

# Bayesian constraints and analytic continuation of scattering amplitudes from lattice QCD

Miguel Salg

In collaboration with Fernando Romero-López and William Jay

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BERN

# Background

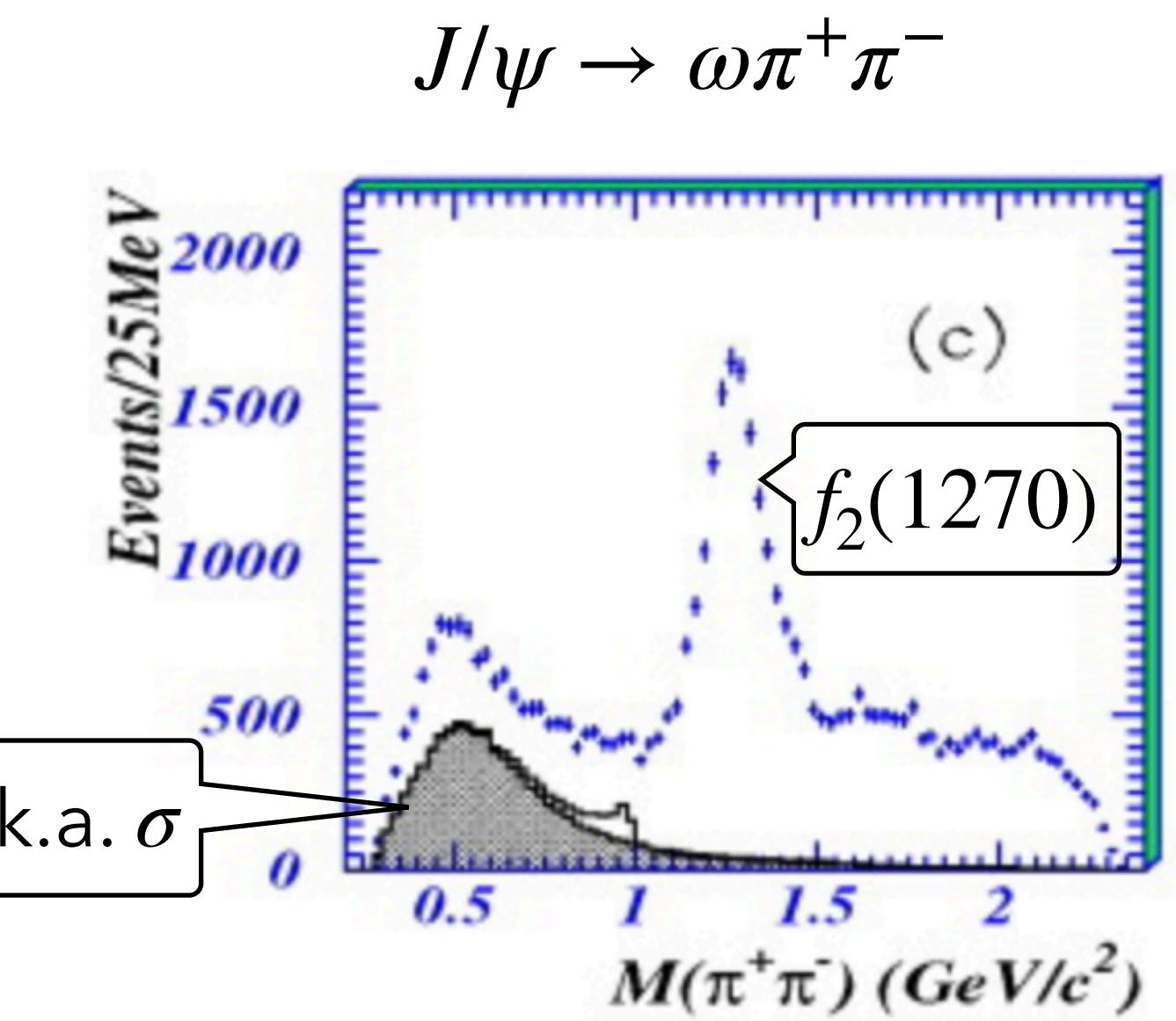
## Hadron spectroscopy

- Goal: obtain **masses** and **decay widths** of hadronic **resonances** from QCD

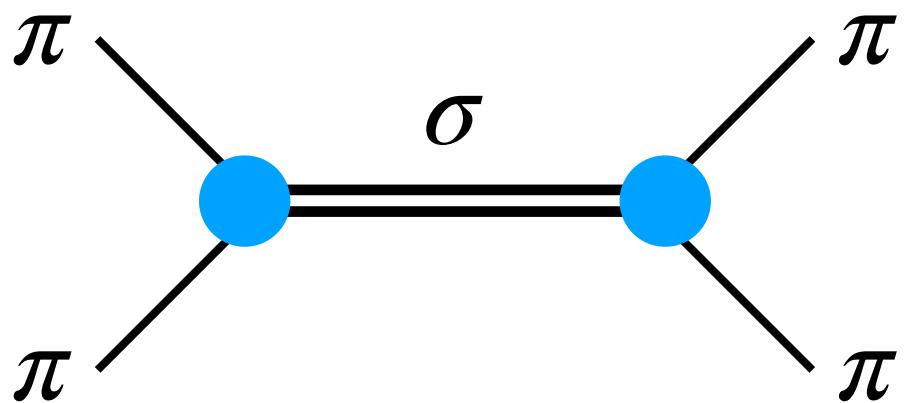
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- Experimentally: resonances as enhancements in **scattering** cross sections



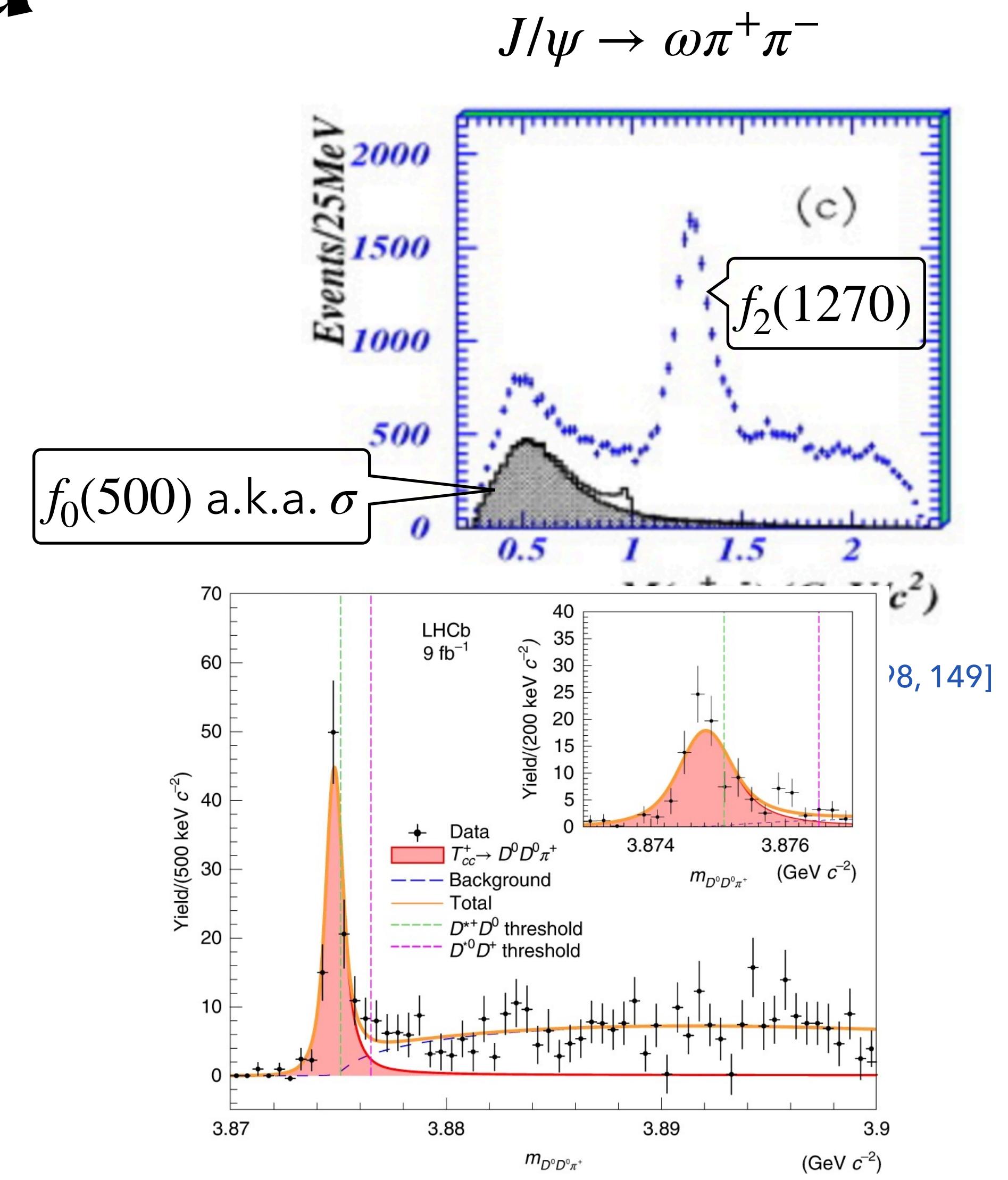
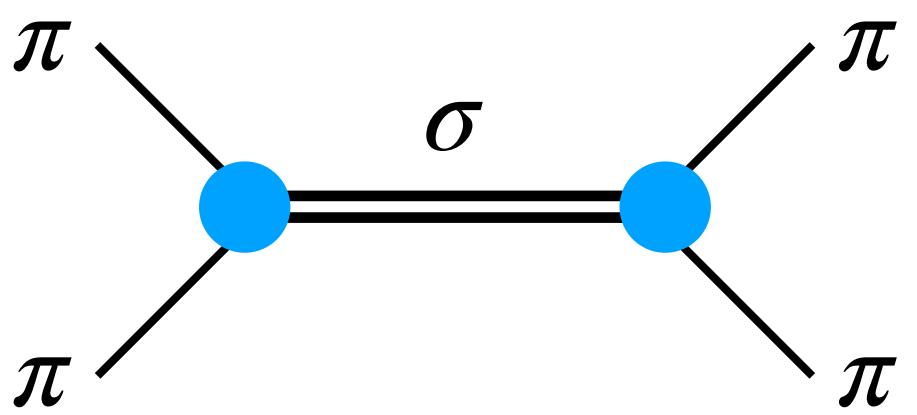
Ablikim et al. (BES) 2004 [Phys. Lett. B 598, 149]



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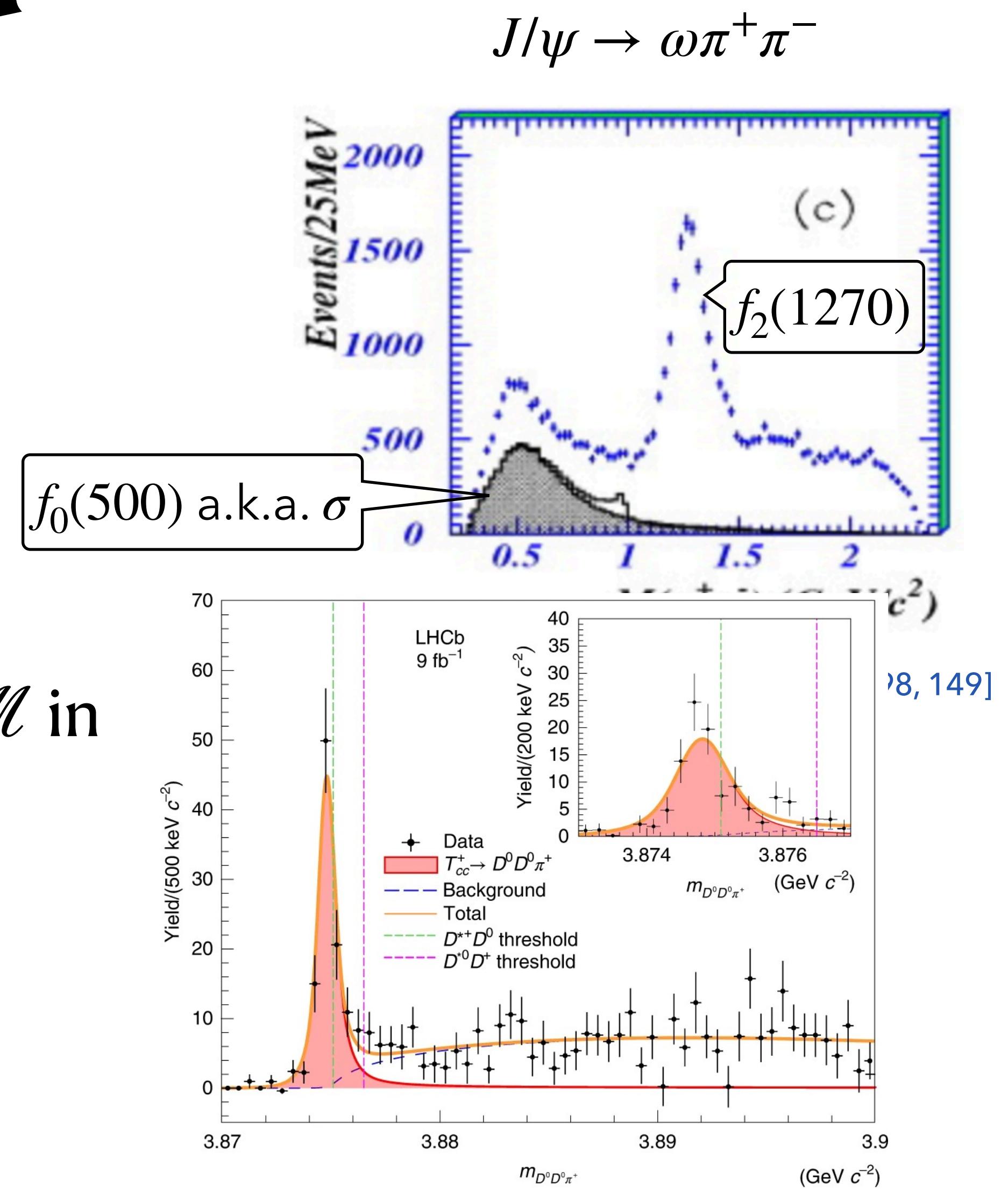
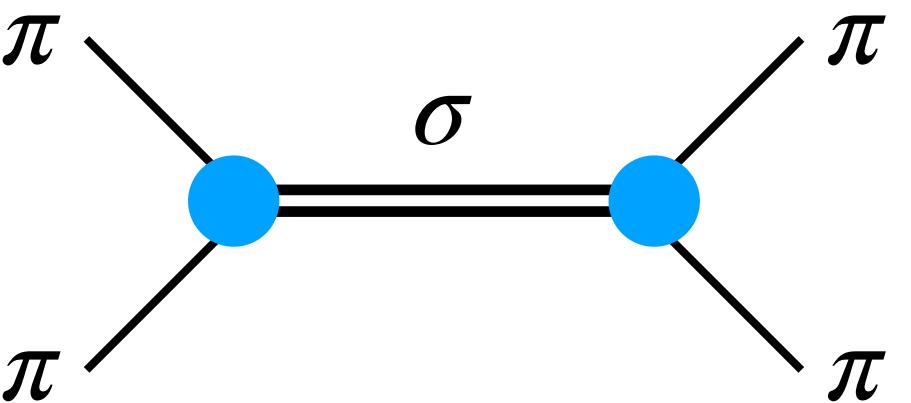


# Background

## Hadron spectroscopy

- Goal: obtain **masses** and **decay widths** of hadronic **resonances** from QCD
- Experimentally: resonances as enhancements in **scattering** cross sections
- Rigorous definition: **poles of scattering amplitude**  $\mathcal{M}$  in complex plane

$$\mathcal{M} \sim -\frac{g^2}{E^2 - E_{\text{pole}}^2}$$



# Background

## Hadron spectroscopy

- Nature of the state depends on the **location** of the pole in the **complex plane**

# Background

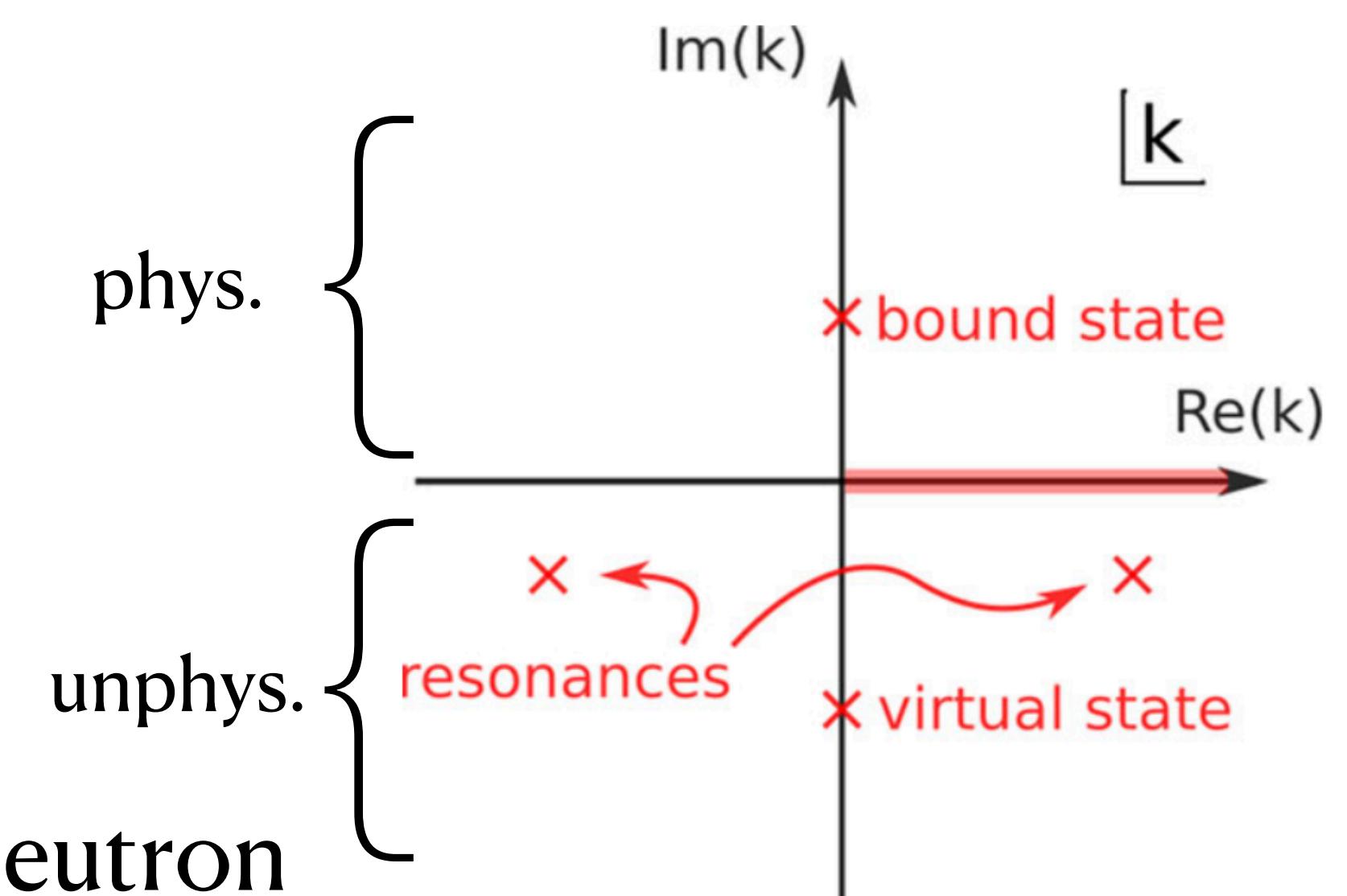
## Hadron spectroscopy

- Nature of the state depends on the **location** of the pole in the **complex plane**
- No poles allowed:
  - Real energies above threshold (**unitarity**)
  - Complex energies on the physical sheet (**causality**)

