

Mapping the QCD phase diagram : Connecting EoS to observables

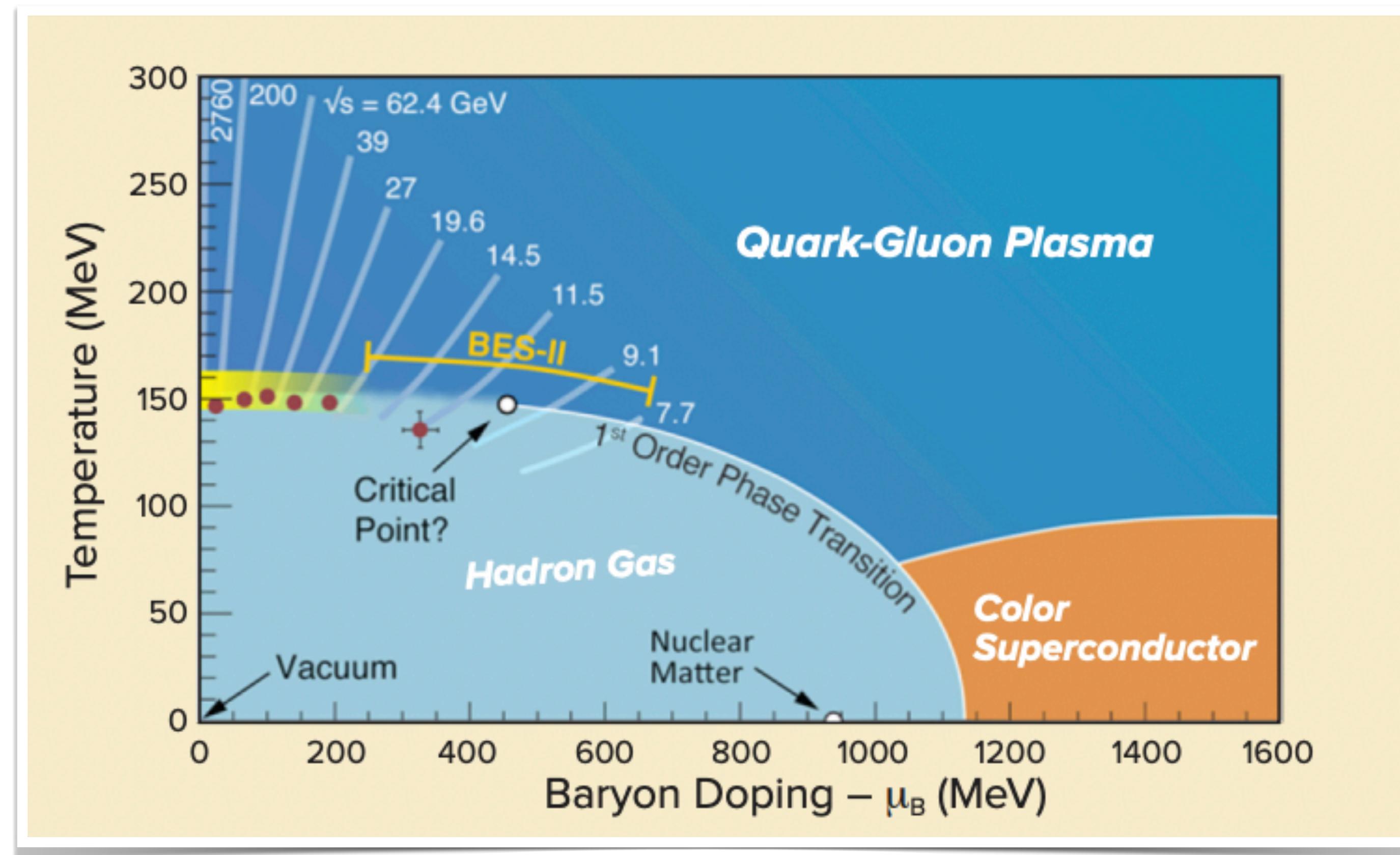
Workshop on “Confinement and Symmetry from vacuum to QCD phase diagram”

Maneesha Pradeep
University of Maryland at College Park

Collaborators : Misha Stephanov, Jamie Karthein, Krishna Rajagopal, Yi Yin

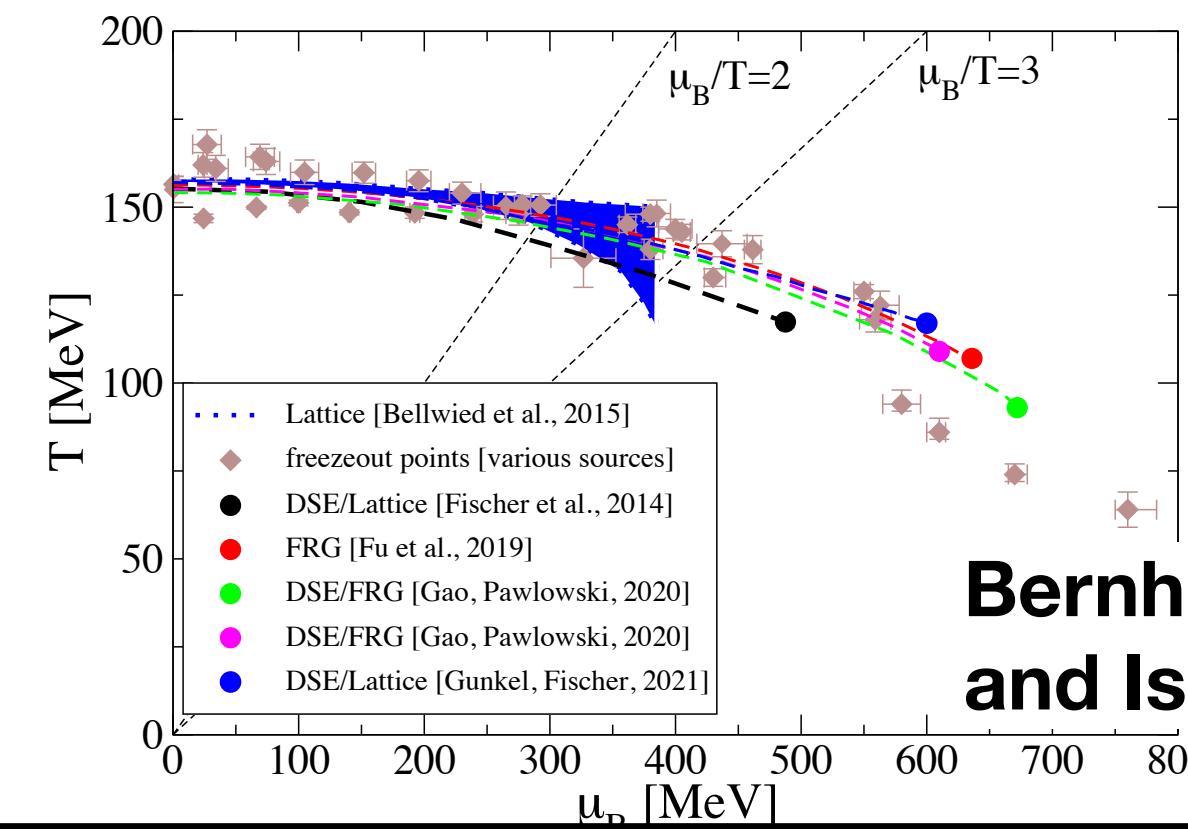
Phys.Rev.Lett. 130 (2023) 16, 16, arXiv 2409.16249 + Work in progress

Is there a critical point on the QCD phase diagram?

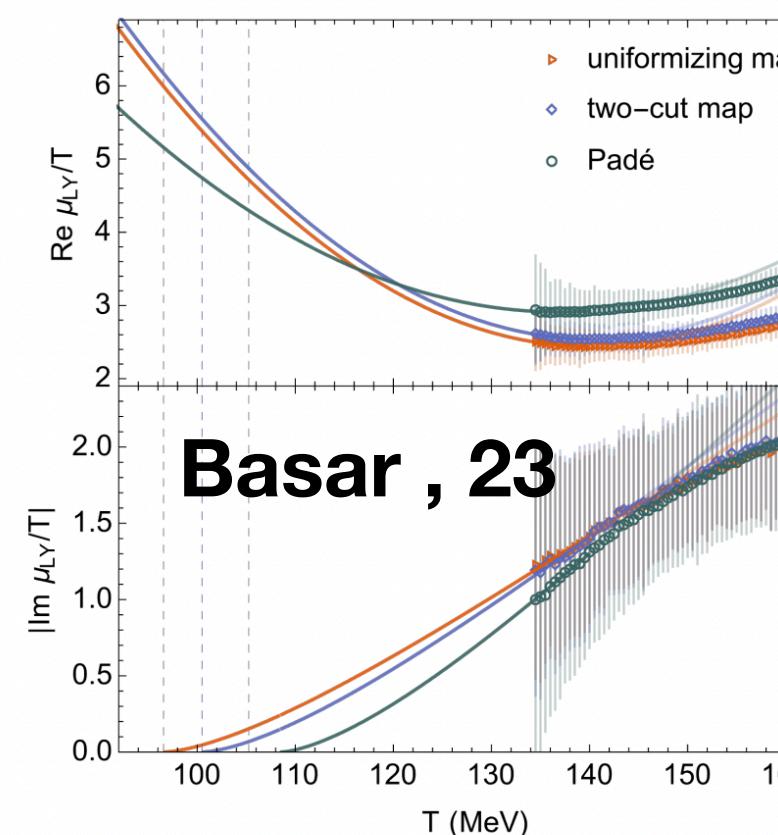


- At $\mu_B = 0$: It is a cross-over. — Lattice QCD simulations
- At T= 0 : Models predict a first order phase transition — Not known from first principles

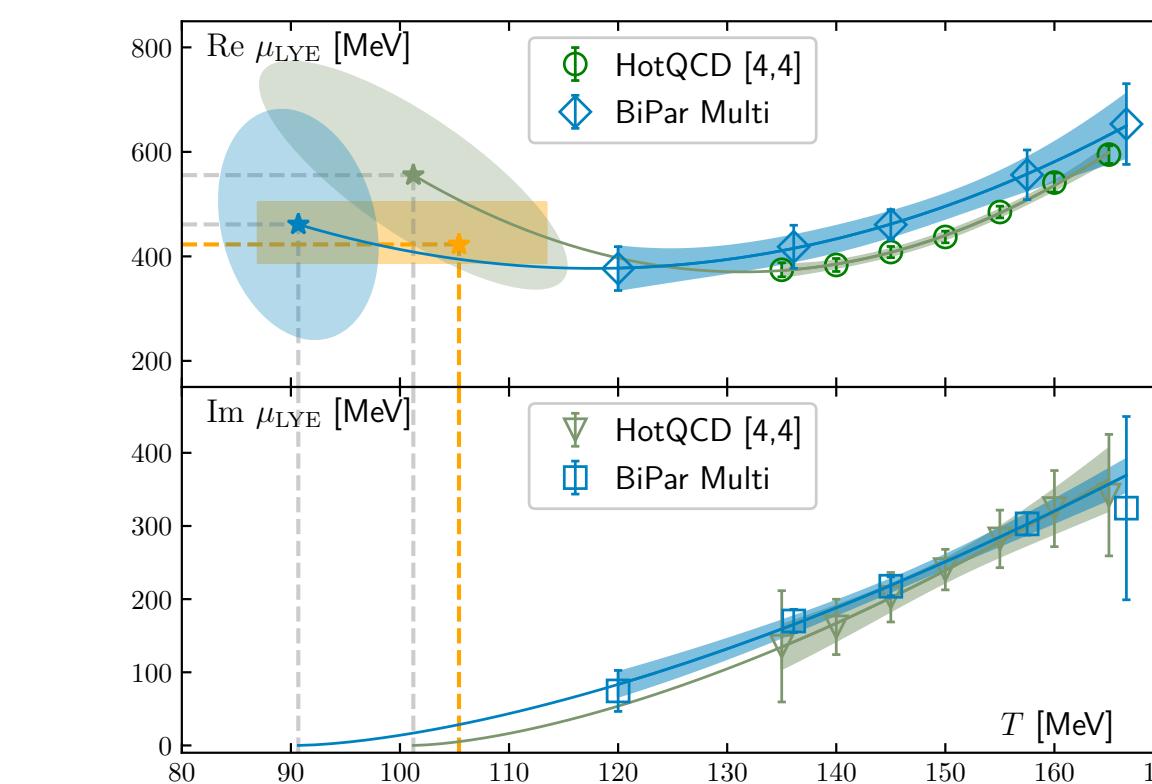
Recent theory guidances for the location of the critical point



$$(\mu_{BC}, T_c) = (495 - 654, 108 - 119) \text{ MeV}$$



$$(\mu_{BC}, T_c) \approx (580, 100) \text{ MeV}$$



**Clarke et al.,
24**

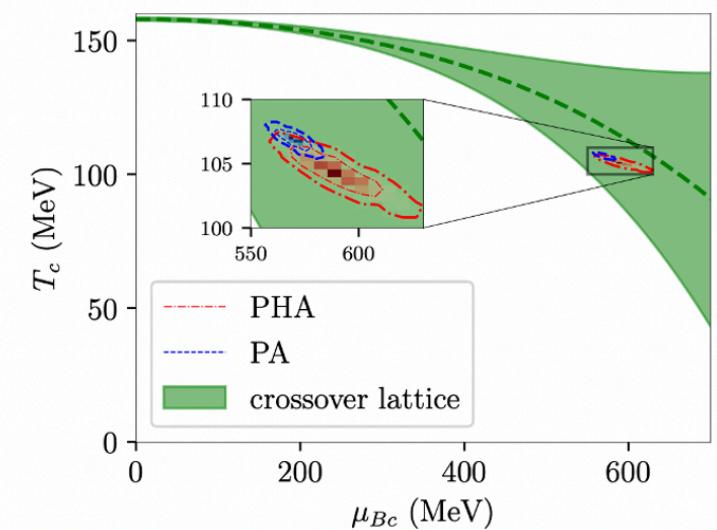
$$(\mu_{BC}, T_c) = (422^{+80}_{-35}, 105^{+8}_{-18}) \text{ MeV}$$

