

# The Nonlinear Dynamics of Axion Inflation on the Lattice

#### Nicolás Loayza Romero

In collaboration with: Daniel G. Figueroa, Joanes Lizarraga, Ander Urio & Jon Urrestilla Based on *Phys.Rev.D* 111 (2025) 6 2411.16368 [astro-ph.CO]

The Dawn of Gravitational Wave Cosmology - Benasque

VniverSitat de València

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## Inflation



# Slow roll period $\phi \rightarrow \phi + c$

#### Shift symmetry

#### We want to protect from being spoiled



#### Symmetry protected flatness of potential









#### Slow roll períod

 $\phi \rightarrow \phi + c$ 

#### Shift symmetry

## **Axion Inflation**

 $\mathscr{L} \supset \frac{\psi}{\kappa} F\tilde{F}$ 

#### $A_{\mu} \equiv U(1)$ gauge field

[K. Freese, J. A. Frieman, A. V. Olinto (PRL 65,3233 1990)] ...















[M. M. Anber, L. Sorbo (0908.4089)] [J. Cook, L. Sorbo (1109.0022)] [N. Barnaby, E. Pajer, M. Peloso (1110.3327)]





# $\mathscr{L} \supset \frac{\phi}{1} F \tilde{F}$ $\xi \equiv \frac{|\dot{\phi}|}{2H\Lambda} \qquad \checkmark$ $\left(\partial_{\tau}^{2} + k^{2} + \operatorname{sign}(\dot{\phi})\frac{2k\xi}{|\tau|}\right)\mathscr{A}^{+}(\tau, \mathbf{k}) = 0$

[M. M. Anber, L. Sorbo (0908.4089)] [J. Cook, L. Sorbo (1109.0022)] [N. Barnaby, E. Pajer, M. Peloso (1110.3327)]





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## Axion Inflation phenomenology

#### Chiral GWBs





## **Axion Inflation** phenomenology

#### Chiral GWBs

#### Primordial Black Holes







## **Axion Inflation** phenomenology

#### Chiral GWBs

#### Primordial Black Holes

#### Non-gaussianities

- [N. Barnaby, M. Peloso (1011.1500)]
- [A. Linde, S. Mooji, E. Pajer (1212.1693)]
- [E. Bugaev, P. Klimai (1312.7435)]
- [N. Bartolo et al (LISA) (1610.06481)]
- [J. G. Bellido, M. Peloso, C. Unal (1610.03763)] .....













#### Dynamical equations in FLRW

$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}$$
$$\dot{\vec{E}} = -H\vec{E} - \frac{1}{a^2}\vec{\nabla}\times\vec{B} - \frac{\alpha_{\Lambda}}{am_p}\left(\dot{\phi}\vec{B} - \vec{\nabla}\phi\times\vec{E}\right)$$
$$\vec{\nabla}\cdot\vec{E} = -\frac{\alpha_{\Lambda}}{am_p}\vec{\nabla}\phi\cdot\vec{B}$$







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$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E} \cdot \vec{E}$$
$$\dot{\vec{E}} = -H\vec{E} - \frac{1}{a^2}\vec{\nabla}\times\vec{B} - \frac{\alpha_{\Lambda}}{am_p}\left(\phi\vec{B} - \vec{\nabla}\phi\times\vec{D}\right)$$
$$\vec{\nabla}\cdot\vec{E} = -\frac{\alpha_{\Lambda}}{am_p}\vec{\nabla}\phi\cdot\vec{B}$$
$$\ddot{a} = -\frac{a}{3m_p^2}\left(2\rho_{\rm K} - \rho_{\rm V} + \rho_{\rm EM}\right)$$
$$H^2 = \frac{1}{3m_p^2}\left(\rho_{\rm K} + \rho_{\rm G} + \rho_{\rm V} + \rho_{\rm EM}\right)$$





#### Dynamical equations in FLRW

$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}$$
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#### Dynamical equations in FLRW

$$\begin{split} \ddot{\phi} &= -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E} \\ \dot{\vec{E}} &= -H\vec{E} - \frac{1}{a^2}\vec{\nabla}\times\vec{B} - \frac{\alpha_{\Lambda}}{am_p}\left(\dot{\phi}\vec{B} - \vec{\nabla}\phi\times\vec{E}\right) \\ \vec{\nabla}\cdot\vec{E} &= -\frac{\alpha_{\Lambda}}{am_p}\vec{\nabla}\phi\cdot\vec{B} \\ \ddot{a} &= -\frac{a}{3m_p^2}\left(2\rho_{\rm K} - \rho_{\rm V} + \rho_{\rm EM}\right) \\ H^2 &= \frac{1}{3m_p^2}\left(\rho_{\rm K} + \rho_{\rm G} + \rho_{\rm V} + \rho_{\rm EM}\right) \end{split}$$