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## Hadron spectroscopy

- Nature of the state depends on the **location** of the pole in the **complex plane**
- No poles allowed:
  - Real energies above threshold (**unitarity**)
  - Complex energies on the physical sheet (**causality**)
- Classification of allowed poles
  - **Bound states:** stable particles, e.g. deuteron
  - **Virtual states:** „non-normalizable“ QM states, e.g. dineutron
  - **Resonances:** unstable hadrons, e.g.  $\rho, \sigma$ ;  $E_{\text{pole}} = M - i\Gamma/2$



Matuschek et al. 2021 [EPJA 57, 101]

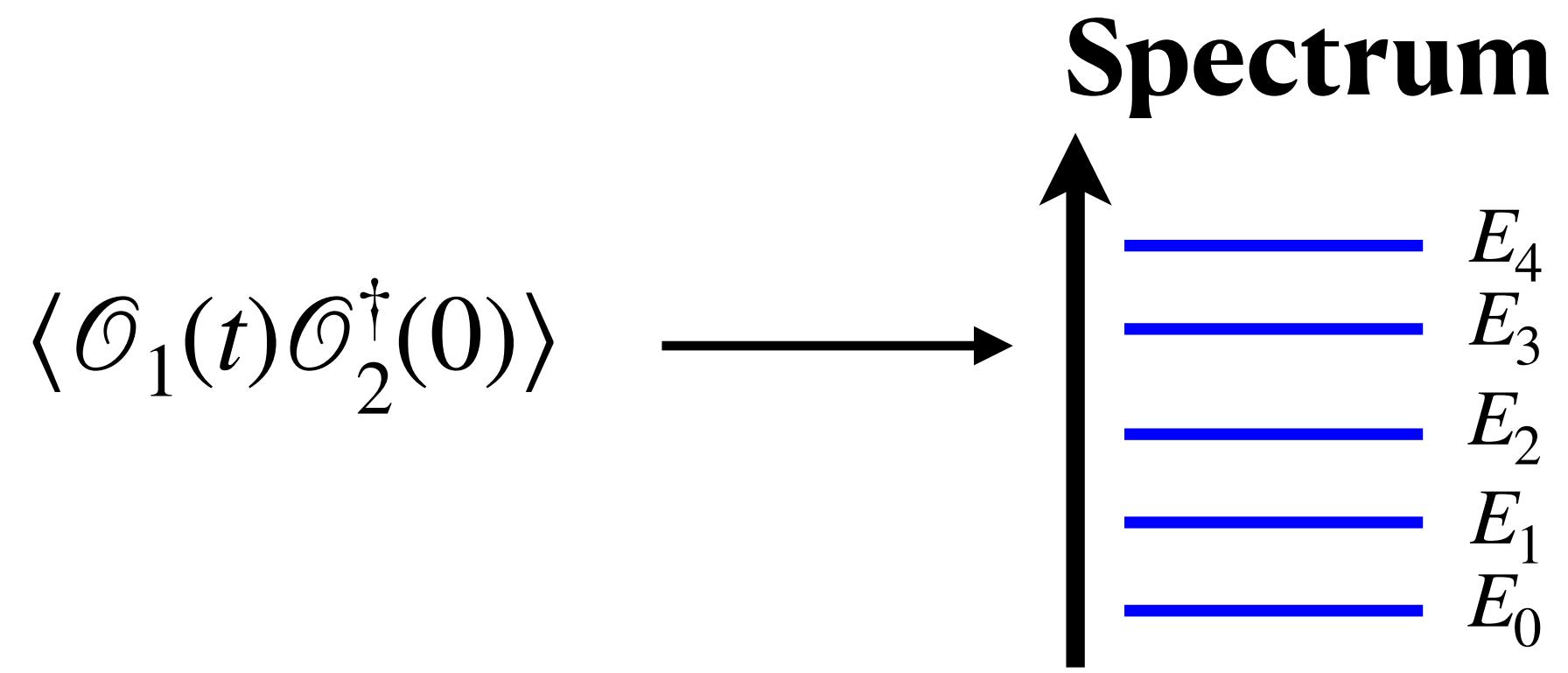
# Background

## Finite-volume formalism for hadron spectroscopy

$$\langle \mathcal{O}_1(t) \mathcal{O}_2^\dagger(0) \rangle$$

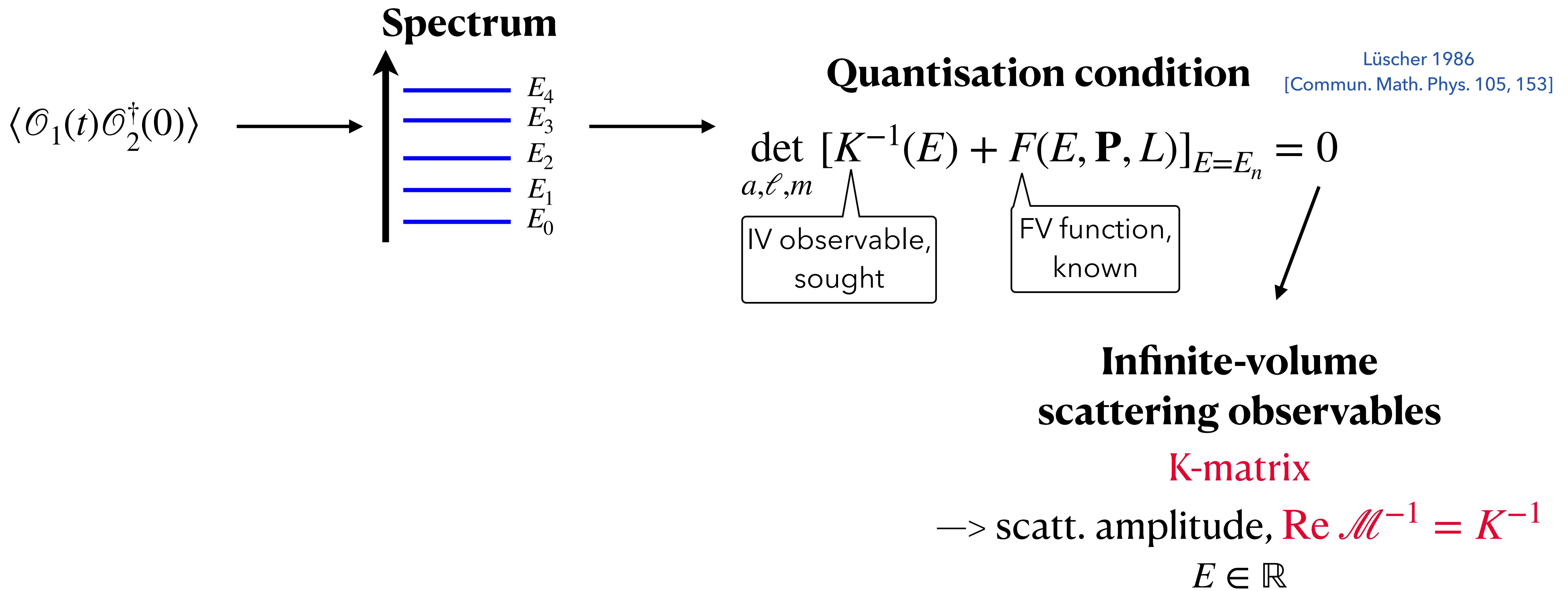
# Background

Finite-volume formalism for hadron spectroscopy



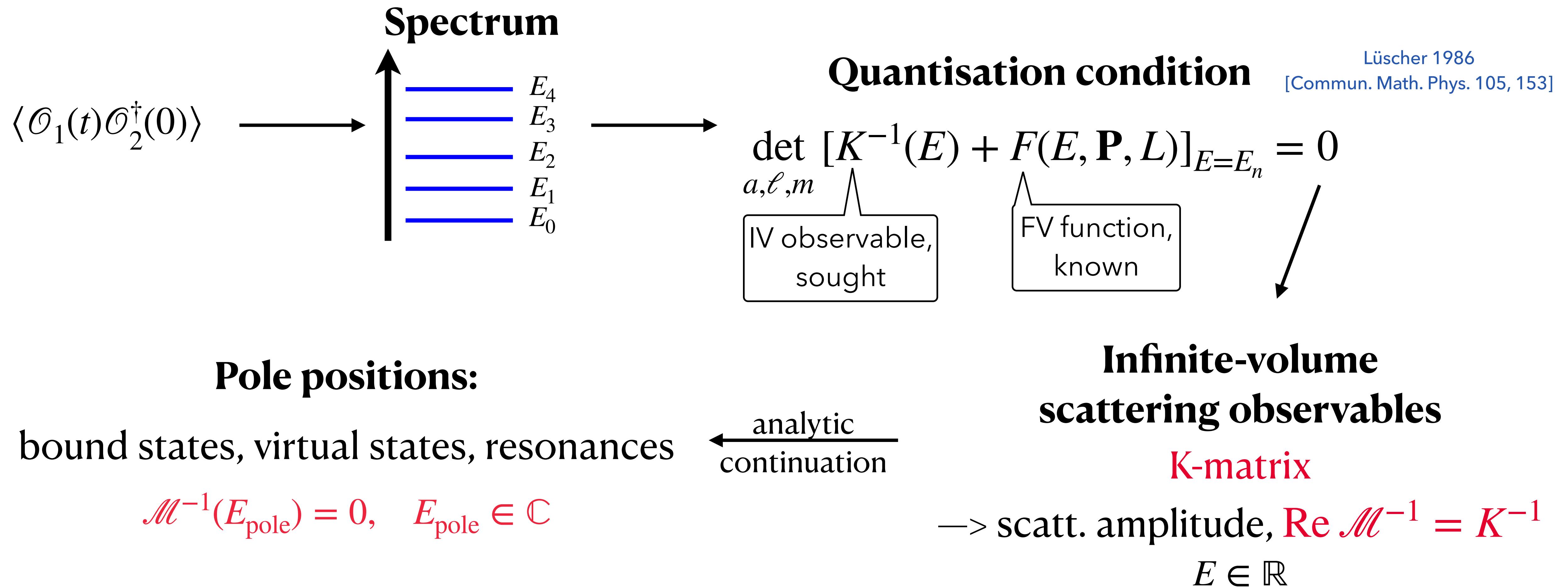
# Background

## Finite-volume formalism for hadron spectroscopy



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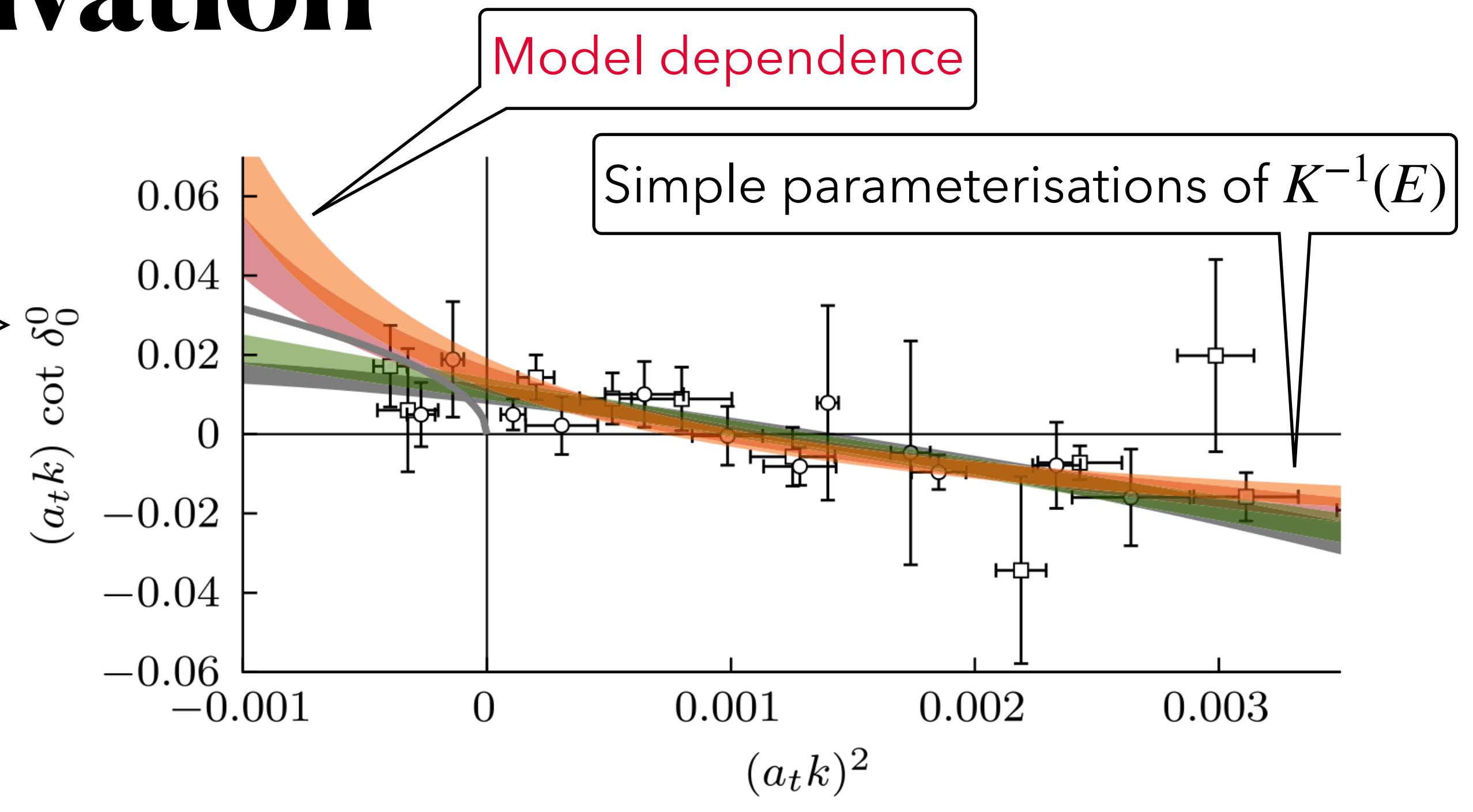
## Finite-volume formalism for hadron spectroscopy



# Motivation

- Conventional approach: **spectrum fits**

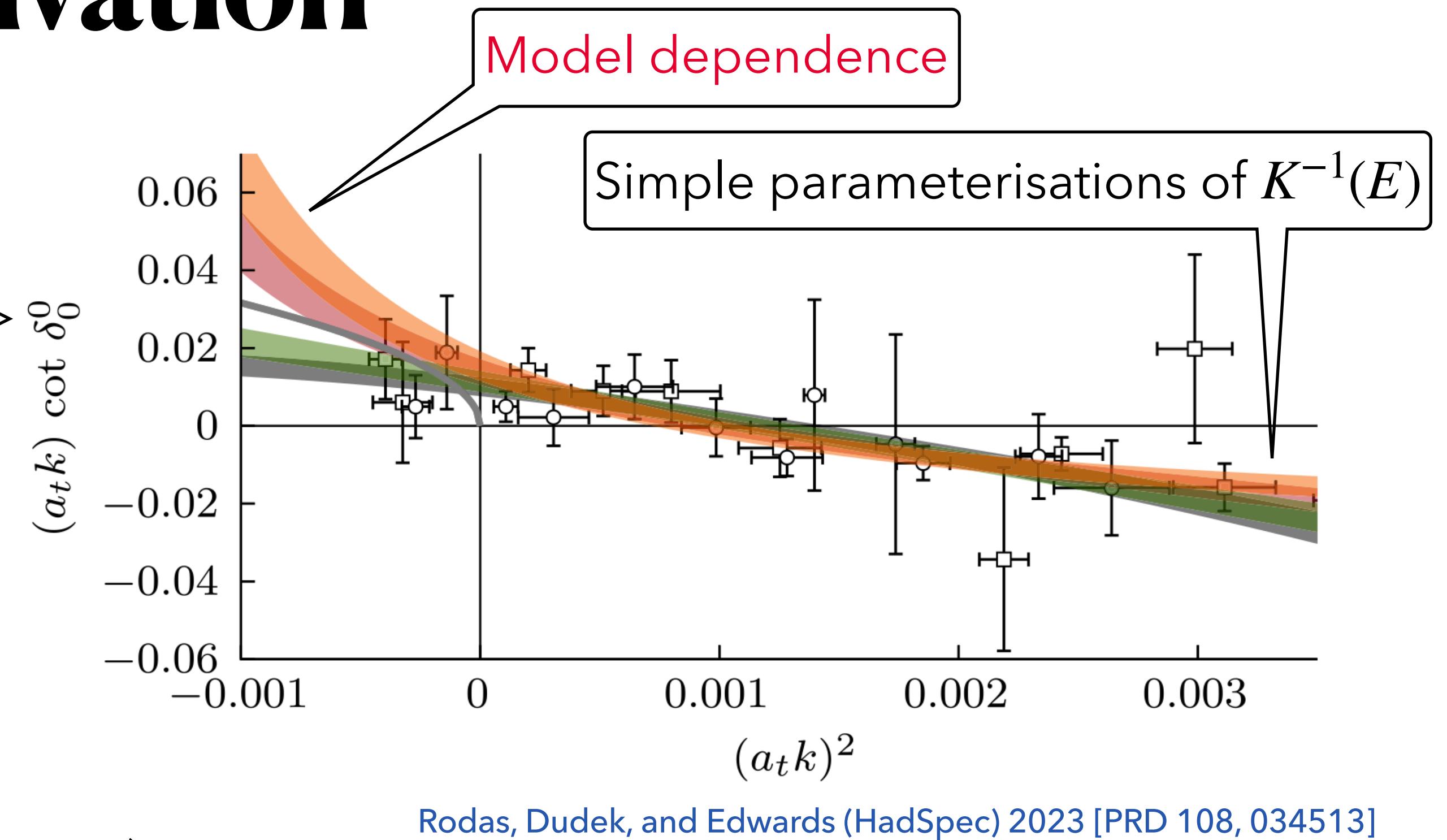
Phase shift,  $K^{-1} \propto k \cot \delta$   
for 1 channel and 1  $\ell$