Font Color
Bayesian Holography + Lattice input at $\mu = 0$

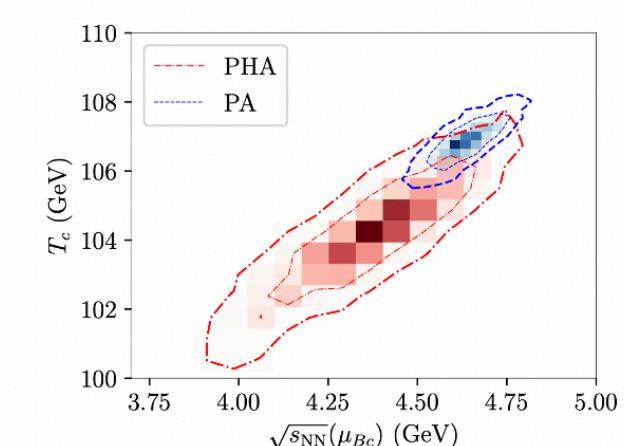
[Hippert et al, e-Print: 2309.00579 \[nucl-th\]](#)

Predict CFP (95% confidence level):

$$T_c = 101 - 108 \text{ MeV} \quad \mu_c = 560 - 625 \text{ MeV}$$



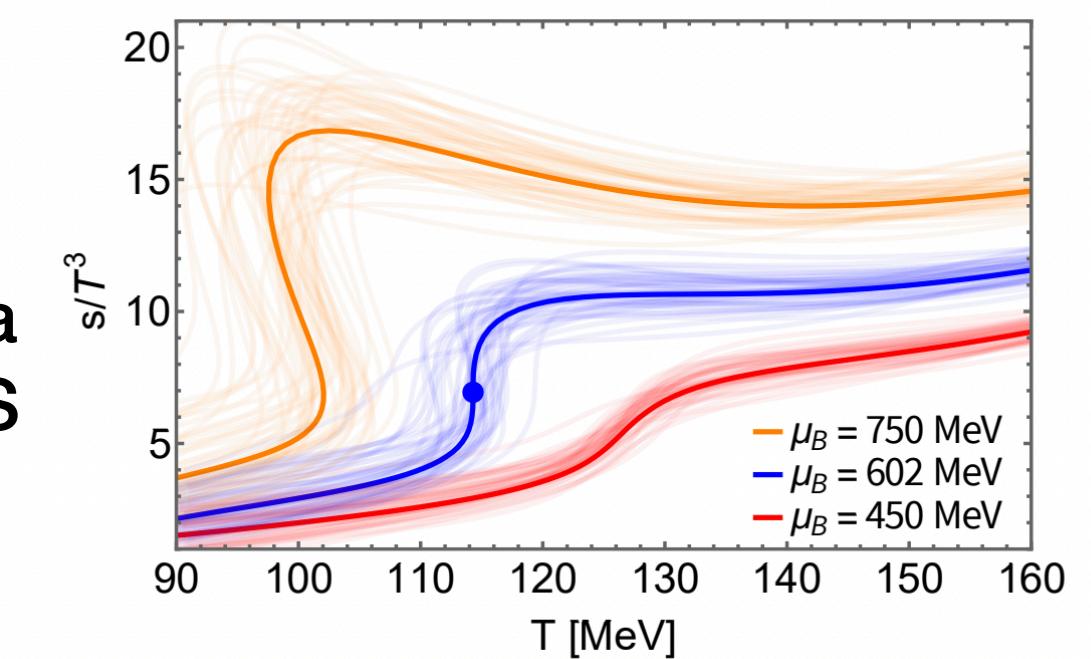
$$\sqrt{s} = 4.0 - 4.8 \text{ GeV}$$



Extrapolations of Lee-Yang edge singularities to real axis

Thermodynamic analysis of a lattice QCD extrapolated EoS

Shah et al., 24



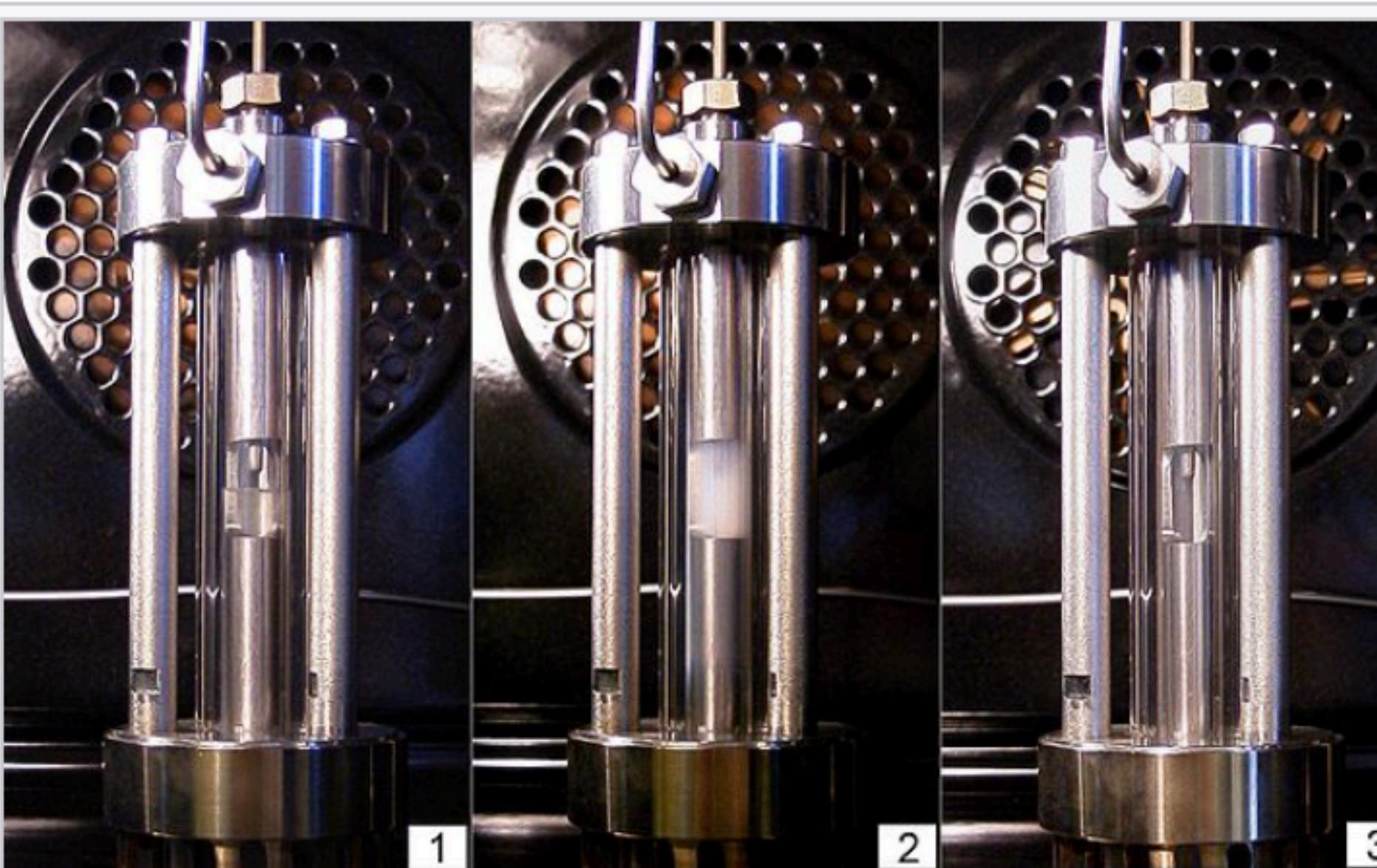
$$(\mu_{BC}, T_c) = (602.1 \pm 62.1, 114.3 \pm 6.9)$$

Probing critical points experimentally

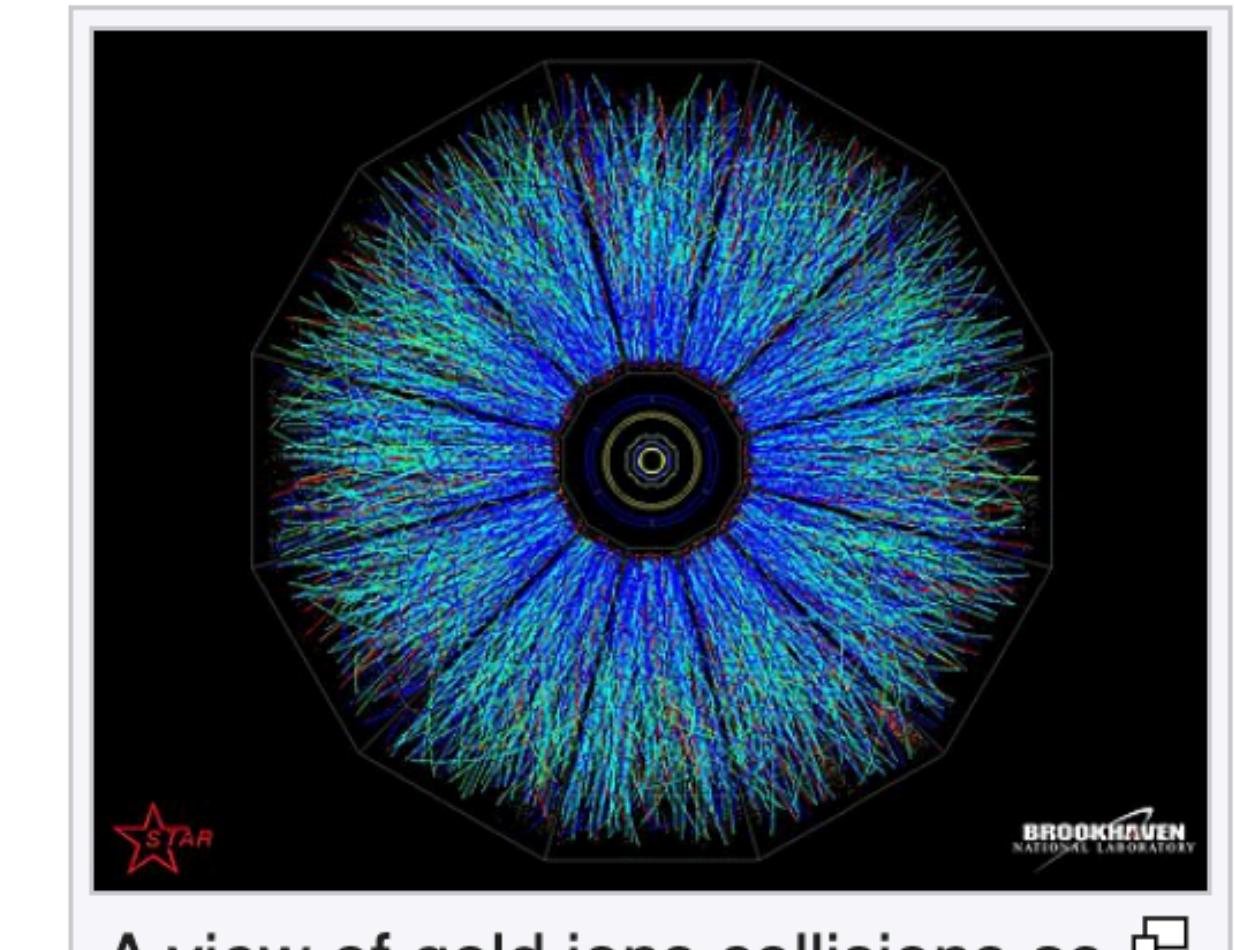


Baron Cagniard de la Tour's experiment

Ethane (C ₂ H ₆)	31.17 °C (304.32 K)	48.077 atm (4,871.4 kPa)
Ethanol (C ₂ H ₅ OH)	241 °C (514 K)	62.18 atm (6,300 kPa)
Fluorine	-128.85 °C (144.30 K)	51.5 atm (5,220 kPa)
Helium	-267.96 °C (5.19 K)	2.24 atm (227 kPa)
Hydrogen	-239.95 °C (33.20 K)	12.8 atm (1,300 kPa)
Krypton	-63.8 °C (209.3 K)	54.3 atm (5,500 kPa)
Methane (CH ₄)	-82.3 °C (190.8 K)	45.79 atm (4,640 kPa)
Neon	-228.75 °C (44.40 K)	27.2 atm (2,760 kPa)
Nitrogen	-146.9 °C (126.2 K)	33.5 atm (3,390 kPa)
Oxygen (O ₂)	-118.6 °C (154.6 K)	49.8 atm (5,050 kPa)
Carbon dioxide (CO ₂)	31.04 °C (304.19 K)	72.8 atm (7,380 kPa)
Nitrous oxide (N ₂ O)	36.4 °C (309.5 K)	71.5 atm (7,240 kPa)
Sulfuric acid (H ₂ SO ₄)	654 °C (927 K)	45.4 atm (4,600 kPa)
Xenon	16.6 °C (289.8 K)	57.6 atm (5,840 kPa)
Lithium	2,950 °C (3,220 K)	652 atm (66,100 kPa)
Mercury	1,476.9 °C (1,750.1 K)	1,720 atm (174,000 kPa)
Sulfur	1,040.85 °C (1,314.00 K)	207 atm (21,000 kPa)
Iron	8,227 °C (8,500 K)	
Gold	6,977 °C (7,250 K)	5,000 atm (510,000 kPa)
Aluminium	7,577 °C (7,850 K)	
Water (H ₂ O) ^{[3][20]}	373.946 °C (647.096 K)	217.7 atm (22,060 kPa)



1. Subcritical ethane, liquid and gas phase coexist.
2. Critical point (32.17 °C, 48.72 bar), displaying critical opalescence.
3. Supercritical ethane, fluid.^[1]



A view of gold ions collisions as captured by the STAR detector.

