



Universität
Münster

Gravitational waves from (low-scale) cosmic strings

Based on: Kai Schmitz & Tobias Schröder, Phys. Rev. D **110** (2024) 6, 063549, 2405.10937.

Kai Schmitz

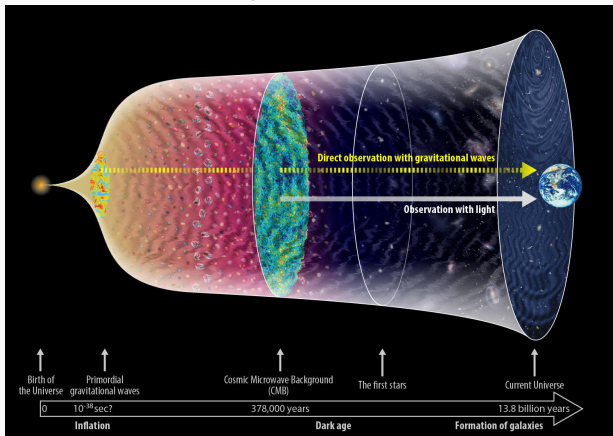
University of Münster | Münster, Germany

The Dawn of Gravitational Wave Cosmology

Benasque Science Center | Benasque, Spain | 30 April 2025

Gravitational-wave echo from the Big Bang

[National Astronomical Observatory of Japan, gwpo.nao.ac.jp]



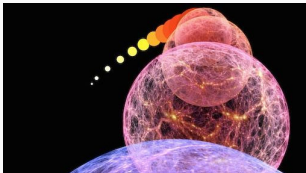
Primordial gravitational waves (GWs): Chance to peek behind the veil of the CMB

- Probe cosmology of the primordial Universe at **very early times**
- Probe particle physics at **extremely high energies** → **New physics!?**

Beyond-the-Standard-Model (BSM) options

① Inflationary tensor perturbations

- Accelerated expansion before the Hot Big Bang
- Complementarity: GWs + CMB observations

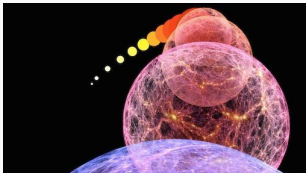


Abbreviations: GW: gravitational wave; CMB: cosmic microwave background; QFT: quantum field theory; EW: electroweak; QCD: quantum chromodynamics; PBH: primordial black hole; GUT: grand unified theory

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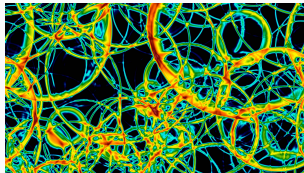
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② Cosmological phase transition

- First-order transition in the QFT vacuum structure
- Complementarity: GWs + EW / QCD / dark sector

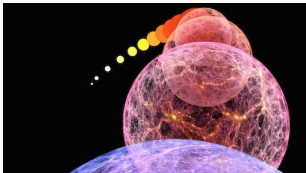


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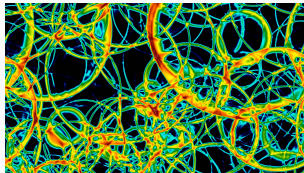
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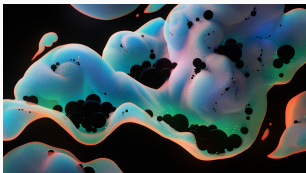
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- Overdensities that emit GWs and collapse to PBHs
- Complementarity: GWs + primordial black holes

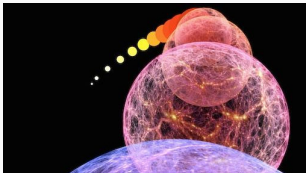


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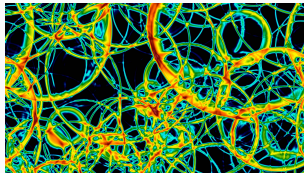
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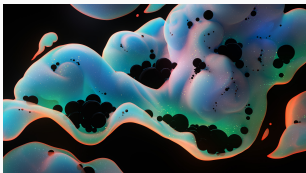
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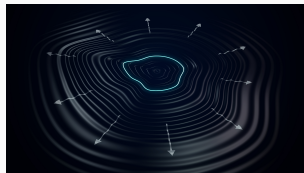
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④ Cosmic defects

- Phase transition remnants preserving the old vacuum
- Complementarity: GWs + grand unified theories



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GWs from cosmic strings

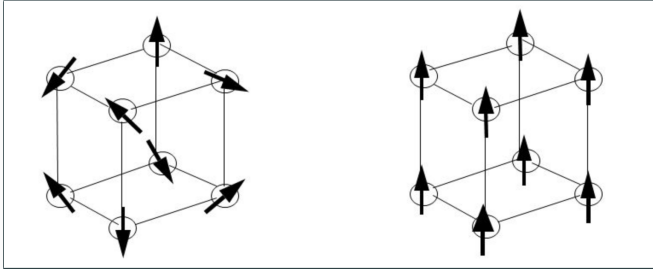
Low-scale cosmic strings

From VOS to BOS

Conclusions

Magnetic domains in a ferromagnet

[wikimedia.org]

