement

relevant corner of the
Hilbert space



local gapped 1D Hamiltonians
have ground states
with area law of entanglement

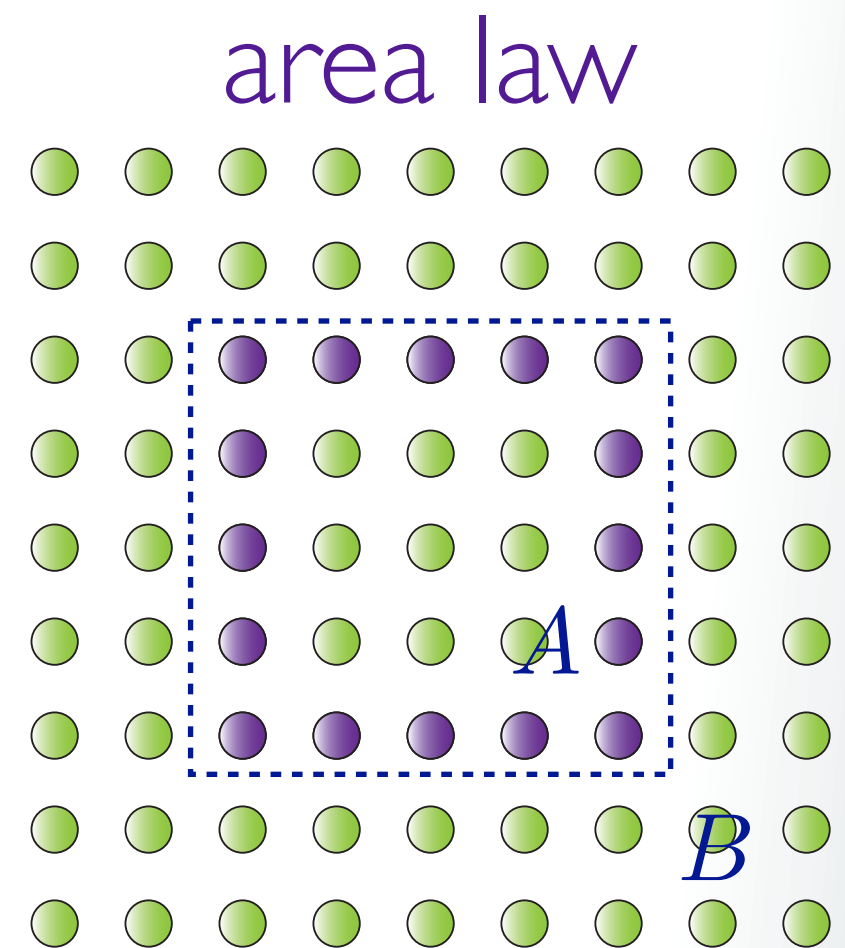
$$S_{A_{\max}} \propto |\delta A| \quad \text{Hastings 2007}$$

in 1D critical systems,
logarithmic corrections

$$S_{A_{\max}} \propto |\delta A| \log A \quad \text{Calabrese, Cardy 2004}$$

satisfied at finite temperature

Wolf, Verstraete, Hastings, Cirac, PRL 2008

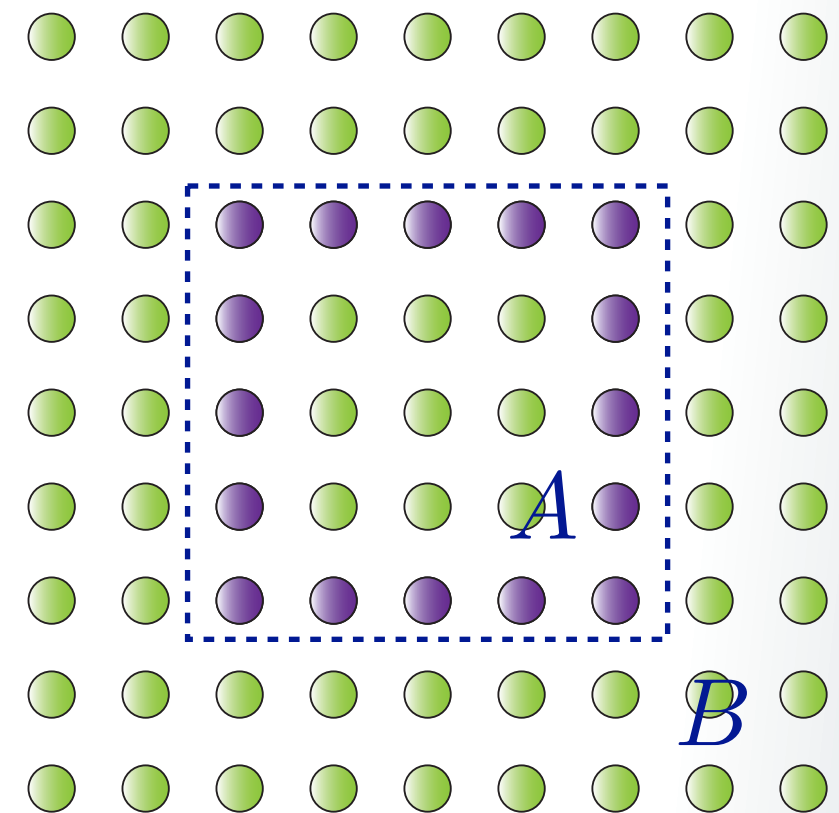


TNS come in different forms...

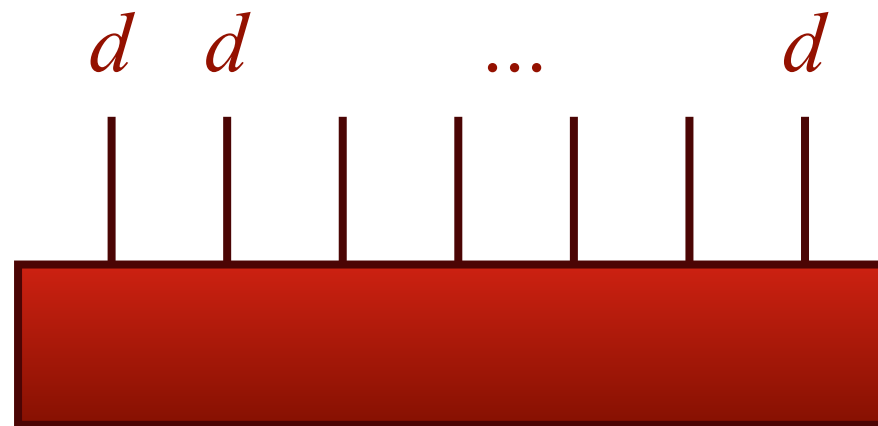
Ansätze satisfying
the area law
by construction

MPS & PEPS

Area law

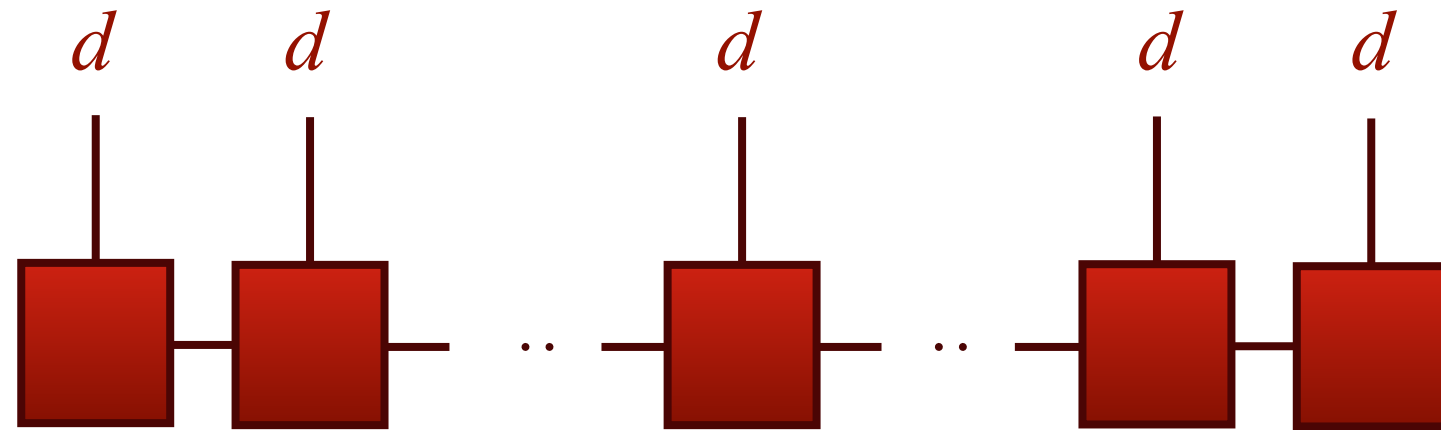


MPS



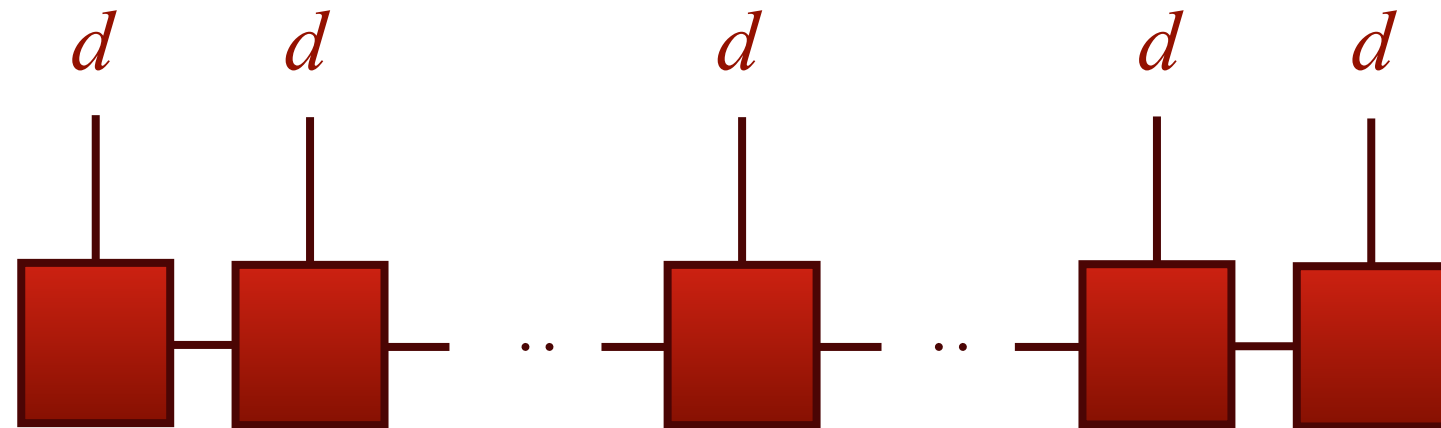
$$|\Phi\rangle = \sum_{s_1, \dots, s_N=1}^d c_{s_1, \dots, s_N} |s_1, \dots, s_N\rangle$$

MPS

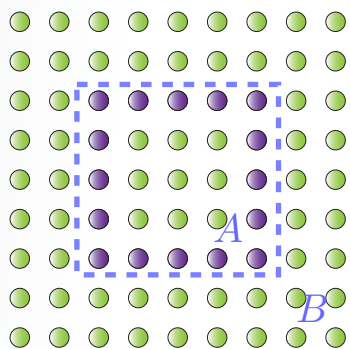


$$|\Phi_D\rangle = \sum_{s_1, \dots, s_N=1}^d \text{tr} (A^{s_1} [1] \dots A^{s_1} [N]) |s_1, \dots, s_N\rangle$$

MPS



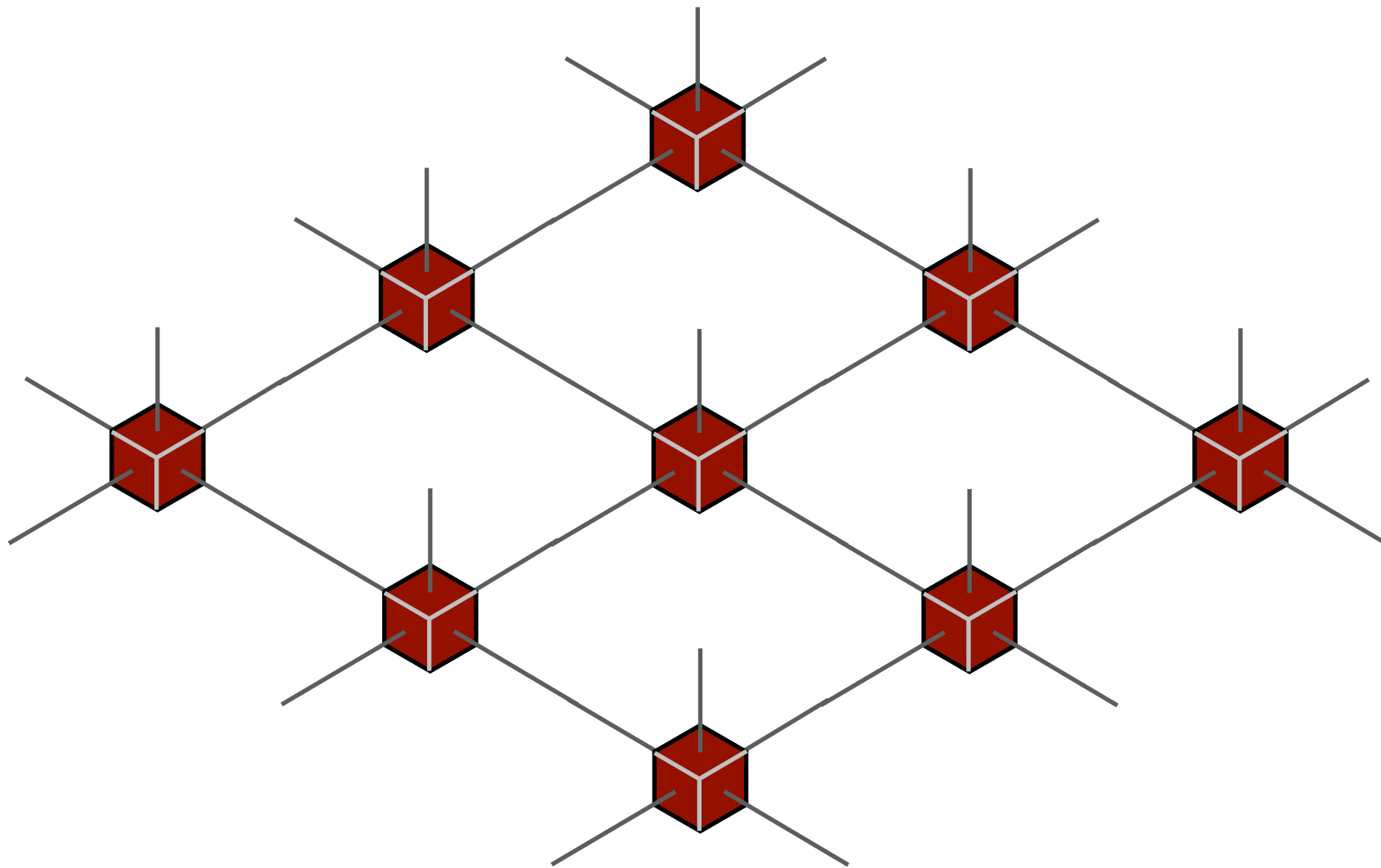
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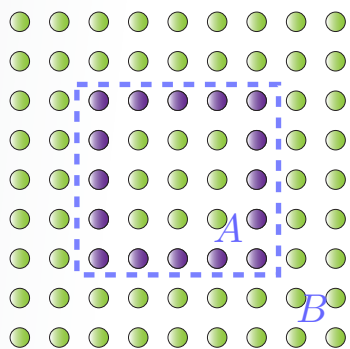
area law by construction

$$S(L/2) \leq \log D$$

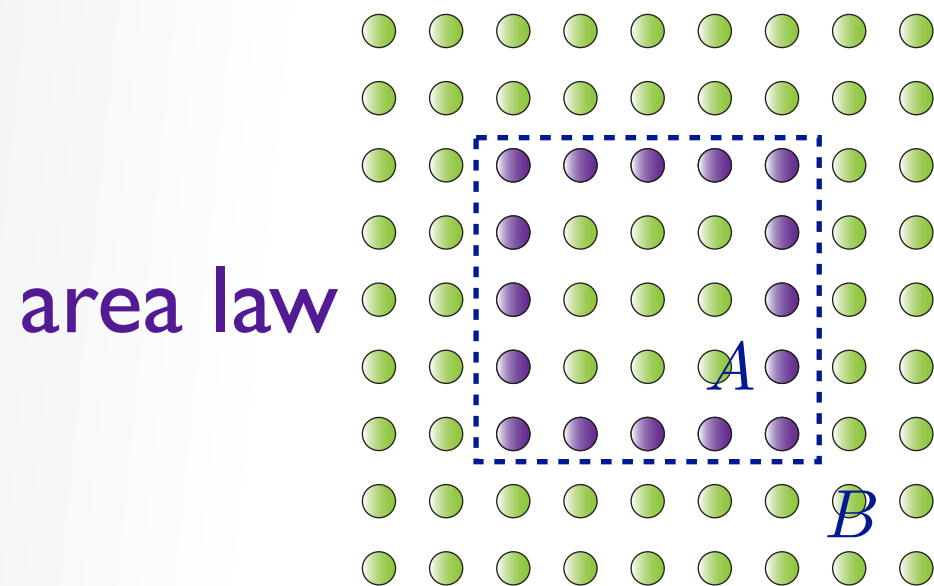
peps



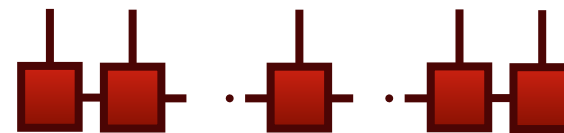
area law by construction



TNS = entanglement based ansatz

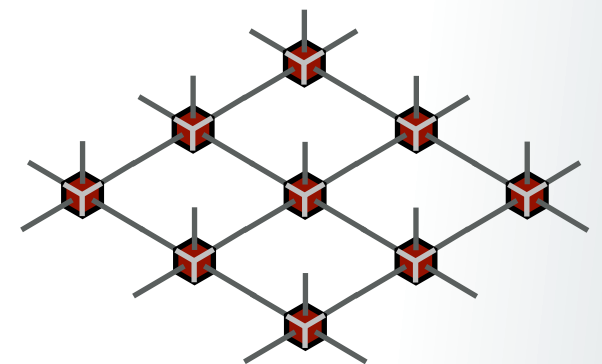


MPS



Schollwöck Ann.Phys.2011

PEPS

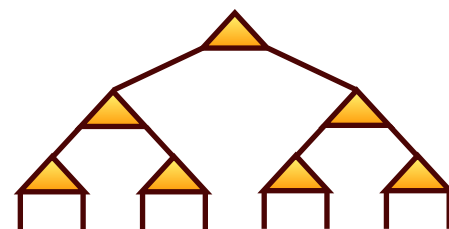


Verstraete et al. Adv. Phys. 2008

other TNS

critical ID
correlations

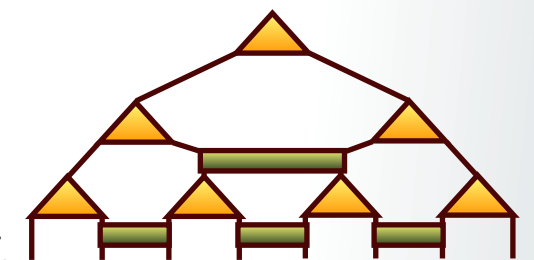
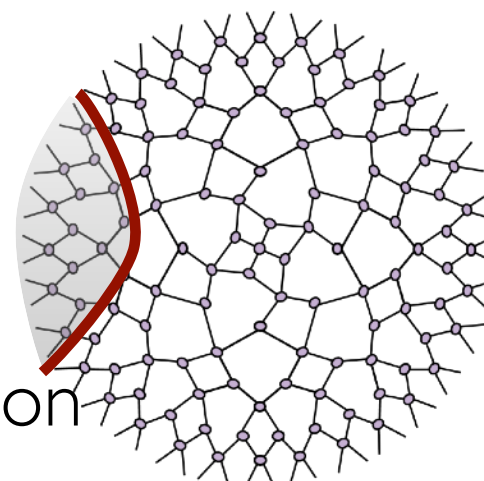
TTN