Colliding heavy-ions at varying center of mass energy

Probing the EoS via heavy-ion collisions

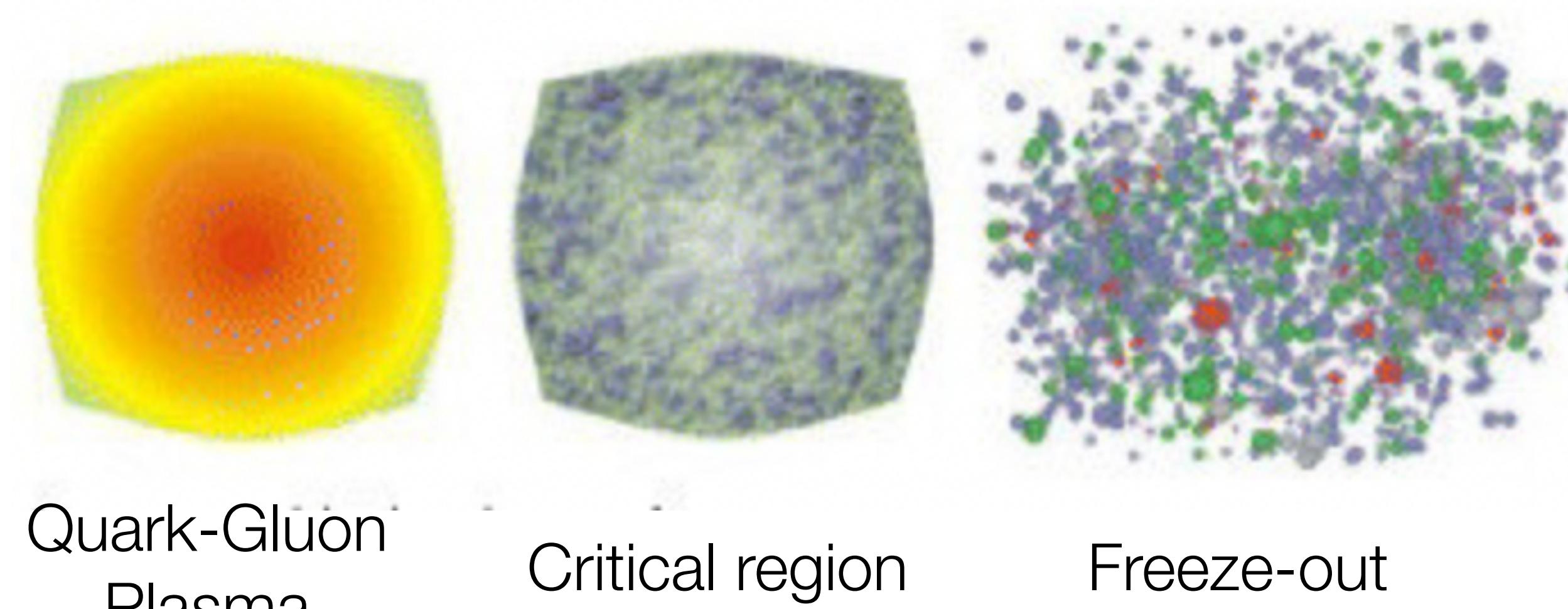
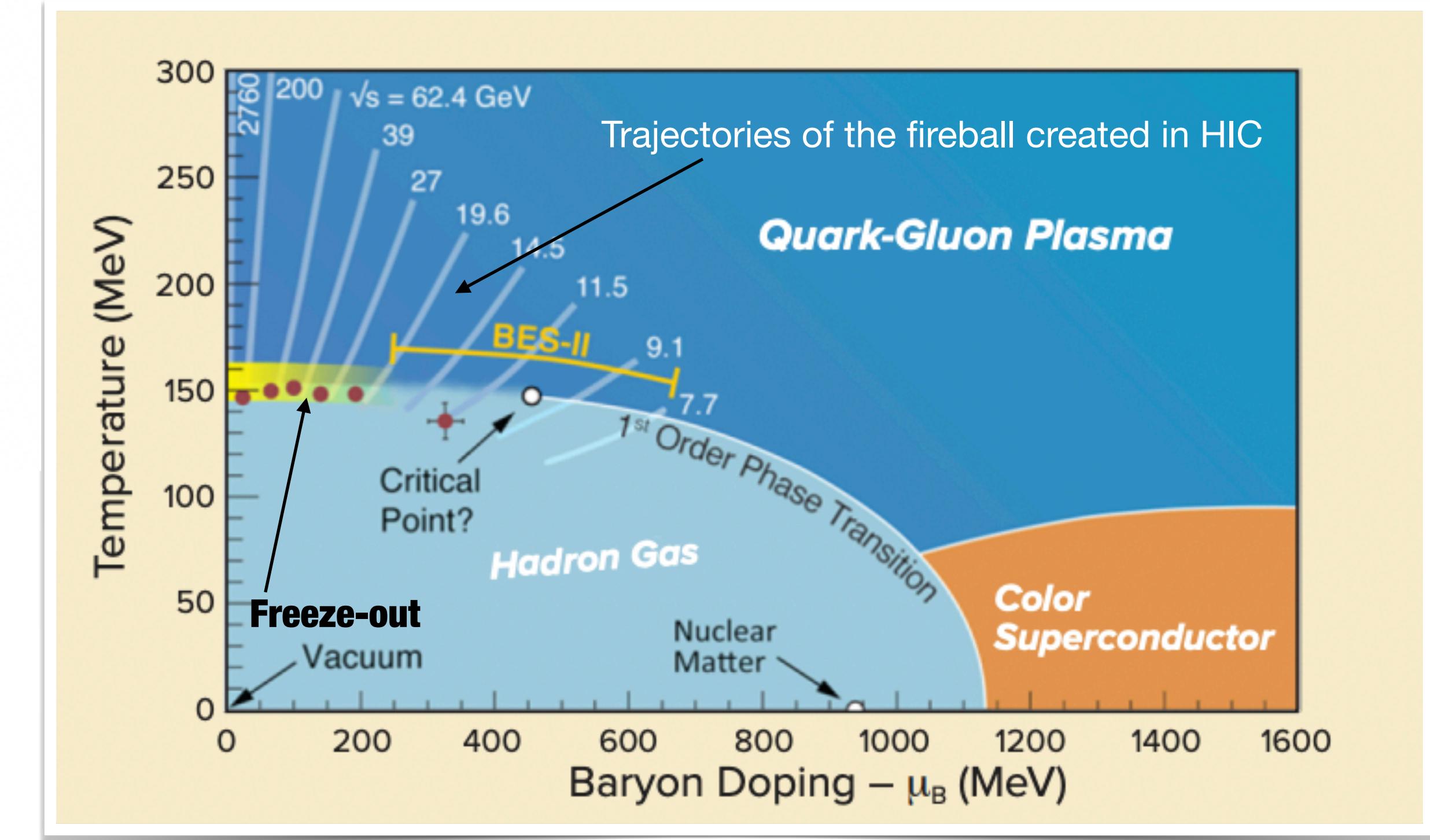
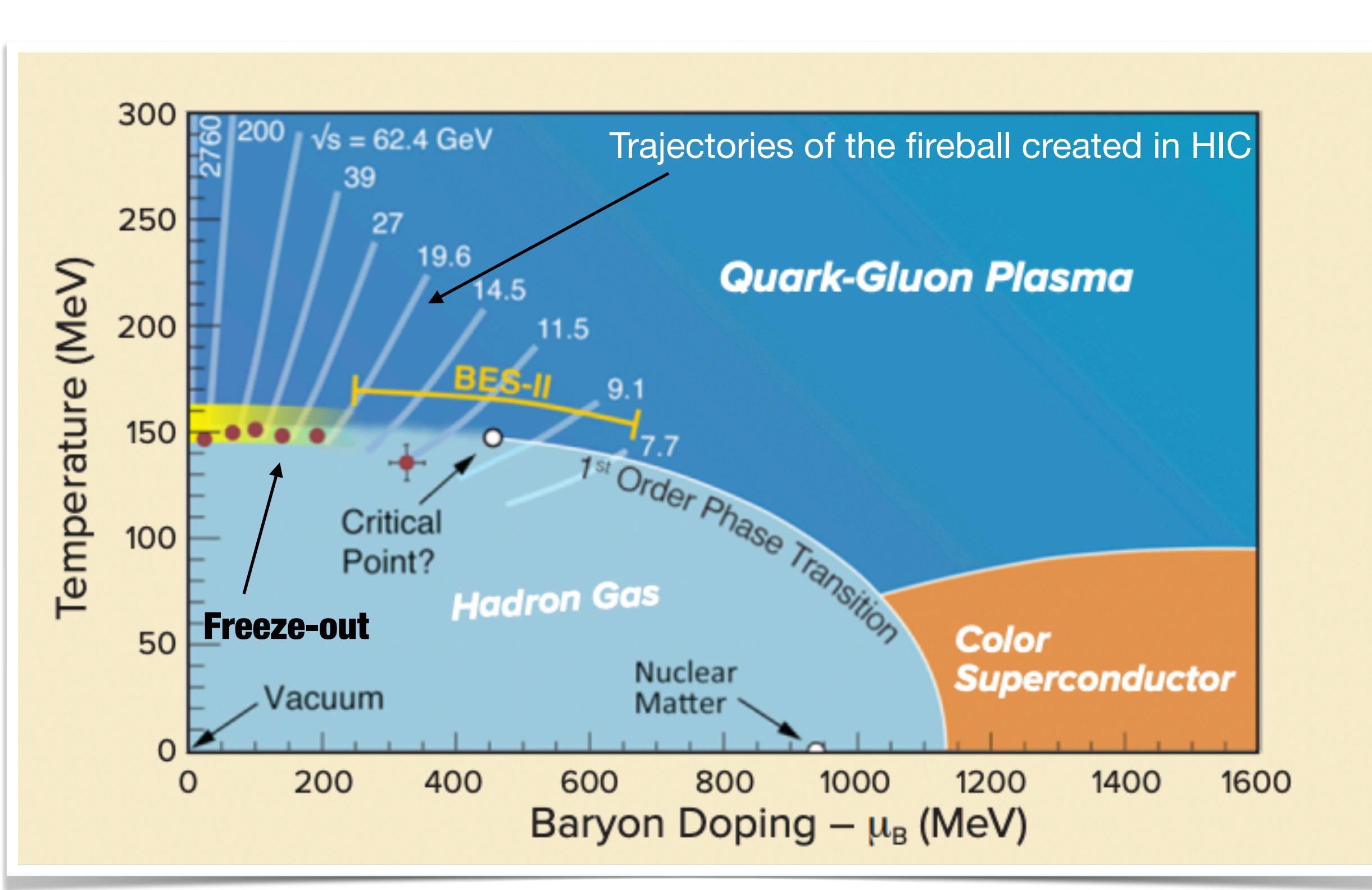


Figure by
S.A.Bass



- QGP thermalizes in $\sim 1\text{fm}/c$
- Freeze-out at about $10\text{ fm}/c$
- The event by event distribution of the particle multiplicities are measured at the detectors.