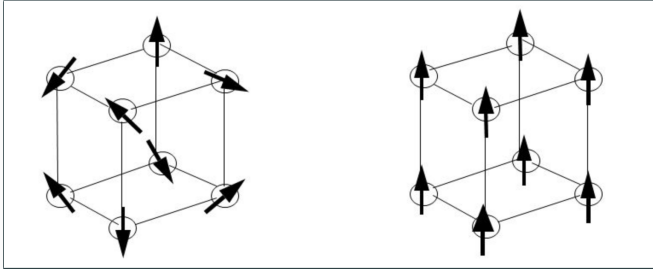


Magnetization in a ferromagnet

- Phase transition at the Curie temperature: paramagnet \rightarrow ferromagnet
- Magnetic dipoles align spontaneously due to exchange interaction
- Translation and rotation invariance spontaneously broken
- Magnetic domains, regions of uniform magnetization, separated by domain walls
- Domain walls are stable, unless an external force (magnetic field) is applied

Magnetic domains in a ferromagnet

[wikimedia.org]



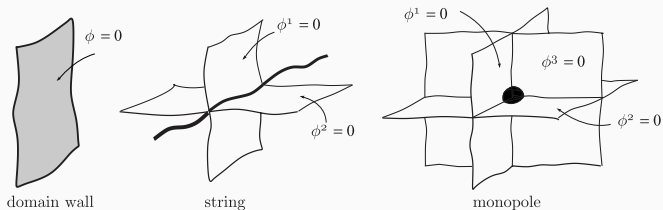
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Similar phenomenology after phase transitions in the early Universe!

Topological defects in the early Universe

[Viatcheslav Mukhanov: Physical Foundations of Cosmology, Cambridge University Press (2005)]

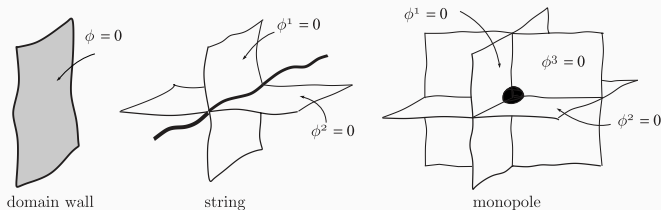


Consider spontaneous symmetry breaking in an N -dimensional scalar field space:

$$V(\Phi) = \frac{\lambda}{4} (\Phi^2 - v^2)^2, \quad \Phi = \frac{1}{\sqrt{N}} (\phi_1, \phi_2, \dots, \phi_N)^T$$

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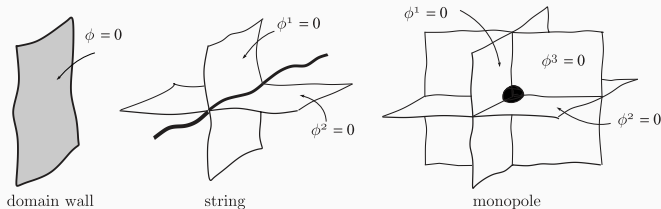
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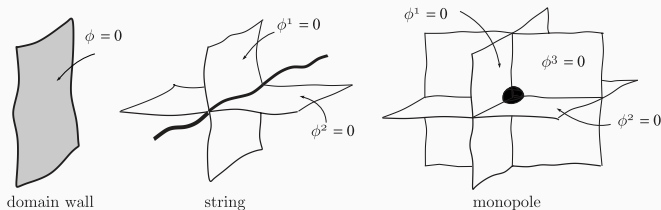
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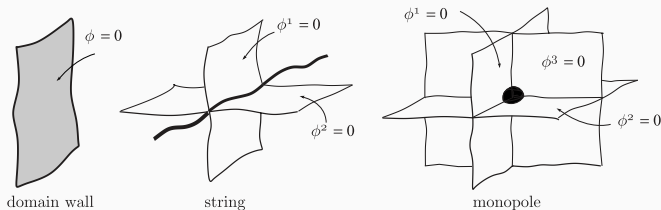
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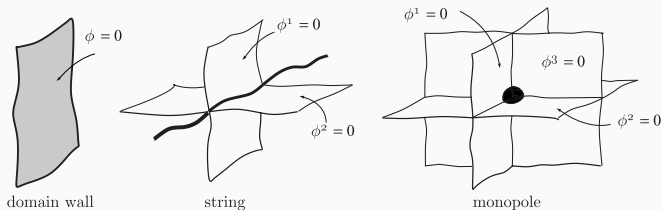
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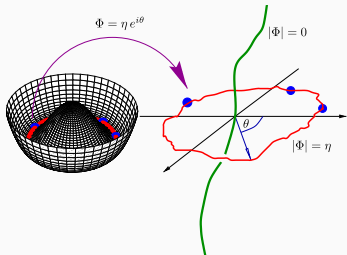
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- Solitonic solutions of classical equations of motion for the gauge and Higgs fields
- Formal description in terms of topology of vacuum manifold $\mathcal{M}(\Phi) \rightarrow$ stability
- In addition, whole zoo of composite defects, non-topological defects, etc.