Rodas, Dudek, and Edwards (HadSpec) 2023 [PRD 108, 034513]

# Motivation

- Conventional approach: **spectrum fits**
  - Phase shift,  $K^{-1} \propto k \cot \delta$  for 1 channel and 1  $\ell$
- Our approach: avoid explicit model dependence
  - Bayesian analysis** of FV spectrum to constrain K-matrix for real energies (inspired by Gaussian processes)
  - Nevanlinna-Pick interpolation** for analytic continuation to complex energies with strict uncertainty quantification —> **Wertevorrat**



Rodas, Dudek, and Edwards (HadSpec) 2023 [PRD 108, 034513]

# Outline of our method

# Bayesian analysis

## General framework

Evaluate  $K^{-1}$  at  $n_p$  discrete „nodes“  $\mathbf{E} \rightarrow$  values  $\mathbf{u} = K^{-1}(\mathbf{E})$

# Bayesian analysis

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Evaluate  $K^{-1}$  at  $n_p$  discrete „nodes“  $\mathbf{E} \rightarrow$  values  $\mathbf{u} = K^{-1}(\mathbf{E})$

**Prior**

$$p(\mathbf{u} | \mathbf{u}_0) = \frac{1}{\sqrt{(2\pi)^{n_p} \det \Sigma_p}} \exp \left[ -\frac{1}{2} (\mathbf{u} - \mathbf{u}_0)^T \Sigma_p^{-1} (\mathbf{u} - \mathbf{u}_0) \right]$$

# Bayesian analysis

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Evaluate  $K^{-1}$  at  $n_p$  discrete „nodes“  $\mathbf{E} \rightarrow$  values  $\mathbf{u} = K^{-1}(\mathbf{E})$

Prior covariance matrix,  
$$\Sigma_p(E, E') = \sigma(E, E')^2 \exp \left[ -\frac{(E - E')^2}{2\ell_c^2} \right]$$

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Prior central values

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Prior central values

### Likelihood

$$p(\mathbf{E}_d | \mathbf{u}) \propto \exp \left[ -\frac{1}{2} \left( \mathbf{E}_d - \mathbf{E}_{QC}(\mathbf{u}) \right)^T \Sigma_d^{-1} (\mathbf{E}_d - \mathbf{E}_{QC}(\mathbf{u})) \right]$$

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LQCD data

Covariance of  $\mathbf{E}_d$

Solution of QC (nonlinear!)

# Bayesian analysis

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LQCD data

**Posterior (Bayes' theorem)**

$$p(\mathbf{u} | \mathbf{E}_d, \mathbf{u}_0) = p(\mathbf{E}_d | \mathbf{u}) p(\mathbf{u} | \mathbf{u}_0)$$

# Bayesian analysis

Computing expectation values and errors

- **Expectation values** given by **Bayesian integrals**, e.g.  $\langle \mathbf{u} \rangle_{\text{post}} = \frac{1}{Z} \int d^{n_p} u \mathbf{u} p(\mathbf{u} | \mathbf{E}_d, \mathbf{u}_0)$
- **Uncertainties** given by **width** of posterior distribution, e.g.  $(\delta \mathbf{u})^2 = \langle \mathbf{u}^2 \rangle_{\text{post}} - \langle \mathbf{u} \rangle_{\text{post}}^2$

# Bayesian analysis

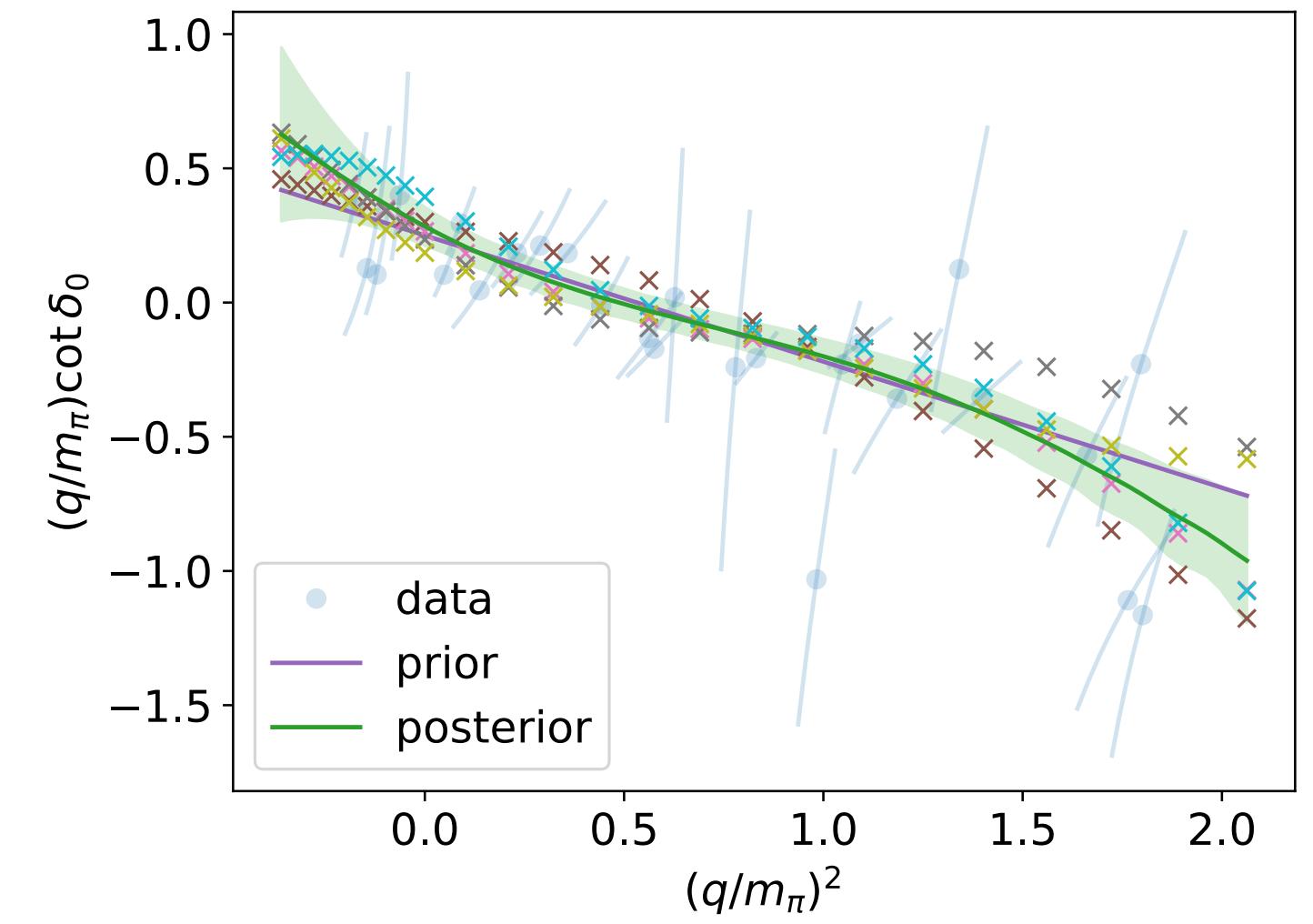
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# Bayesian analysis

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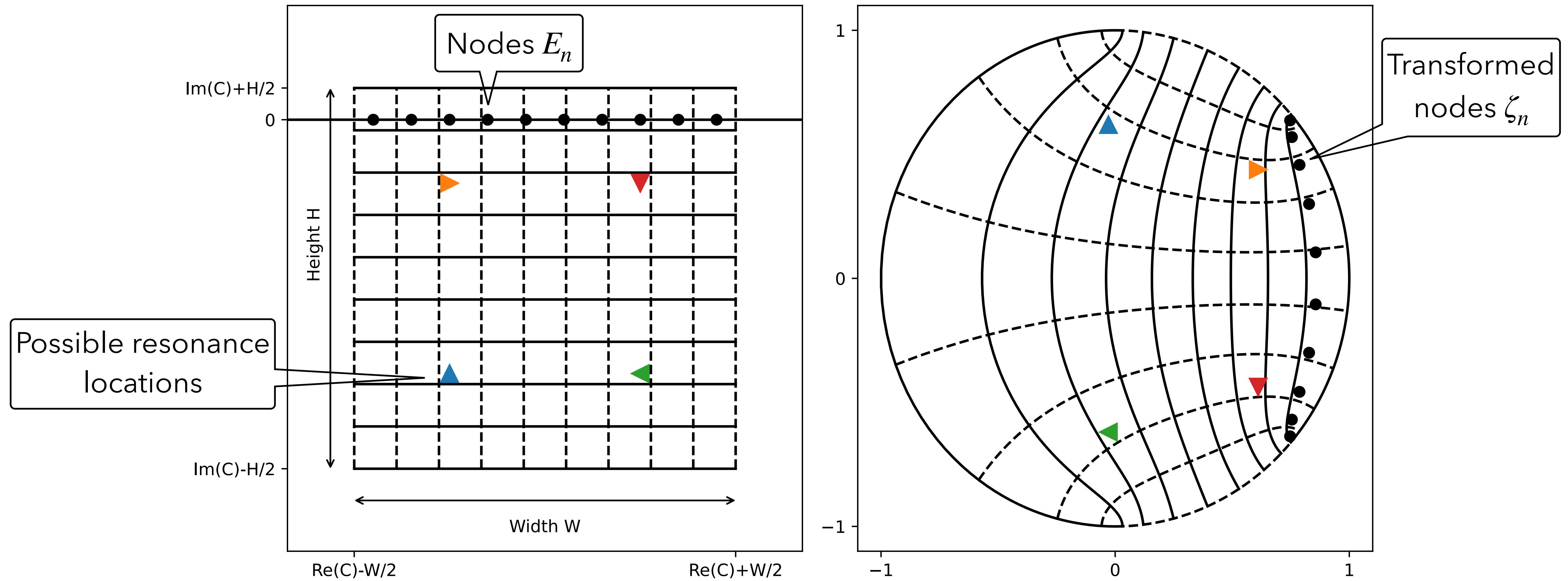
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- Quantify prior dependence by Bayes factor,  $B_{12} = \frac{p(\mathbf{E}_d | M^{(1)})}{p(\mathbf{E}_d | M^{(2)})} = \frac{Z_{\text{post}}^{(1)}}{Z_{\text{post}}^{(2)}}$   
—> ratio of average reweighting factors

Model  $M^{(2)}$  with prior  $\mathbf{u}_0^{(2)}$

Jeffreys 1935 [Math. Proc. Cambridge Philos. Soc. 31, 203]  
Kass and Raftery 1995 [J. Am. Stat. Assoc. 90, 773]