Non-monotonic dependence of cumulants as a function of collision energy



$$\langle \delta N_B^2 \rangle_c = \langle (N_B - \langle N_B \rangle)^2 \rangle, \langle \delta N_B^3 \rangle_c = \langle (N_B - \langle N_B \rangle)^3 \rangle$$

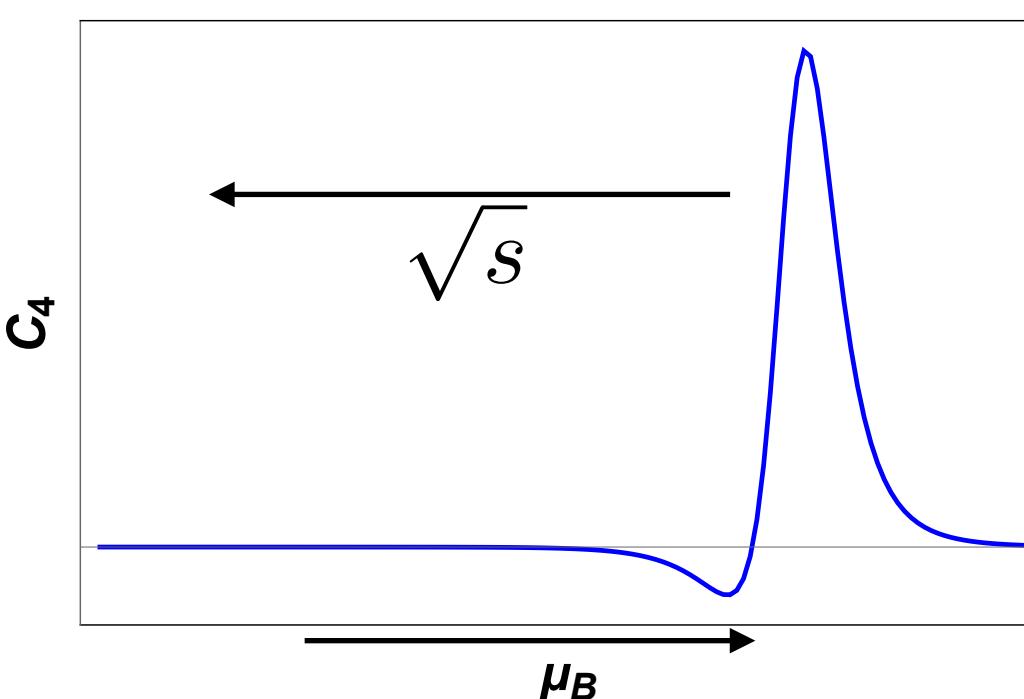
$$\langle \delta N_B^4 \rangle_c = \langle (N_B - \langle N_B \rangle)^4 \rangle - 3 \langle (N_B - \langle N_B \rangle)^2 \rangle$$

Diverges at CP

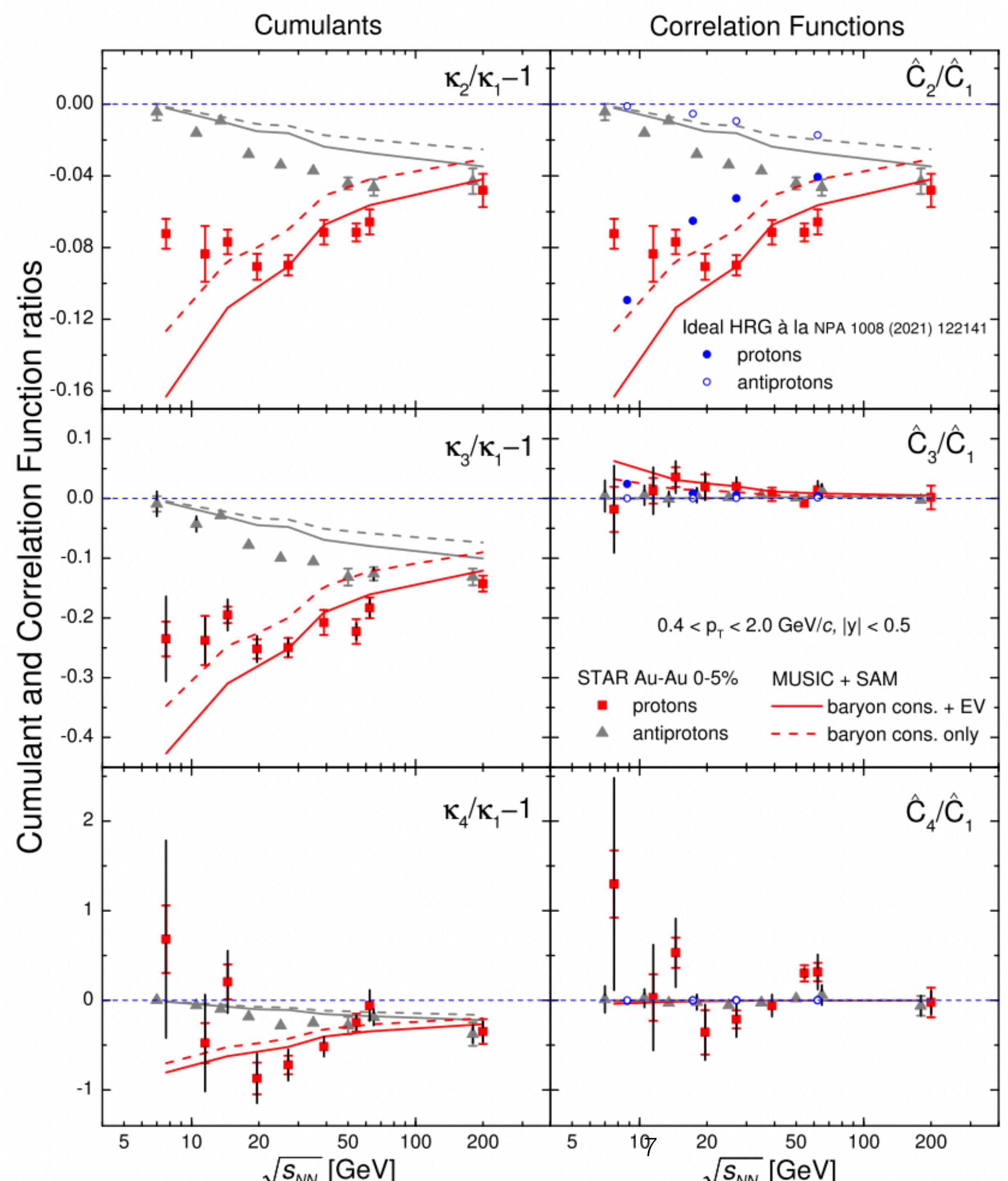
$$C_k \equiv \langle \delta N_B^k \rangle_c \text{ in eq. } VT^{k-1} \frac{\partial^k P}{\partial \mu^k}$$

Fluctuations are enhanced near CP

Similar relations between cumulants of entropy density and derivatives of pressure with respect to temperature



Rajagopal, Shuryak, Stephanov, 98,
99, Stephanov, 09



2021, STAR Collaboration

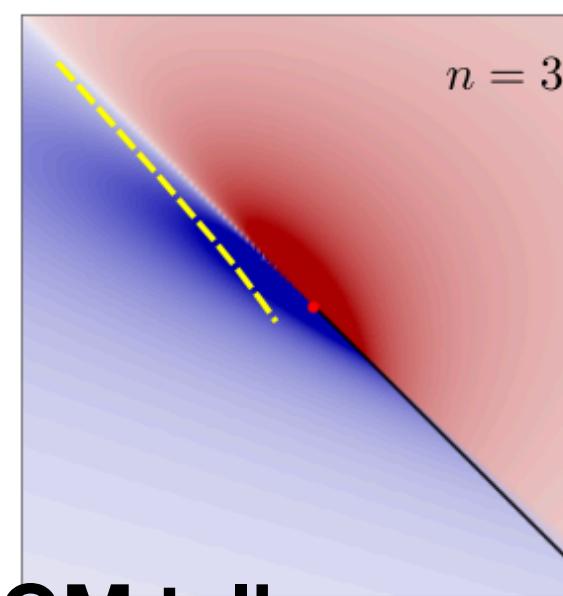
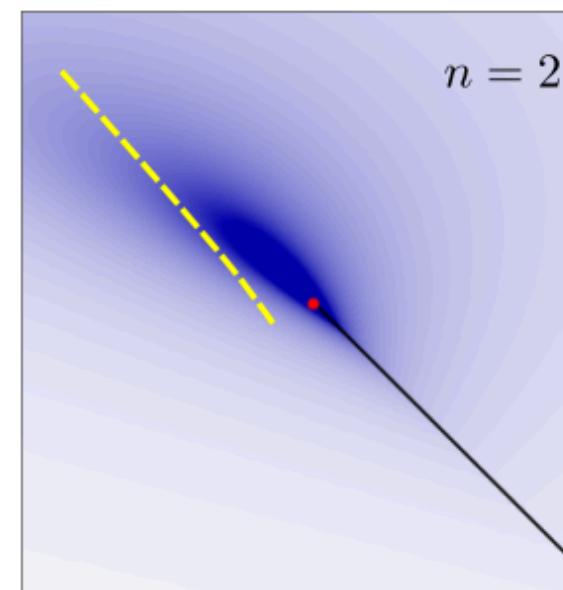
BES-I data for proton multiplicity cumulants

A **clear excess** of scaled proton-number variance from non-critical baseline reported for $\sqrt{s_{NN}} \leq 10 \text{ GeV}$

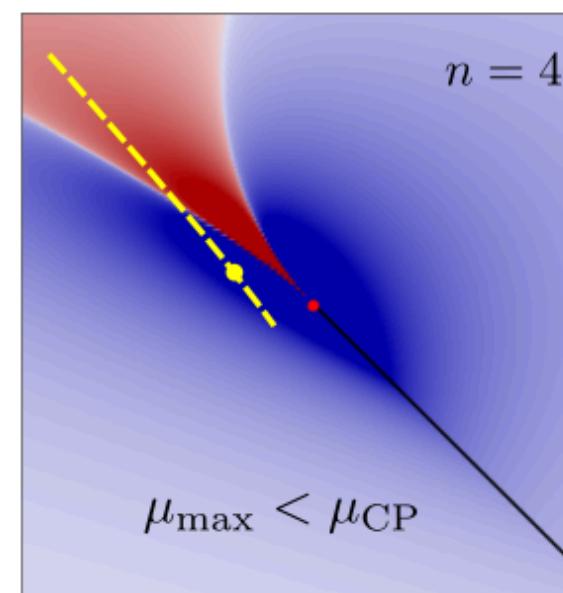
Vovchenko, Koch, Shen, 22

Theory vs BES-II data

(universal EOS) critical χ_n :

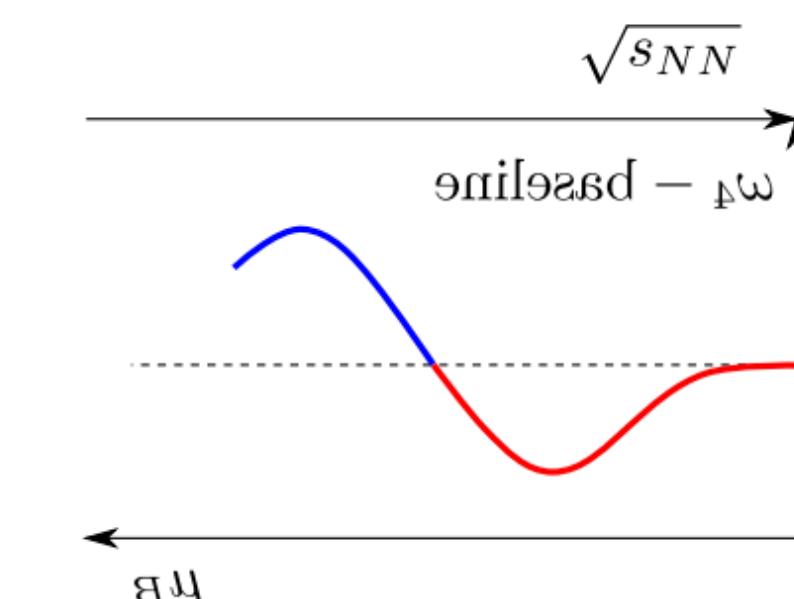
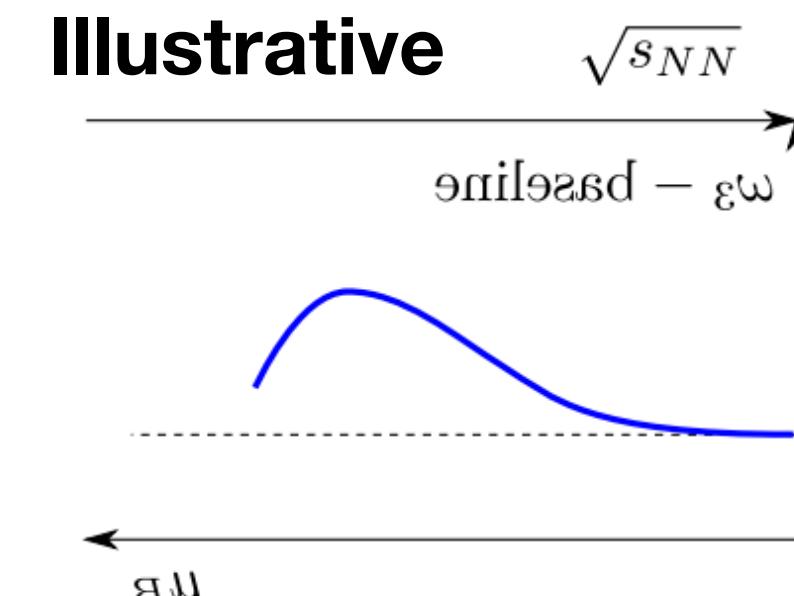
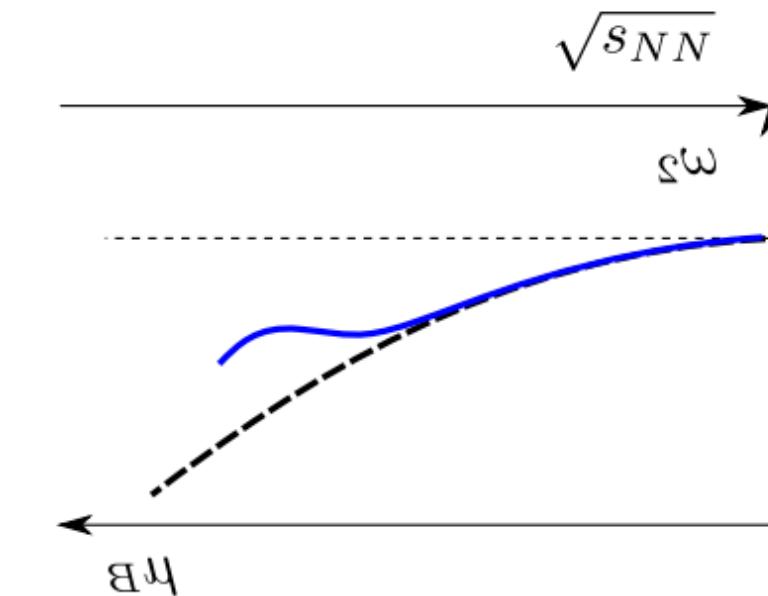