Cosmic strings

[Ringeval: 1005.4842]

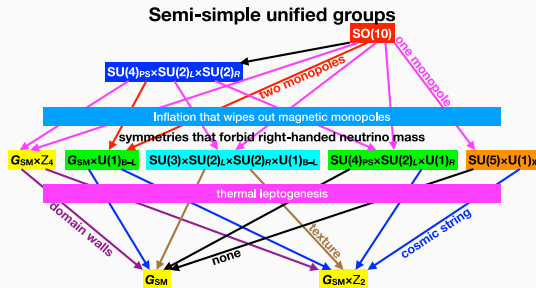


- **Topological defects** after spontaneous symmetry breaking, s.t. $\pi_1(\mathcal{M})$ nontrivial
- For instance, breaking of global / local $U(1)$; symmetry restored at string cores
- Condensed matter: Magnetic field vortices (quantum vortices) in a superconductor

Relevant parameters

- $G\mu$: String tension = energy per unit length, in units of $G = 1/M_{\text{P}}^2$
- α : Size of string loops at time of formation, in units of the horizon $d_h \sim t \sim H^{-1}$

Cosmic strings in grand unified theories



Cosmic-string tension: Controlled by energy scale of spontaneous symmetry breaking

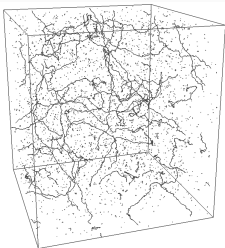
$$\mu \sim 2\pi v^2, \quad G\mu \sim 4 \times 10^{-8} \left(\frac{v}{10^{15} \text{ GeV}} \right)^2$$

Interesting possibilities

$$v \sim \Lambda_{\text{GUT}} \sim 10^{15 \cdots 16} \text{ GeV}, \quad v \sim \Lambda_{\text{intermediate}} \sim 10^{9 \cdots 10} \text{ GeV}$$

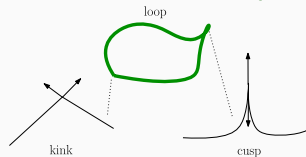
Gravitational waves from cosmic strings

[Allen, Martins, Shellard: ctc.cam.ac.uk/outreach]



Infinitely long strings and string loops;
scaling regime: $\rho_{\text{CS}} \propto \rho_{\text{crit}} \propto H^2$

[Gouttenoire, Servant, Simakachorn: 1912.02569]

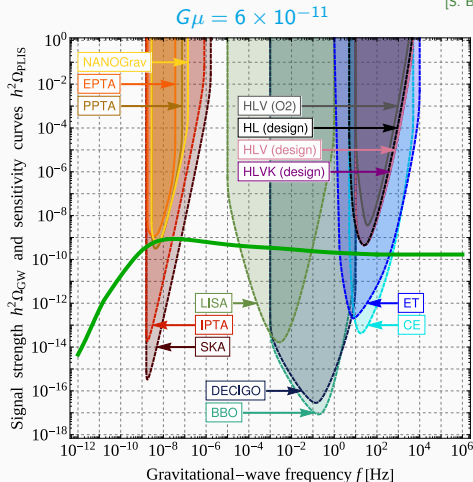


Gravitational waves from

- Cusps
- Kinks
- Kink–kink collisions

-
- **Nambu–Goto strings:** Infinitely thin, particle emission irrelevant at late times
 - **Abelian-Higgs strings:** Short-lived loops, decay into massive particles

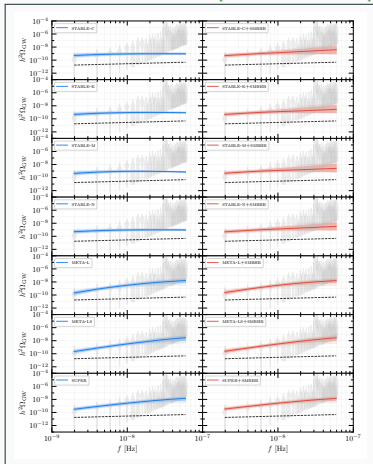
[Vachaspati, Vilenkin: PRD 31 (1985) 3052] [LISA Cosmology Working Group, Auclair et al.: 1909.00819]



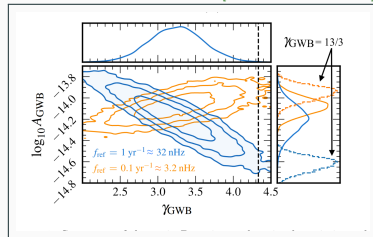
Broadband signal

- Reflects scaling regime, GW emission during radiation and matter domination
- Interesting target for future GW experiments. Source of the **PTA** signal?