Shi et al PRA 2006

suggested connection
to AdS/CFT

Vidal PRL 2007 MERA



Swingle PRD 2012
Molina JHEP 2013
Nozaki et al JHEP 2012
Bao et al PRD 2015

some interesting properties of MPS & PEPS

complete families

good approximation of ground/thermal states of
local Hamiltonians

Verstraete, Cirac, PRB 2006

Hastings J. Stat. Phys 2007

Hastings PRB 2006

Molnar et al PRB 2015

can be defined directly in the thermodynamic limit

ground states of local parent Hamiltonians

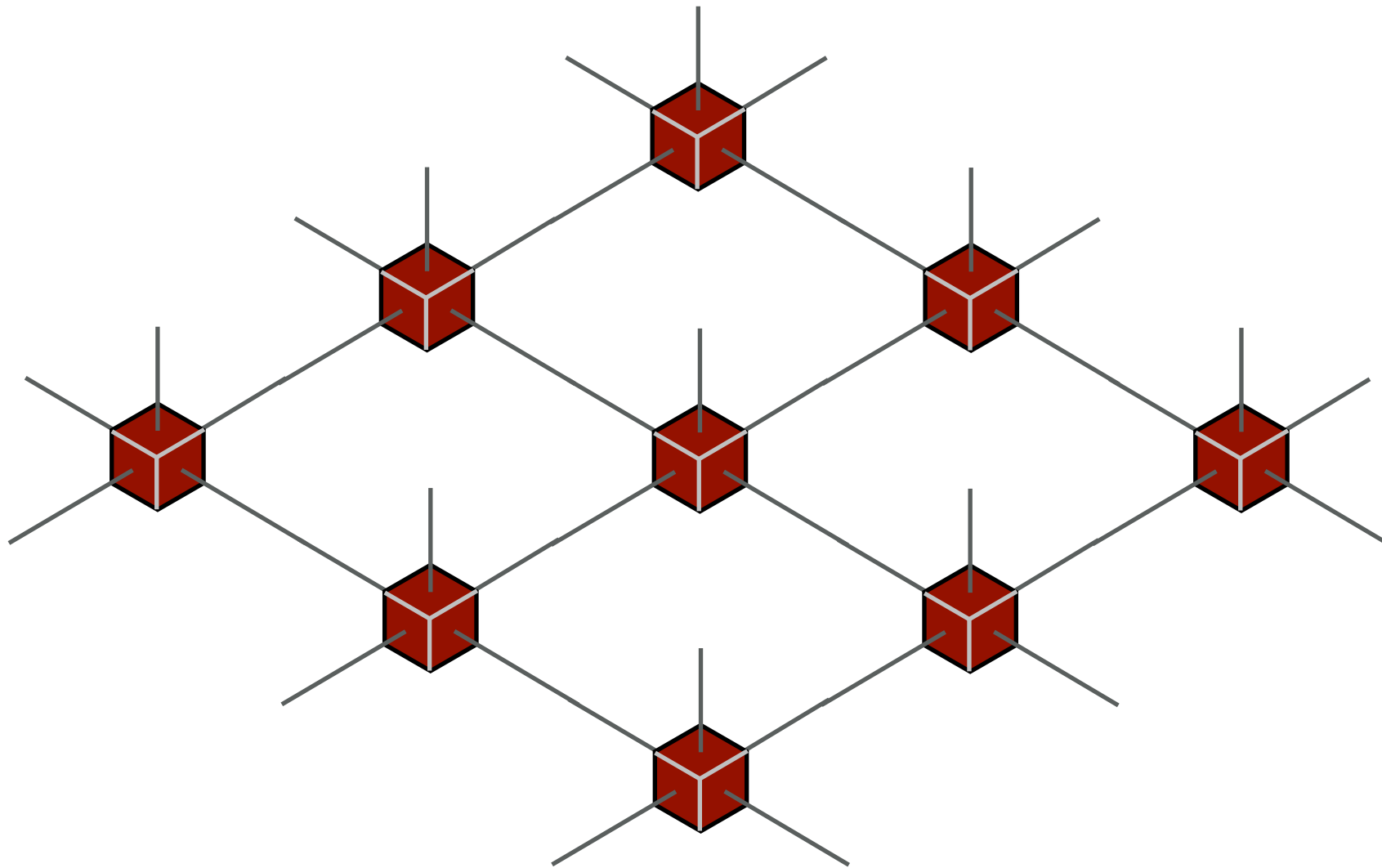
Pérez-García et al., QIC 2007

Schuch et al., Ann. Phys. 2010

Cirac et al RMP 2021

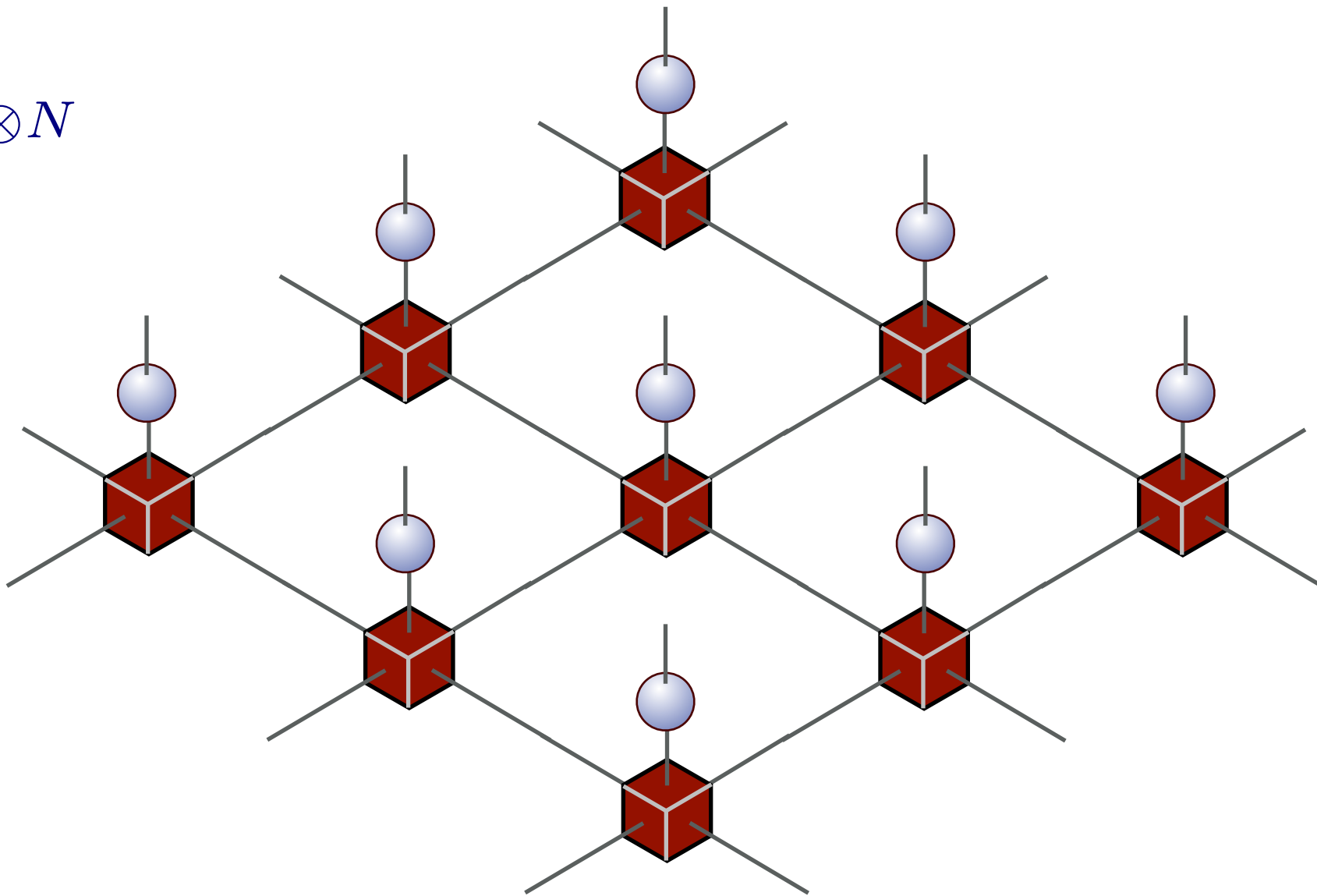
TNS play well with symmetries

state invariant under global symmetry

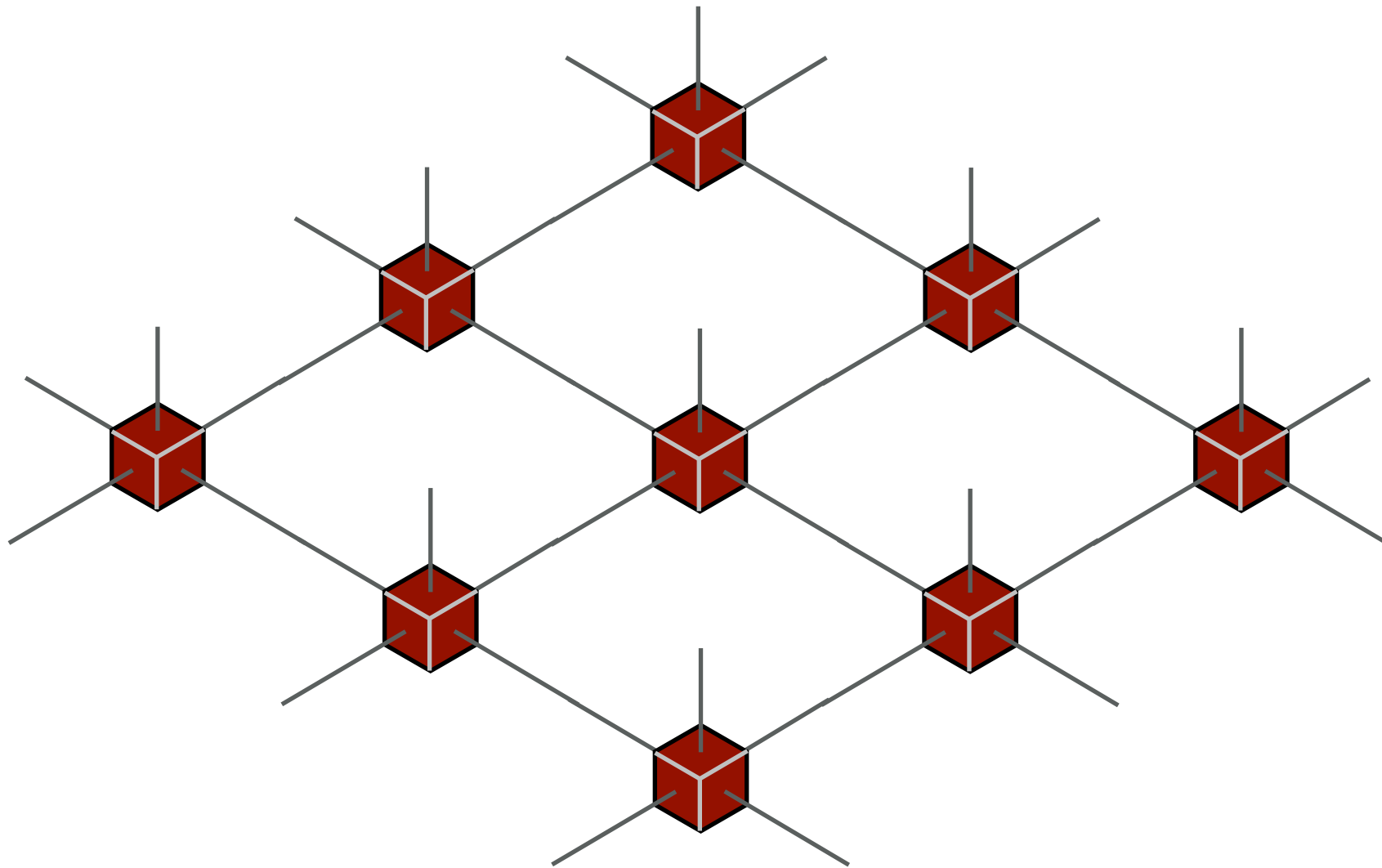


state invariant under global symmetry

$$U_g^{\otimes N}$$

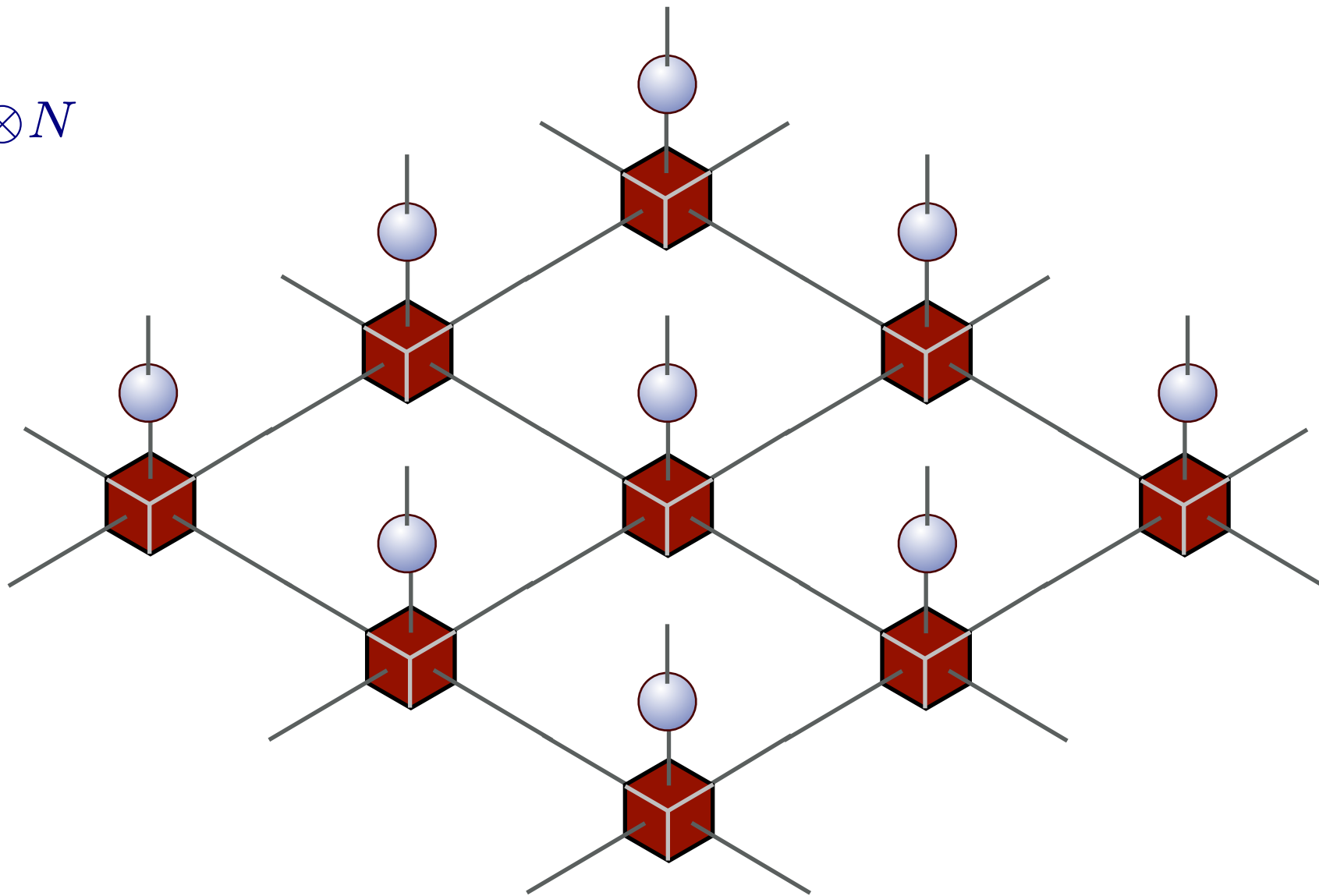


state invariant under global symmetry

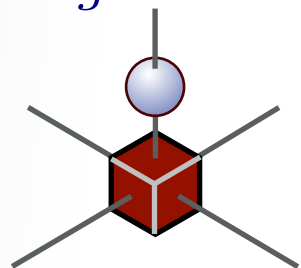


state invariant under global symmetry

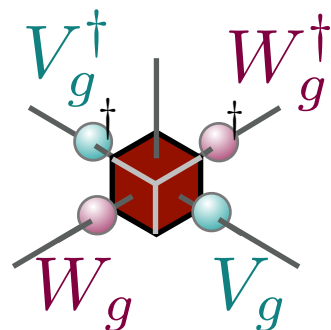
$U_g^{\otimes N}$



U_g



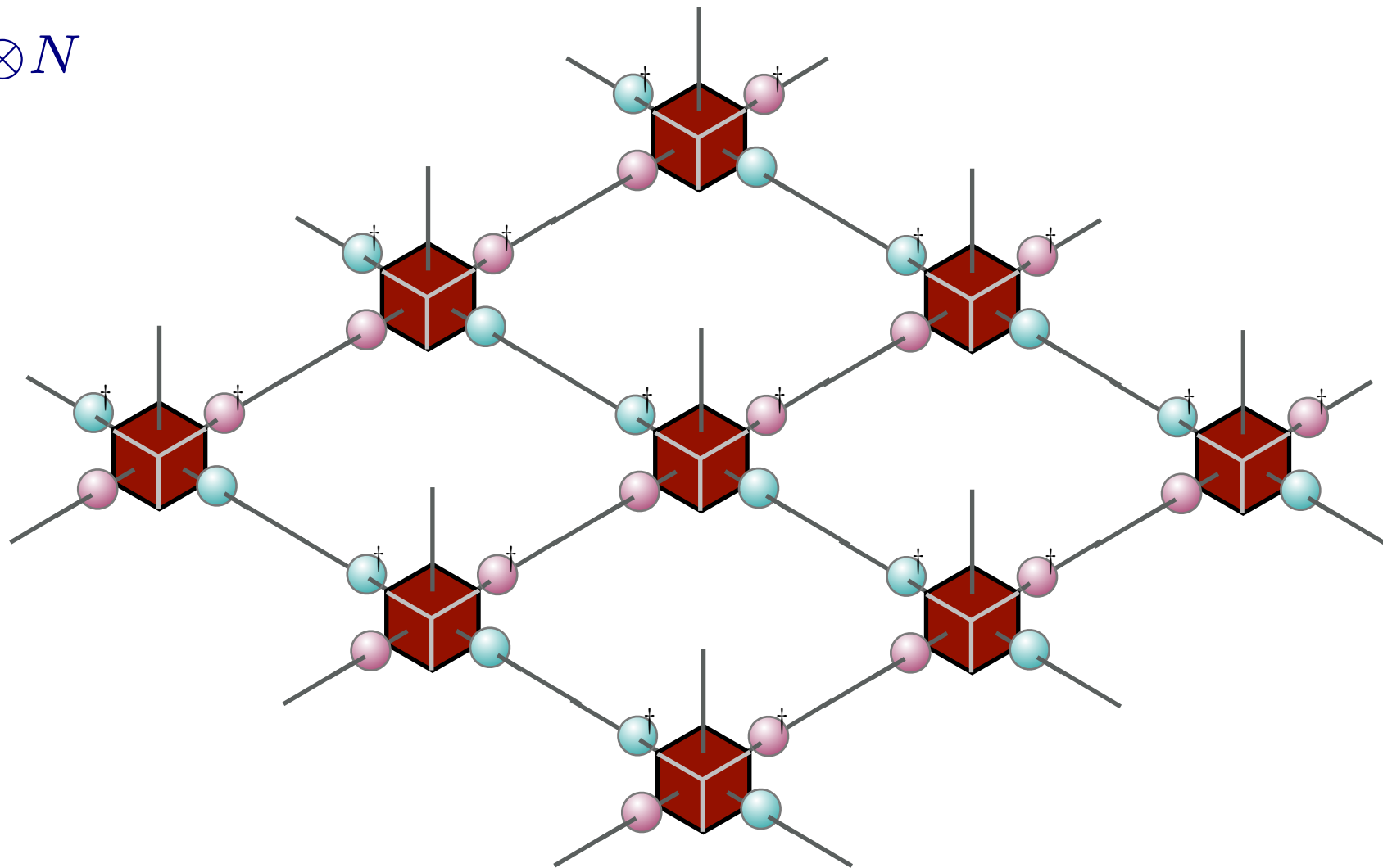
=



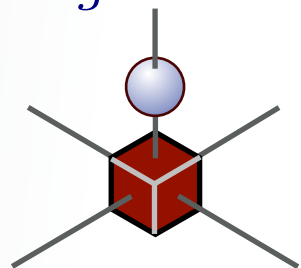