## Linear Regime

$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2} \sqrt[3]{\phi} - m^2 \phi + \frac{\alpha_N}{a^3 m_p}$$
$$\dot{\vec{E}} = -H\vec{E} - \frac{1}{a^2} \vec{\nabla} \times \vec{B} - \frac{\alpha_N}{am_p} \left( \dot{\phi}\vec{B} - \vec{\nabla}\phi \right)$$
$$\ddot{a} = -\frac{a}{3m_p^2} (2\rho_{\rm K} - \rho_{\rm V} + \phi_{\rm V})$$



$$\left(\partial_{\tau}^{2} + k^{2} + \operatorname{sign}(\dot{\phi})\frac{2k\xi}{|\tau|}\right) \mathscr{A}^{+}(\tau, \mathbf{k}) = 0$$
$$\mathscr{A}^{+}(\tau, \mathbf{k}) \simeq \frac{1}{k|\tau| \ll 2\xi} \frac{1}{\sqrt{2k}} \left(\frac{k}{2\xi aH}\right)^{1/4} e^{\pi\xi - 2\sqrt{2\xi k/(aH)}}$$



## **Homogeneous Backreaction Regime**

$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2} \sqrt[3]{\phi} - m^2 \phi + \frac{\alpha_{\Lambda}}{a^3 m_p} \langle \vec{B} \rangle$$
$$\dot{\vec{E}} = -H\vec{E} - \frac{1}{a^2} \vec{\nabla} \times \vec{B} - \frac{\alpha_{\Lambda}}{a m_p} \left( \dot{\phi} \vec{B} - \vec{\nabla} \phi \right)$$
$$\ddot{a} = -\frac{a}{3m_p^2} (2\rho_{\rm K} - \rho_{\rm V} + \rho_{\rm EM})$$

- Integrate  $\langle \vec{E} \cdot \vec{B} \rangle$ every time step

[Cheng, Lee, Ng (1508.00251)] [Notari, Tywoniuk, (1608.06223)] [Dall'Agata, González-Martín, Papageorgiu, Peloso (1912.09950)] [Domcke, Guidetti, Welling, Westphal (2002.02952)] ...

> -Gradient Expansion Formalism

[Sobol, Gorbar, Vilchinskii (1907.10443)] [Gorbar, Schmitz, Sobol, Vilchinskii (2109.01651)] [Durrer, Sobol, Vilchinskii (2303.04583)] [Durrer, von Eckardstein, Garg, Schmitz, Sobol (2404.19694)]



$$\begin{split} \ddot{\phi} &= -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}\cdot\vec{B} \\ \dot{\vec{E}} &= -H\vec{E} - \frac{1}{a^2}\vec{\nabla}\times\vec{B} - \frac{\alpha_{\Lambda}}{am_p}\left(\dot{\phi}\vec{B} - \vec{\nabla}\phi\times\vec{E}\right) \\ \vec{\nabla}\cdot\vec{E} &= -\frac{\alpha_{\Lambda}}{am_p}\vec{\nabla}\phi\cdot\vec{B} \\ \ddot{a} &= -\frac{a}{3m_p^2}\left(2\rho_{\rm K} - \rho_{\rm V} + \rho_{\rm EM}\right) \\ H^2 &= \frac{1}{3m_p^2}\left(\rho_{\rm K} + \rho_{\rm G} + \rho_{\rm V} + \rho_{\rm EM}\right) \end{split}$$

[Figueroa, Lizarraga, Urio, Urrestilla (2303.17436)] [Figueroa, Lizarraga, NL, Urio, Urrestilla (2411.16368)]

$$\begin{split} \ddot{\phi} &= -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}\cdot\vec{B} \\ \dot{\vec{E}} &= -H\vec{E} - \frac{1}{a^2}\vec{\nabla}\times\vec{B} - \frac{\alpha_{\Lambda}}{am_p}\left(\dot{\phi}\vec{B}\right)\cdot\vec{\nabla}\phi\times\vec{E} \\ \vec{\nabla}\cdot\vec{E} &= -\frac{\alpha_{\Lambda}}{am_p}\vec{\nabla}\phi\cdot\vec{B} \\ \ddot{a} &= -\frac{a}{3m_p^2}(2\rho_{\rm K} - \rho_{\rm V} + \rho_{\rm EM}) \\ H^2 &= \frac{1}{3m_p^2}(\rho_{\rm K} + \rho_{\rm G} + \rho_{\rm V} + \rho_{\rm EM}) \end{split}$$