[NANOGrav: 2306.16219]



[NANOGrav: 2306.16213]

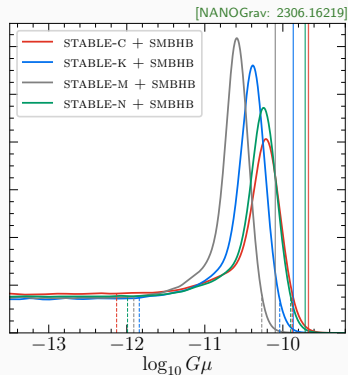
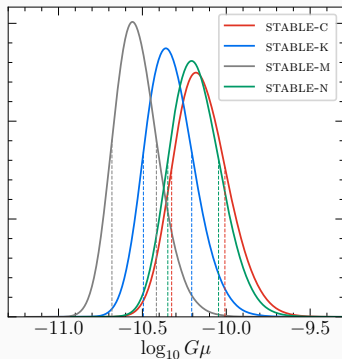


Recent PTA results

- Evidence for nHz GWB signal
- Stable cosmic strings do not yield a good fit (spectrum too flat)
- Alternatives doing a better job: metastable strings, superstrings

PTA upper limit on the tension of stable Nambu–Goto strings

$$G\mu \lesssim 10^{-10} \quad \longleftrightarrow \quad v \lesssim 5 \times 10^{13} \text{ GeV}$$

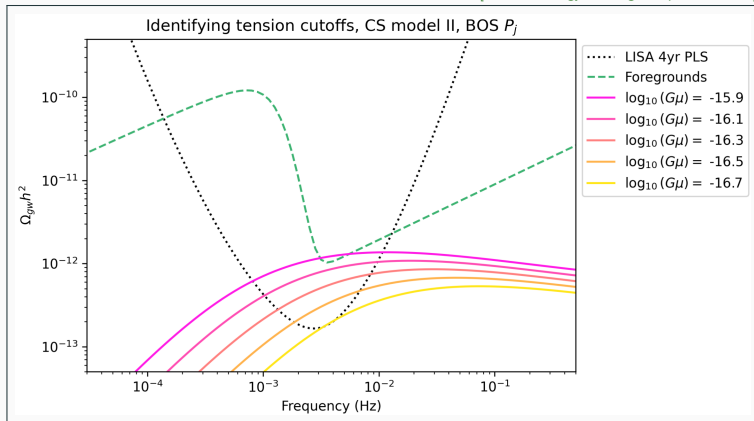


[NANOGrav: 2306.16219]

Different models

- GW emission dominated by cusps (c), kinks (k), fundamental mode (m); numerical result (n)
- GWs from cosmic strings only or in combination with GWs from supermassive BH binaries

[LISA Cosmology Working Group: 2405.03740]



Expected sensitivity: $G\mu \sim 10^{-(16 \dots 17)}$ \longleftrightarrow $\nu \sim \text{few} \times 10^{10} \text{ GeV}$

- GW signal from cosmic strings competes with galactic and extragalactic foregrounds

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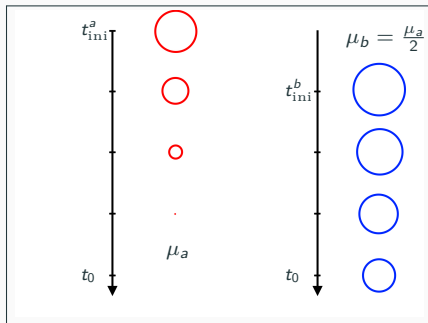
Two facts about string loops

- **Shrink** because of GW emission

$$\frac{d\ell}{dt} = -\Gamma G\mu, \quad \Gamma \simeq 50$$

- **Characteristic length at birth**

$$\ell_* = 2\alpha t_*, \quad \alpha \simeq 0.05$$



Phenomenology at low string tension

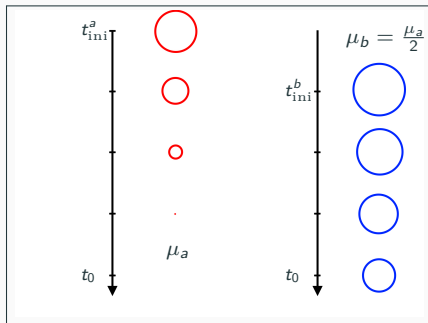
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Loop length decreases linearly in time between birth and today

$$\ell(t_0) = \ell_* - \Gamma G\mu(t_0 - t_*), \quad t_* \in [t_{\text{ini}}, t_0]$$