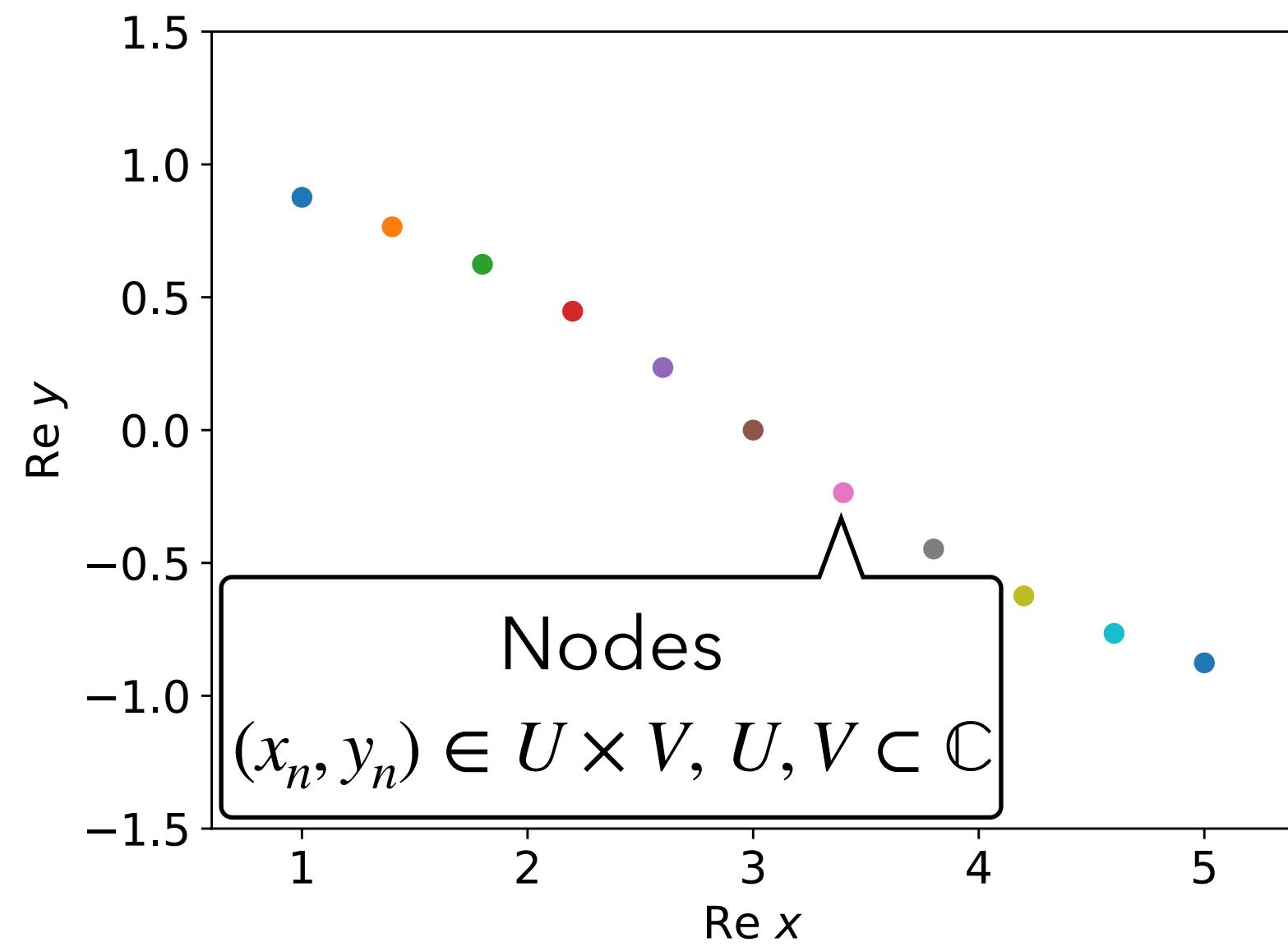
# Nevanlinna-Pick interpolation

## Mapping to the unit disk



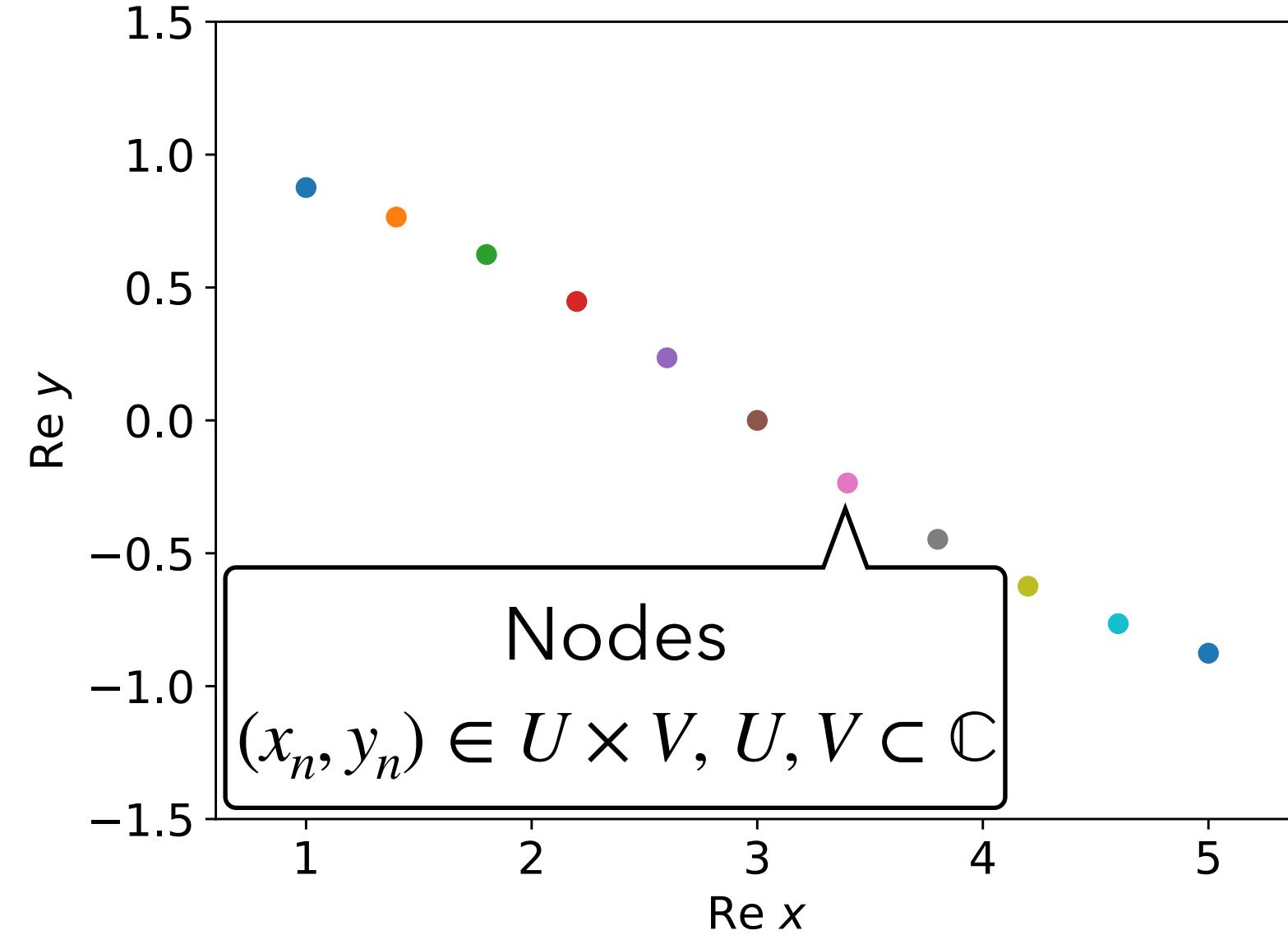
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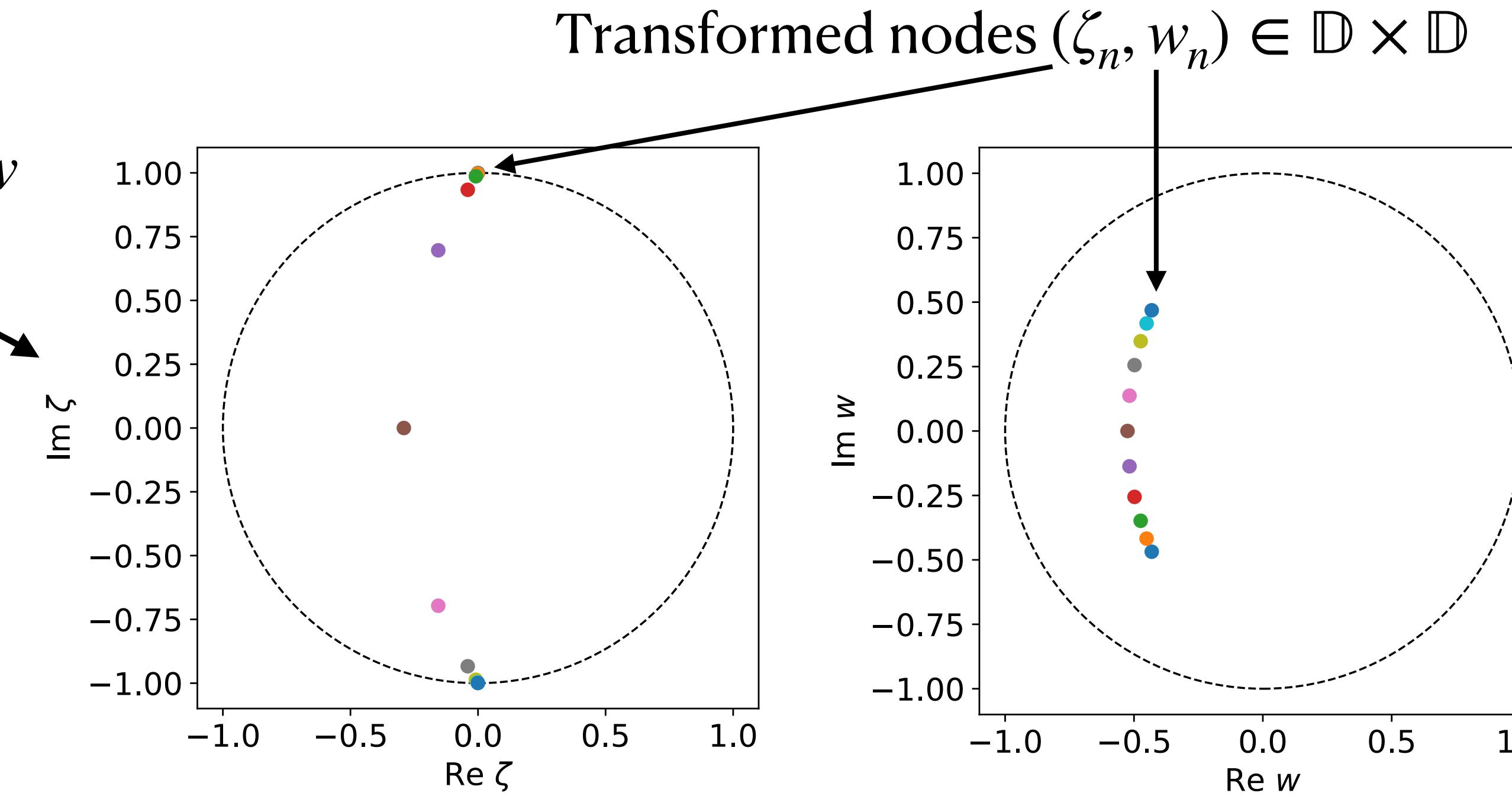


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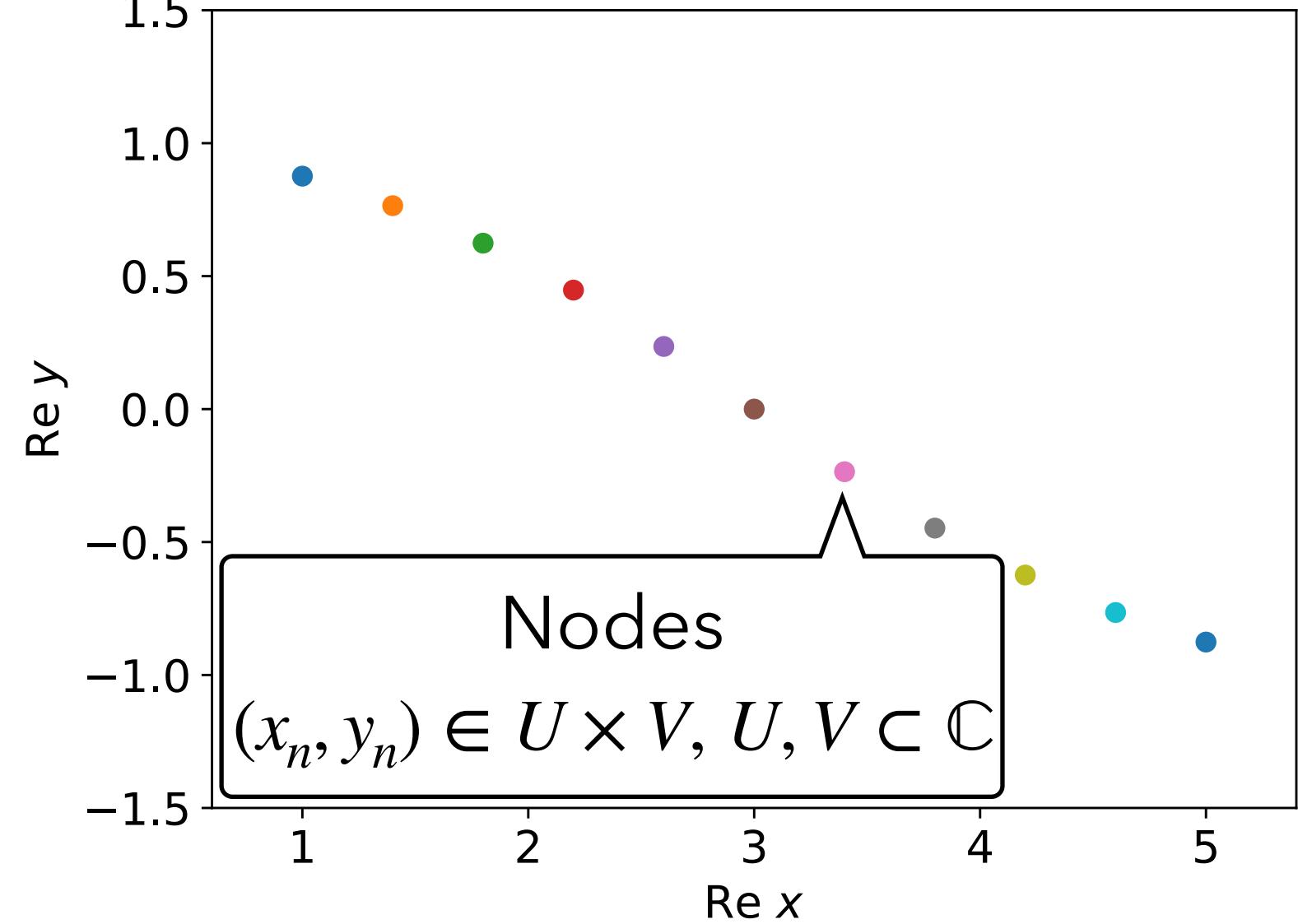


$x \mapsto \zeta, y \mapsto w$   
Map to disk  
**Riemann mapping theorem**

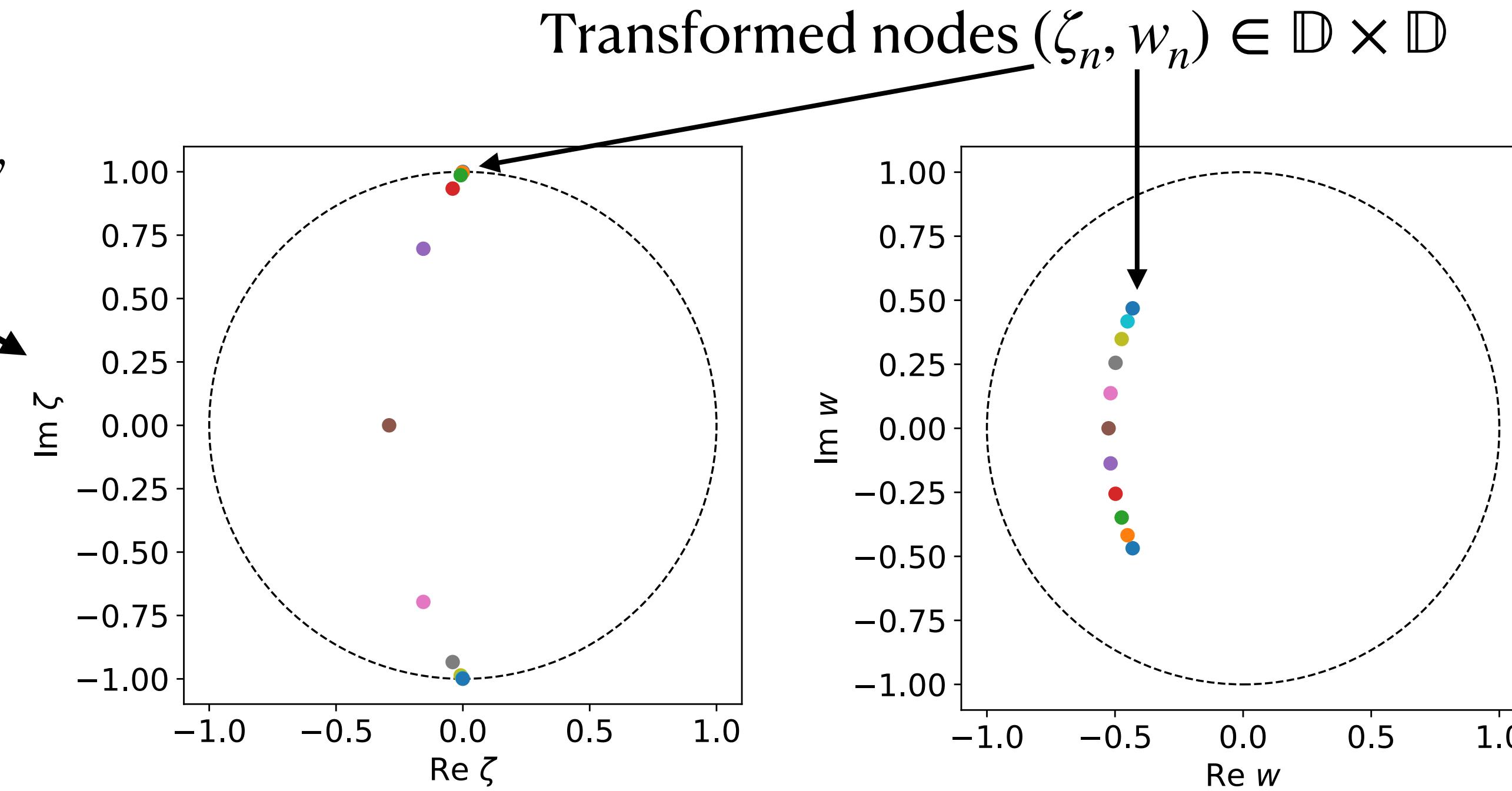


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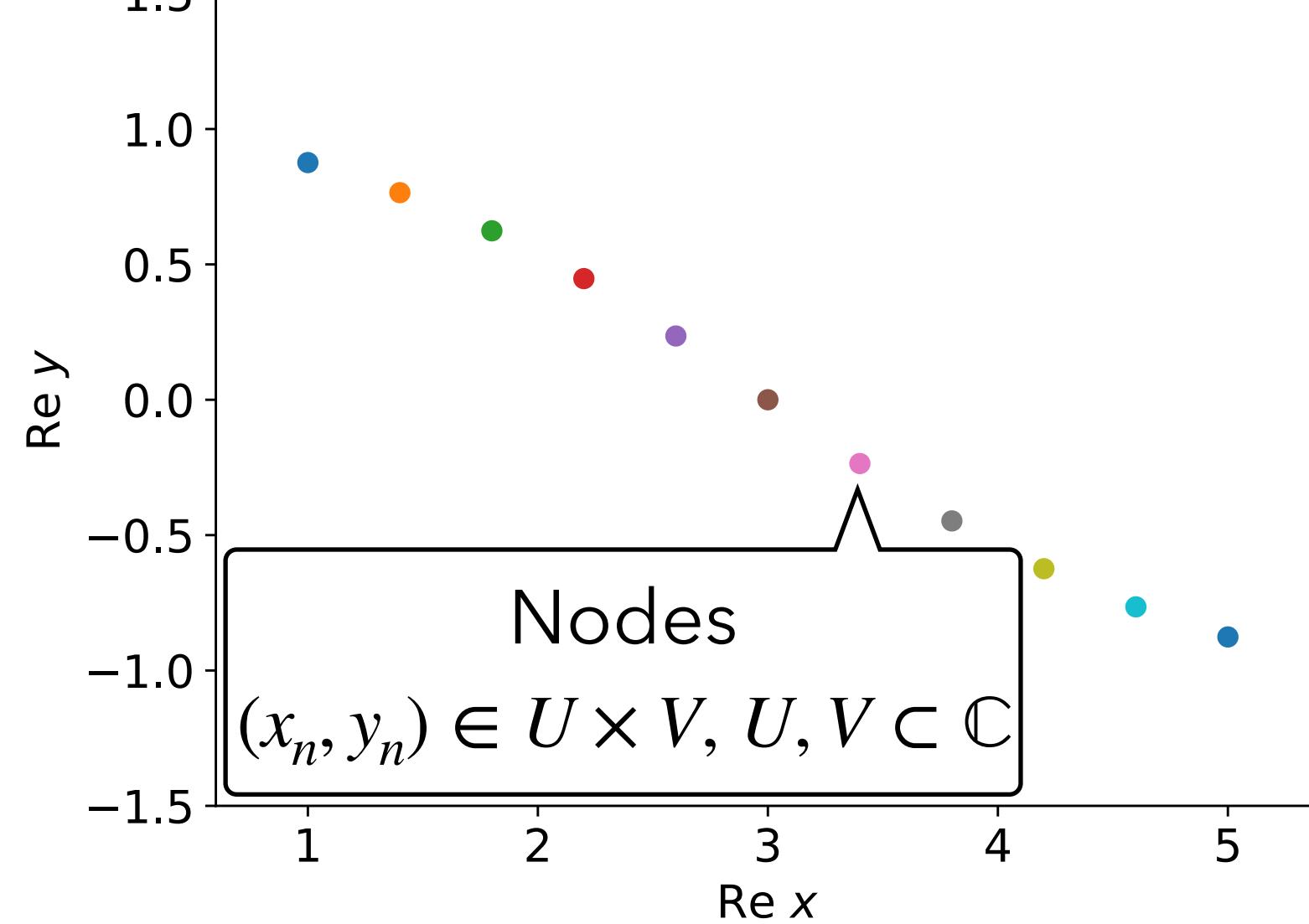


- Any solution can be written as  $f(\zeta) = \frac{P(\zeta)g(\zeta) + Q(\zeta)}{R(\zeta)g(\zeta) + S(\zeta)}$  with  $g : \mathbb{D} \rightarrow \mathbb{D}$  analytic, but arb.
- For fixed  $\zeta$ , all possible solutions fall within a disk  $\Delta(\zeta) \subset \mathbb{D} \rightarrow$  Wertevorrat (WV)