Misha Stephanov's SQM talk

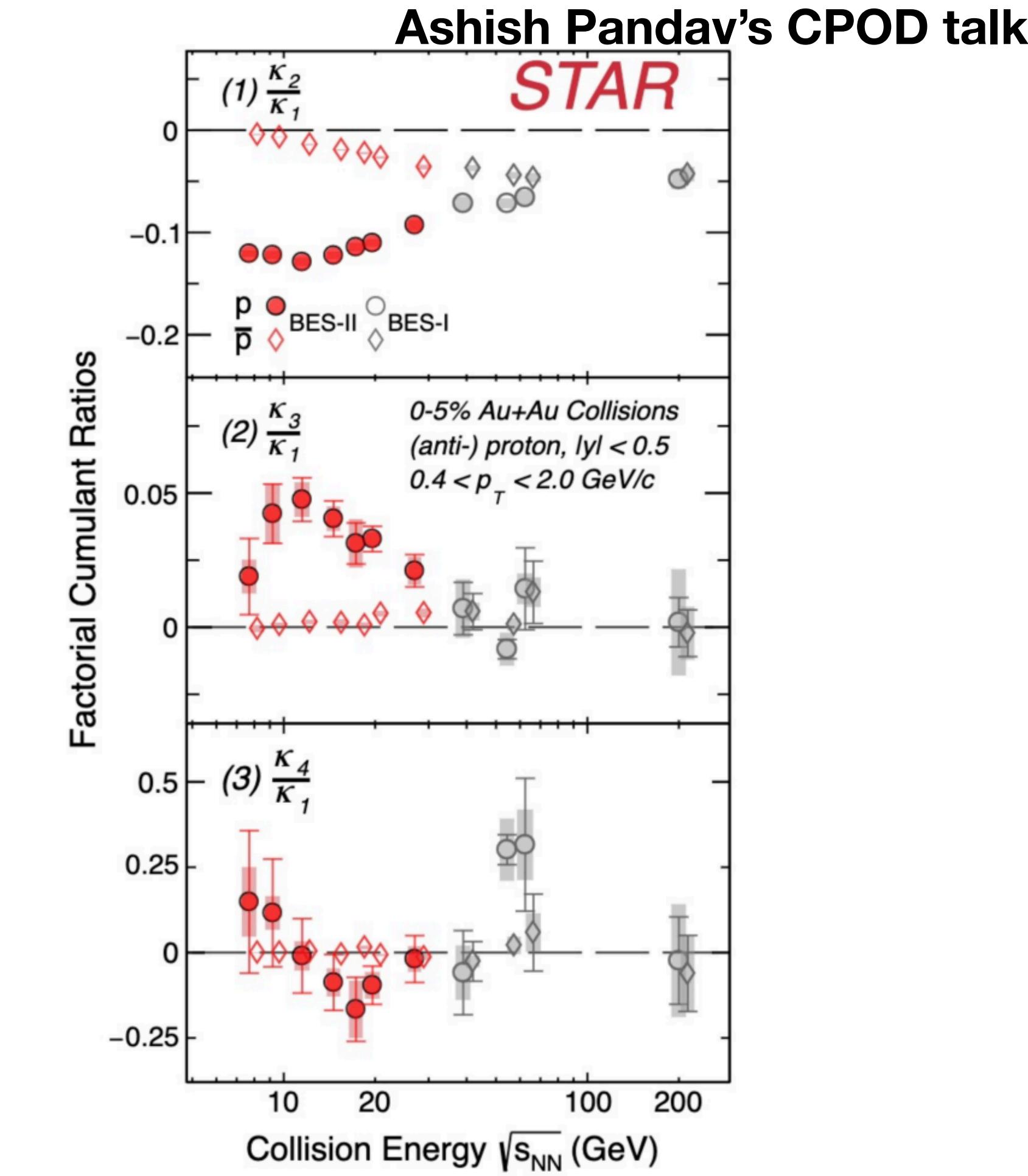


Bzdak et al review 1906.00936

(irreducible correlations) $FC_n [N_p] \sim \chi_n$ (Pradeep, MS 2211.09142), $\omega_n \equiv FC_n / FC_1$



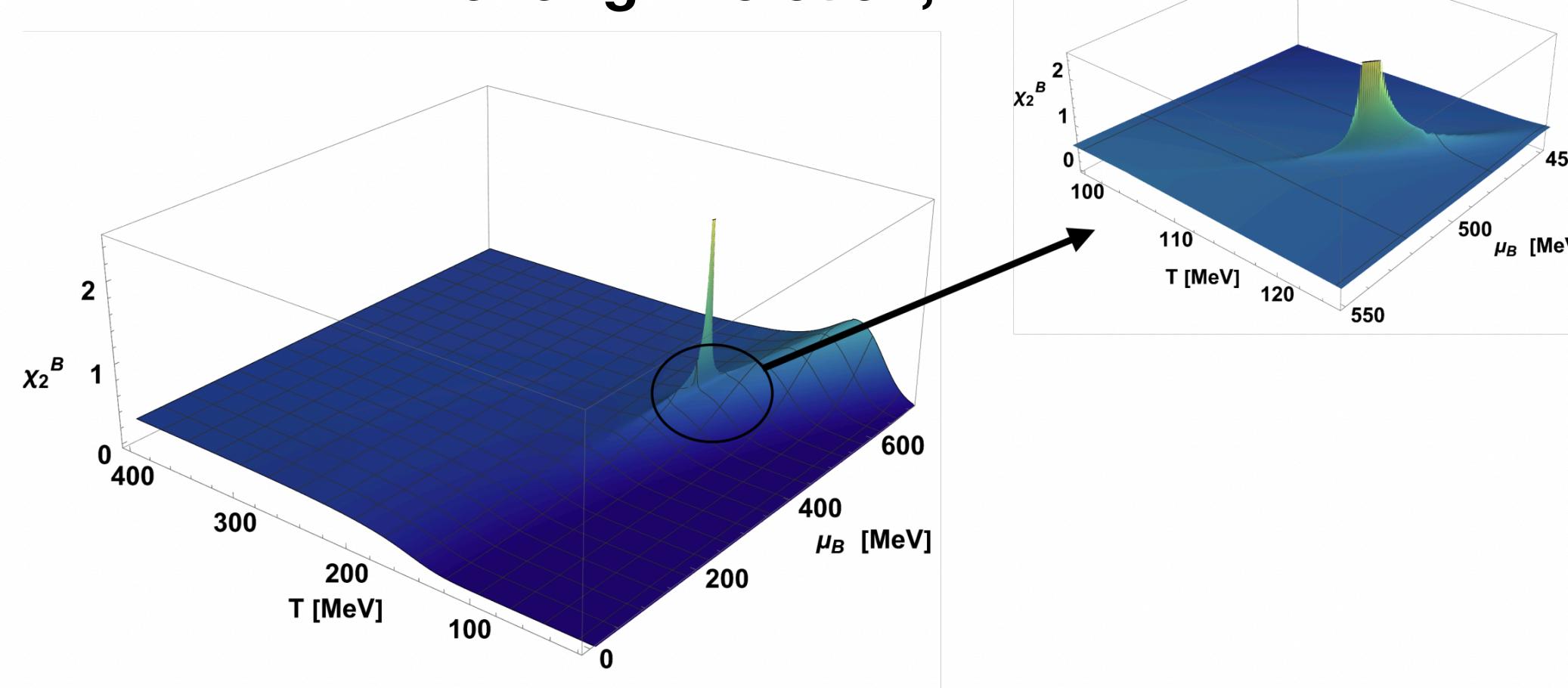
Illustrative



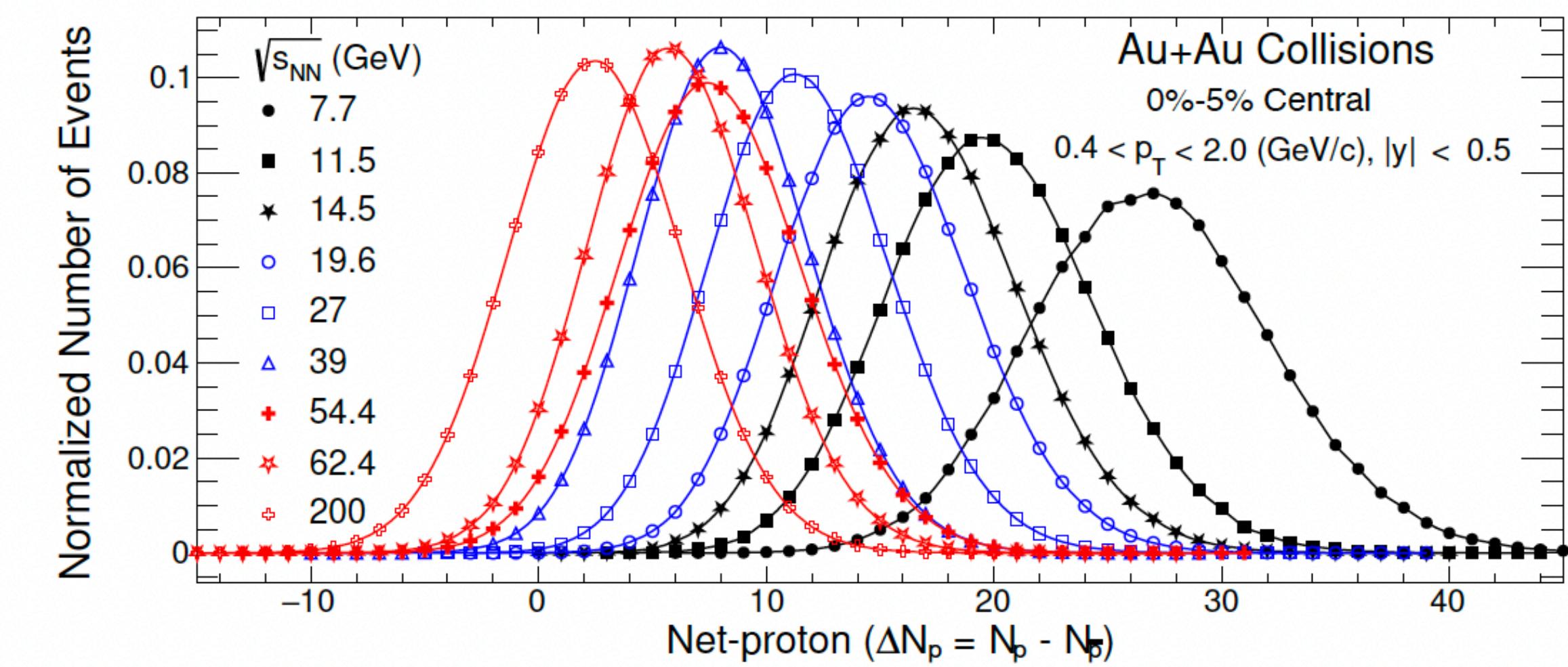
Expected signatures: **bump** in ω_2 and ω_3 , **dip** then **bump** in ω_4 for CP at $\mu_B > 420$ MeV