Pérez-García et al., PRL 2008
 Sanz et al., PRA 2009
 Schuch et al., Ann. Phys. 2010
 Singh et al., NJP 2007, PRA 2010

state invariant under global symmetry

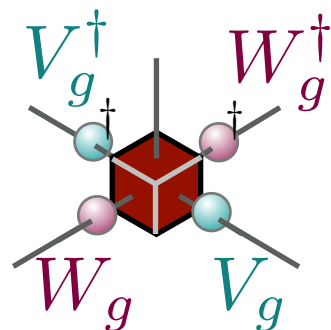
$$U_g^{\otimes N}$$



$$U_g$$



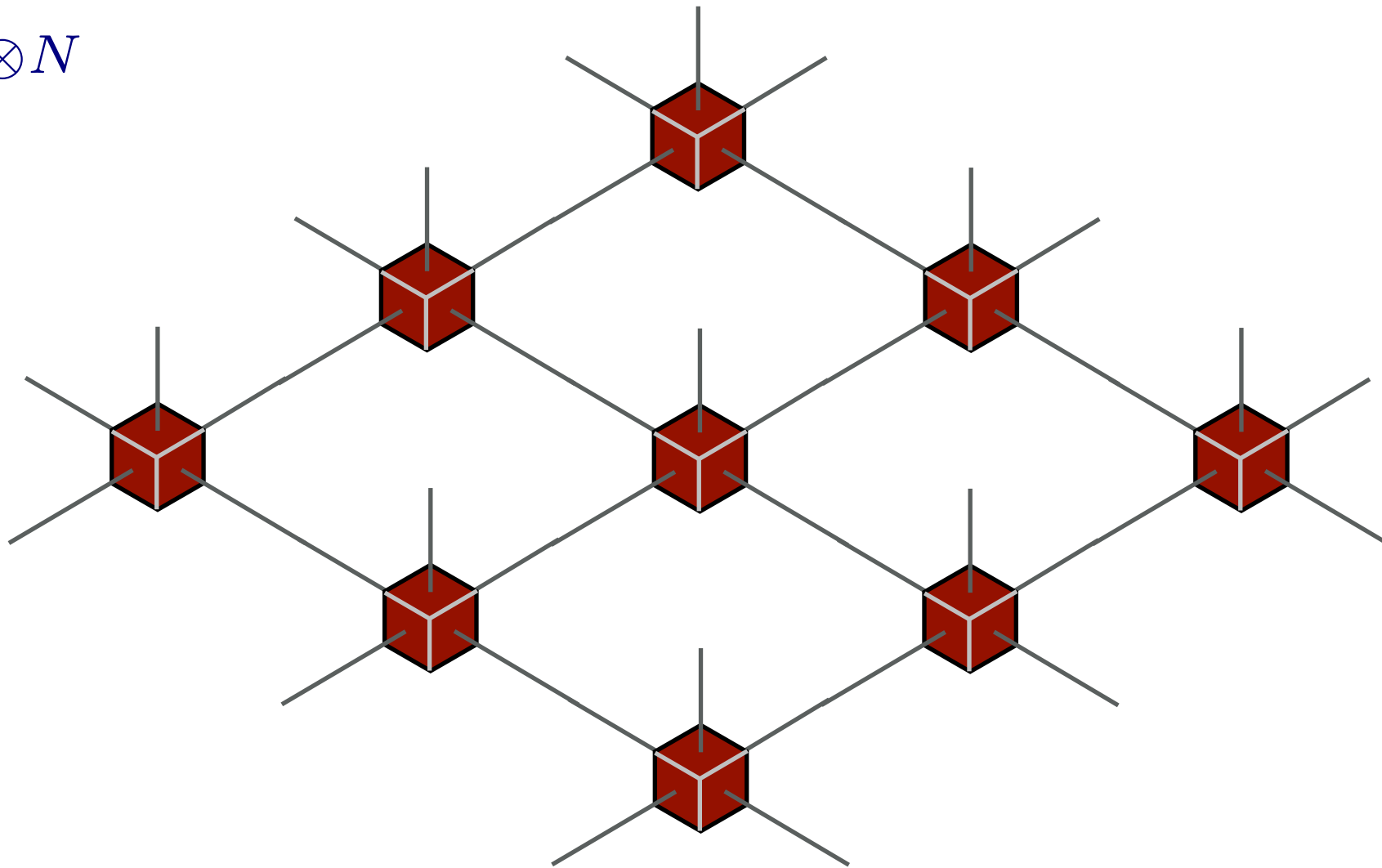
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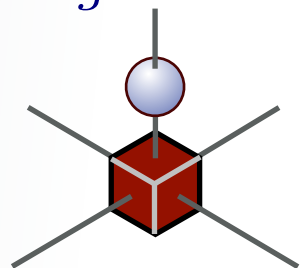
Pérez-García et al., PRL 2008
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 Singh et al., NJP 2007, PRA 2010

state invariant under global symmetry

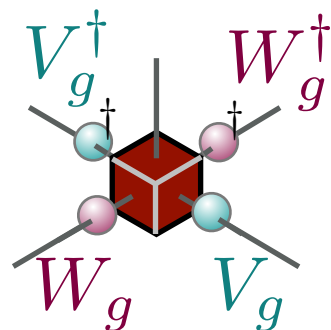
$$U_g^{\otimes N}$$



$$U_g$$



=

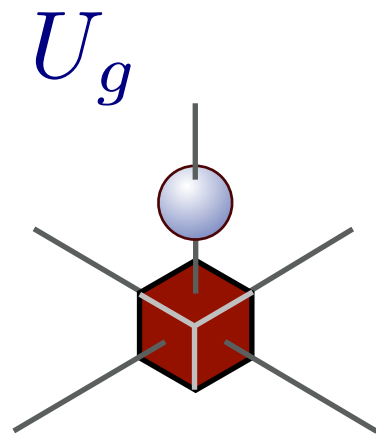


Pérez-García et al., PRL 2008
 Sanz et al., PRA 2009
 Schuch et al., Ann. Phys. 2010
 Singh et al., NJP 2007, PRA 2010

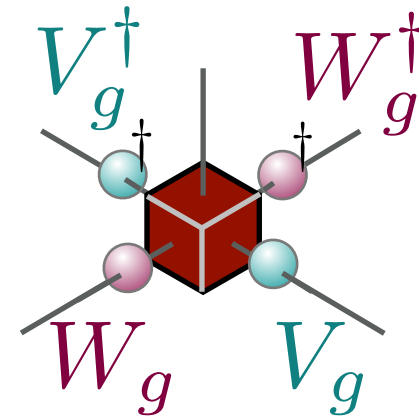
state invariant under global symmetry
constructed with invariant tensors

for MPS & PEPS

state (globally) invariant $\Leftrightarrow$



=



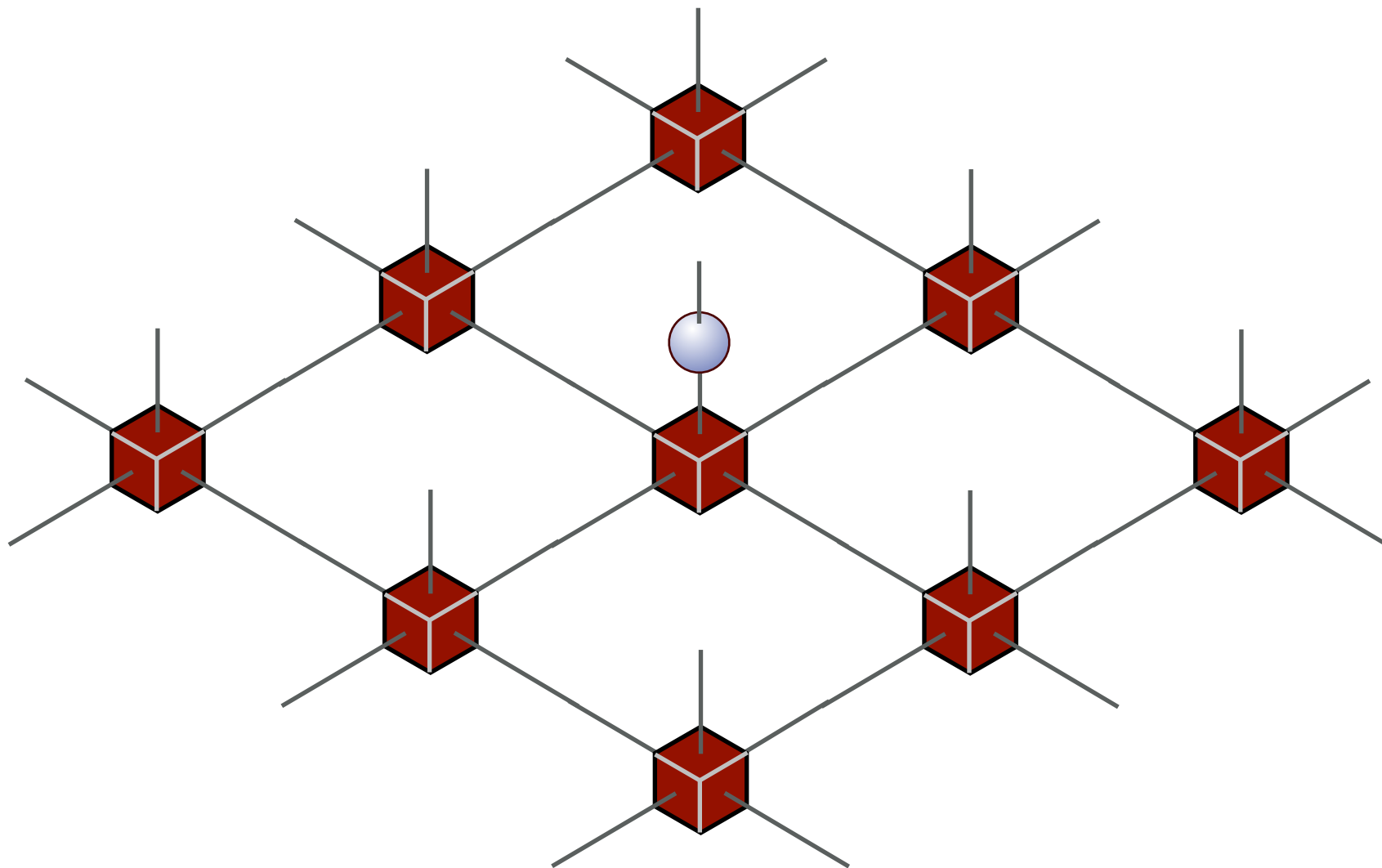
structure
of
tensors

also gauge symmetries

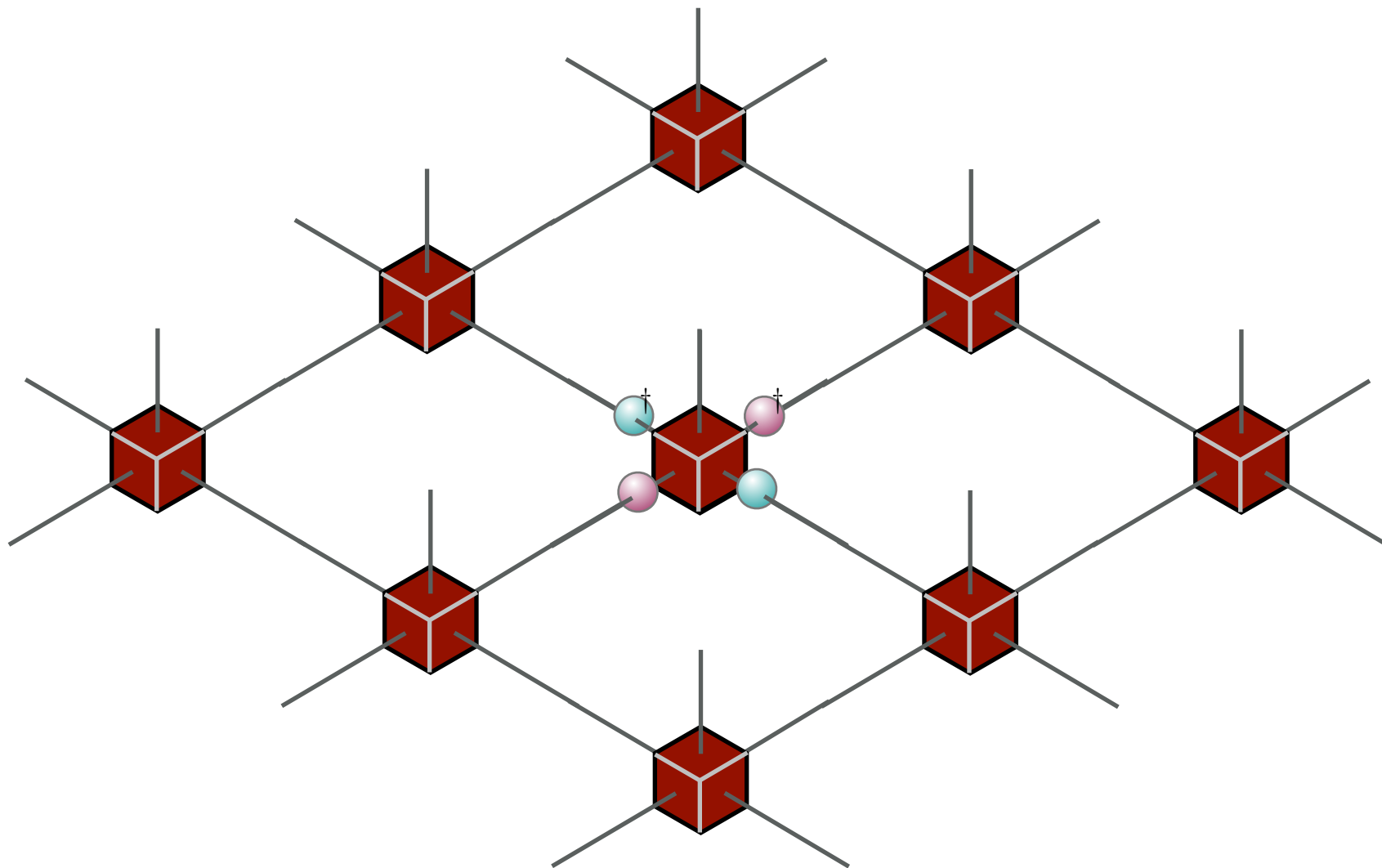
Tagliacozzo et al PRX 2014
Haegeman et al PRX 2014
Zohar et al Ann Phys 2015