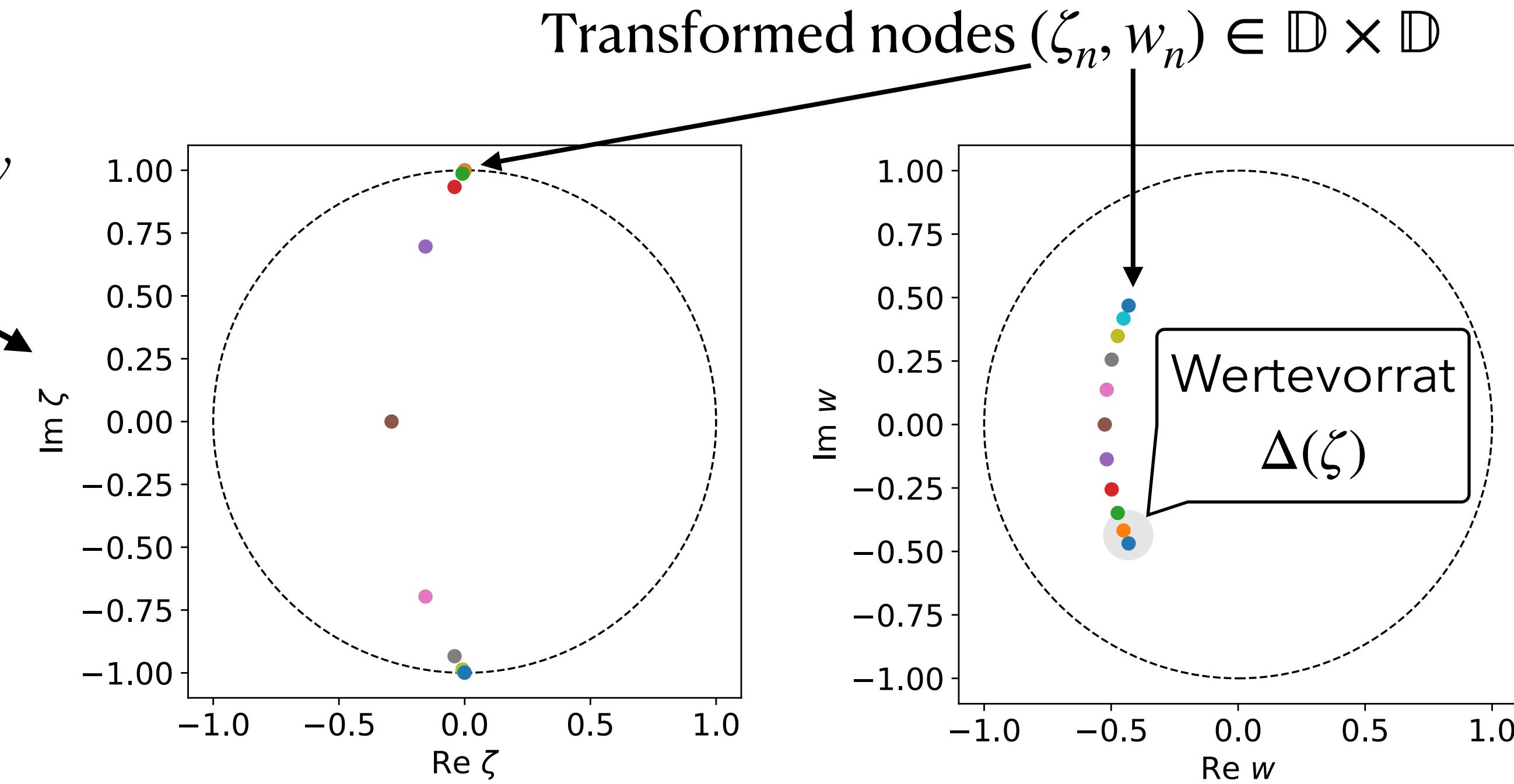
Recursively calculable coefficients

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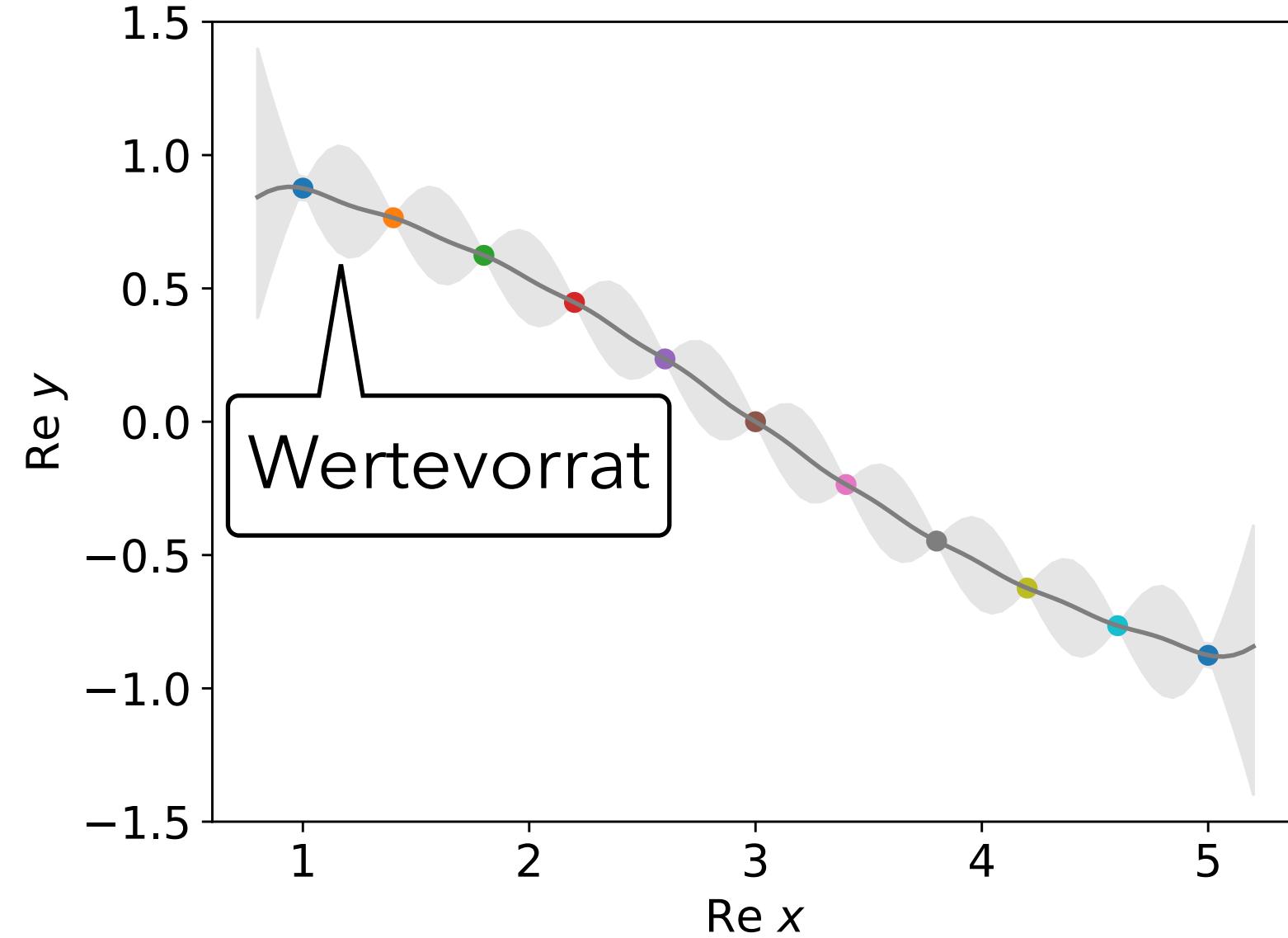


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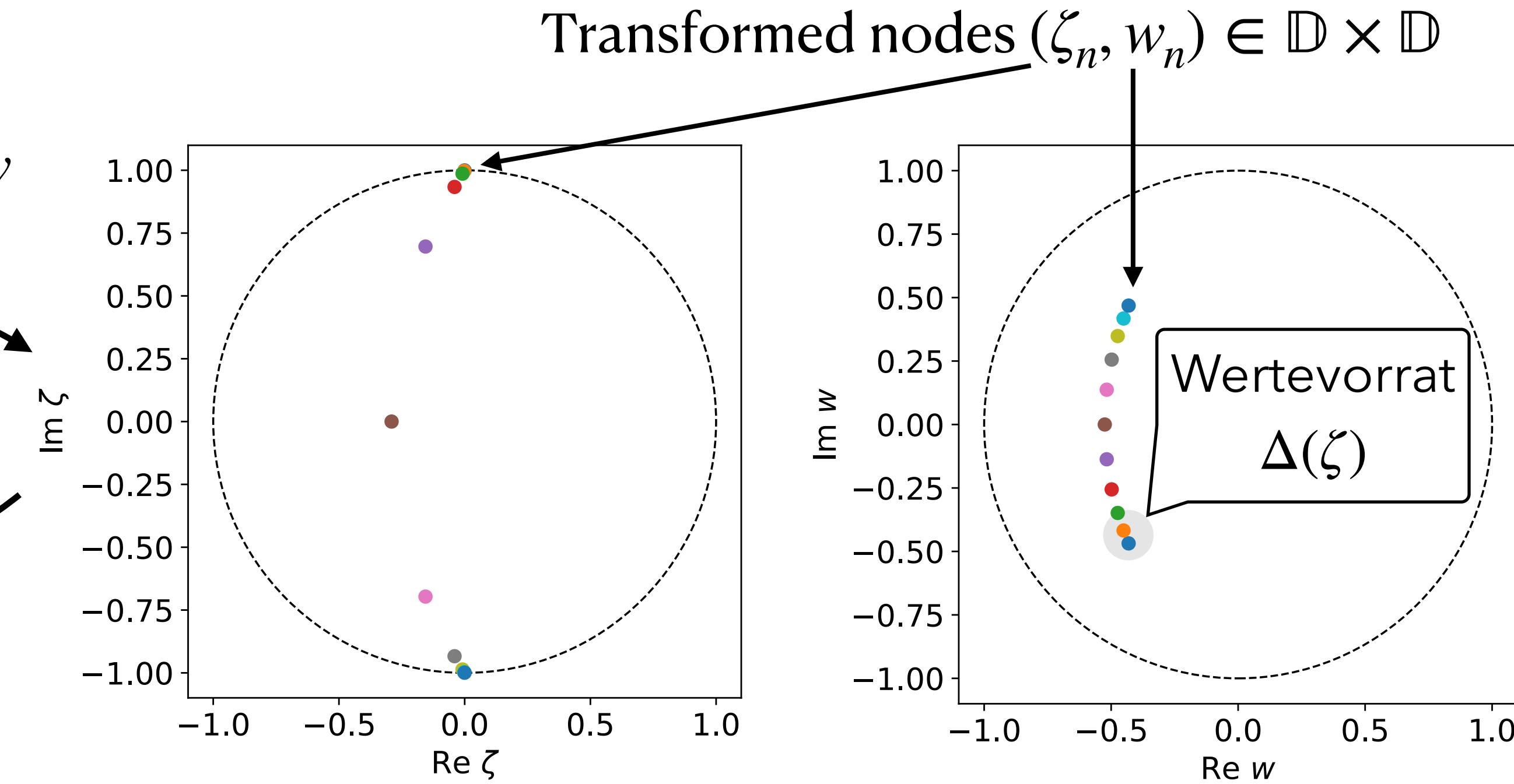
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## General framework



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Map to disk  
**Riemann mapping theorem**  
Map back



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# Nevanlinna-Pick interpolation

Application to analytic continuation of scattering amplitudes

1. Transform domain (energies  $\mathbf{E}$ ) and codomain (values  $\mathbf{u}$ ) of  $K^{-1}$  to unit disk using Schwarz-Christoffel mapping of a rectangle to the disk [Liu, Chen, and Karageorghis 2017 \[J. Sci. Comput. 71, 1035\]](#)

# Nevanlinna-Pick interpolation

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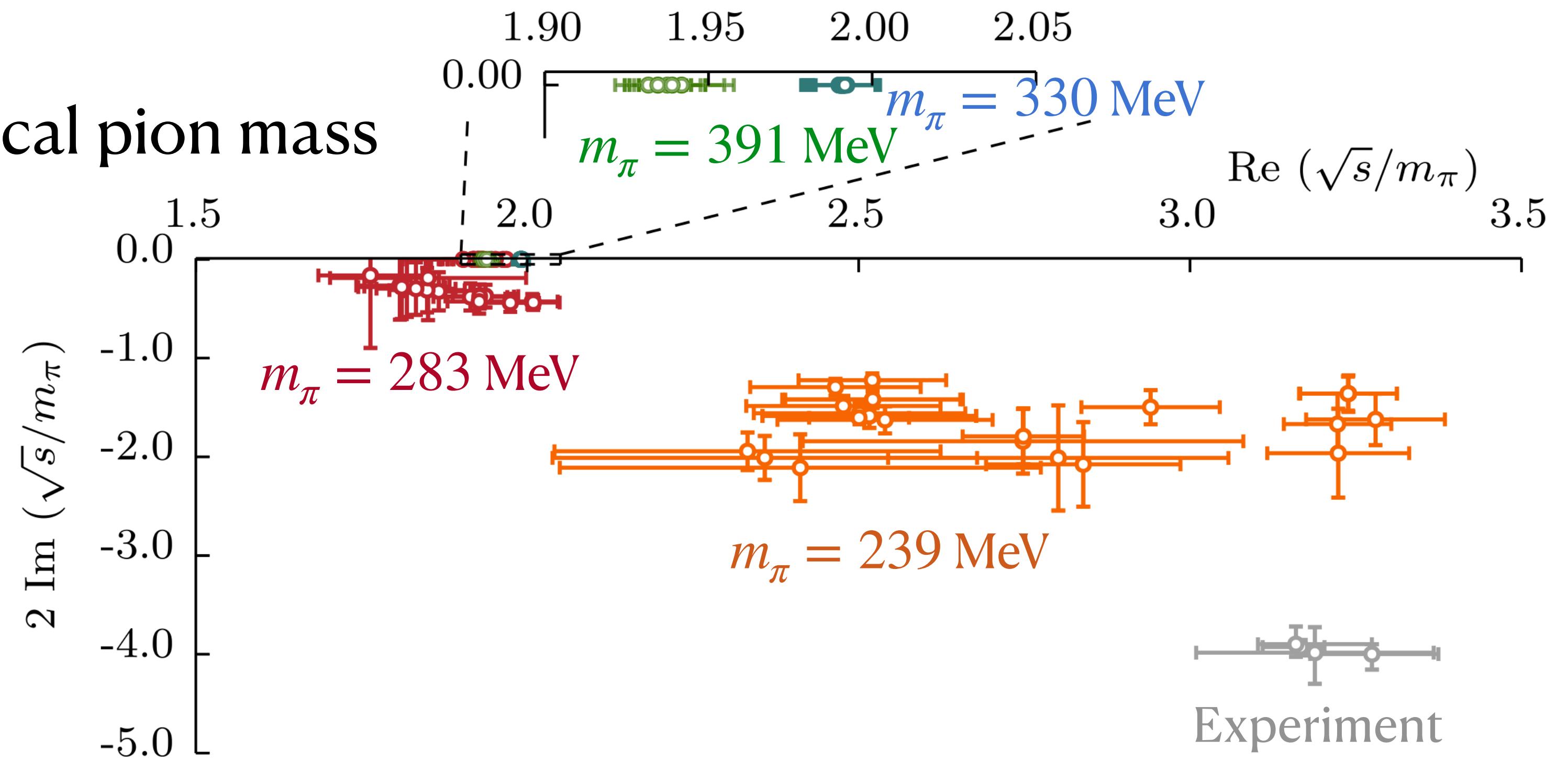
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2. For each sample from posterior:
  - 2.1. Compute Nevanlinna coefficients  $P, Q, R, S$
  - 2.2. Evaluation of interpolating function + uncertainty (Wertevorrat) for complex energies furnishes desired analytic continuation (after mapping back)
  - 2.3. Solve for zeros of  $\mathcal{M}^{-1}(E)$  in the complex energy plane, incorporating WV uncertainty („sampling from WV“) —> resonances

# Applications to lattice QCD analyses

# Isospin-0 $\pi\pi$ system on HadSpec ensembles

Pion-mass dependence of the  $\sigma$  pole position

Broad  $\sigma$  meson resonance at physical pion mass



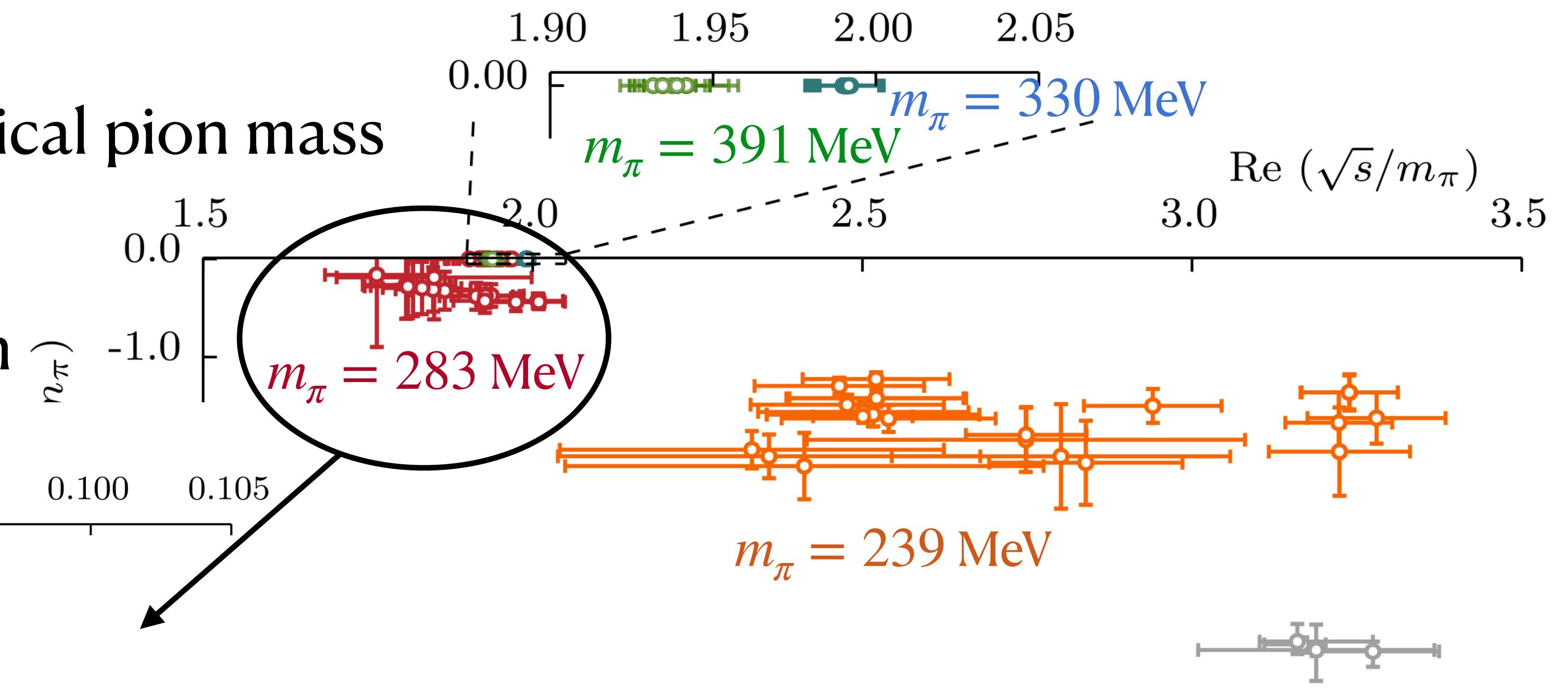
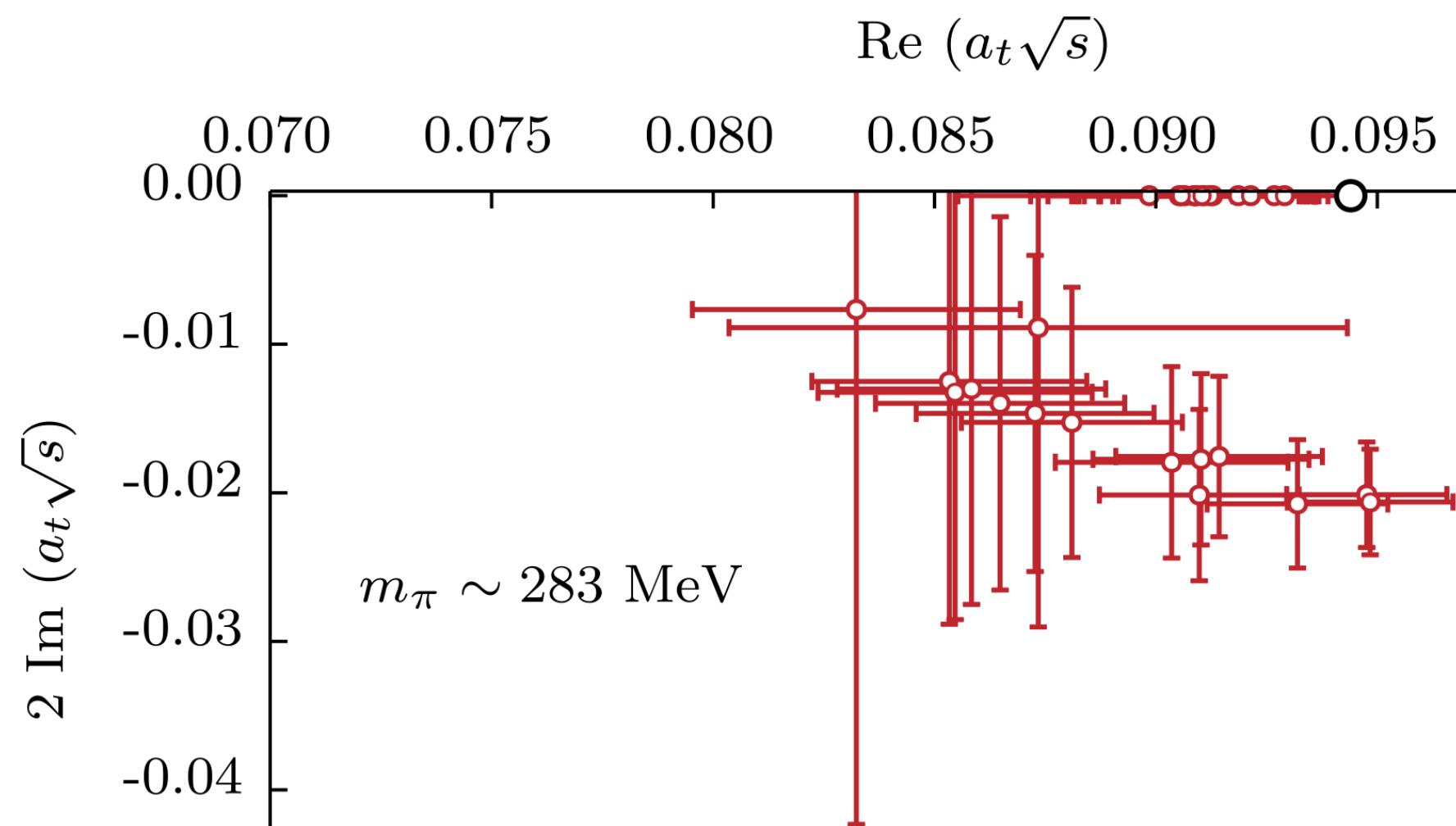
Rodas, Dudek, and Edwards (HadSpec) 2023 [PRD 108, 034513]