EoS \longleftrightarrow ? Particle multiplicity distributions

Kahangirwe et al., 24



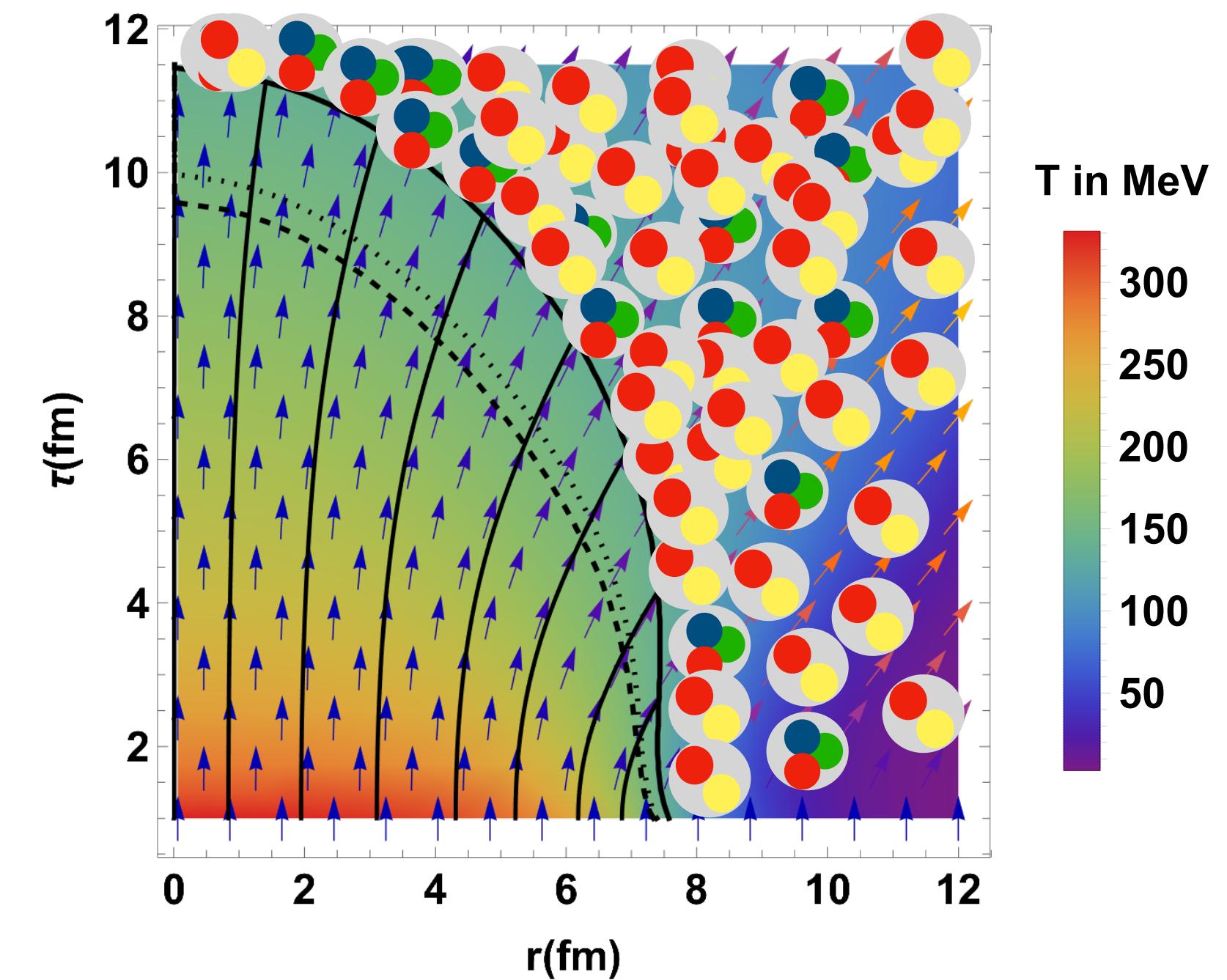
Susceptibilities that diverge at CP



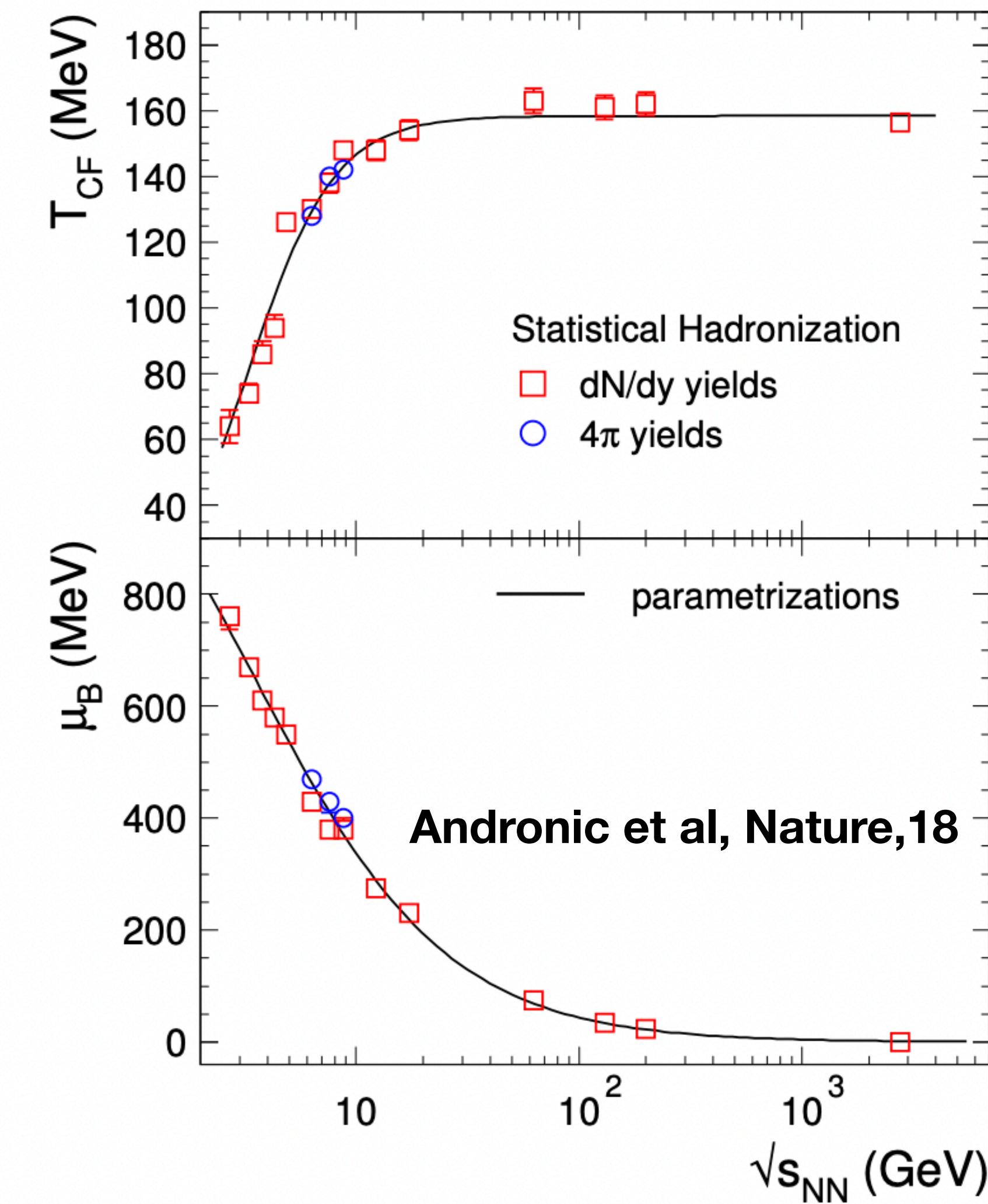
Imprints on Particle Distribution Functions?

STAR Collaboration, PRL 2014, PRL 2021

Freeze-out fits for mid-rapidity yields of hadrons



Freeze-out hypersurface : Both the hydrodynamics and kinetic description in terms of a hadron resonance gas apply



Freeze-out in heavy-ion collisions

Variables at freeze-out

Hydrodynamic mean densities

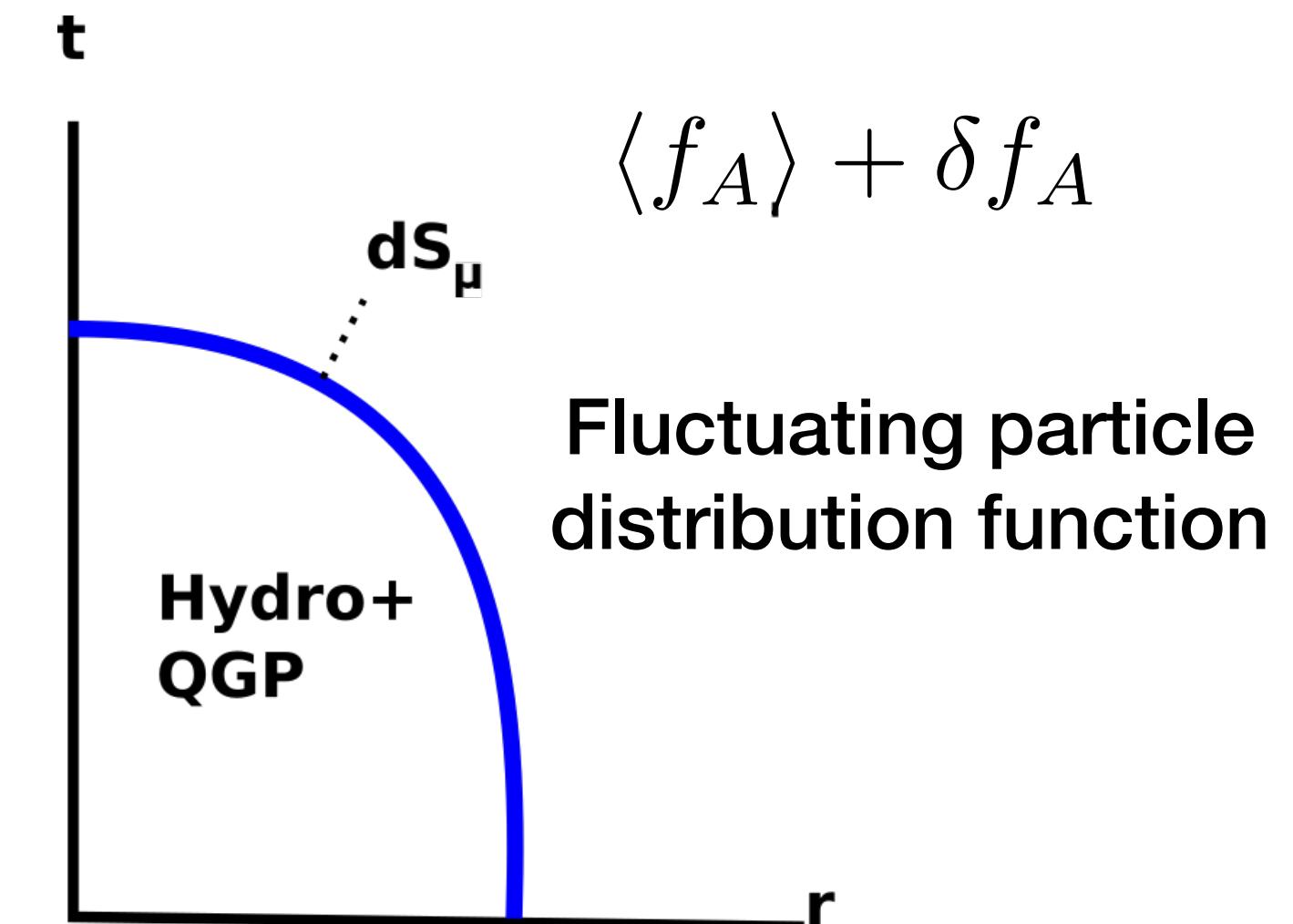
$$\{\langle \epsilon u^\mu \rangle, \langle n \rangle\} \equiv \Psi^a$$

Hydrodynamic correlations

$$\Psi^a, \langle \delta \Psi^a \delta \Psi^b \rangle \equiv H^{ab}, \dots H^{abc\dots}$$

Particle distribution function at freeze-out

$$\langle f_A \rangle = \bar{f}_A, \langle \delta f_A \delta f_B \rangle = G_{AB}, \langle \delta f_A \delta f_B \delta f_C \rangle = G_{ABC} \dots$$



$$\langle f_A \rangle + \delta f_A$$

Fluctuating particle distribution function

Matching conditions at freeze-out

$$\langle \epsilon u^\mu \rangle = \sum_A \int_{p_A} \bar{f}_A p_A^\mu, \quad \langle n \rangle = \sum_A q_A \int_{p_A} \bar{f}_A$$

$$\Psi^a = \sum_A \int_{p_A} \bar{f}_A P_A^a$$

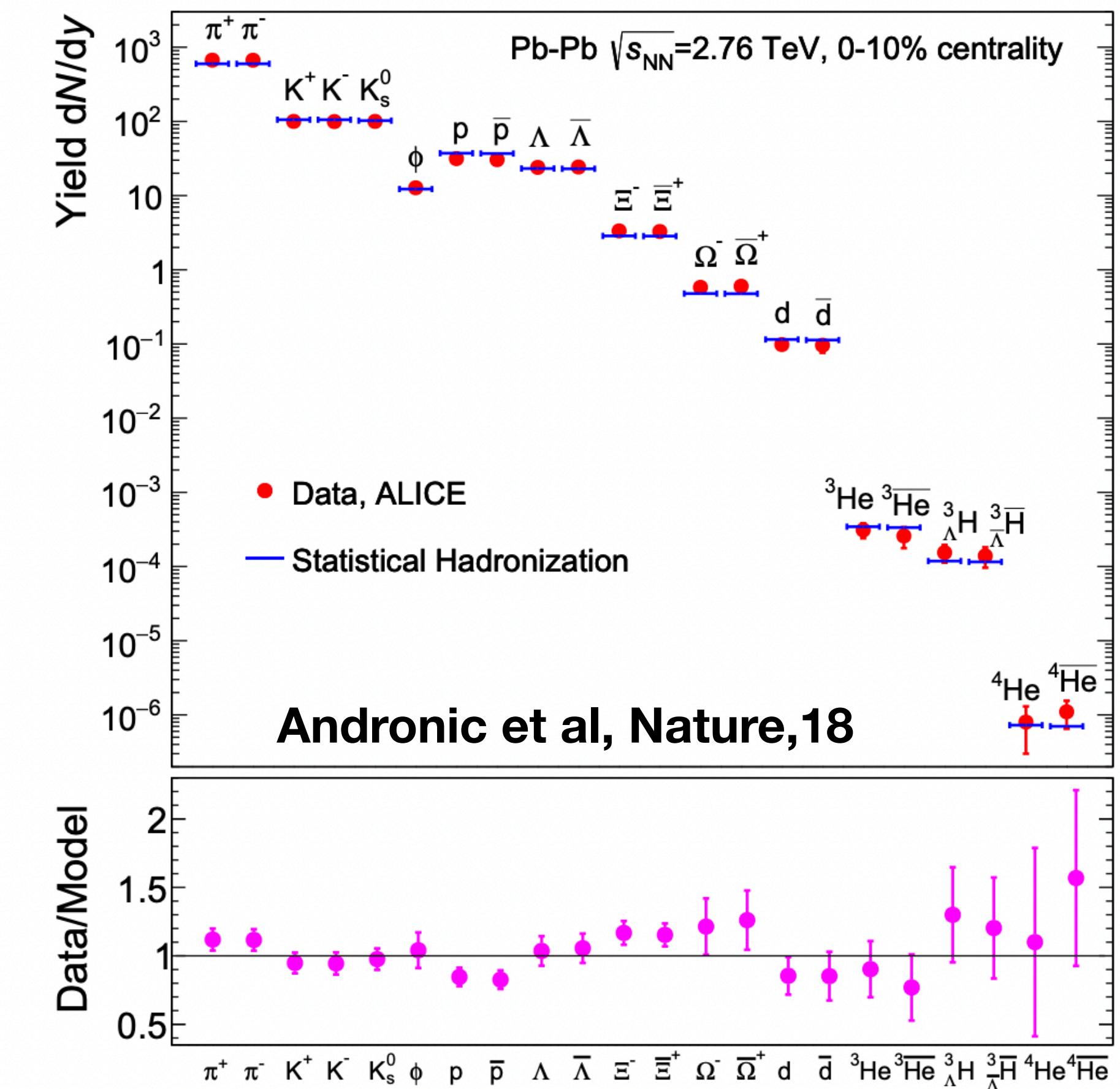
$$P_A = \begin{bmatrix} p_A^\mu \\ q_A \end{bmatrix}$$

$$H^{abc\dots} = \sum_{A,B,C,\dots} \int_{p_A p_B p_C \dots} G_{ABC\dots} P_A^a P_B^b P_C^c \dots$$