# - local description of $\overrightarrow{E} \cdot \overrightarrow{B}$ and $\phi \overrightarrow{B}$

[Figueroa, Lizarraga, Urio, Urrestilla (2303.17436)] [Figueroa, Lizarraga, NL, Urio, Urrestilla (2411.16368)]

$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}\cdot\vec{B}$$
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- local description of  $\overrightarrow{E} \cdot \overrightarrow{B}$  and  $\phi \overrightarrow{B}$ 



$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}\cdot\vec{B}$$
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- local description of  $\overrightarrow{E} \cdot \overrightarrow{B}$  and  $\phi \overrightarrow{B}$ 

Turn on gradients  $\nabla^2 \phi$  and  $\overrightarrow{\nabla} \phi$ 





CosmoLattice

[D. G. Figueroa, A. Florio, F. Torrenti & W. Valkenburg (2006.15122)] [D. G. Figueroa, A. Florio, F. Torrenti & W. Valkenburg (2102.01031)]



$$\ddot{\phi} = -3H\dot{\phi} + \frac{1}{a^2}\nabla^2\phi - m^2\phi + \frac{\alpha_{\Lambda}}{a^3m_p}\vec{E}\cdot\vec{B}$$
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Alternative lattice descriptions

# - local description of $\overrightarrow{E} \cdot \overrightarrow{B}$ and $\phi \overrightarrow{B}$

Turn on gradients  $\nabla^2 \phi$  and  $\nabla \phi \times \vec{E}$ 





CosmoLattice

[D. G. Figueroa, A. Florio, F. Torrenti & W. Valkenburg (2006.15122)][D. G. Figueroa, A. Florio, F. Torrenti & W. Valkenburg (2102.01031)]

[Sharma, Brandenburg, Subramanian, Vikman (2411.04854)] [Caravano, Komatsu, Lozanov, Weller (2204.12874)] [Caravano, Komatsu, Lozanov, Weller (2110.10695)]

# 04854)]

## Lattice implementation of axion coupling



## Lattice implementation of axion coupling

$$\frac{1}{\Lambda} \overrightarrow{E} \cdot \overrightarrow{B}$$





Satisfy all conditions

> Satisfies - Gauge transformations  $A_{\mu} \rightarrow A_{\mu} + \partial_{\mu} \alpha(\mathbf{x})$

Lattice

 $\sum \frac{\phi}{\Lambda} E_i^{(2)} B_i^{(4)}$ 

- Bianchi identities

 $\overrightarrow{\nabla} \times \overrightarrow{E} = \overrightarrow{B} \quad \nabla \cdot \overrightarrow{B} = 0$ 

- Topological term as a total derivative

$$F_{\mu\nu}\tilde{F}^{\mu\nu}=\partial_{\mu}K^{\mu}$$

## Lattice implementation of axion coupling

$$\frac{1}{\Lambda} \overrightarrow{E} \cdot \overrightarrow{B}$$







## Lattice formulation of axion inflation

$$\pi_{\phi}' = -3\pi_{\phi} + \frac{1}{H} \left( \frac{1}{a^2} \sum_{i} \Delta_i^{-} \Delta_i^{+} \phi - m^2 \phi + \frac{\alpha_{\Lambda}}{a^3 m_p} \sum_{i} E_i^{(2)} B_i^{(4)} \right)$$

$$E_i' = -E_i - \frac{1}{H} \left( \frac{1}{a^2} \sum_{jk} \epsilon_{ijk} \Delta_j^{-} B_k + \frac{\alpha_{\Lambda}}{2am_p} \left( \pi_{\phi} B_i^{(4)} + \pi_{\phi,+i} B_{i,+i}^{(4)} \right) - \frac{\alpha_{\Lambda}}{4am_p} \sum_{\pm} \sum_{j,k} \epsilon_{ijk} \left\{ \left[ (\Delta_j^{\pm} \phi) E_{k,\pm j}^{(2)} \right]_{+i} + \left[ (\Delta_j^{\pm} \phi) E_{k,\pm j}^{(2)} \right\} \right]$$

$$\sum_i \Delta_i^{-} E_i = -\frac{\alpha_{\Lambda}}{2am_p} \sum_{\pm} \sum_i \left( \Delta_i^{\pm} \phi \right) B_{i,\pm i}^{(4)}$$