# Isospin-0 $\pi\pi$ system on HadSpec ensembles

Pion-mass dependence of the  $\sigma$  pole position

Broad  $\sigma$  meson resonance at physical pion mass

Pole position hard to pin down



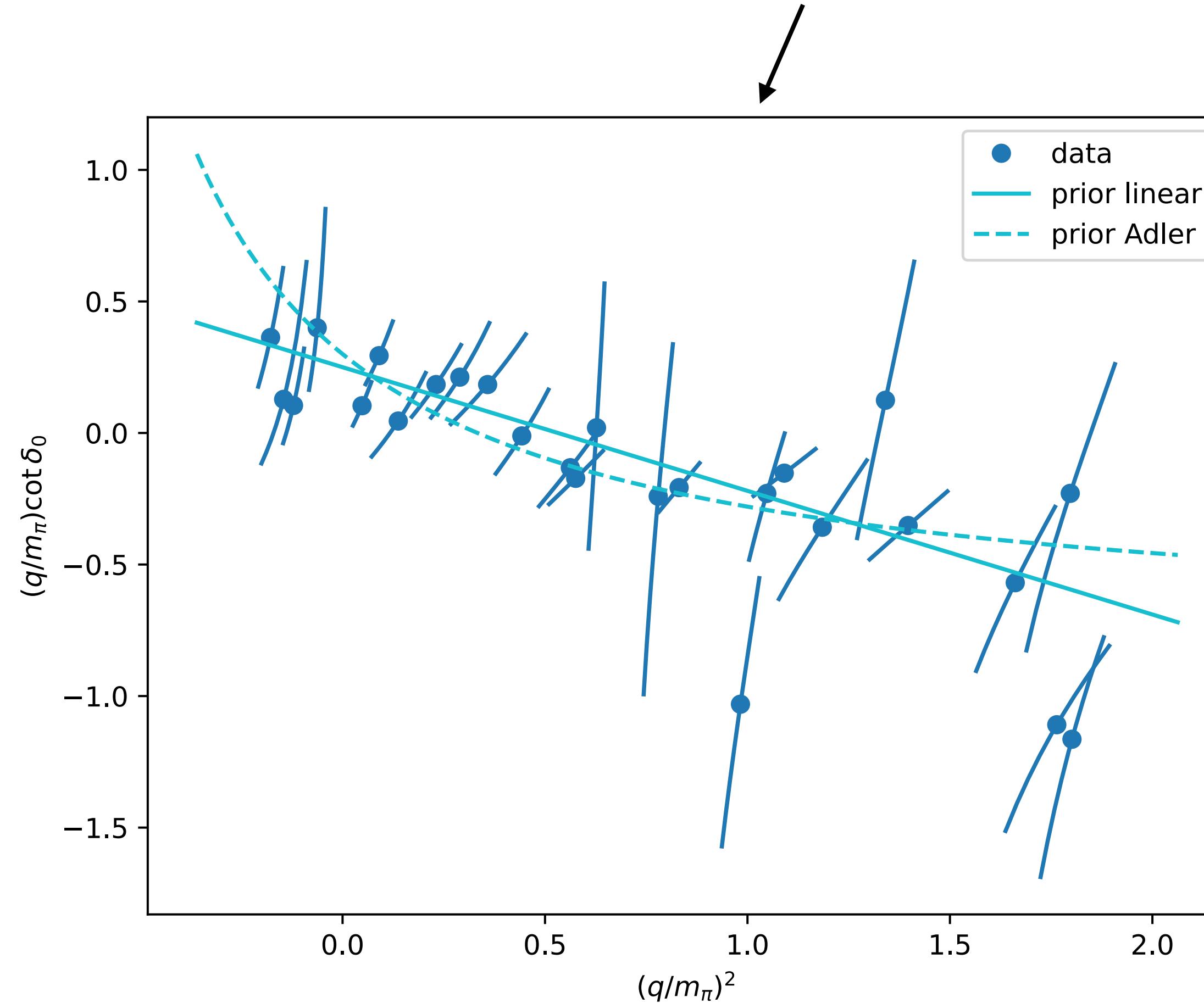
Switches from being a resonance to a virtual bound state  
at  $m_\pi = 283$  MeV

Rodas, Dudek, and Edwards (HadSpec) 2023 [PRD 108, 034513]

Big thanks for sharing data!

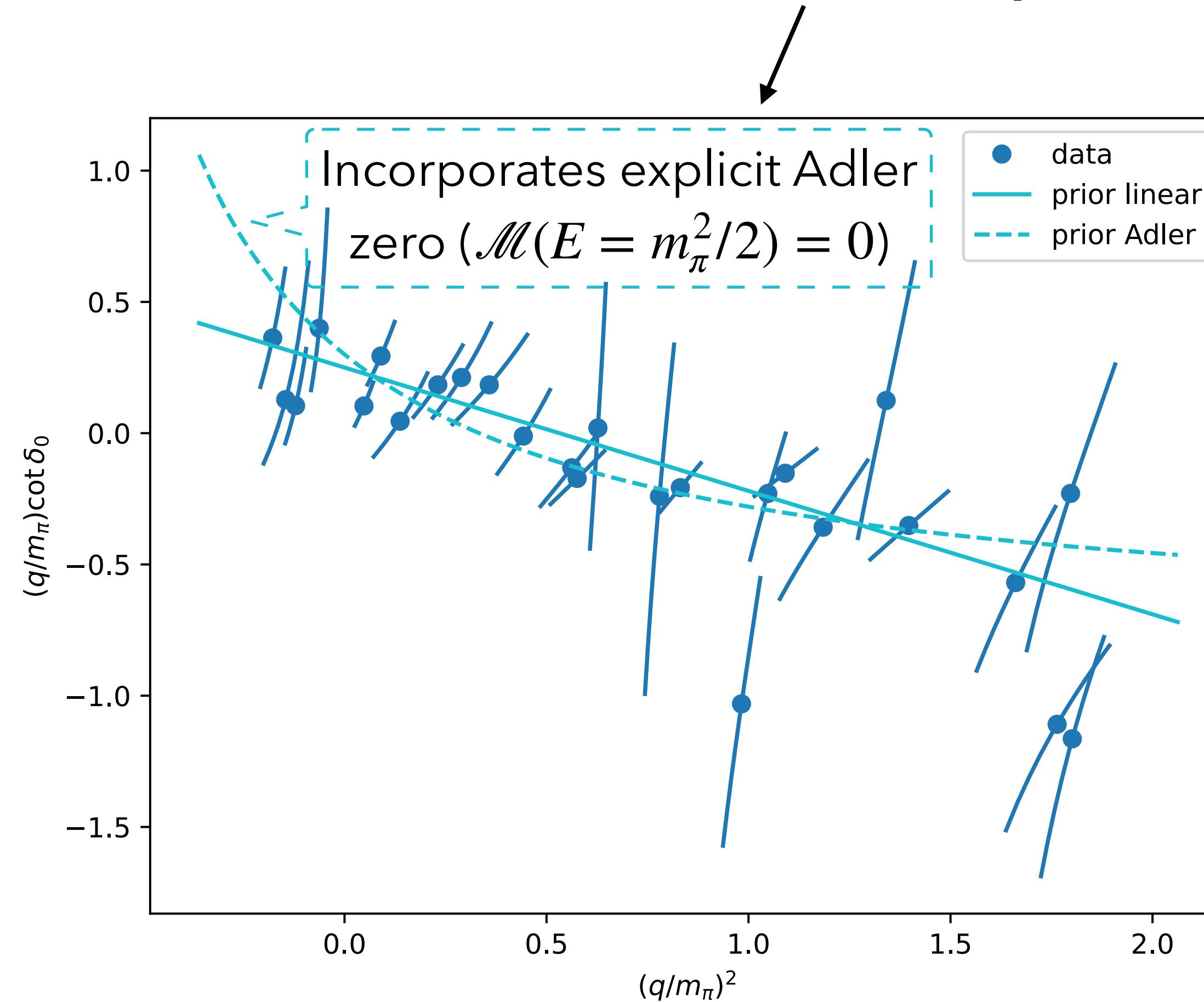
# Isospin-0 $\pi\pi$ system on HadSpec ensembles

Priors and posteriors on the real energy line



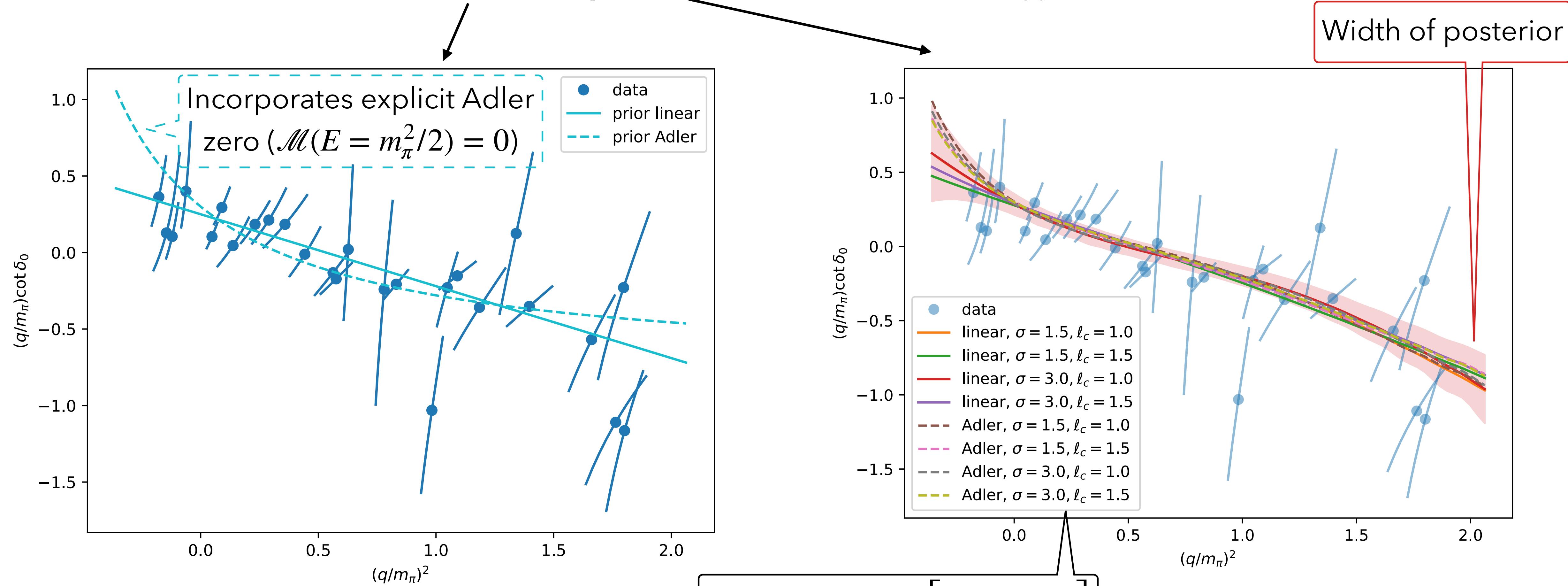
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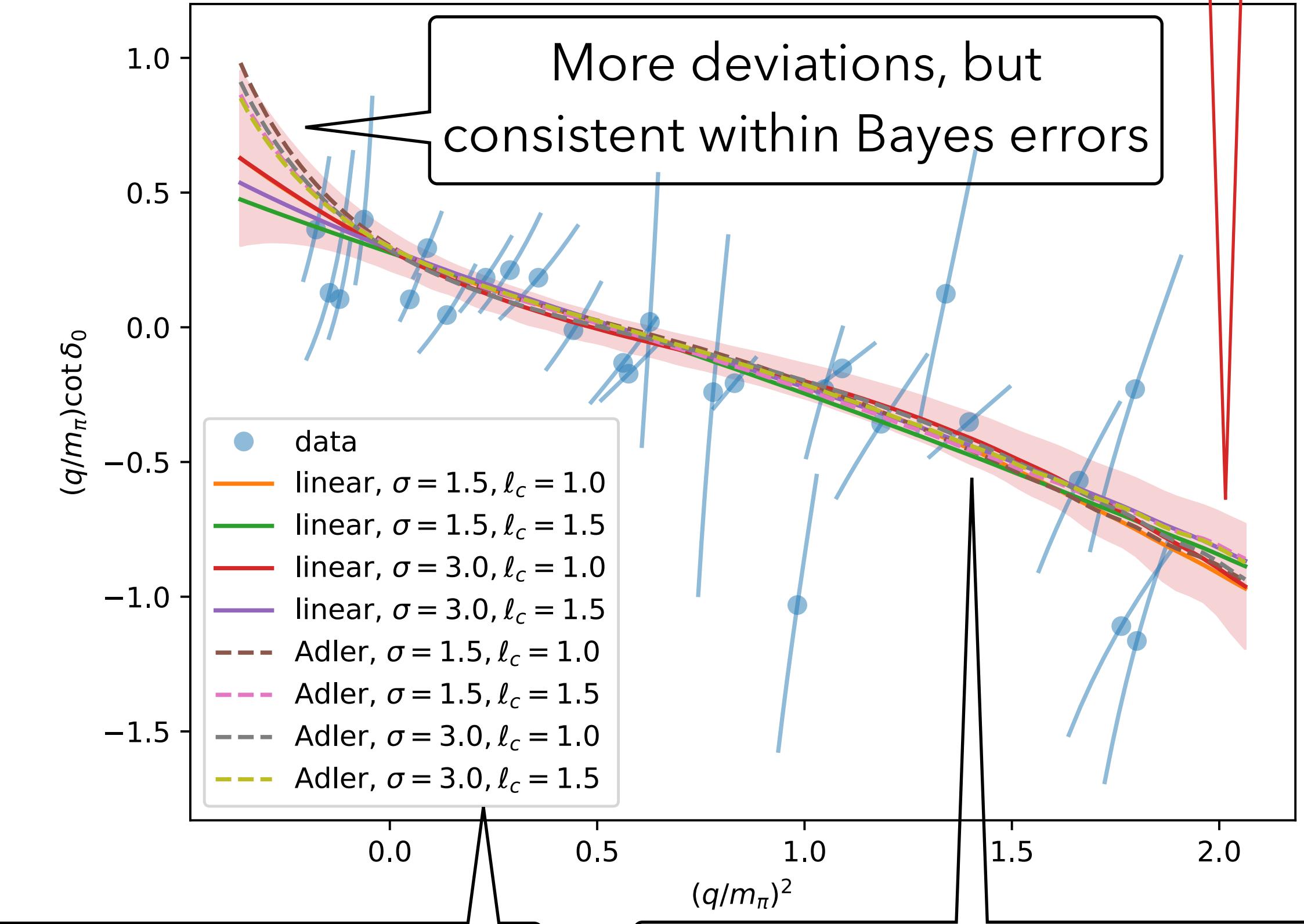
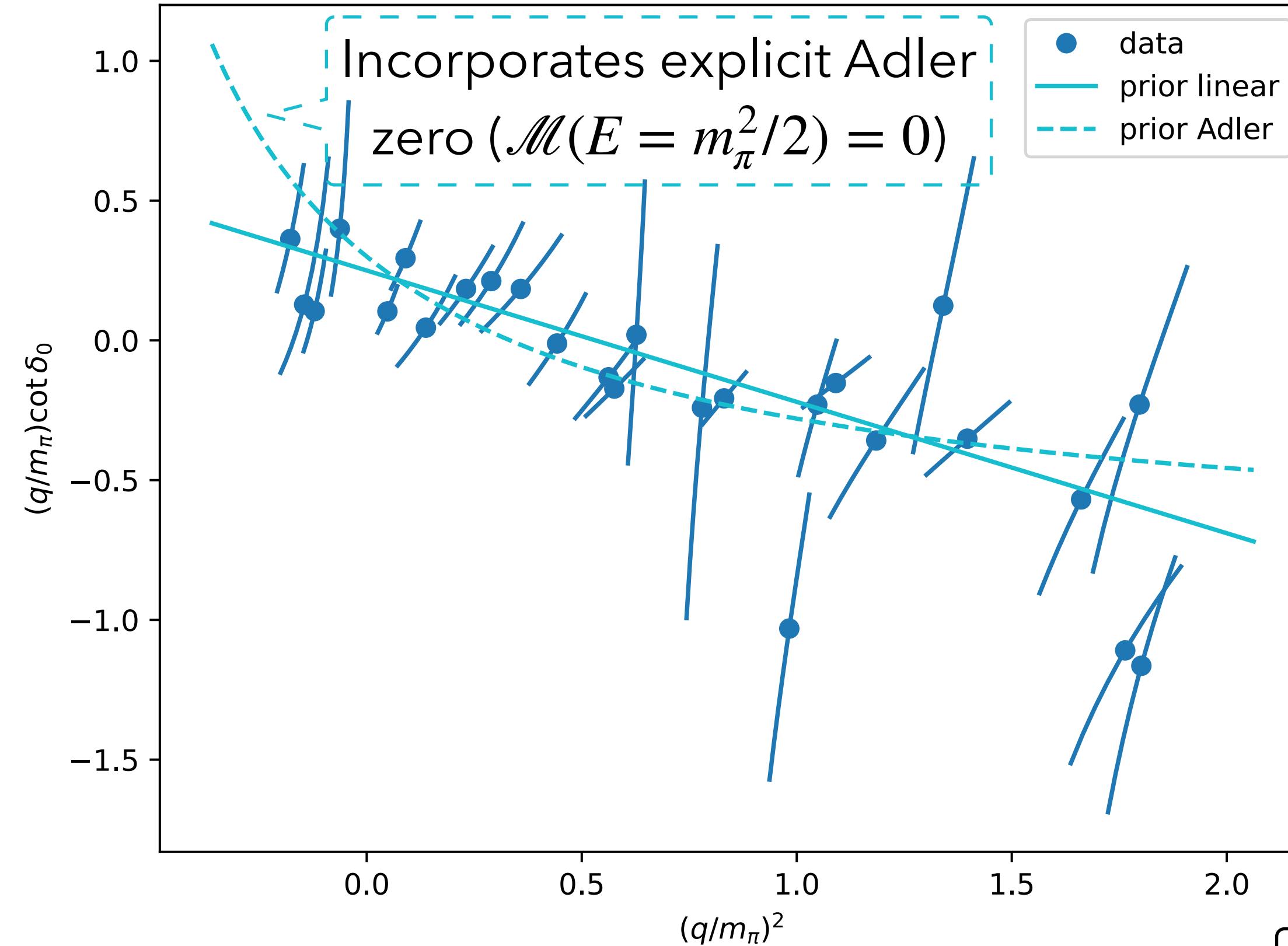
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# Isospin-0 $\pi\pi$ system on HadSpec ensembles