- Matching conditions for averages of conserved densities
- Infinitely many sets of distribution functions that satisfy these matching conditions
- Freeze-out prescription corresponds to choosing one of these sets - **How to choose?**

Maximum entropy freeze-out without fluctuations

- Results in a thermal gas of hadrons specified by the local temperature and chemical potential. Coincides with Cooper-Frye, 74
- Recent work on extension to viscous hydrodynamics : Everett-Heinz-Chattopadhyay, 21



Maximum entropy freeze-out

Given all the information about the hydrodynamic densities on the freeze-out hypersurface, what is the ***least biased*** ensemble of free streaming particles after freeze-out that obeys the matching conditions?

It is the one which maximizes the entropy of the fluctuating particle distribution function:

$$S[\bar{f}, G_2, G_3, \dots]$$

subject to the constraints of the matching conditions.

Equilibrium

$$S[\bar{f}]$$

Which entropy to maximize?

Generalized

$$S[P(f)]$$

$$S[P(f)] = S_0 [\bar{f}] + \int_f P(f) \log \frac{P_{\text{eq}}(f)}{P(f)}$$

$$S[\bar{f}, G_2, G_3, \dots]$$

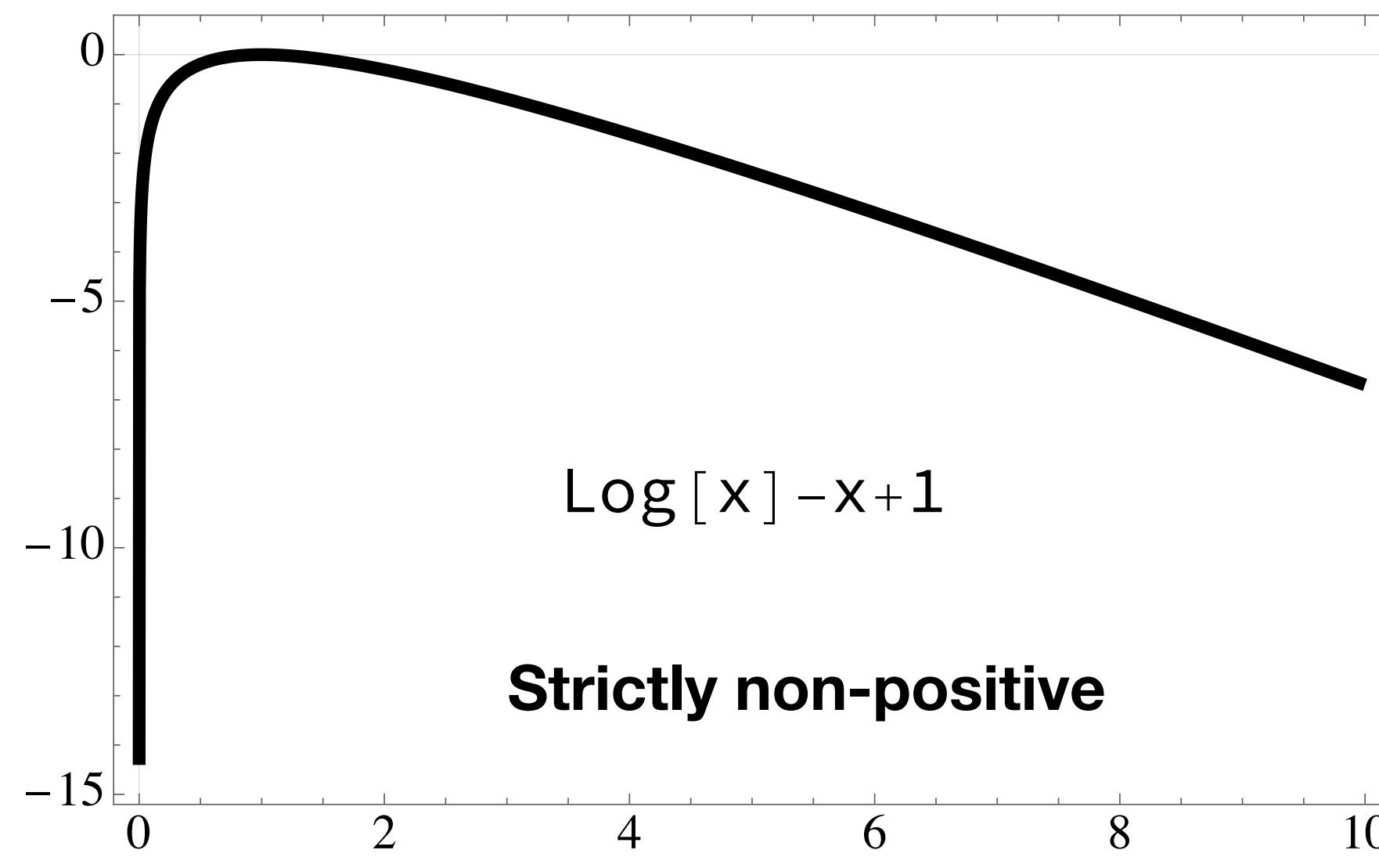
G s are the correlation functions in the Hadron Gas description

- Maximize the **relative entropy** when correlations are out of equilibrium
- Constraints from matching conditions

Entropy to describe out-of equilibrium two-point correlations in ideal HRG

$$S_2 = S_0 + \frac{1}{2} \text{Tr} [\log G \bar{G}^{-1} - G \bar{G}^{-1} + 1] ,$$

Equilibrium



$\text{Log}[x] - x + 1$

Strictly non-positive

Similar to 2-PI action

Berges, 04, Stephanov, Yin, 17...

2-PI entropy

MP, Stephanov, PRL, 23

Two-point correlation function of proton multiplicities in phase space,

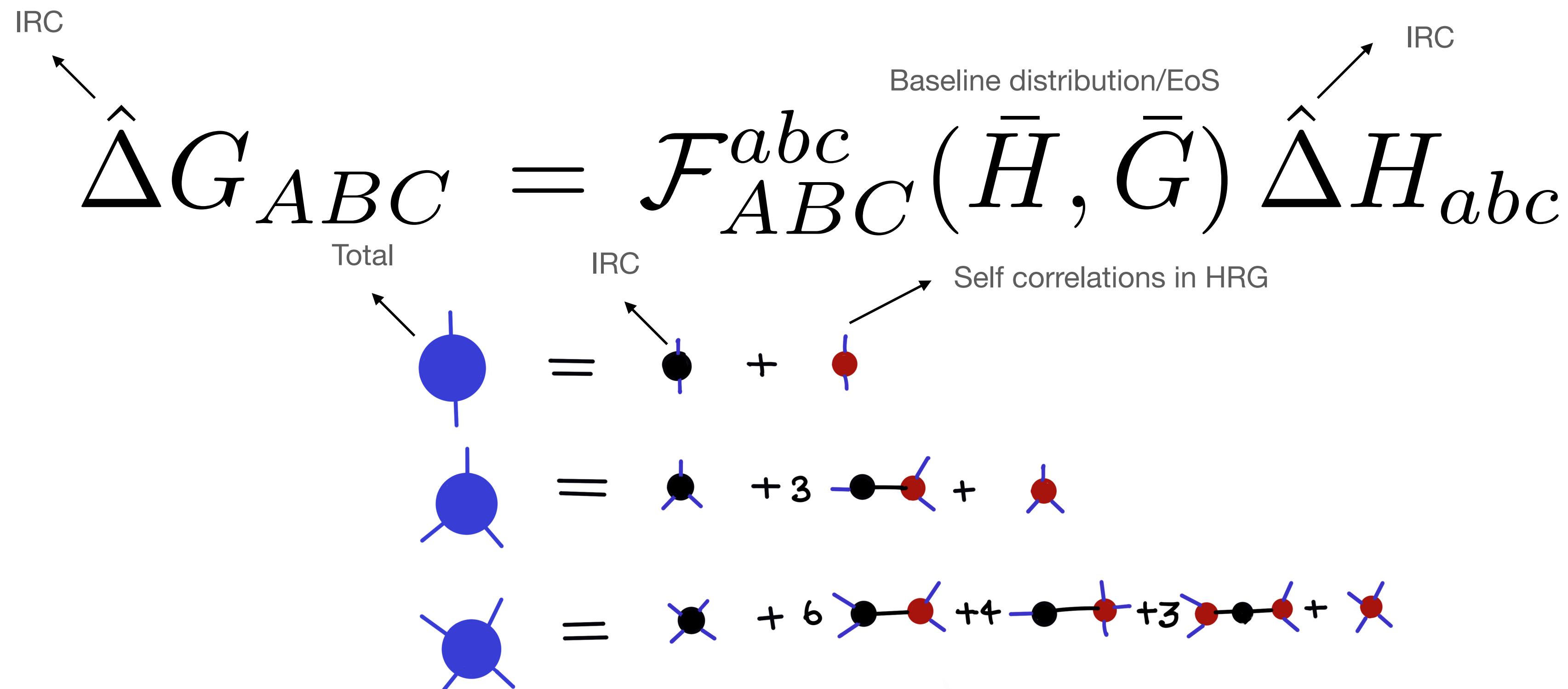
$$\Delta G_{pp}(k, k') = \#_{pp}^{\epsilon\epsilon}(k, k') \Delta H_{\epsilon\epsilon} + \#_{pp}^{n\epsilon}(k, k') \Delta H_{n\epsilon} + \#_{pp}^{nn}(k, k') \Delta H_{nn} + \dots$$

Contributions from baryon and energy density correlations

Δ denotes deviation of the hydrodynamic correlation function from its value in Ideal Hadron Resonance Gas

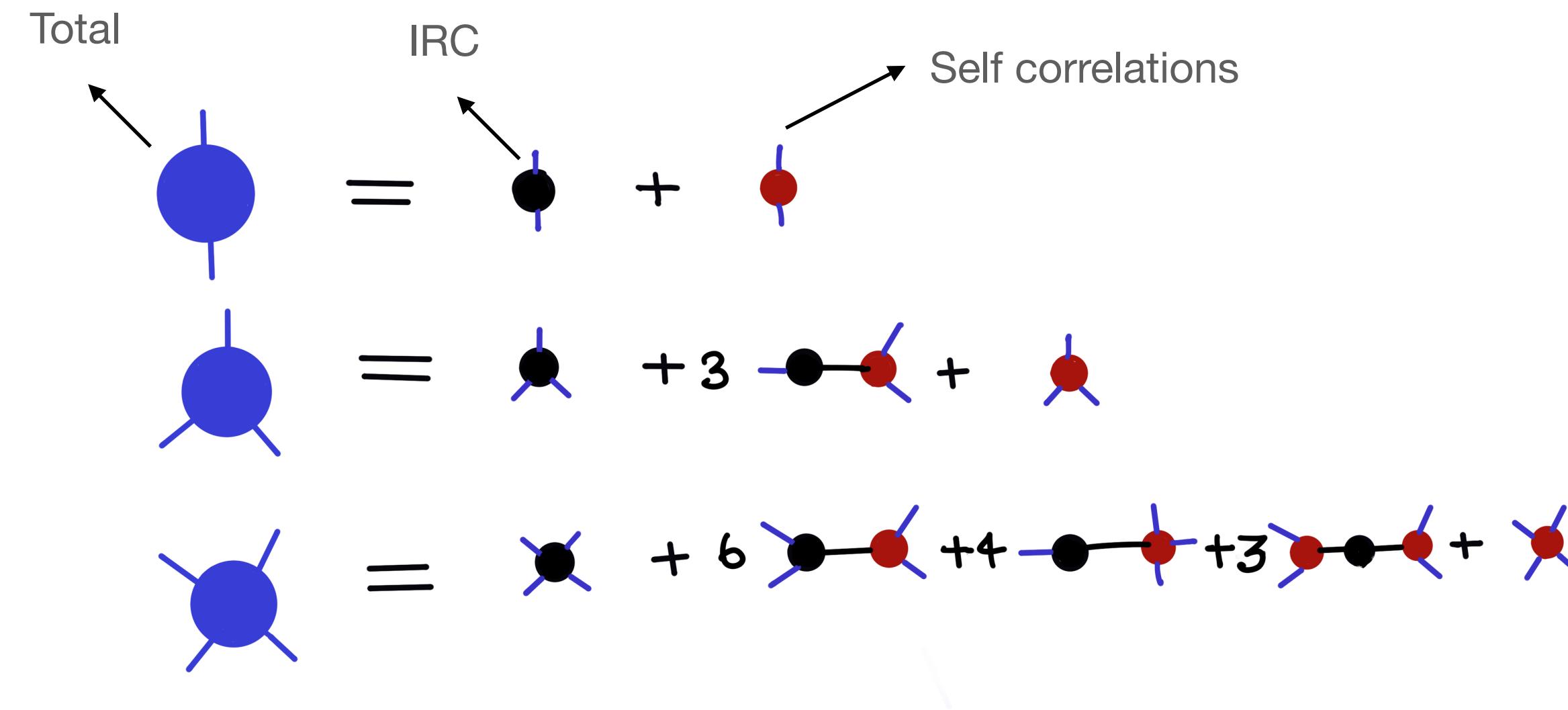
Organizational Scheme for Non-Gaussian Correlations

$$G_{nnn}(p, p', p'') - \bar{G}_{nnn}(p, p', p'') = \sum_{a,b,c=\epsilon,n} \#_{pp'p''}^{abc} \Delta H_{abc} + \sum_{a,b=\epsilon,n} \#_{pp'p''}^{ab} \Delta H_{ab} + \dots$$



Irreducible relative cumulants

For classical gas, irreducible relative cumulants (IRCs) reduce to so called “factorial cumulants”.