Computation of GW signal only valid starting from some early initial time t_{ini}

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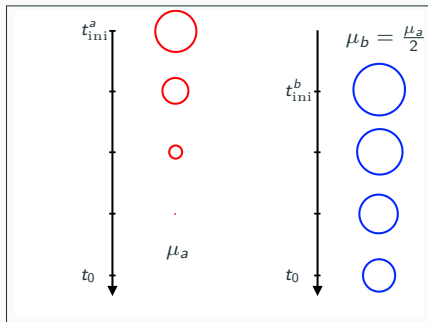
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Computation of GW signal only valid starting from some early initial time t_{ini}

Observation: For late t_{ini} and small enough $G\mu$, no loop ever reaches zero length!

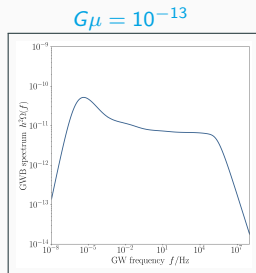
Shortest loops = earliest loops today. Length: $\ell_{min} = \ell_*(t_{ini}) - \Gamma G\mu(t_0 - t_{ini})$

Sharp cutoff frequency

Present-day frequencies of GWs emitted by strings

$$f = \frac{a(t)}{a_0} \frac{2k}{\ell(t)}$$

$2k/\ell(t)$ Frequency at emission
 $\ell(t)$ Loop length at emission
 k Mode number ($k = 1, 2, \dots$)
 $a(t)/a_0$ Cosmological redshift factor

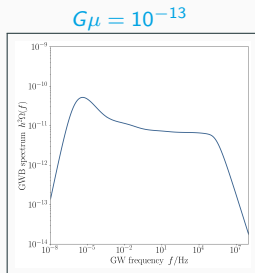


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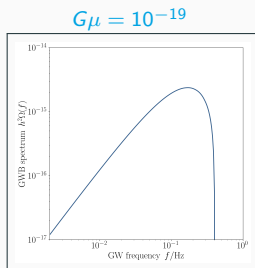
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| k | Mode number ($k = 1, 2, \dots$) |
| $a(t)/a_0$ | Cosmological redshift factor |



Minimal length ℓ_{\min} implies frequency cutoff

- Focus on fundamental mode ($k = 1$) for now
- Shortest loops today: minimal length, minimal redshift \rightarrow highest possible frequency

$$f_{\text{cut}} = \frac{2}{\ell_{\min}} = \frac{2}{2\alpha t_{\text{ini}} - \Gamma G\mu (t_0 - t_{\text{ini}})}$$



Cutoff frequency is positive and finite if: $2\alpha t_{\text{ini}} > \Gamma G\mu (t_0 - t_{\text{ini}}) \approx \Gamma G\mu t_0$

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Four reasonable options

① $t_{\text{ini}} = t_{\text{form}}$ Network formation

$$\rho_{\text{tot}} = 3H^2 M_{\text{Pl}}^2 \sim \mu^2$$

Initial loop length

Cutoff frequency is positive and finite if: $2\alpha t_{\text{ini}} > \Gamma G\mu (t_0 - t_{\text{ini}}) \approx \Gamma G\mu t_0$

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② $t_{\text{ini}} = t_{\text{fric}}$ End of friction regime

$$\beta T^3 / \mu \sim 2H$$

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| ② | $t_{\text{ini}} = t_{\text{fric}}$ | End of friction regime | $\beta T^3 / \mu \sim 2H$ |
| ③ | $t_{\text{ini}} = t_{\text{kink}}$ | Particles from kinks subdominant | $P_{\text{kink}} \sim \frac{N_k \mu^{1/2}}{\ell} \sim \Gamma G\mu^2$ |