# Priors and posteriors on the real energy line



$$\Sigma_p(E, E') = \sigma^2 \exp \left[ -\frac{(E - E')^2}{2\ell_c^2} \right]$$

Little variation above threshold

# Isospin-0 $\pi\pi$ system on HadSpec ensembles

Prior	Probability of v.b.s.	$\sigma$ pole location		Pole position under assumption of...
		$E_0/m_\pi$	$E_0/m_\pi$ (v.b.s.)	
Linear	0.502(57)	$1.890(90) - 0.08(10)i$	1.88(10)	$1.904(77) - 0.164(86)i$
Adler-zero	0.328(52)	$1.901(97) - 0.099(91)i$	1.86(11)	$1.921(82) - 0.147(73)i$

( $\sigma = 3, \ell_c = 1$ )

Full error: integration + Bayes + Wertevorrat

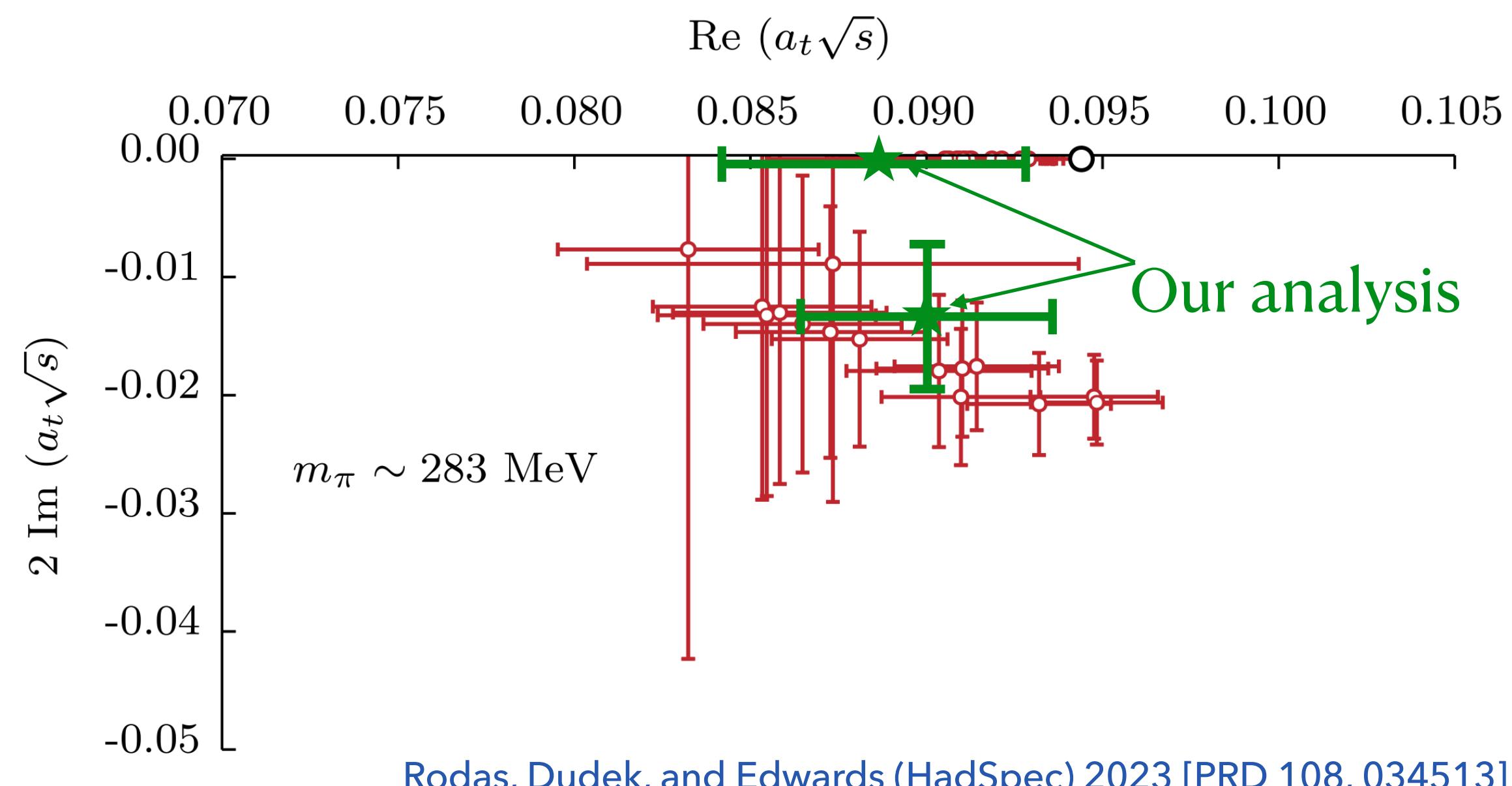
- Pole locations **consistent** within errors between different priors
- Adler-zero prior prefers resonance

# Isospin-0 $\pi\pi$ system on HadSpec ensembles

Prior	Probability of v.b.s.	$\sigma$ pole location		E0/m_pi (res.)
		$E_0/m_\pi$	$E_0/m_\pi$ (v.b.s.)	
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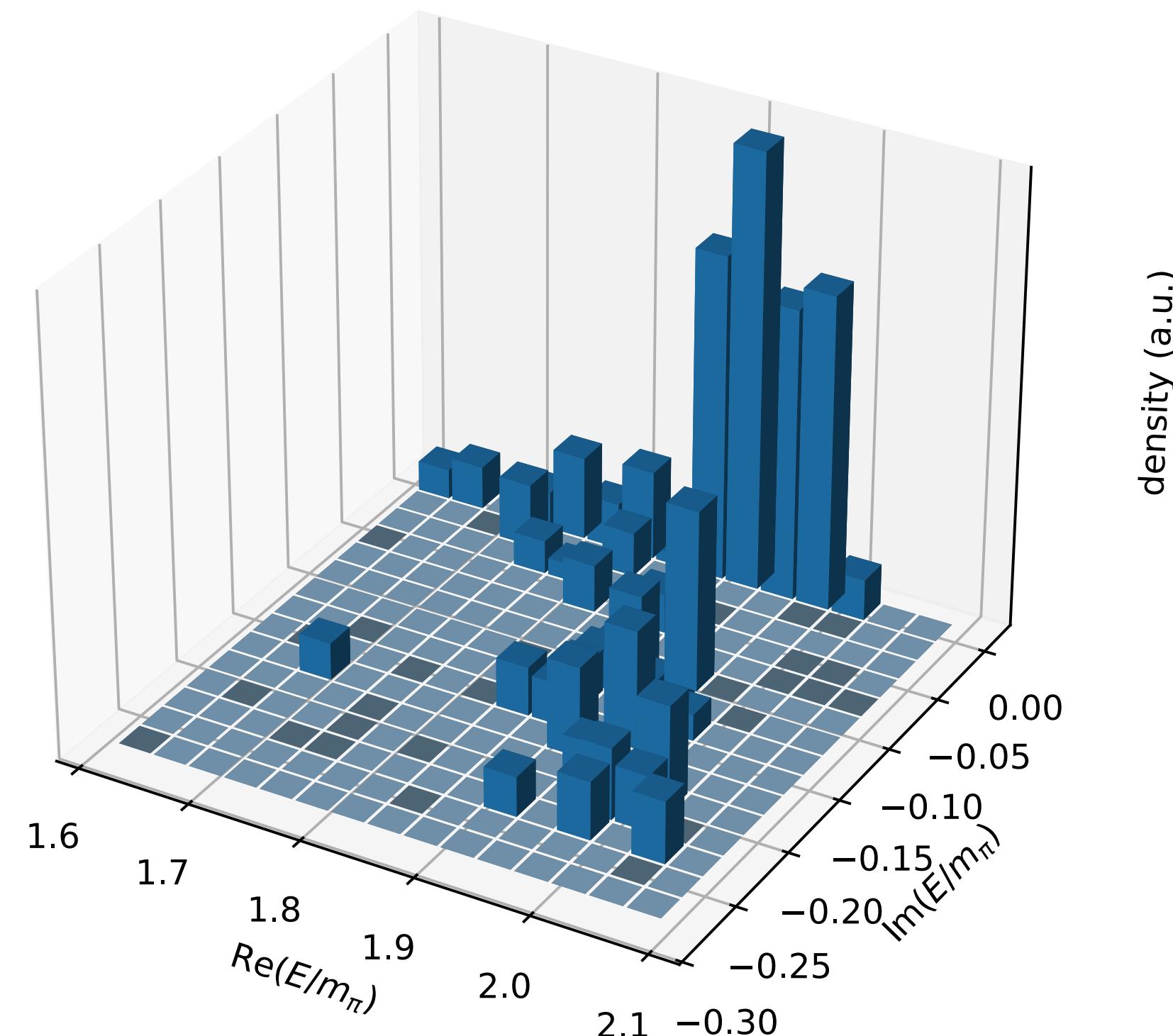
$(\sigma = 3, \ell_c = 1)$

- Pole locations **consistent** within errors between different priors
- Adler-zero prior prefers resonance
- Compares well** with original analysis, but quantifies model dependence

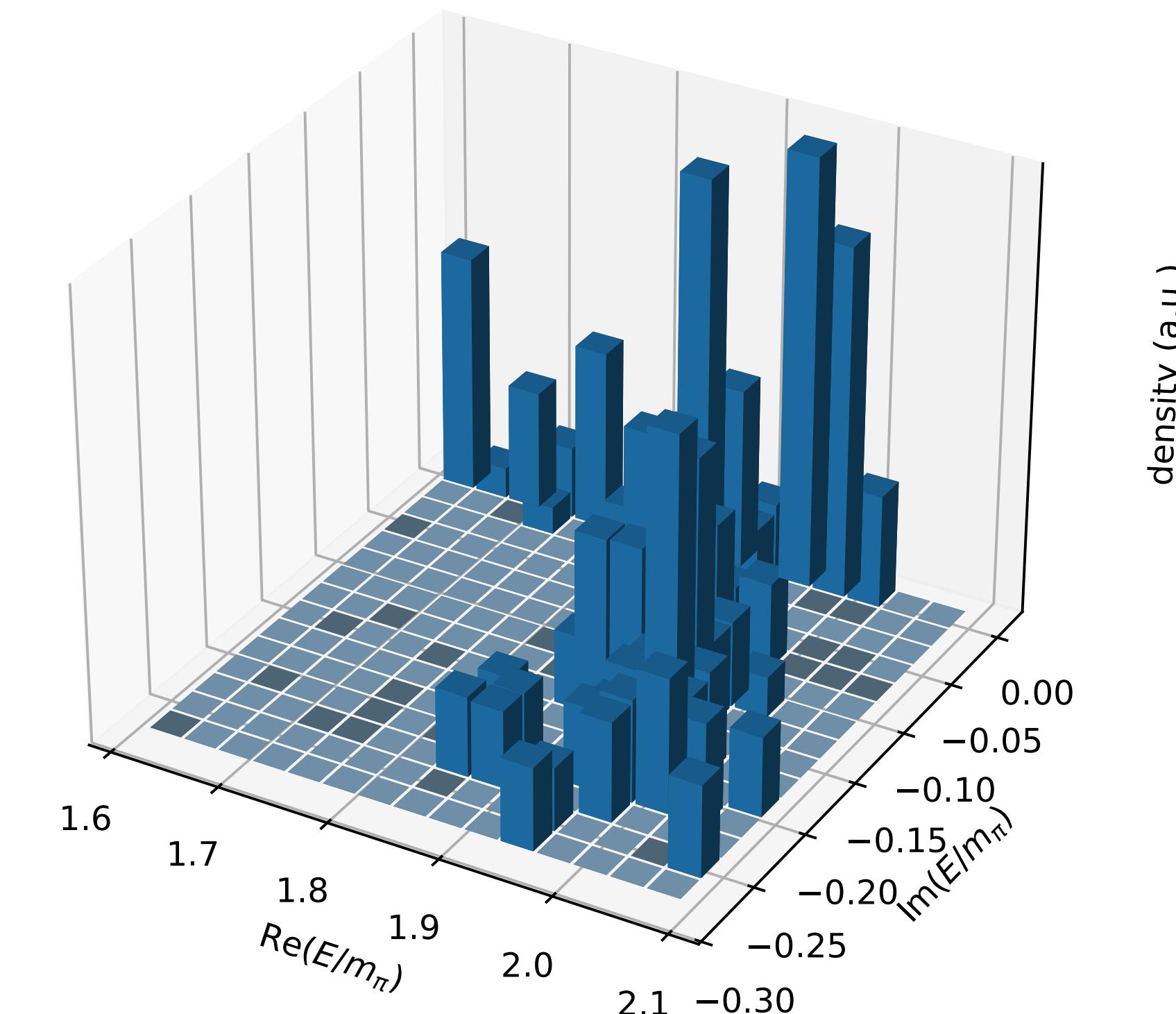


# Isospin-0 $\pi\pi$ system on HadSpec ensembles

$\sigma$  pole location



Linear prior



Adler-zero prior

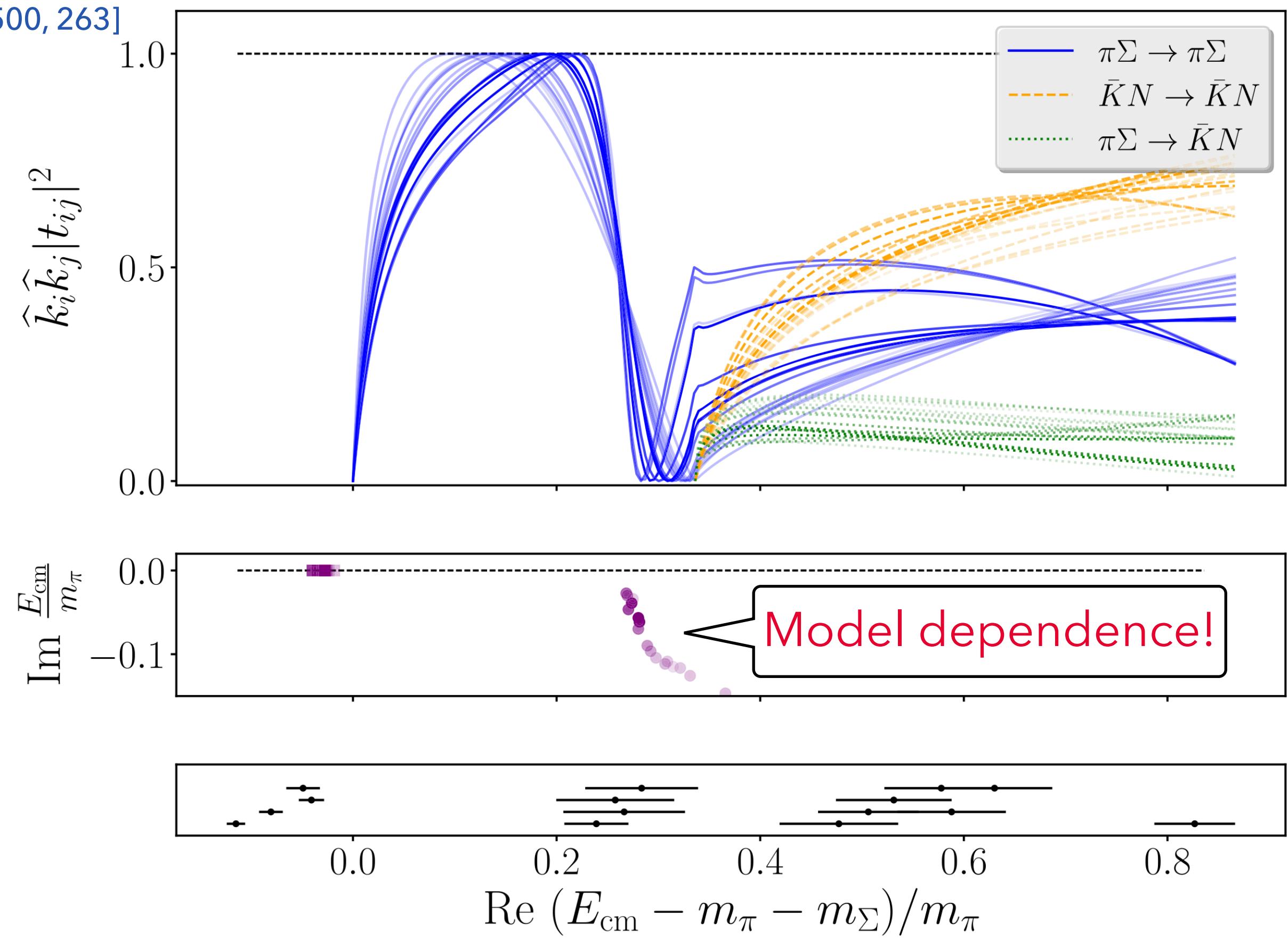
# Coupled-channel $\pi\Sigma - \bar{K}N$ scattering on D200