Pérez-García et al., PRL 2008
Sanz et al., PRA 2009
Schuch et al., Ann. Phys. 2010
Singh et al., NJP 2007, PRA 2010

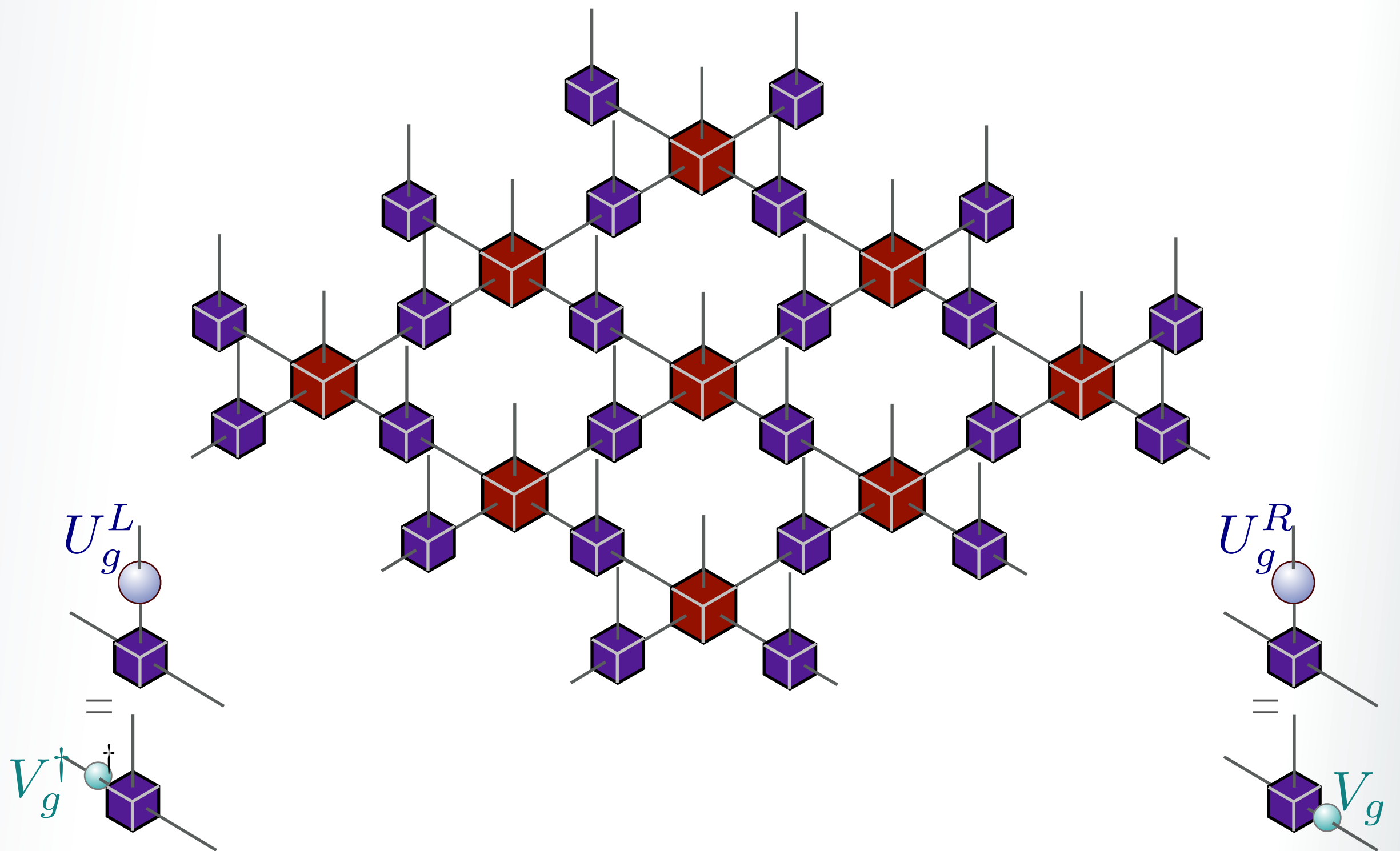
gauging PEPS



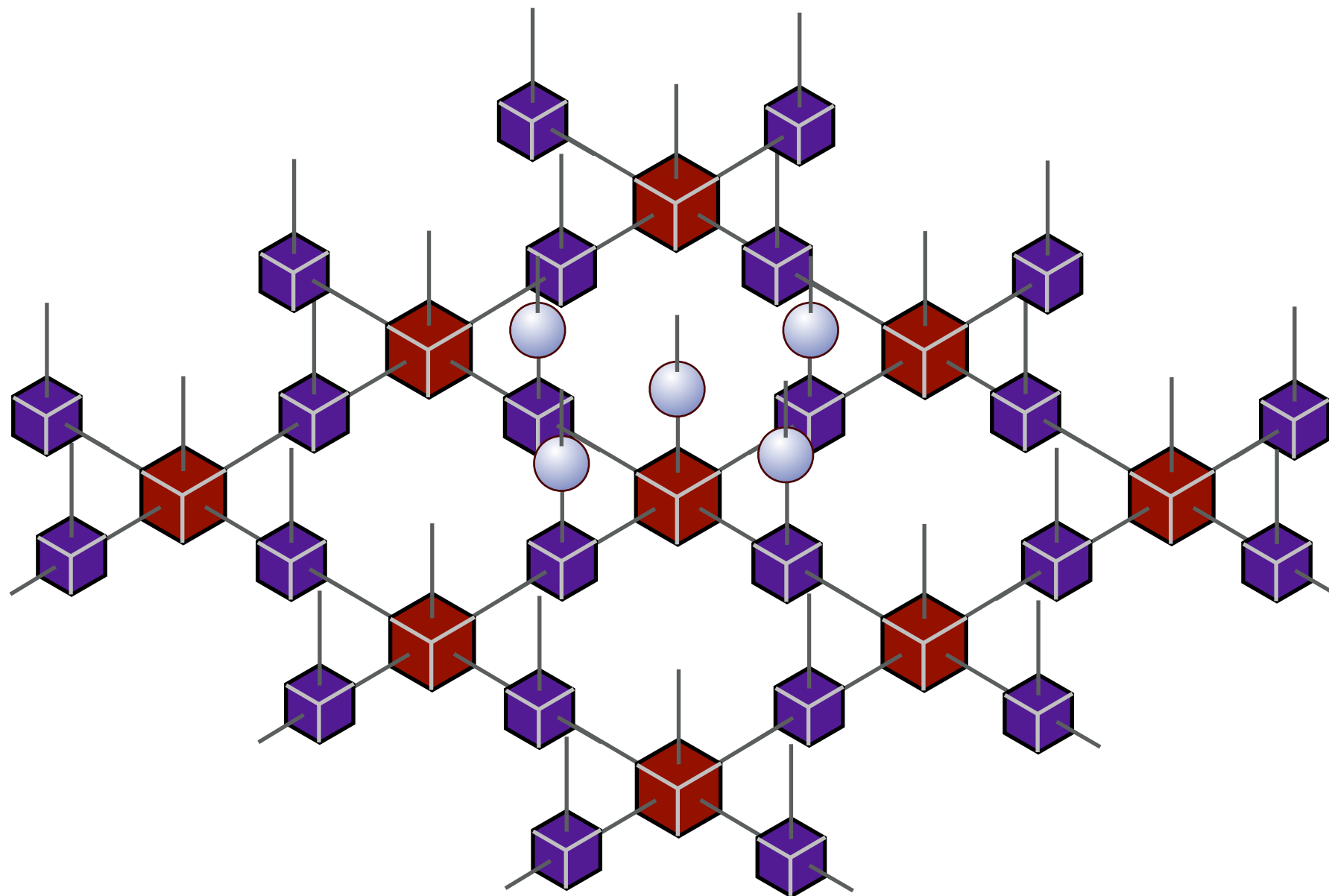
gauging PEPS



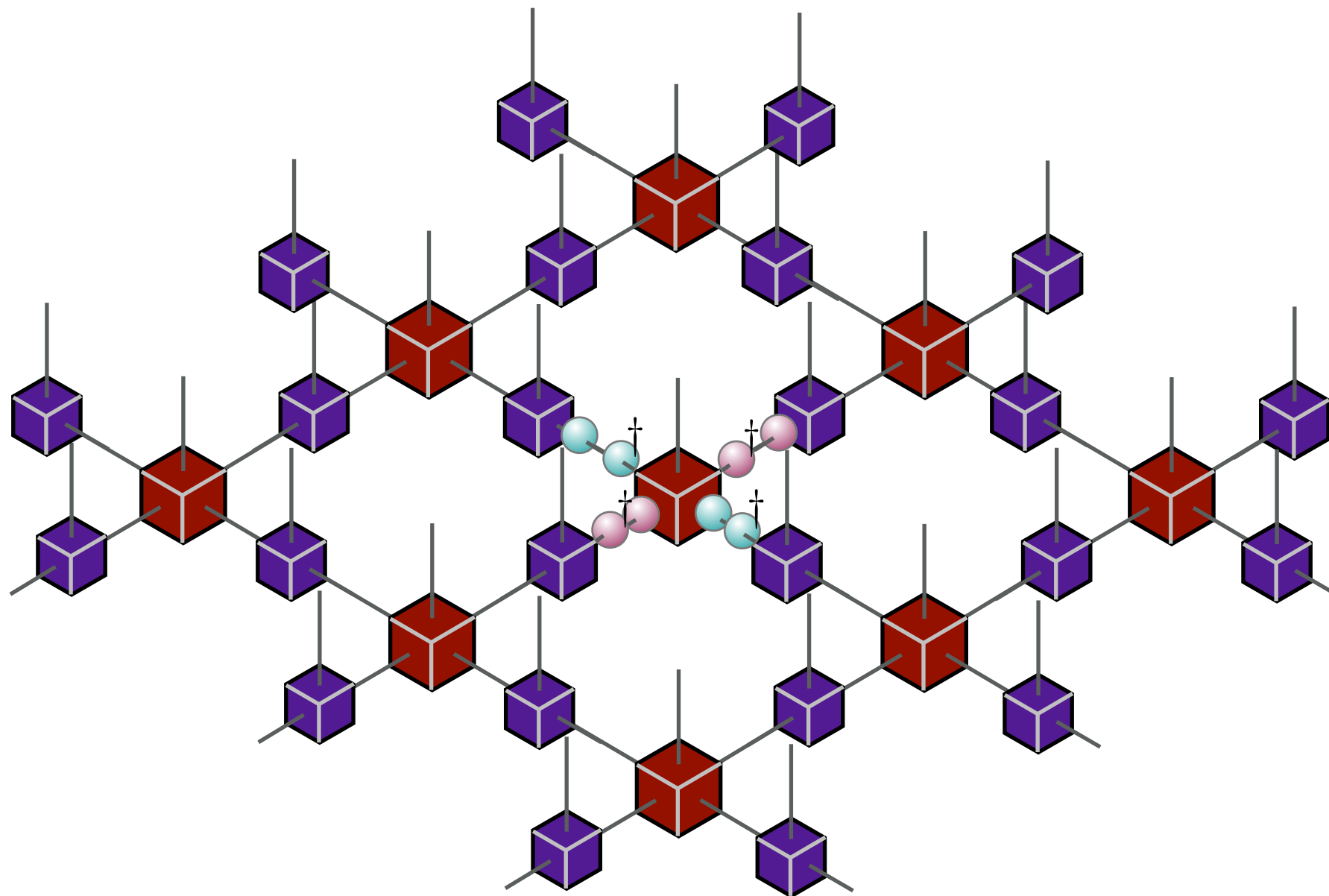
gauging PEPS



gauging PEPS



gauging PEPS

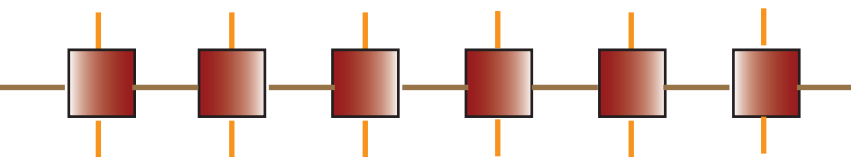


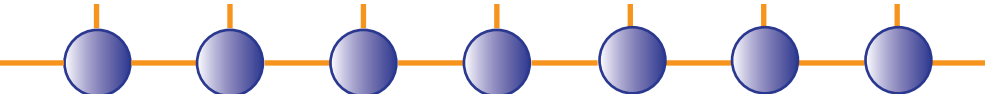
numerics with TNS

basic algorithms

variational minimization of energy

local
Hamiltonian

$$H =$$


$$|E_0\rangle \simeq$$


ground state
excitations

apply local operators $\rightarrow$ simulate time evolution

imaginary time $\rightarrow$ ground state
thermal state



also: evolve using TDVP

White, PRL 1992; Schollwöck, Ann. Phys. 2011;
Vidal, PRL 2003, 2004; Verstraete et al., PRL 2004;
Haegeman et al., PRL 2011

basic algorithms

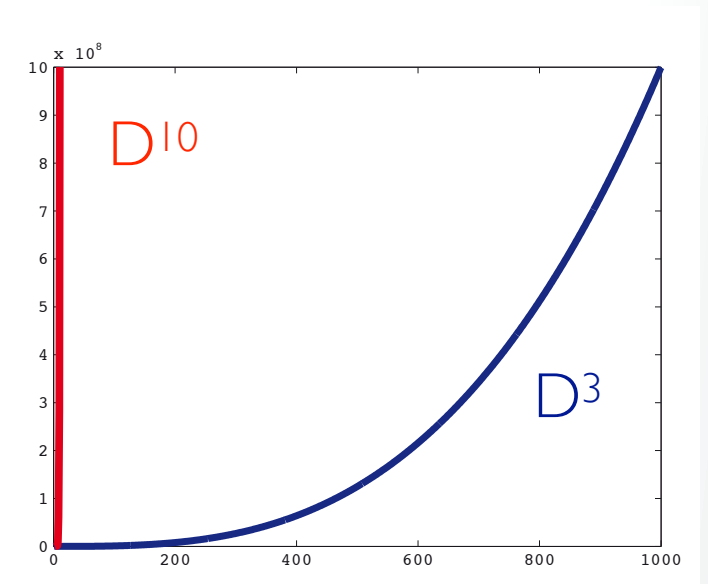
basically work in 1D or 2D

different computational costs

$$1D \rightarrow O(D^3)$$

$$2D \rightarrow O(D^{10})$$

no exact contractions

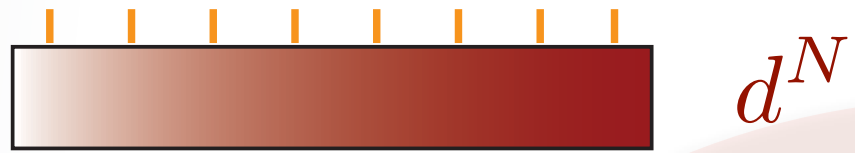


regarding real time

entanglement may grow fast in non-equilibrium scenarios! up to exponentially growing D

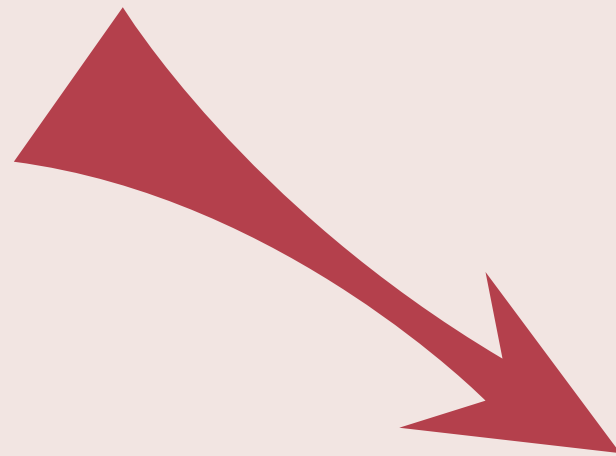
works for close to equilibrium or moderate times

arbitrary many-
body state



$$|\Psi\rangle = \sum_{i_j} c_{i_1 \dots i_N} |i_1 \dots i_N\rangle$$

exponential



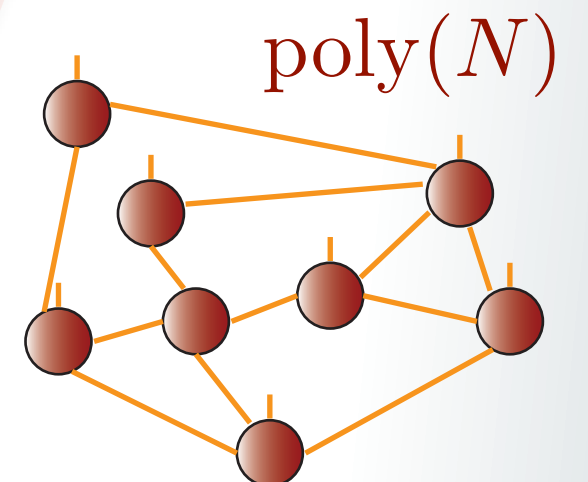
polynomial

TNS: restricted
family

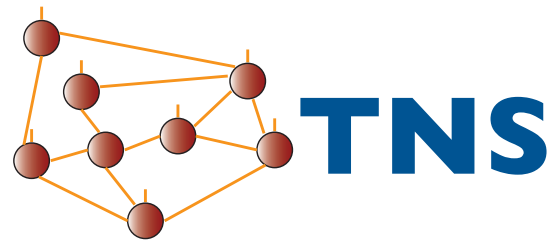
good ansatz for ground states
and thermal equilibrium: area
law

entanglement hierarchy

efficient numerics



TNS for QFT



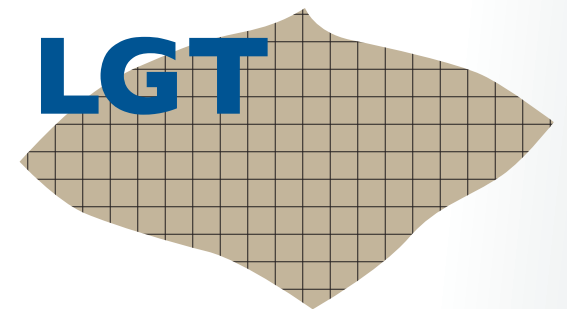
Non-perturbative for Hamiltonian systems