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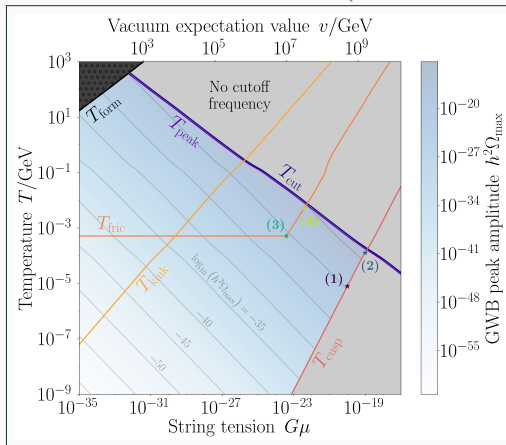
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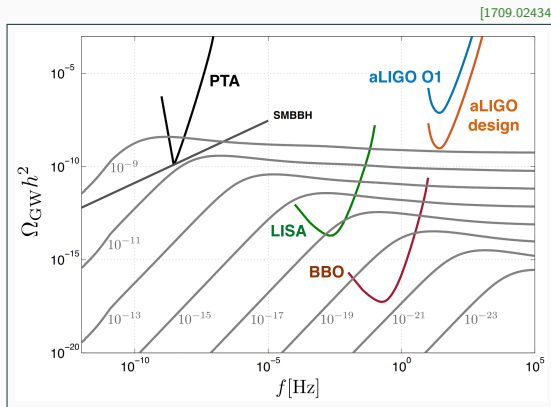
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| ② | $t_{\text{ini}} = t_{\text{fric}}$ | End of friction regime | $\beta T^3 / \mu \sim 2H$ |
| ③ | $t_{\text{ini}} = t_{\text{kink}}$ | Particles from kinks subdominant | $P_{\text{kink}} \sim \frac{N_k \mu^{1/2}}{\ell} \sim \Gamma G\mu^2$ |
| ④ | $t_{\text{ini}} = t_{\text{cusp}}$ | Particles from cusps subdominant | $P_{\text{cusp}} \sim \frac{N_c \mu^{3/4}}{\ell^{1/2}} \sim \Gamma G\mu^2$ |

[KS, Schröder: 2405.10937]



- $G\mu$ and T_{ini} values resulting in a cutoff frequency in the $k = 1$ GWB spectrum
- Hierarchy of temperature scales for $G\mu \sim 10^{-20}$: $T_{\text{cusp}} \ll T_{\text{fric}} \ll T_{\text{kink}} \ll T_{\text{form}}$

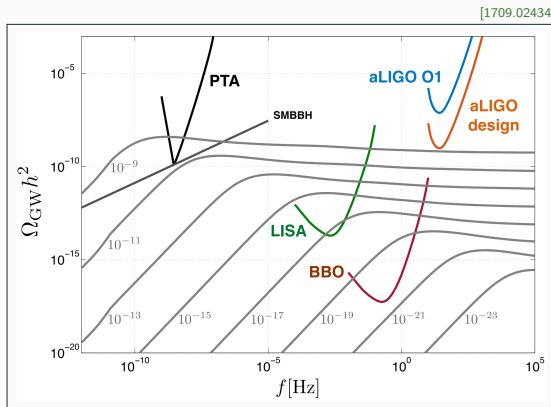
Earlier results in the literature



What's new?

$$\Omega_{\text{GW}}(f) = \frac{8\pi}{3H_0^2} (G\mu)^2 \sum_{k=1}^{k_{\text{max}}} \frac{\Gamma}{H_{k_{\text{max}}}^q} \frac{1}{k^q} \frac{2k}{f} \int_{t_{\text{ini}}}^{t_0} dt \left(\frac{a(t)}{a_0} \right)^5 n \left(\frac{2k}{f} \frac{a(t)}{a_0}, t \right)$$

Earlier results in the literature



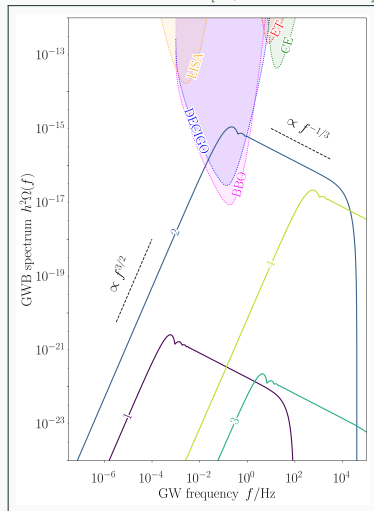
What's new? Nothing! Our Ω_{GW} is standard, but we do not integrate from $t_{\text{ini}} = 0$

$$\Omega_{\text{GW}}(f) = \frac{8\pi}{3H_0^2} (G\mu)^2 \sum_{k=1}^{k_{\text{max}}} \frac{\Gamma}{H_{k_{\text{max}}}^q} \frac{1}{k^q} \frac{2k}{f} \int_{t_{\text{ini}}}^{t_0} dt \left(\frac{a(t)}{a_0} \right)^5 n \left(\frac{2k}{f} \frac{a(t)}{a_0}, t \right)$$

Example spectra

Numerical spectra based on VOS loop
number densities

[KS, Schröder: 2405.10937]