## Lattice formulation of axion inflation

$$\pi_{\phi}' = -3\pi_{\phi} + \frac{1}{H} \left( \frac{1}{a^2} \sum_{i} \Delta_{i}^{-} \Delta_{i}^{+} \phi - m^2 \phi + \frac{\alpha_{\Lambda}}{a^3 m_p} \sum_{i} E_{i}^{(2)} B_{i}^{(4)} \right)$$

$$E_{i}' = -E_{i} - \frac{1}{H} \left( \frac{1}{a^2} \sum_{jk} \epsilon_{ijk} \Delta_{j}^{-} B_{k} + \frac{\alpha_{\Lambda}}{2am_p} \left( \pi_{\phi} B_{i}^{(4)} + \pi_{\phi,+i} B_{i,+i}^{(4)} \right) - \frac{\alpha_{\Lambda}}{4am_p} \sum_{\pm} \sum_{j,k} \epsilon_{ijk} \left\{ \left[ (\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)} \right]_{+i} + \left[ (\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)} \right] \right\}$$

$$\sum_{i} \Delta_{i}^{-} E_{i} = -\frac{\alpha_{\Lambda}}{2am_p} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi \right) B_{i,\pm i}^{(4)}$$

- kernel of 
$$\mathcal{K}_{A}[a, \dot{a}, \phi, \pi_{\phi}, A_{i}, E_{i}]$$



## Lattice formulation of axion inflation

$$\pi_{\phi}' = -3\pi_{\phi} + \frac{1}{H} \left( \frac{1}{a^2} \sum_{i} \Delta_{i}^{-} \Delta_{i}^{+} \phi - m^2 \phi + \frac{\alpha_{\Lambda}}{a^3 m_p} \sum_{i} E_{i}^{(2)} B_{i}^{(4)} \right)$$

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$$\sum_{i} \Delta_{i}^{-} E_{i} = -\frac{\alpha_{\Lambda}}{2am_p} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi \right) B_{i,\pm i}^{(4)}$$

- kernel of  
$$\mathscr{K}_{A}[a, \dot{a}, \phi, \pi_{\phi}, A_{i}, E_{i}]$$

Non-symplectic integrators like Runge-Kutta















#### Run simulations until end of inflation $\epsilon_H = 1$



$$\begin{array}{c} \alpha_{\Lambda} = 10\\ \alpha_{\Lambda} = 11\\ \alpha_{\Lambda} = 12\\ \alpha_{\Lambda} = 12\\ \alpha_{\Lambda} = 13\\ \alpha_{\Lambda} = 13\\ \alpha_{\Lambda} = 14\\ \alpha_{\Lambda} = 15\\ \alpha_{\Lambda} = 15\\ \alpha_{\Lambda} = 16\\ \alpha_{\Lambda} = 17\\ \alpha_{\Lambda} = 18\\ \alpha_{\Lambda} = 19\\ \alpha_{\Lambda} = 20\end{array}$$











$$\Delta {\cal N}_{
m br}$$



![](_page_35_Figure_1.jpeg)

![](_page_35_Picture_3.jpeg)

![](_page_36_Figure_1.jpeg)

Extrapolations  $\alpha_{\Lambda} = 25 \rightarrow \Delta \mathcal{N}_{br} \sim 10 - 12$   $\alpha_{\Lambda} = 30 \rightarrow \Delta \mathcal{N}_{br} \sim 15 - 18$   $\alpha_{\Lambda} = 35 \rightarrow \Delta \mathcal{N}_{br} \sim 18 - 25$ 

![](_page_36_Picture_4.jpeg)

![](_page_37_Figure_1.jpeg)

![](_page_37_Picture_4.jpeg)

![](_page_38_Figure_1.jpeg)

![](_page_38_Picture_4.jpeg)

![](_page_39_Figure_1.jpeg)

![](_page_39_Picture_4.jpeg)

![](_page_40_Figure_0.jpeg)

![](_page_40_Picture_5.jpeg)

![](_page_41_Figure_0.jpeg)

![](_page_41_Picture_5.jpeg)

![](_page_42_Figure_0.jpeg)

![](_page_42_Picture_5.jpeg)

![](_page_43_Figure_0.jpeg)

![](_page_43_Picture_5.jpeg)