Extremely practical (and successful) for 1D systems (MPS)

Powerful, but more costly for higher dimensions

ground states
low-lying excitations
thermal states
time evolution

apply to
LGT



using TNS for LGT

formal approach

gauging the symmetry

explicitly invariant states

general prescriptions, $U(1)$, $SU(2)$

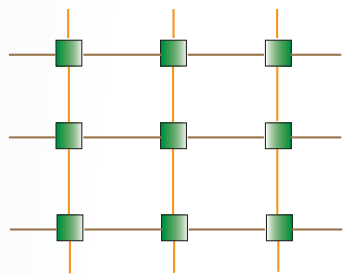
Tagliacozzo et al PRX 2014

Haegeman et al PRX 2014

Zohar et al Ann Phys 2015

numerical simulations

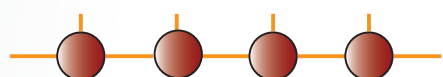
no sign problem



TN describe partition functions (observables)

TRG approaches to classical and quantum models

Liu et al PRD 2013; Shimizu, Kuramashi, PRD 2014; Kawauchi, Takeda 2015;
review Meurice et al. 2010.06539



TN describe states

general strategy

Hamiltonian formulation

acting on a Hilbert space

→ choose proper basis

Finite dimensional degrees of freedom

fermions

→ ✓ no sign problem

gauge bosons require attention

→ truncating, integrating out (also QLinks)

★ Common ingredients for quantum simulation

Zohar et al. PRL 2010, 2012 ,
Tagliacozzo et al., Nat. Comm. 2013
Banerjee et al., PRL 2012

Rico et al. PRL 2014
Pichler et al, PRX 2016
Zohar, Burrello, PRD 2015

there is long way to go until LQCD

journey begins with $I + ID$ steps

early works with DMRG/TNS

Byrnes PRD2002; Sugihara NPB2004
Tagliacozzo PRB2011; Sugihara JHEP2005
Meurice PRB2013

Schwinger model $U(1)$ in 1D precise equilibrium simulations, feasibility of QSim

MCB et al JHEP11(2013)158;
Rico et al PRL 2014; Buyens et al. PRL 2014;
Kühn et al., PRA 90, 042305 (2014);
MCB et al PRD 2015, Buyens et al. PRD 2016;
Pichler et al. PRX 2016;
review Dalmonte, Montangero, Cont. Phys. 2016
MCB, Cichy, Cirac, Jansen, Kühn, arXiv:1810.12838

MCB, K. Cichy 1910.00257
QTFLAG Collab.1911.00003

3+1 dimensions

Magnifico et al. Nat. Com.12, 3600 (2021)

2+1 dimensions

Felser et al. PRX10, 041040 (2020)
Robaina et al. PRL126, 050401 (2021)
Emonts et al. PRD102, 074501 (2020)

Non-Abelian in 1D string breaking dynamics

S. Kühn et al., JHEP 07 (2015) 130;
Silvi et al., Quantum 2017
S. Kühn et al. PRX 2017

$SU(3)$ QLM

Silvi et al, PRD 2019

finite density

S. Kuehn et al, PRL118 (2017) 071601

spectrum

finite density

entropy

I + I D RESULTS

thermal equilibrium

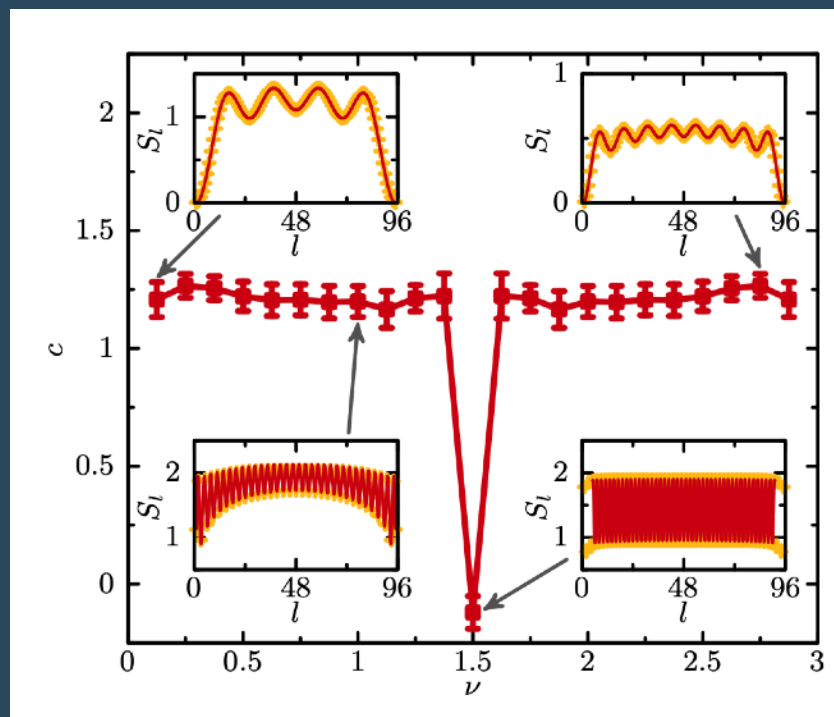
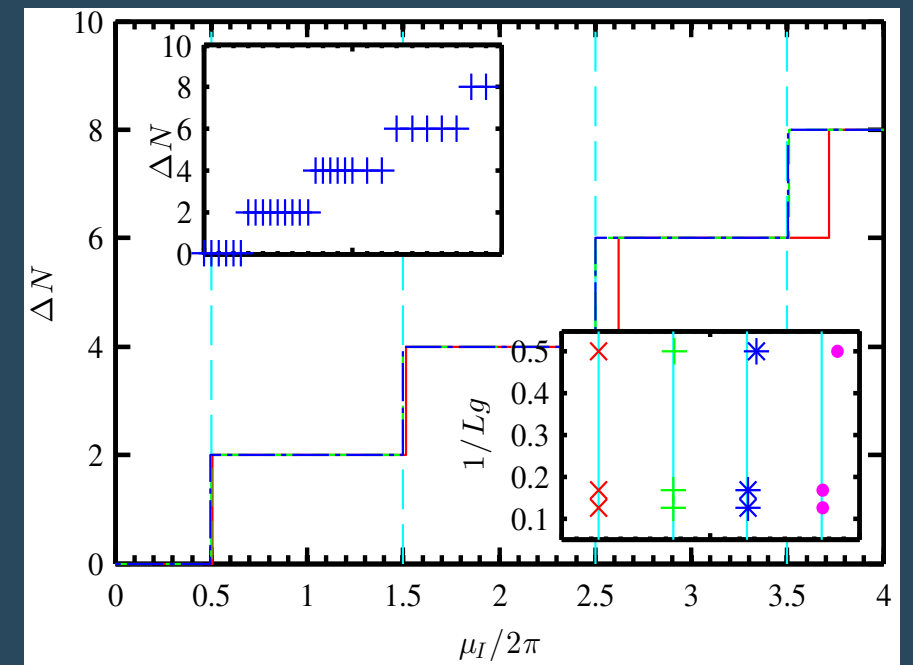
time evolution



finite density

Schwinger model with several
flavours and chemical potentials
(continuum)

S. Kühn et al, PRL 118 (2017) 071601



phase diagram of SU(2) and
SU(3) QLM at finite density

Silvi et al, Quantum 1, 9 (2017)

PRD 100, 074512 (2019)

some results in 2+1, 3+1
U(1)

spectrum

finite density

entropy

I + ID RESULTS

thermal equilibrium

time evolution

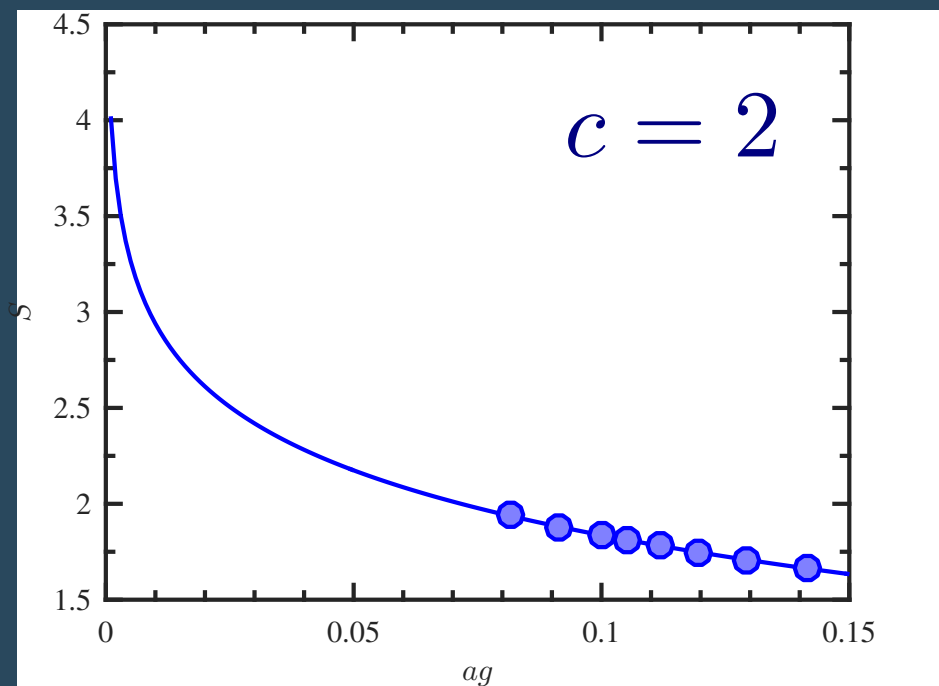


gauge constraints not purely local $\Rightarrow$ not all entropy physical

Casini et al 2014; Gosh et al JHEP 2015

Soni, Trivedi JHEP 2016; van Acoleyen et al PRL 2016

SU(2)



Kühn PRX7, 041046 (2017)

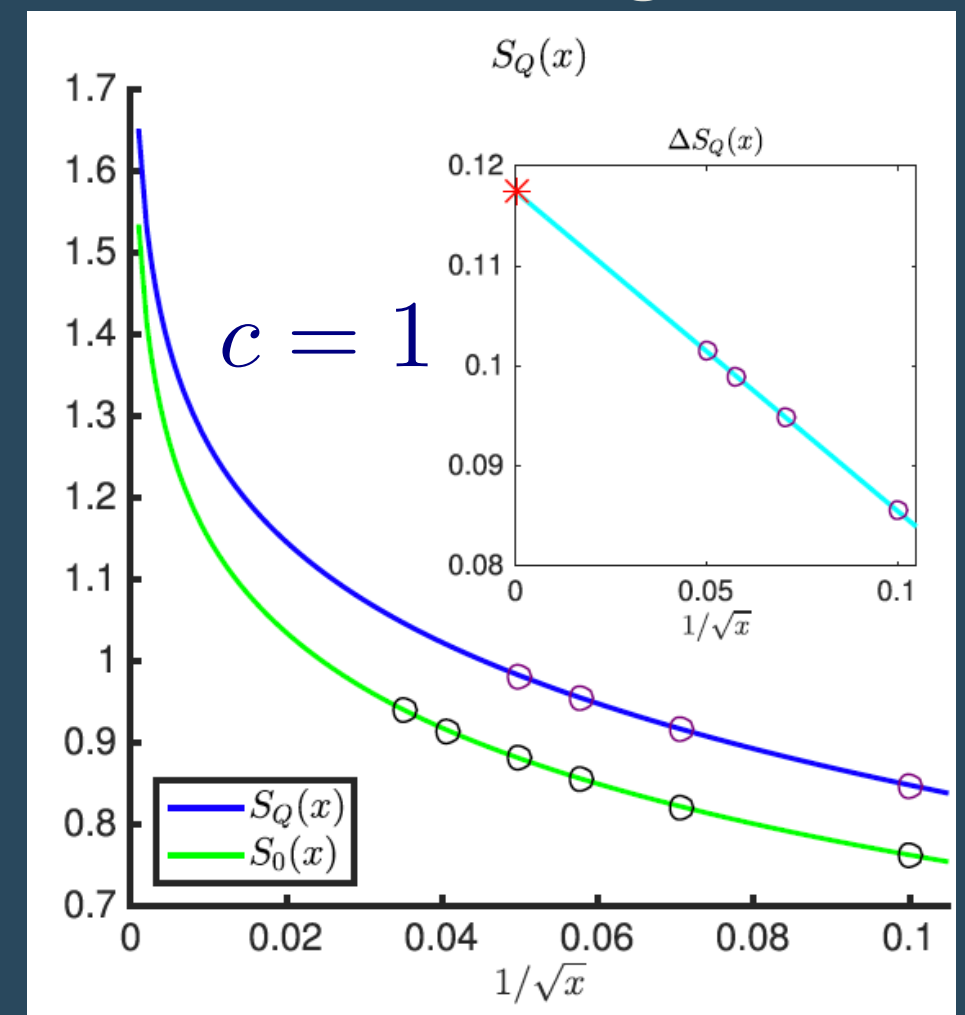
divergence in the continuum limit

$$S \propto \frac{c}{6} \log_2 \frac{\xi}{a}$$

Calabrese, Cardy JStatMech 2004

entropy

Schwinger



Buyens PRX6, 041040 (2016)

spectrum

finite density

entropy

I + ID RESULTS

thermal equilibrium

time evolution



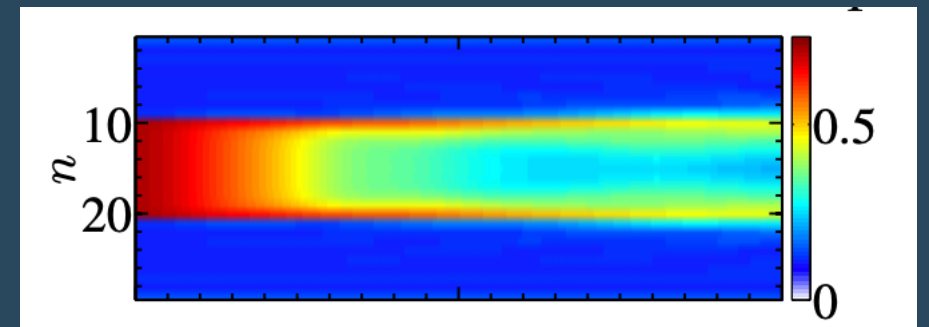
string breaking in lattice models: SU(2), U(1)

Buyens et al. PRL 113, 091601 (2014)

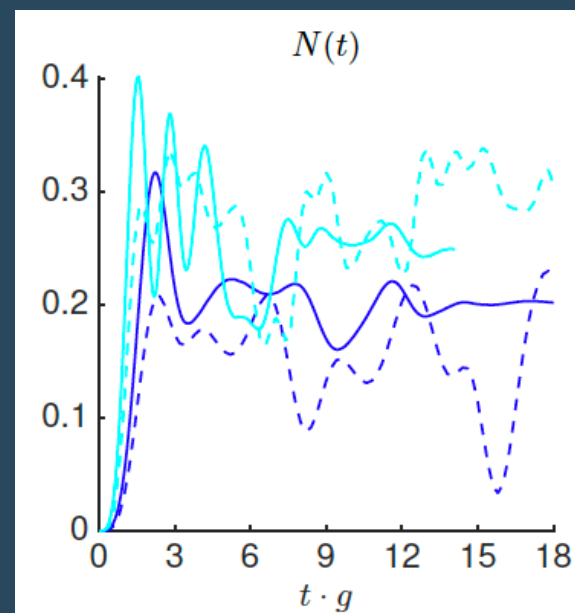
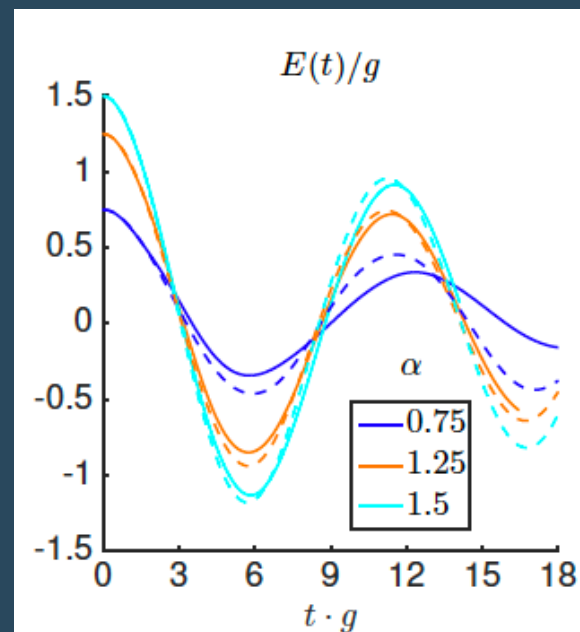
Kühn et al, JHEP 07 (2015) 130