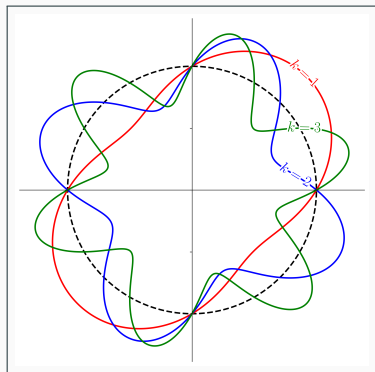


| | $\log_{10}(G\mu)$ | $\log_{10}(T_{\text{ini}}/\text{GeV})$ | |
|---|-------------------|--|---------------------|
| 1 | -20.0 | -5.1 | (T_{cusp}) |
| 2 | -19.0 | -3.9 | (T_{cusp}) |
| 3 | -23.4 | -3.3 | (T_{fric}) |
| 4 | -22.4 | -2.3 | (T_{fric}) |

- No fine-tuning required
- Sweet spot where signal even observable by BBO and DECIGO
- Power-law behavior can be understood analytically

Challenge: Subtraction of galactic and extragalactic foregrounds. Impossible?

Summation of oscillation modes



Total GWB spectrum

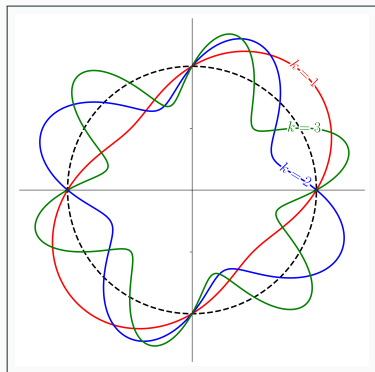
$$\Omega_{\text{GW}}(f) = \sum_{k=1}^{k_{\text{max}}} \frac{\Omega_{\text{GW}}^{(1)}(f/k)}{H_{k_{\text{max}}}^q k^q}$$

can be written in terms of $\Omega_{\text{GW}}^{(1)}$, i.e., spectrum from the fundamental mode

Simple approximation for $\Omega_{\text{GW}}^{(1)}$

$$\Omega_{\text{GW}}^{(1)} \approx \Theta(f_{\text{cut}} - f) \mathcal{A} f^{3/2}$$

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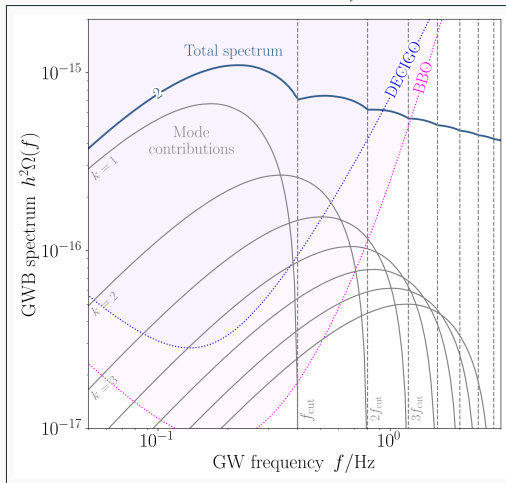
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Power-law behavior of the total GWB spectrum at low and high frequencies

$$h^2 \Omega_{\text{low}} \propto f^{3/2}, \quad h^2 \Omega_{\text{high}} \propto \frac{1}{q + 1/2} \left(\frac{f_{\text{cut}}}{f} \right)^{q-1}$$

Features in the GW spectrum

[KS, Schröder: 2405.10937]



Novel features in the spectrum: Series of peaks and dips at integer multiples of f_{cut} on top of a broken power law ($f^{3/2} \rightarrow f^{-1/3}$) → **Clear target for GW experiments**

Loop number and energy densities

Assumption: $G\mu$ is so low that no loop has fully decayed yet because of GW emission

Consequence: All loops produced since t_{ini} still exist in our present Universe

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Present-day loop number density

$$N(t_0) \approx \int_0^\infty d\ell n_{\text{RM}}(\ell, t_0) \sim \frac{50}{\text{kpc}^3} \left(\frac{10^2 \text{ s}}{t_{\text{ini}}} \right)^{3/2}$$

Present-day loop energy density

$$h^2 \Omega(t_0) \approx \frac{1}{\rho_{\text{crit}}} \int_0^\infty d\ell \mu \ell n_{\text{RM}}(\ell, t) \simeq 10^{-13} \left(\frac{G\mu}{10^{-19}} \right) \left(\frac{10^2 \text{ s}}{t_{\text{ini}}} \right)^{1/2}$$

- Cosmologically harmless
- Signatures from nearby loops? **Microlensing, GW bursts?**