$\Lambda(1405)$  and  $\Lambda(1380)$  resonances

$M_\pi = 200$  MeV

- One or two poles?
- Lattice QCD favours two poles

Dalitz and Tuan 1959 [Annals Phys. 8, 100]  
Oller and Mei  ner 2001 [Phys. Lett. B 500, 263]  
Anisovich et al. 2020 [EPJA 56, 139]  
Mai 2021 [EPJ Spec. Top. 230, 1593]



Bulava et al. (BaSc) 2024 [PRD 109, 014511]

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$\Lambda(1405)$  and  $\Lambda(1380)$  resonances

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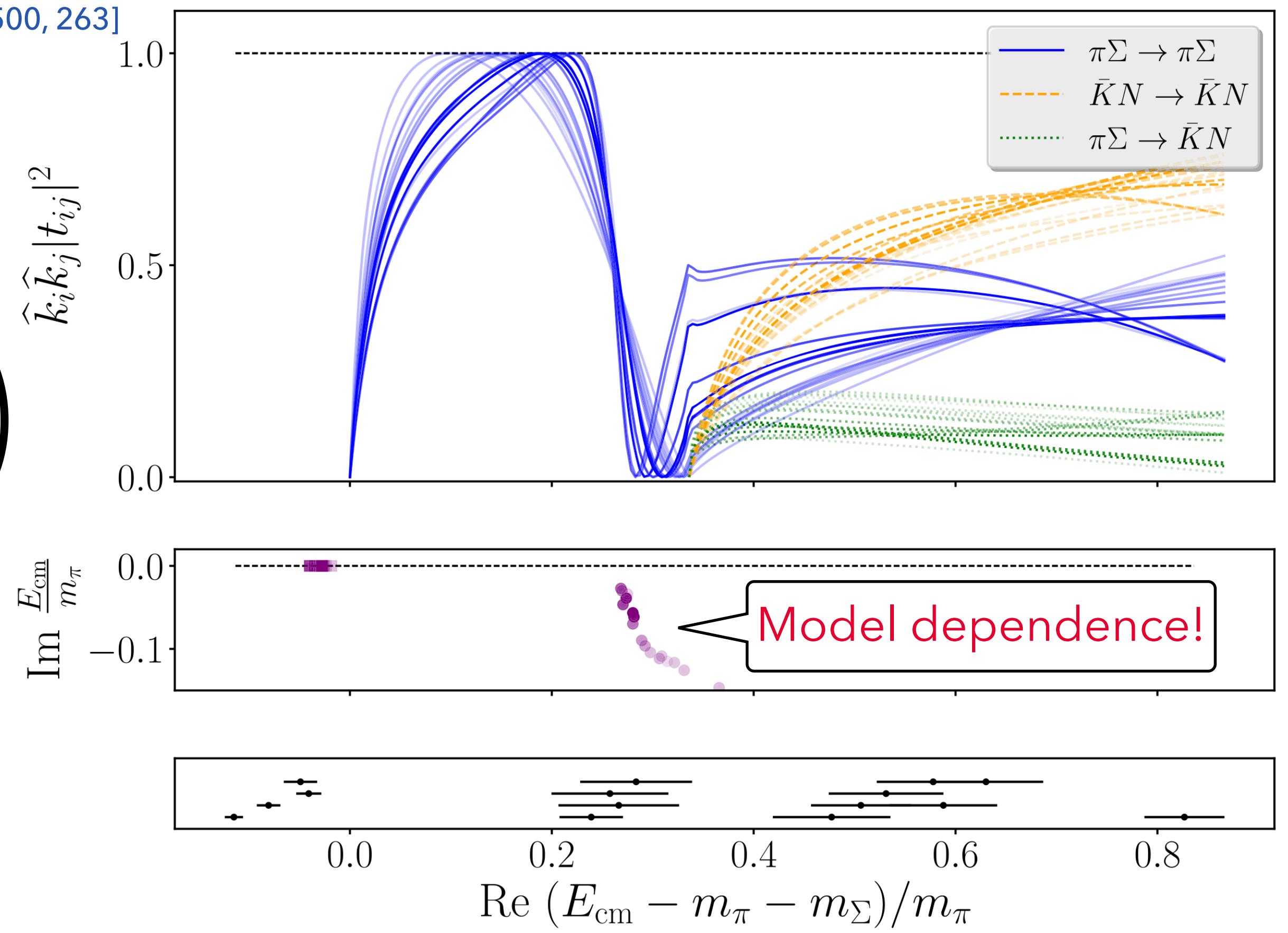
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Dalitz and Tuan 1959 [Annals Phys. 8, 100]  
 Oller and Mei  ner 2001 [Phys. Lett. B 500, 263]  
 Anisovich et al. 2020 [EPJA 56, 139]  
 Mai 2021 [EPJ Spec. Top. 230, 1593]

- Lattice QCD favours two poles

- Coupled-channel scattering:

$$K^{-1} = \begin{pmatrix} K^{-1}(\pi\Sigma \rightarrow \pi\Sigma) & K^{-1}(\pi\Sigma \rightarrow \bar{K}N) \\ K^{-1}(\bar{K}N \rightarrow \pi\Sigma) & K^{-1}(\bar{K}N \rightarrow \bar{K}N) \end{pmatrix}$$



# Coupled-channel $\pi\Sigma - \bar{K}N$ scattering on D200

$\Lambda(1405)$  and  $\Lambda(1380)$  resonances

$M_\pi = 200$  MeV

- One or two poles?

Dalitz and Tuan 1959 [Annals Phys. 8, 100]  
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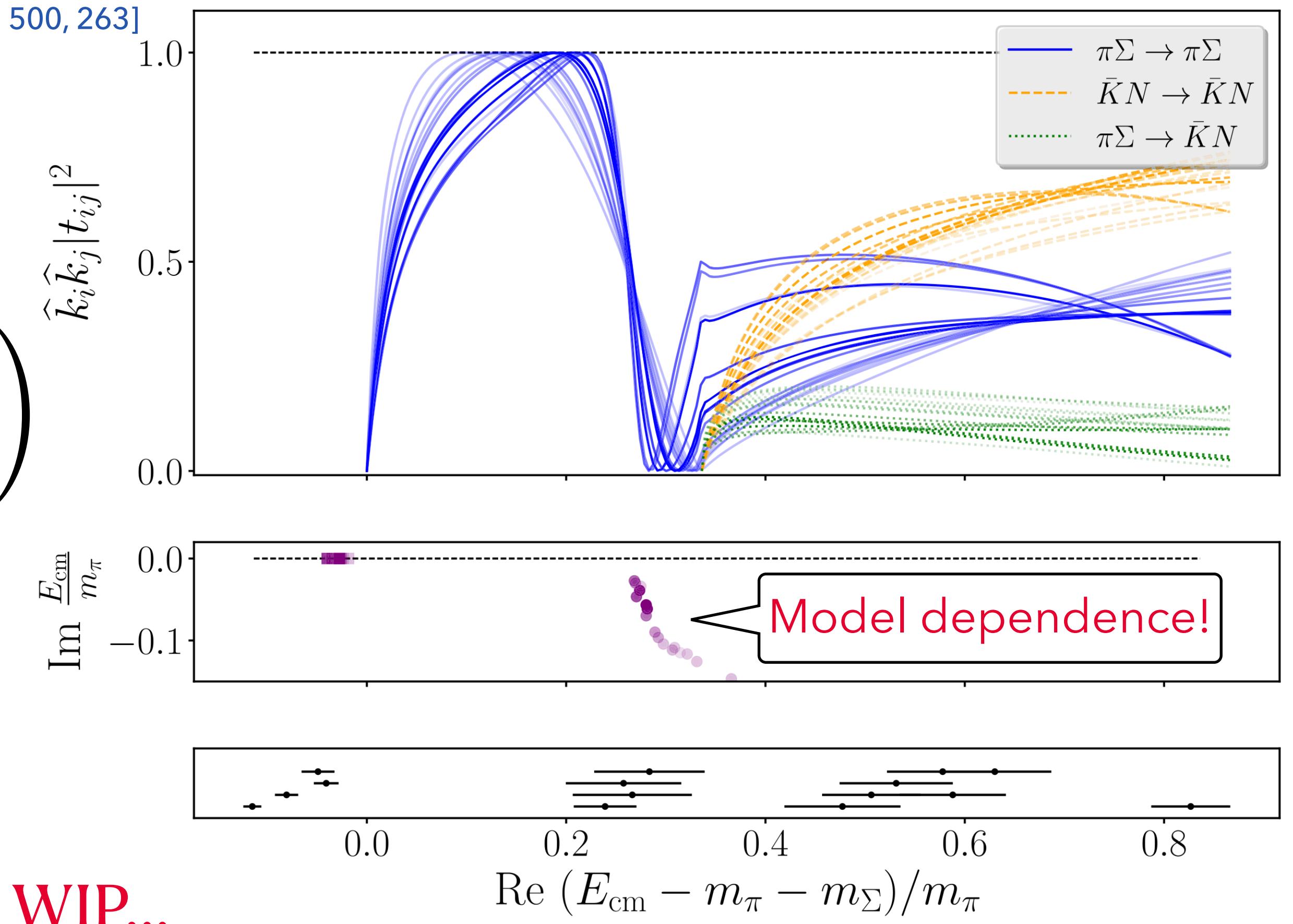
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- Symmetries:

- $(K^{-1})^T = K^{-1} \rightarrow K_{00}, K_{01}, K_{11}$

- $K_{01} \rightarrow -K_{01}$  at each  $E$  separately



WIP...

# Summary and outlook

# Summary and outlook

- Novel analysis procedure for hadron spectroscopy from LQCD using FV formalism  
—> incorporate systematic error without need for explicit choice of parametrisation
- Key ingredients
  - Bayesian analysis of FV spectrum
  - Nevanlinna-Pick interpolation for analytic continuation
- Single-channel case works very well —> promising results in analysing  $\sigma$  resonance
- Coupled-channel analysis ongoing —>  $\Lambda(1405)$
- Future directions: three-hadron systems —> parameterise  $K_3$  in addition to  $K_2$

Hansen and Sharpe 2014 [PRD 90, 116003], 2015 [PRD 92, 114509]  
Mai and Döring 2017 [EPJA 53, 240]

# Backup

# Bayesian analysis

## Sampling from posterior distribution

- Linearisation of QC around  $\mathbf{u}_l$ :  $E_{QC}(\mathbf{u}) - E_l \simeq J(\mathbf{u} - \mathbf{u}_l)$

- Approximate posterior distribution:

$$p^*(\mathbf{u}) = \frac{1}{\sqrt{(2\pi)^{n_p} \det \Sigma^*}} \exp \left[ -\frac{1}{2} (\mathbf{u} - \mathbf{u}^*)(\Sigma^*)^{-1} (\mathbf{u} - \mathbf{u}^*) \right] \text{ with}$$
$$(\Sigma^*)^{-1} = J^T \Sigma_d^{-1} J + \Sigma_p^{-1} \text{ and } \mathbf{u}^* = \Sigma^* \left[ \Sigma_p^{-1} \mathbf{u}_0 + J^T \Sigma_d^{-1} (E_d - E_l + J \mathbf{u}_l) \right]$$

- Reweighting factors:  $w_* = p(\mathbf{u} | E_d, \mathbf{u}_0) / p^*(\mathbf{u} | E_d, \mathbf{u}_0)$

- Probabilities entering Bayes factors:

$$p(E_d | M^{(k)}) = \int d^{n_p} u p(E_d | \mathbf{u}) p(\mathbf{u} | \mathbf{u}_0^{(k)}) = \int d^{n_p} u p(\mathbf{u} | E_d, \mathbf{u}_0^{(k)}) = \langle w_*^{(k)} \rangle_*^{(k)}$$

# Nevanlinna-Pick interpolation

## General framework and employed mappings

- **Schwarz-Christoffel mapping** from rectangle to unit disk:  $\zeta = \frac{s - i}{s + i}$  with

$$s = \operatorname{sn} \left[ \left( z - C + i \frac{H}{2} \right) \frac{2K(k)}{W}, k \right], \text{ where } K(k)/K(\sqrt{1 - k^2}) = W/2H$$

Rectangle height ( $\in \mathbb{R}$ )

Jacobi elliptic function

$$\bullet \text{ Inverse mapping: } z = F \left[ \arcsin \left( i \frac{1 + \zeta}{1 - \zeta} \right), k \right] \frac{W}{2K(k)} - i \frac{H}{2} + C$$

Rectangle width ( $\in \mathbb{R}$ )

Incomplete elliptic integral, 1st kind

$$\bullet \text{ WV center (in disk coords.): } c(\zeta) = \frac{(Q(\zeta)\bar{S}(\zeta) - P(\zeta)\bar{R}(\zeta))}{(|S(\zeta)|^2 - |R(\zeta)|^2)}$$

$$\bullet \text{ WV radius (in disk coords.): } r(\zeta) = \frac{|P(\zeta)S(\zeta) - Q(\zeta)R(\zeta)|}{(|S(\zeta)|^2 - |R(\zeta)|^2)}$$

# Nevanlinna-Pick interpolation

Uncertainty analysis for pole positions

- Use **randomly chosen point within WV** as function estimator,  
 $w(\zeta) = c(\zeta) + \rho r(\zeta) e^{i\theta}$  with **uniform** random numbers  $\rho \in [0,1)$  and  $\theta \in [0,2\pi)$
- Use independent random point(s) for each sample from the posterior (cf. source positions in lattice QCD)
- Effectively, the **expectation value** of the **pole location** takes the form  
$$\langle E_0 \rangle = \frac{1}{\pi Z} \int d^n p u \int_0^1 dr r \int_0^{2\pi} d\theta E_0(\mathbf{u}, r, \theta) p(\mathbf{u} | \mathbf{E}_d, \mathbf{u}_0)$$
- **Uncertainty:**  $(\delta E_0)^2 = \langle E_0^2 \rangle - \langle E_0 \rangle^2$

# Isospin-0 $\pi\pi$ system on X252

$M_\pi = 280 \text{ MeV}$

## Setup and preliminary results

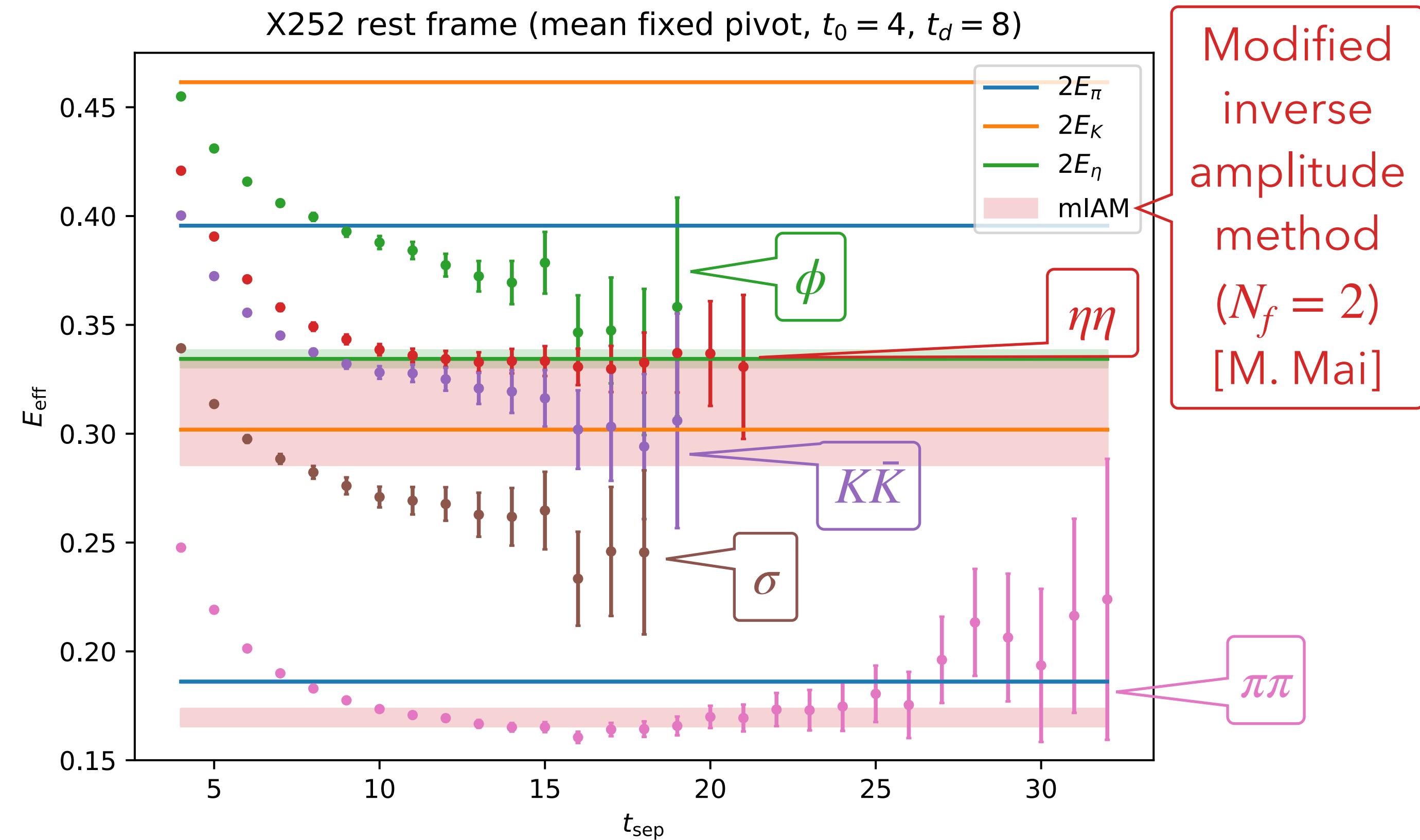
- Exact distillation with 48 eigenvecs, all source times, 1050 configs
- Site-local operators included so far:  
 $\sigma, \phi, \pi\pi, \eta_l\eta_l, \eta_s\eta_s, \eta_l\eta_s, K\bar{K}$
- Outlook
  - Singly-displaced operators
  - Moving frames
  - More ensembles: X253, D251, J250

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WIP in collaboration with H. Alharazin, J. Bulava,  
F. Romero-López, and the BaSc collaboration