Pichler et al., PRX 6, 011023 (2016)

PRX 6, 041040 (2016)

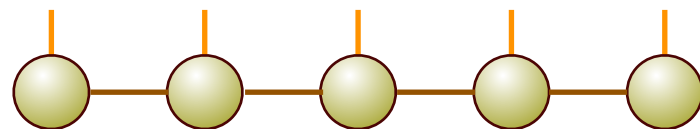


near the continuum limit: Schwinger pair production



time evolution

real time evolution with MPS



TEBD, t-DMRG

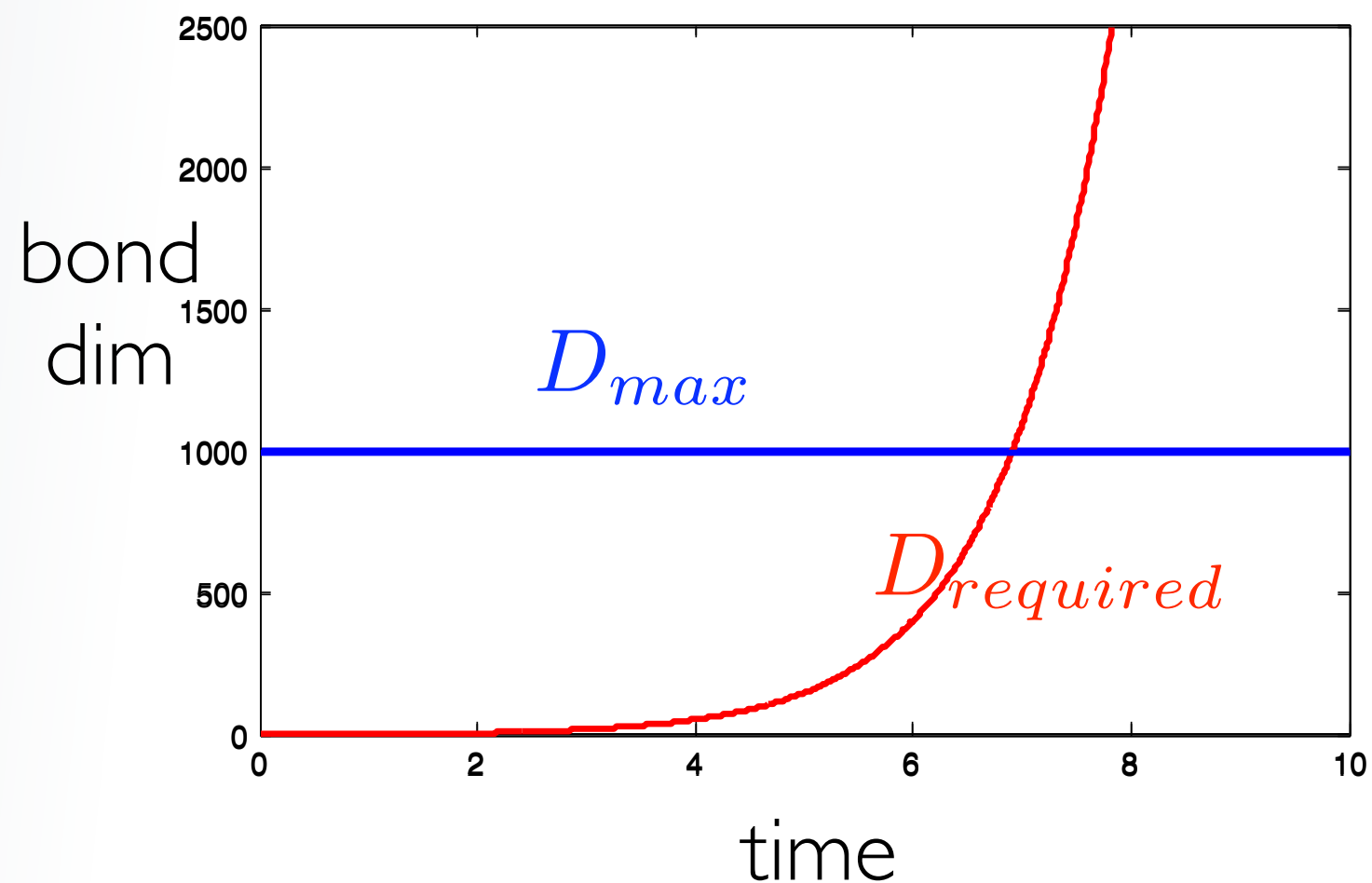
Vidal, PRL 2003, 2004

Verstraete, García-Ripoll, Cirac, PRL 2004

time evolved state
approximated by MPS

but entanglement grows

Osborne, PRL 2006
Schuch et al., NJP 2008



required bond for
fixed precision

$$D \sim e^{\alpha t}$$

yet many physical situations (in closed and open quantum systems) can be successfully studied!

short times, adiabatic, low energy can work well

García-Ripoll, NJP 2006

Wall, Carr NJP 2012

Paeckel et al arXiv:1901.05824

PDF in Schwinger model

with K. Cichy, C.J. David Lin, M. Schneider;
LATTICE'24 arXiv:2409.16996
2504.XXXXXX

fermionic PDF

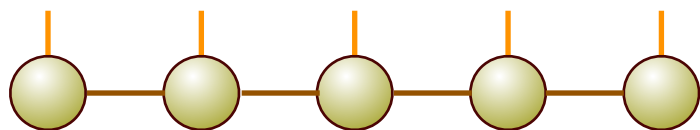
$$f_{\Psi}(\xi) = \int \frac{dz^-}{4\pi} e^{-i\xi P^+ z^-} \langle P | \bar{\Psi}(z^-) W_{(z^-,0)} \gamma^+ \Psi(0) | P \rangle$$

meson state

Wilson line along
lightcone direction

lattice version (staggered fermions)

$$f_{\Psi}(\xi) = \frac{NM}{8\pi x} \sum_{\Delta z=0,2,4,\dots,N} e^{-i\xi \frac{M\Delta z}{2x}} [\mathcal{M}_{ee}(\Delta z) - \mathcal{M}_{eo}(\Delta z) - \mathcal{M}_{oe}(\Delta z) + \mathcal{M}_{oo}(\Delta z)]$$



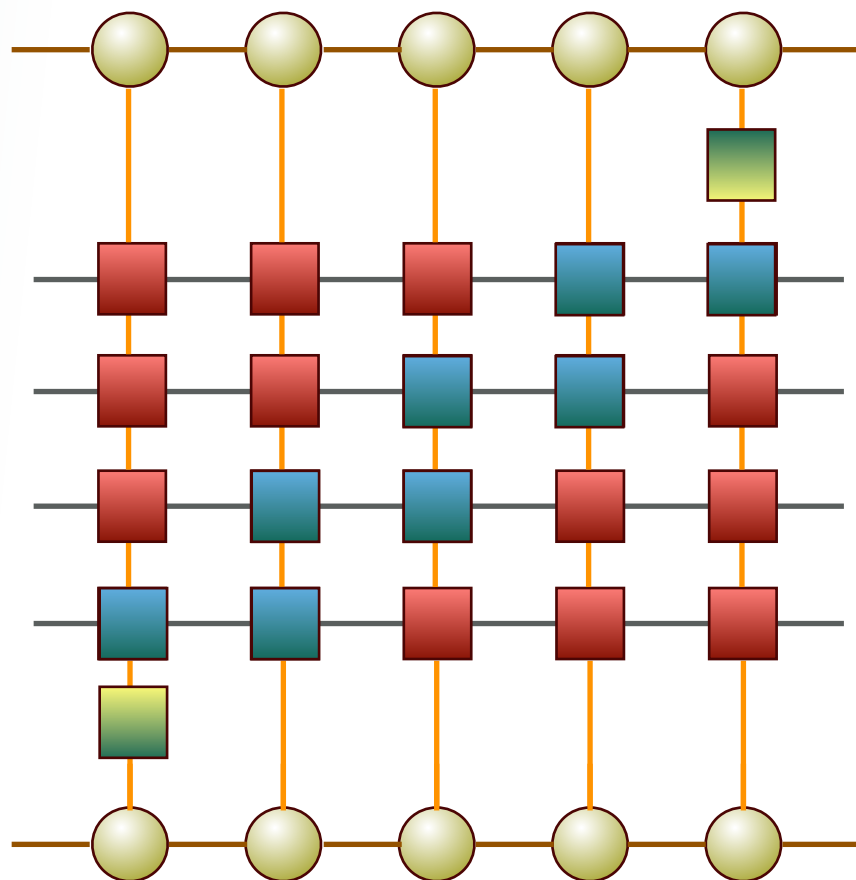
MCB, Cichy, Lin, Schneider, to appear: 2504.XXXXXX
preliminary results presented in LATTICE'24 arXiv:2409.16996

fermionic PDF

$$\langle P | \bar{\Psi}(z^-) W_{(z^-, 0)} \gamma^+ \Psi(0) | P \rangle$$

lattice version (staggered fermions)

$$\langle h | e^{iH \frac{\Delta t}{2} t} \prod_{k < \Delta z} (i\sigma_k^z) \sigma_{\Delta z}^+ e^{-iH_0 \frac{\Delta t}{2} t} \prod_{k' < 0} (-i\sigma_{k'}^z) \sigma_0^- | h \rangle$$



Jordan-Wigner spin mapping

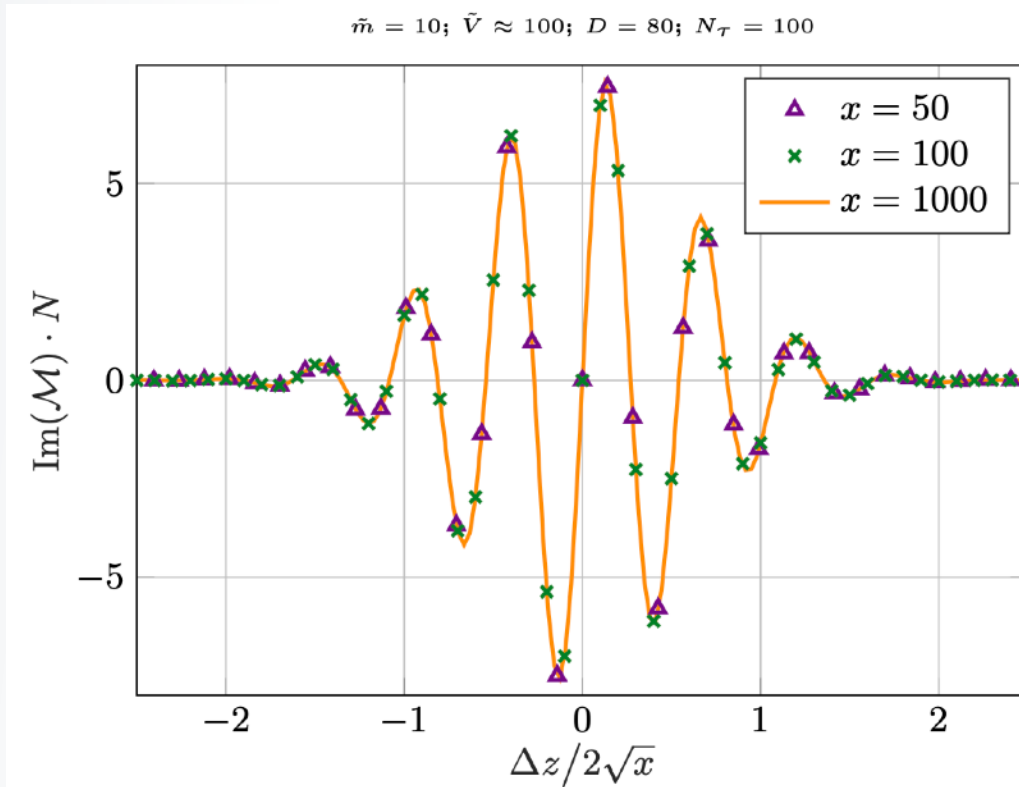
eigenstate

Trotterized evolution

zigzag light cone

MCB, Cichy, Lin, Schneider, to appear: 2504.XXXXXX
preliminary results presented in LATTICE'24 arXiv:2409.16996

fermionic PDF

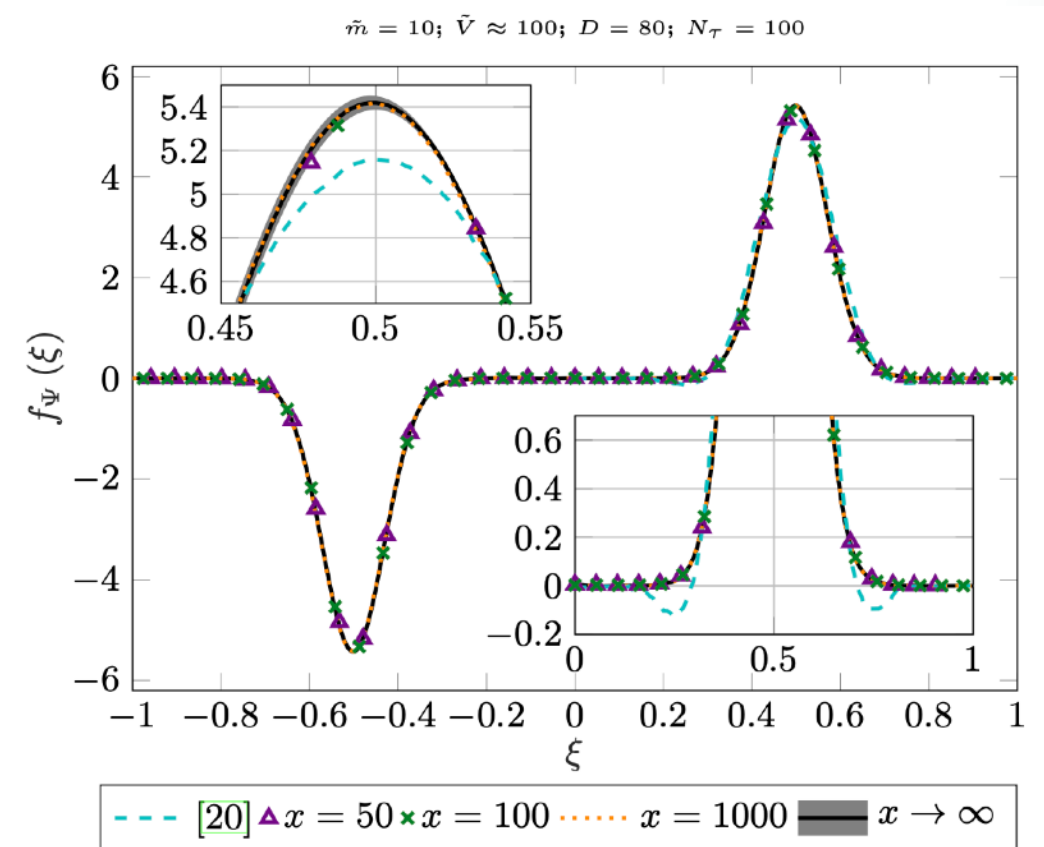


basic matrix element

$$x = \frac{1}{(ga)^2}$$

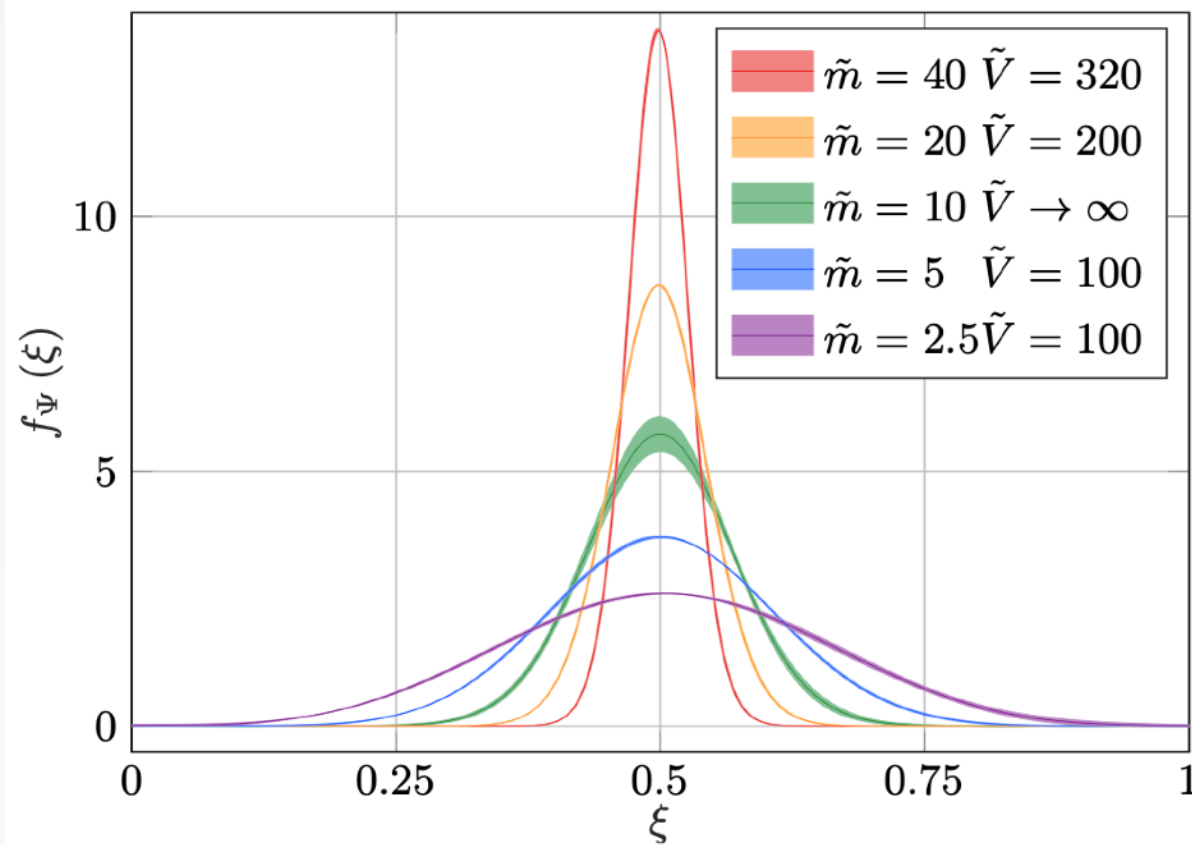
convergence with very moderate size of tensors