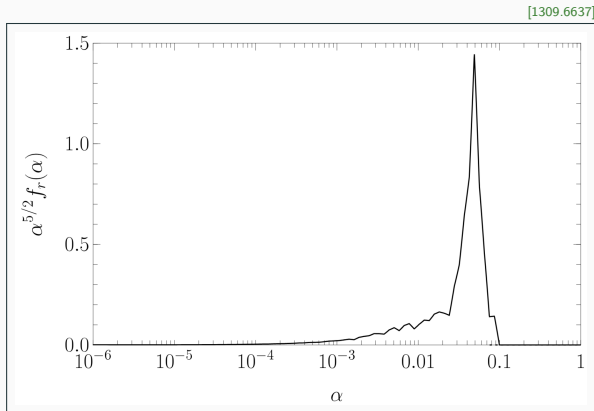
GWs from cosmic strings

Low-scale cosmic strings

From VOS to BOS

Conclusions

Distribution of initial loop lengths



Sharp spectral features follow from the assumption of a unique initial loop length, $\ell_* \simeq 2\alpha t_*$ with $\alpha \simeq 0.05$. But initial loop length deviates from perfect delta peak.

Loop number density in terms of the loop production function

$$n(\ell, t) = \int_{t_{\text{ini}}}^t dt' f(\ell', t') \left(\frac{a(t')}{a(t)} \right)^3, \quad \ell' = \ell + \Gamma G\mu (t - t')$$

Standard choice in the velocity-dependent one-scale (**VOS**) model

$$f(\ell, t) = \frac{\mathcal{F}C_r}{2\alpha t^4} \delta(\ell - 2\alpha t), \quad \alpha \simeq 0.05$$

Numerical simulations by Blanco-Pillado, Olum, and Shlaer (**BOS**) better described by

$$f(\ell, t) = \frac{A_r}{\sqrt{2\pi} \sigma \ell^5 t^{5/2}} \exp \left[-\frac{1}{2\sigma^2} \left(\ln \left(\frac{\ell}{2t} \right) - \nu \right)^2 \right], \quad \nu \simeq -3.0, \quad \sigma \simeq 0.14$$

New number densities

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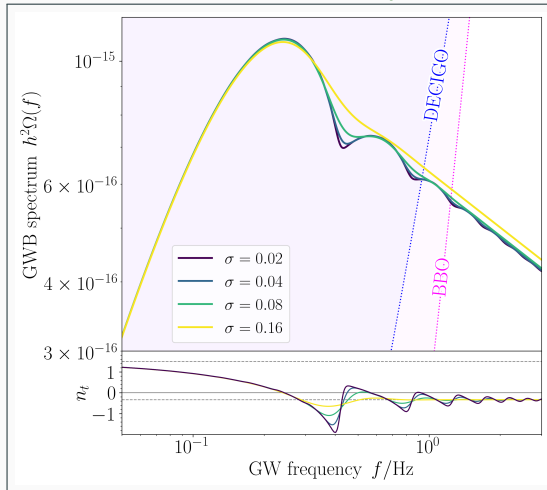
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New loop number densities providing a better description of the BOS results

$$n_{\text{RR}}(\ell, t) \approx \frac{A_r(\text{erf}_{t_{\text{ini}}} - \text{erf}_t)/2}{t^{3/2}(\ell + \Gamma G\mu t)^{5/2}}, \quad n_{\text{RM}}(\ell, t) \approx \left(\frac{a_{\text{eq}}}{a(t)} \right)^3 \frac{A_r(\text{erf}_{t_{\text{ini}}} - \text{erf}_{t_{\text{eq}}})/2}{t_{\text{eq}}^{3/2}(\ell + \Gamma G\mu t)^{5/2}}$$

Smeared GW spectrum

[KS, Schröder: 2405.10937]



Series of peaks and dips washed out for broader distributions of initial loop lengths

Still, even for broad distributions, oscillations in the index n_t may remain detectable

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Summary and outlook

Low-scale cosmic strings

$$G\mu \sim 10^{-33} \dots 10^{-19} \quad \longleftrightarrow \quad v \sim 10^2 \text{ GeV} \dots 10^9 \text{ GeV}$$

Initial time $t_{\text{ini}} \neq 0$: Loop production no longer impeded by thermal friction, GW emission from loops no longer subdominant to particle emission from cusps and kinks,

$$\Omega_{\text{GW}} = \frac{16\pi}{3H_0^2} (G\mu)^2 \sum_{k=1}^{k_{\text{max}}} \frac{kP_k}{f} \int_{t_{\text{ini}}}^{t_0} dt [\dots]$$

- No loop produced at $t \geq t_{\text{ini}}$ ever shrinks to zero length \rightarrow microlensing, bursts?
- Sharp frequency cutoff in $k = 1$ GWB spectrum, series of peaks and dips in Ω_{GW}

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- Nonscaling models where particle emission occurs whenever $\ell \leq \ell_{\text{crit}}$
- Model building: Cosmic strings $v \sim 10^9 \text{ GeV}$, GWs from phase transition?

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Stay tuned! Thanks a lot for your attention