![](_page_44_Figure_1.jpeg)

#### Strong

![](_page_44_Picture_5.jpeg)

![](_page_45_Figure_1.jpeg)

#### Strong

![](_page_45_Picture_5.jpeg)

![](_page_46_Figure_1.jpeg)

![](_page_46_Figure_2.jpeg)

![](_page_47_Figure_1.jpeg)

![](_page_47_Figure_2.jpeg)

![](_page_48_Figure_1.jpeg)

![](_page_48_Figure_2.jpeg)

![](_page_49_Figure_1.jpeg)

![](_page_49_Figure_2.jpeg)

![](_page_50_Figure_1.jpeg)

![](_page_50_Figure_2.jpeg)

![](_page_51_Figure_1.jpeg)

(Electro)Magnetic Slow Roll

![](_page_51_Figure_3.jpeg)

![](_page_52_Figure_0.jpeg)

![](_page_52_Figure_1.jpeg)

![](_page_52_Figure_2.jpeg)

![](_page_52_Picture_3.jpeg)

![](_page_52_Picture_4.jpeg)

![](_page_52_Picture_5.jpeg)

![](_page_53_Figure_0.jpeg)

 $\Delta^{(\pm)}_A(\mathcal{N},k)/m^2$ 

![](_page_54_Figure_0.jpeg)

## UV sensitivity and convergence

![](_page_55_Figure_1.jpeg)

Evolve equally

During strong backreaction, Different UV resolution affects evolution

![](_page_56_Figure_1.jpeg)

#### UV sensitivity and convergence

UV resolution affects the evolution of the self consistent background

![](_page_56_Figure_4.jpeg)

![](_page_57_Figure_0.jpeg)

![](_page_57_Figure_1.jpeg)

#### UV sensitivity and convergence

UV resolution affects the evolution of the self consistent background

![](_page_57_Figure_4.jpeg)

![](_page_58_Figure_1.jpeg)

![](_page_58_Figure_2.jpeg)

## Conclusions

Local backreaction is crucial. 

Inflation is delayed several e-folds.

• Dominance of magnetic energy during well st stage of inflation.

Scale dependent chirality.

![](_page_59_Picture_5.jpeg)

 $\sum \frac{\phi}{\Lambda} E_i^{(2)} B_i^{(4)}$ 

$$\Delta \mathcal{N}_{\rm pr} = b_2 (\alpha_{\Lambda} - 14.31)^{a_2} + 1.44$$

![](_page_59_Figure_8.jpeg)

![](_page_59_Figure_10.jpeg)

 $10^{0}$ 

![](_page_59_Picture_12.jpeg)

![](_page_60_Figure_1.jpeg)

# Excitation of longitudinal and negative $A_{\mu}$ modes

$$(\pi_{\phi}\overrightarrow{B})_{\overrightarrow{k}} = -i\sum_{\lambda=\pm}\lambda\int^{3}q\dot{\phi}^{*}_{(\overrightarrow{k}-\overrightarrow{q})}qA^{\lambda}_{\overrightarrow{q}}\overrightarrow{\varepsilon}^{\lambda}_{\overrightarrow{q}}$$

$$(\pi_{\phi}\overrightarrow{B})_{\overrightarrow{k}} \equiv -i \int d^{3}\overrightarrow{q} \ \dot{\phi}_{(\overrightarrow{k}-\overrightarrow{q})}^{*} q (A_{\overrightarrow{q}}^{+}(\overrightarrow{\varepsilon}_{\overrightarrow{k}}^{+} \cdot \overrightarrow{\varepsilon}_{\overrightarrow{q}}^{+}) - A_{\overrightarrow{q}}^{-}(\overrightarrow{\varepsilon}_{\overrightarrow{k}}^{+})$$

$$\simeq -\frac{i}{2} \int d^2\hat{q} \, dq \, q^3 \, \dot{\phi}^*_{(\vec{k}-\vec{q})} (1-\cos\theta) A^+_{\vec{q}}$$

![](_page_60_Picture_7.jpeg)

![](_page_60_Picture_8.jpeg)

![](_page_61_Figure_1.jpeg)

![](_page_61_Picture_2.jpeg)

#### **Comparison with Homogeneous backreaction**

![](_page_62_Figure_1.jpeg)

#### **Comparison with Homogeneous backreaction**

### Power spectrum of homogeneous back reaction

![](_page_63_Figure_1.jpeg)

![](_page_63_Picture_3.jpeg)