PDF (Fourier transform)



MCB, Cichy, Lin, Schneider, to appear: 2504.XXXXXX
preliminary results presented in LATTICE'24 arXiv:2409.16996

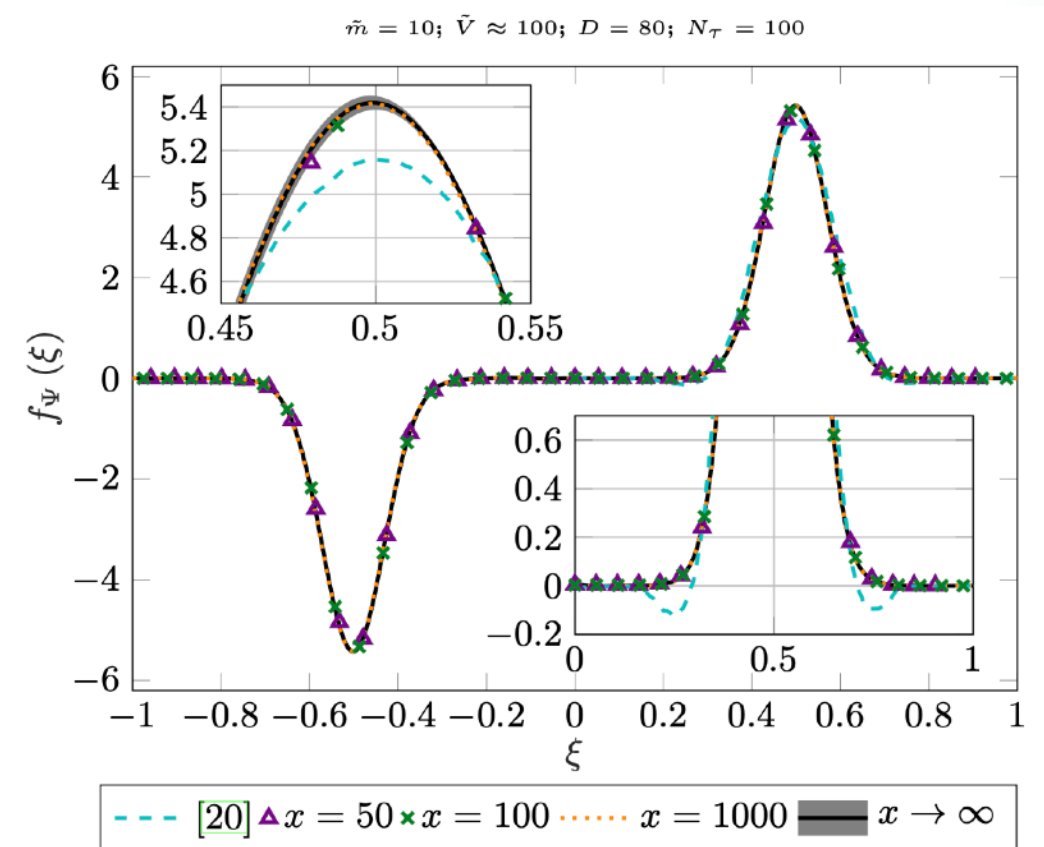
fermionic PDF



different fermion mass, continuum

convergence with very moderate size of tensors

PDF (Fourier transform)



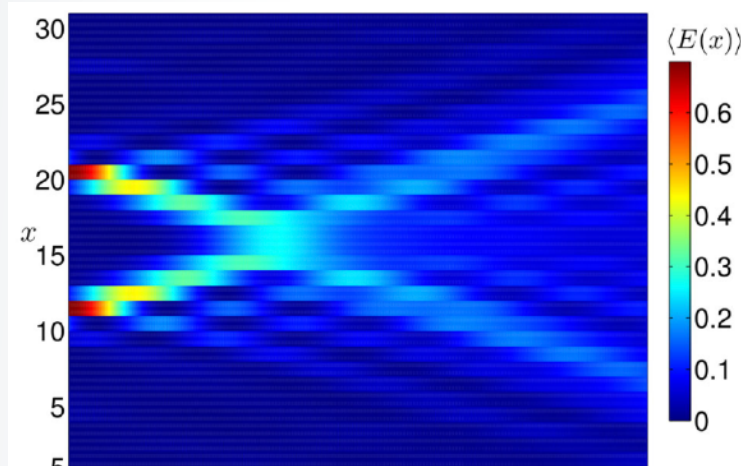
MCB, Cichy, Lin, Schneider, to appear: 2504.XXXXXX
preliminary results presented in LATTICE'24 arXiv:2409.16996

(inelastic) scattering in the Schwinger model

with I. Papaefstathiou, J. Knolle, PRD **111**, 014504 (2025)

earlier simulations: elastic scattering in LGT

PHYSICAL REVIEW X **6**, 011023 (2016)



Real-Time Dynamics in U(1) Lattice Gauge Theories with Tensor Networks

T. Pichler,¹ M. Dalmonte,^{2,3} E. Rico,^{4,5,6} P. Zoller,^{2,3} and S. Montangero¹

PHYSICAL REVIEW D **104**, 114501 (2021)

Entanglement generation in (1+1)D QED scattering processes

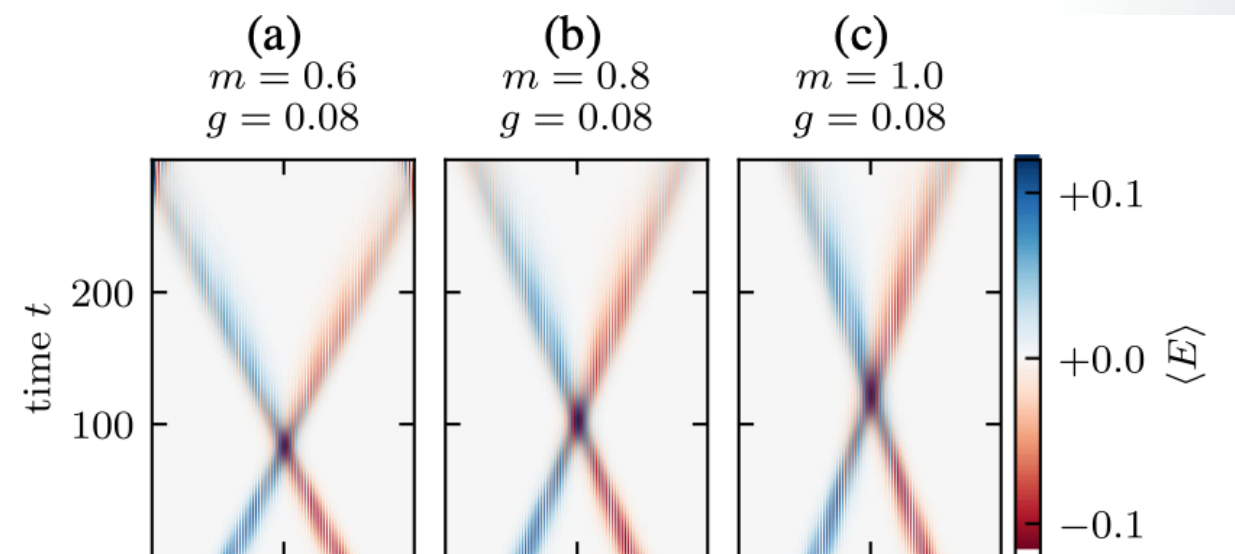
Marco Rigobello^{ID},* Simone Notarnicola^{ID}, Giuseppe Magnifico^{ID}, and Simone Montangero^{ID}

notice also:

Surace, Leroose, New J. Phys. 23 (2021) 062001

Vovrosh et al. PRX Quantum 3, 040309 (2022)

Su, Osborne, Halimeh arXiv:2401.05489



inelastic scattering in the Schwinger model

collision of two vector mesons can produce two scalars

initial setup

$$\sum_n e^{-(n-n_0)^2/(2\sigma^2)} e^{-ikn} (\sigma_n^+ \sigma_{n+1}^- - \sigma_{n+1}^+ \sigma_n^-) |\Omega\rangle$$

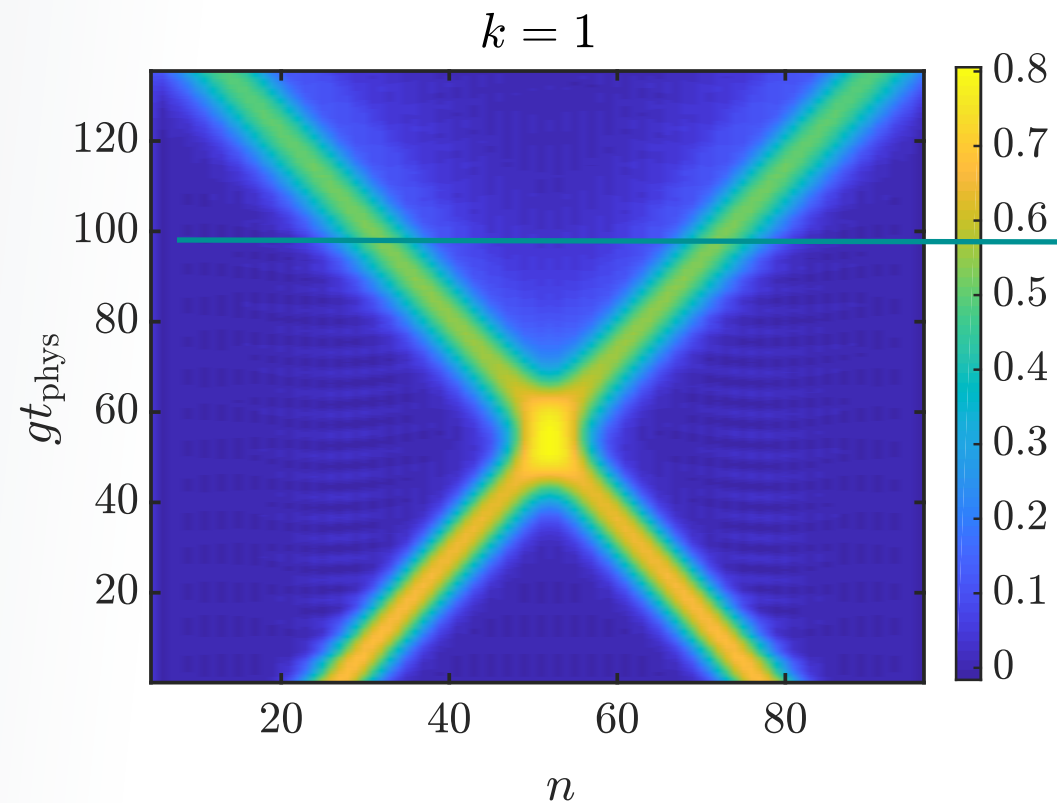
Gaussian wavepacket with momentum k

probe the inelastic threshold

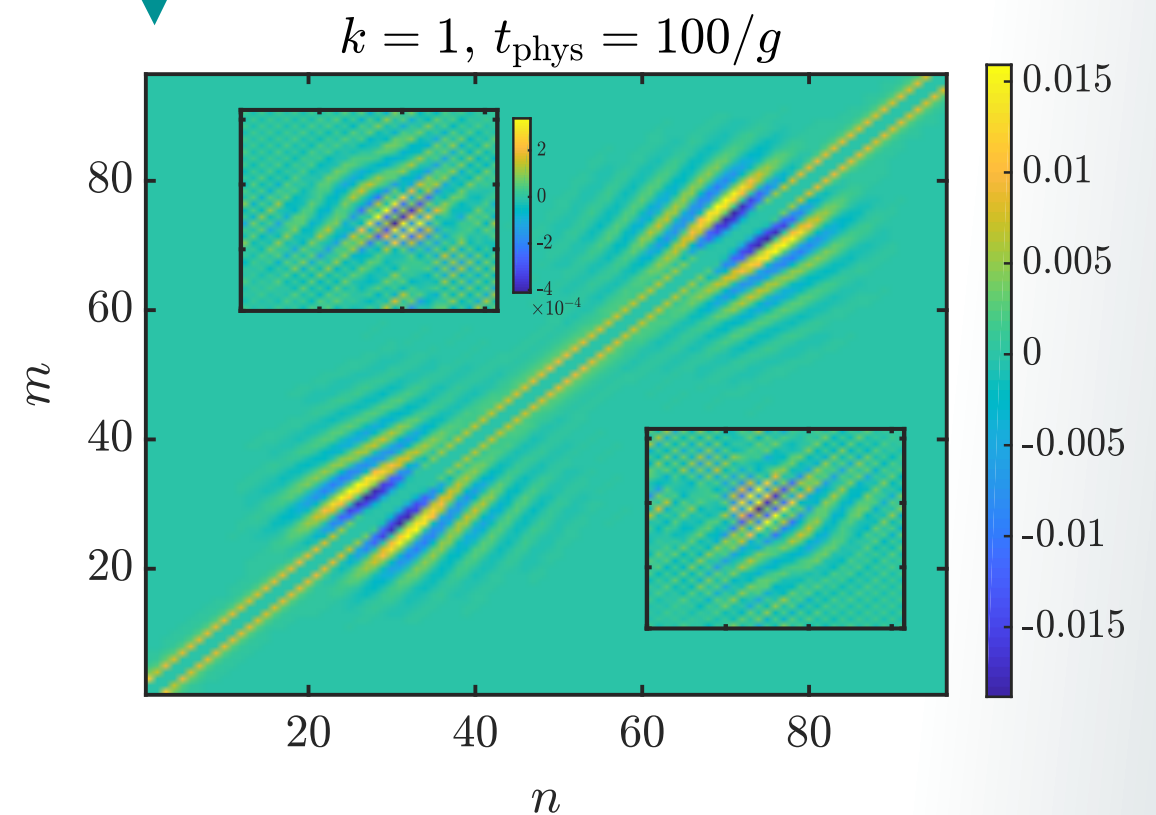
strong coupling regime

below momentum threshold

entropy of two sites

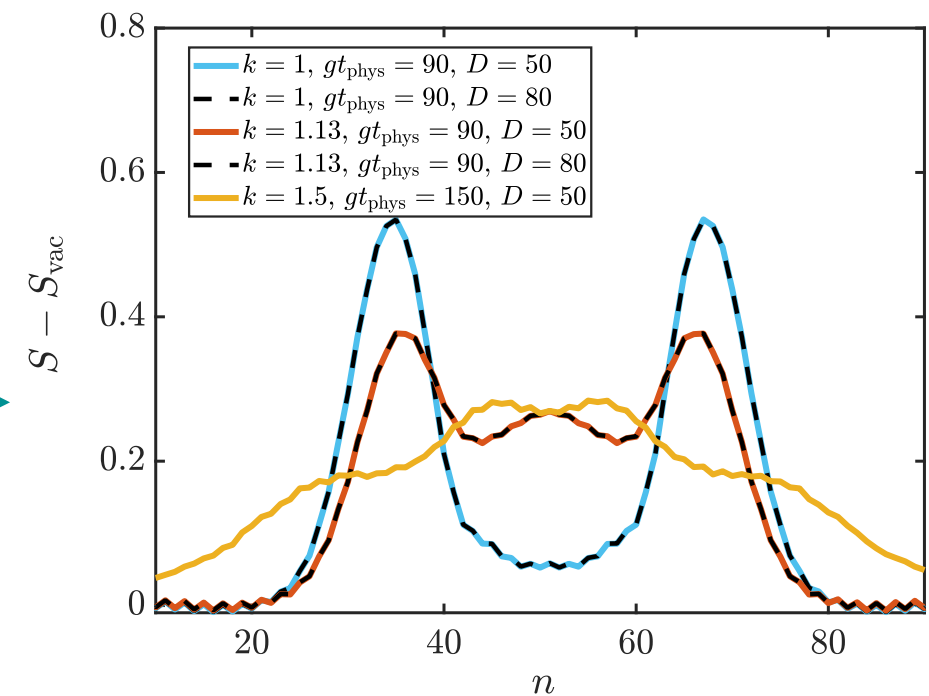
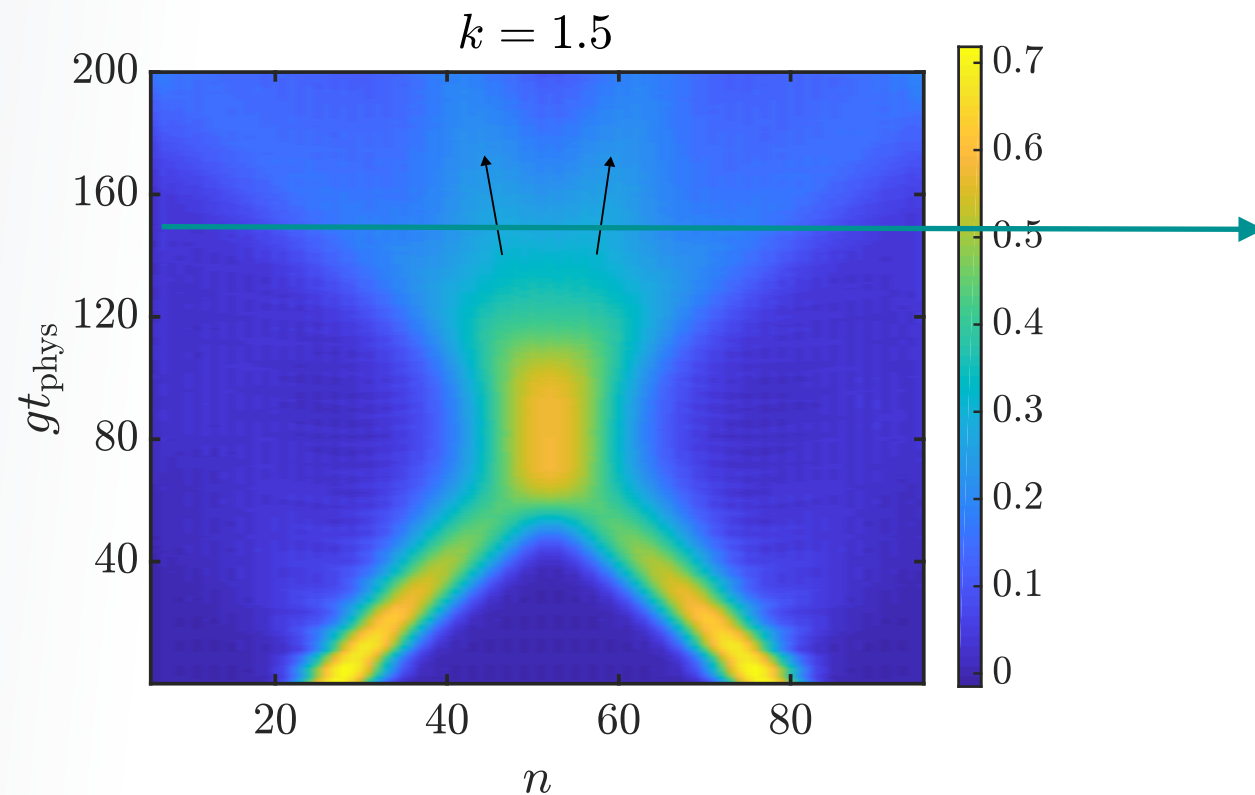


electric flux correlator

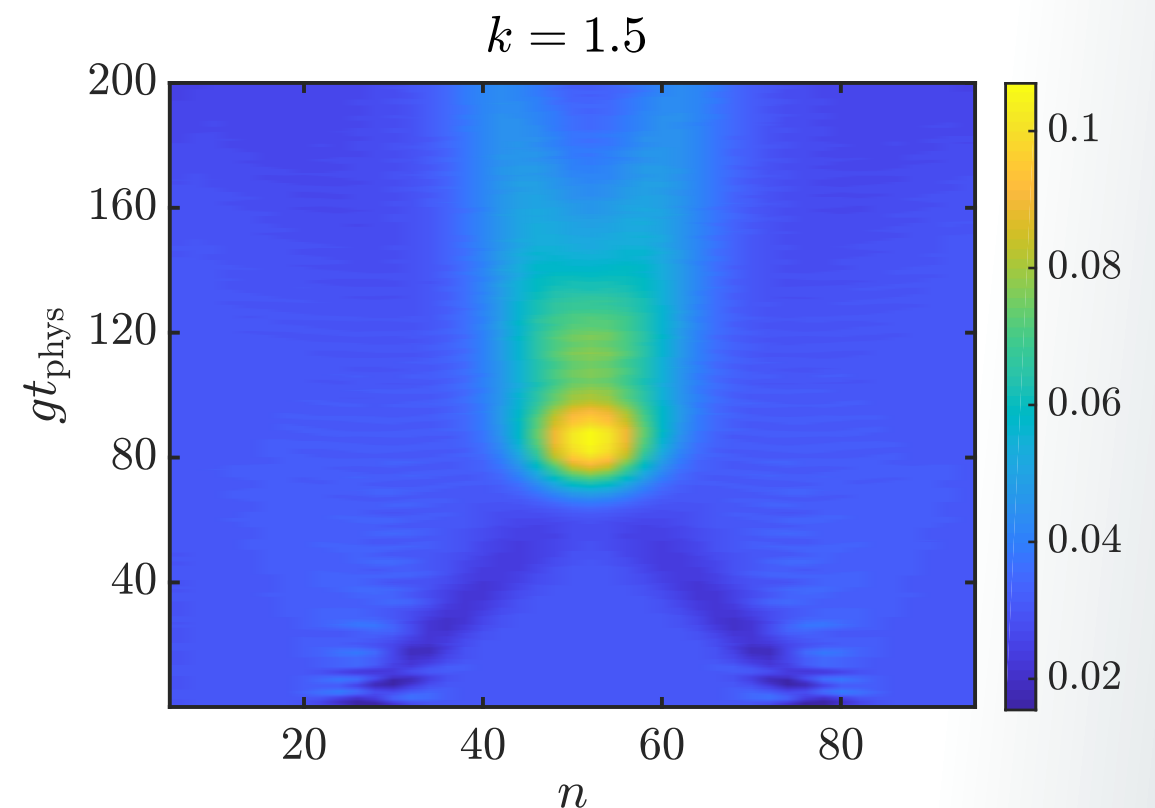


above momentum threshold

entropy of two sites



4-fermion projector



BEYOND ID

active research: PEPS for LGT

explicitly gauge invariant PEPS

restricted ansatz calculations

also fully Gaussian PEPS

Zohar, Cirac PRD 2018

few results with other TNS

Tagliacozzo, Vidal PRB 2011

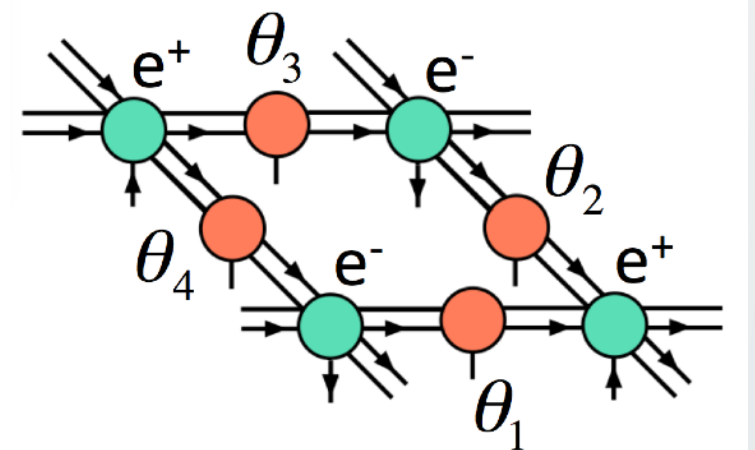
Felser et al. PRX 10, 041040 (2020)

standard PEPS toolbox contains all ingredients

for full variational computation

computational cost, required D

plaquette terms



Zapp, Orús PRD 2017

Tagliacozzo et al PRX 2014

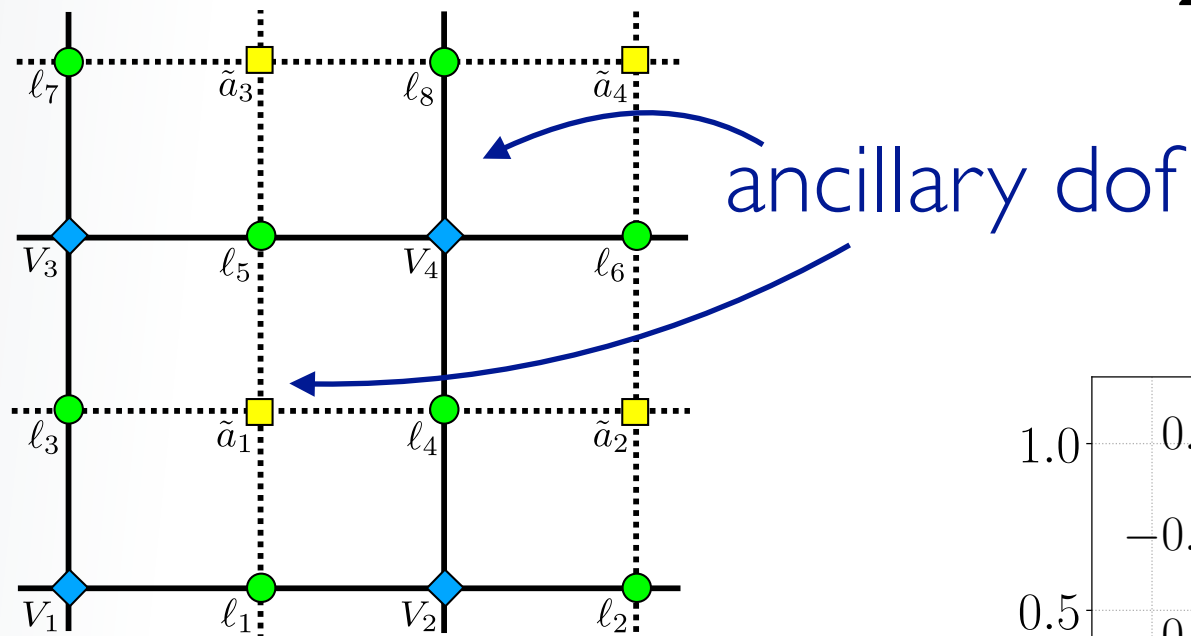
Haegeman et al PRX 2014

Zohar et al Ann Phys 2015

arXiv:1807.01294

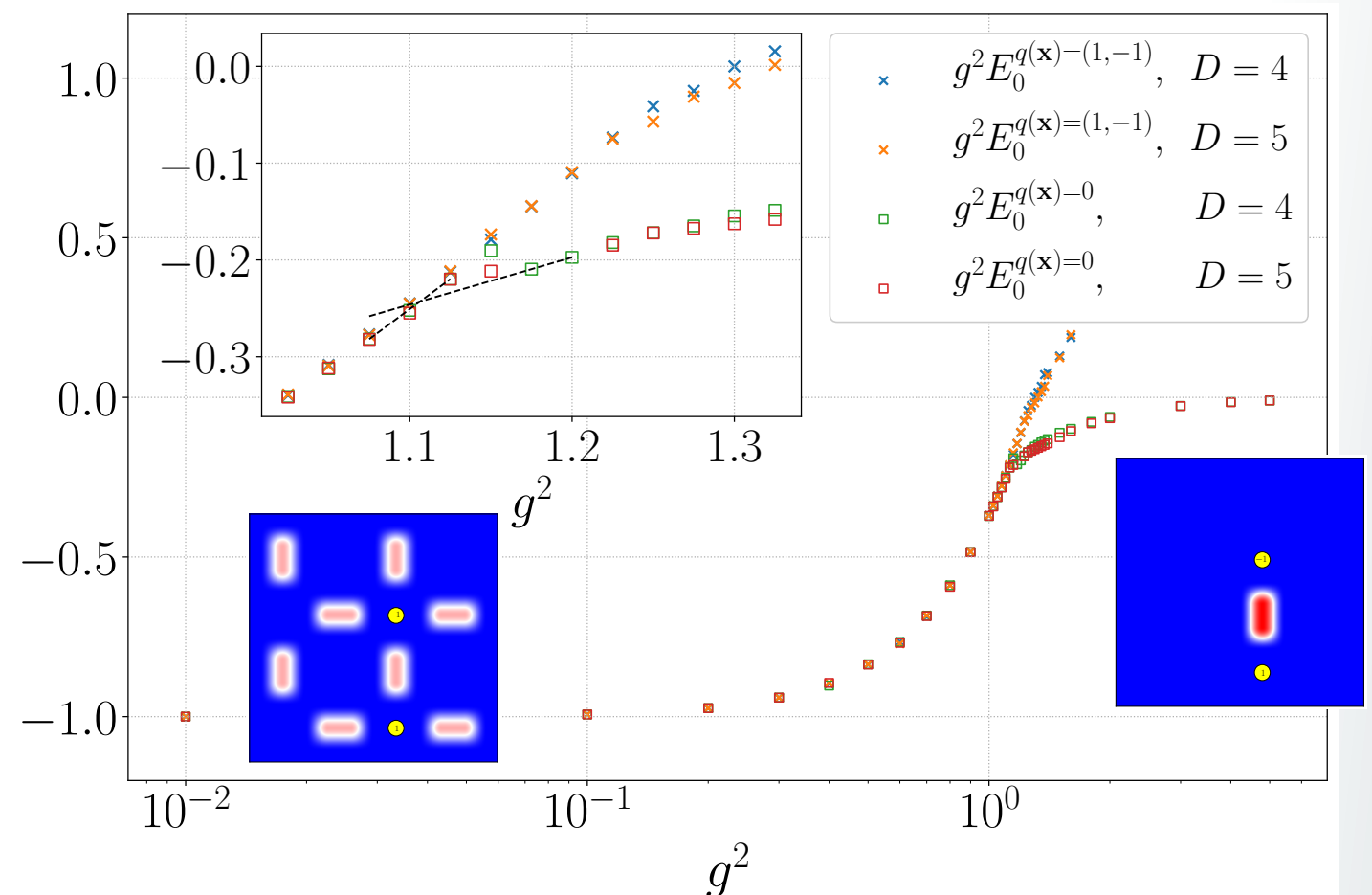
variational iPEPS study of $\mathbb{Z}_3$ in 2+1D

$$H_{\mathbb{Z}_3} = \frac{g^2}{2} \sum_{\ell} E_{\ell}^2 - \frac{1}{2g^2} \sum_P \left(U_P + U_P^{\dagger} \right)$$

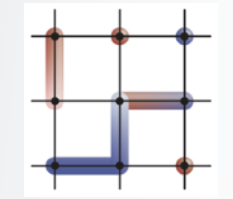


plaquette terms can be implemented as sequence of 2-body gates

Zohar et al, PRA 2017



Robaina, MCB, Cirac, PRL 126, 050401 (2021)



DFG FOR 5522

Thanks for your attention!



MAX PLANCK INSTITUTE
OF QUANTUM OPTICS



TNS: entanglement-based ansatz can be a suitable ansatz also for LGT/ QFT

real time is more challenging than equilibrium

current results: fermion PDF Schwinger model

ab initio calculation in Minkowski space

controlled continuum extrapolation

LATTICE'24 arXiv:2409.16996
to appear 2504.XXXXXX

other dynamical simulations: inelastic scattering

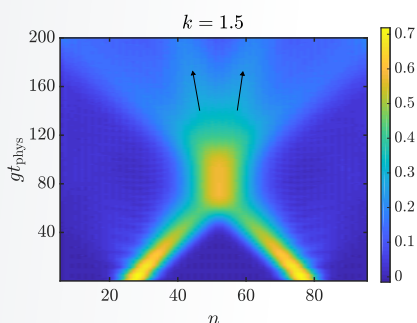
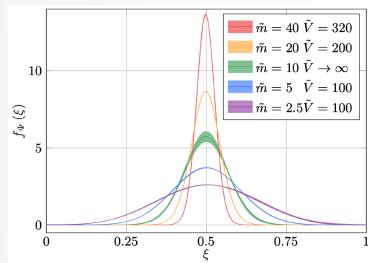
momentum threshold observed

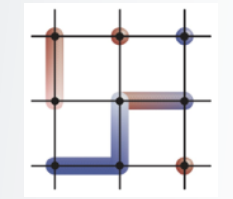
production detected through local observables

Papaefstathiou, arXiv:2402.18429
Belyanski et al PRL 132, 091903 (2024)

ongoing progress in higher dimensions

PEPS, tree TN, MPS





DFG FOR 5522

Thanks for your attention!

TNS: entanglement-based ansatz can be a suitable ansatz also for LGT/ QFT

real time is more challenging than equilibrium

current results: fermion PDF Schwinger model

ab initio calculation
controlled continuation

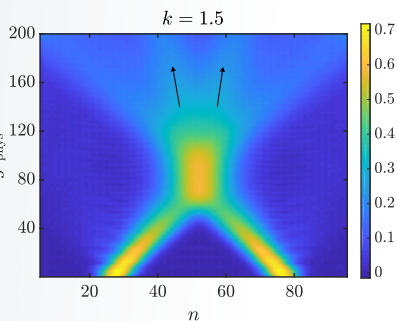
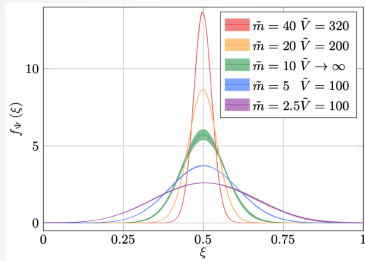
other dynamical simulations

momentum three

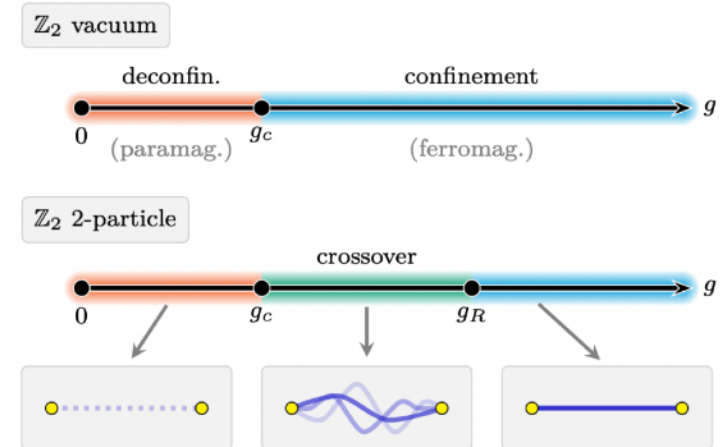
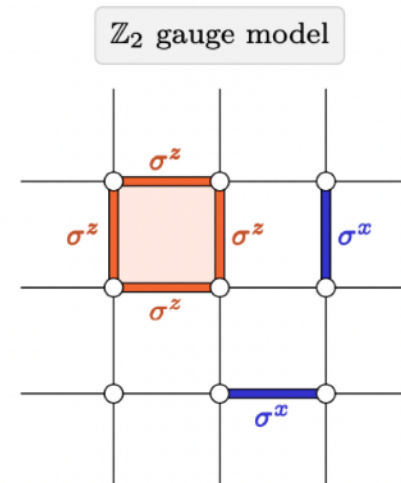
production detection

ongoing progress in

PEPS, tree TN, M



$\mathbb{Z}_2$ roughening transition, and string dynamics



$$H_{\text{LGT}}(g) = -g \sum_{(x,\hat{\mu})} \sigma^x_{(x,\hat{\mu})} - \frac{1}{g} \sum_{\square} (\sigma^z \sigma^z \sigma^z \sigma^z)_{\square}$$

Di Marcantonio, Pradhan, Vallecorsa, Rico, in progress