- For gases obeying different statistics, IRCs quantify the non-trivial correlations
- Non-trivial correlations relative to any specified baseline distribution

Maximum entropy freeze-out procedure

$$\hat{\Delta} \langle \delta N_{A_1 \dots A_k}^k \rangle = X_{A_1}^{b_1} X_{A_2}^{b_2} \dots X_{A_k}^{b_k} \hat{\Delta} H_{b_1 \dots b_k}$$

- Integrating over phase space bins, we get factorial cumulants of particle multiplicities
- Factorial cumulants remove various sources of trivial correlations
- They contain crucial information about criticality in the Equation of State / Hydro

MP, Stephanov, PRL, 23

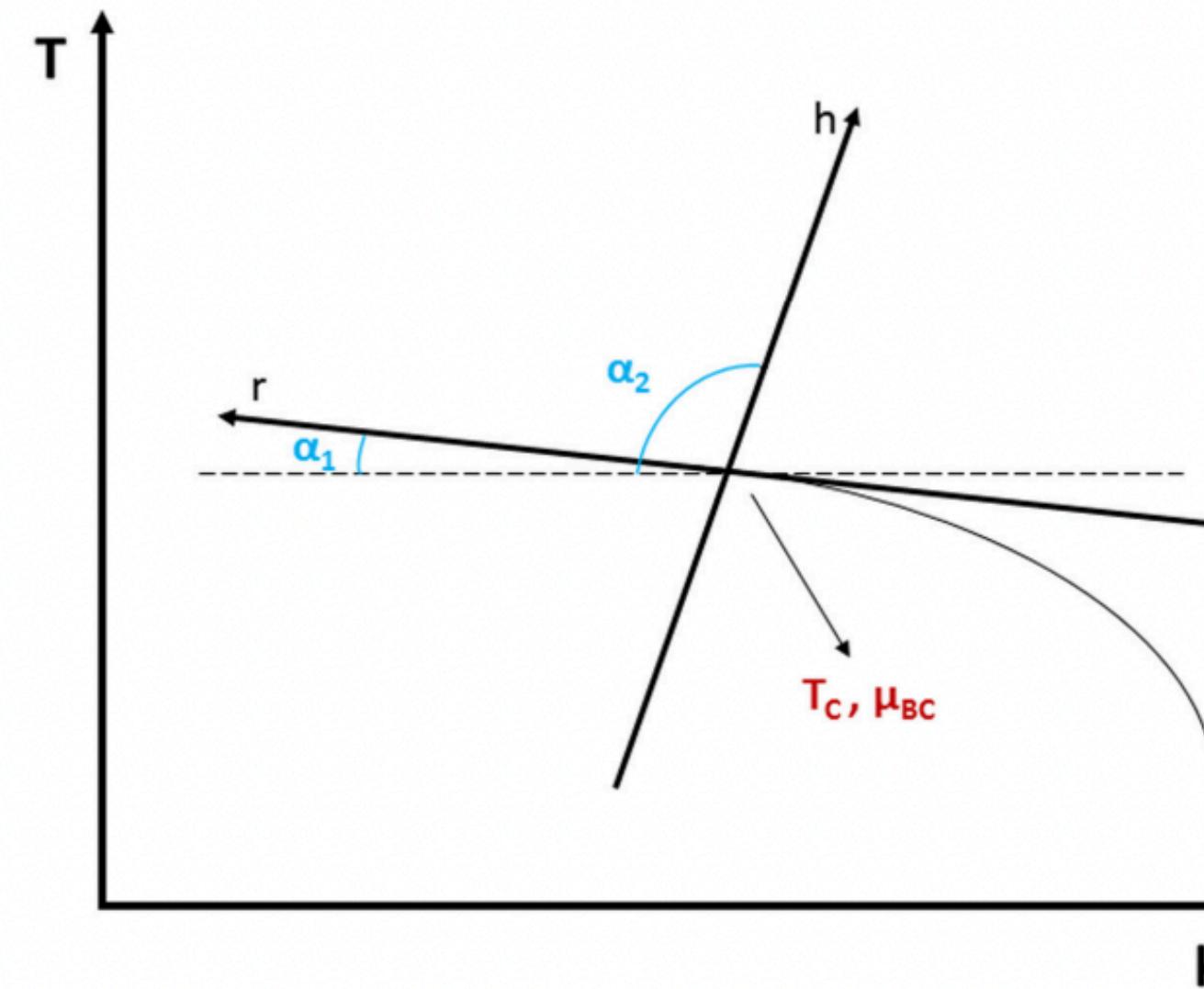
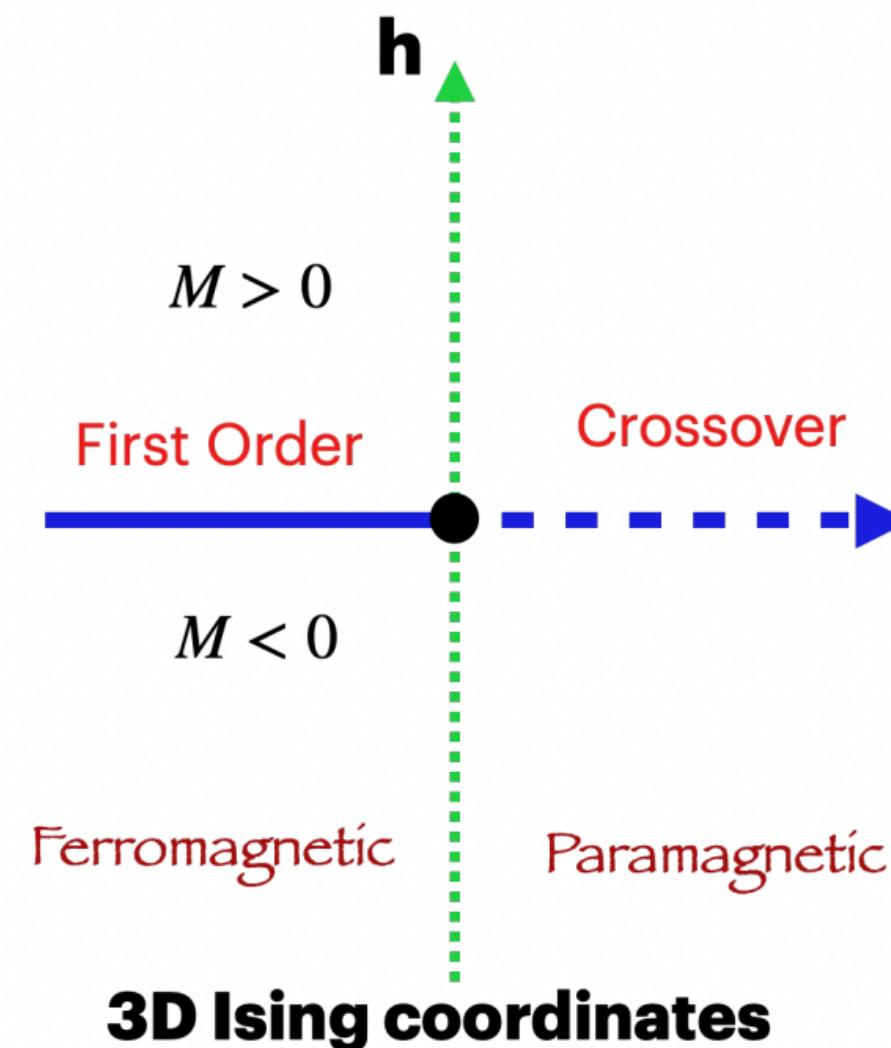
But, the QCD EoS near the critical point is not known from first-principles?!

Equation of States with a QCD critical point

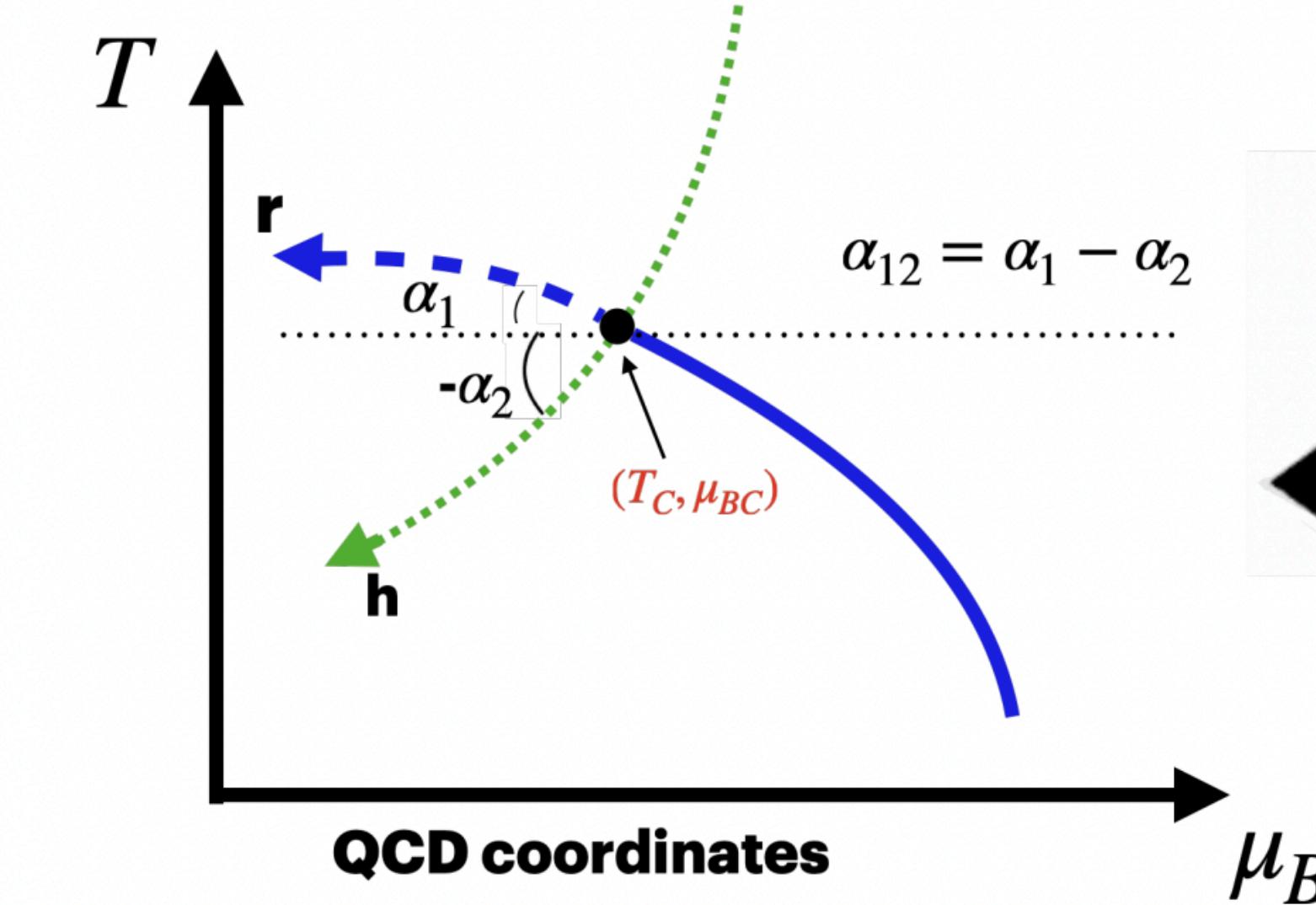
- Must agree with the Taylor Expanded EoS from lattice
- Compatible with other limits : PQCD, HRG
- Critical point in the 3D Ising universality class

Examples of recently developed EoSs that have a CP in the Ising universality class but differ in their implementation:
Parotto et al, 19, Karthein et al.,21, Grefa et al., 21,Kapusta&Welle,22, Kahangirwe et al.,24

QCD EoS near the Critical Point



Parotto et al., 18, Karthein et al., 21



Kahangirwe et al., 24

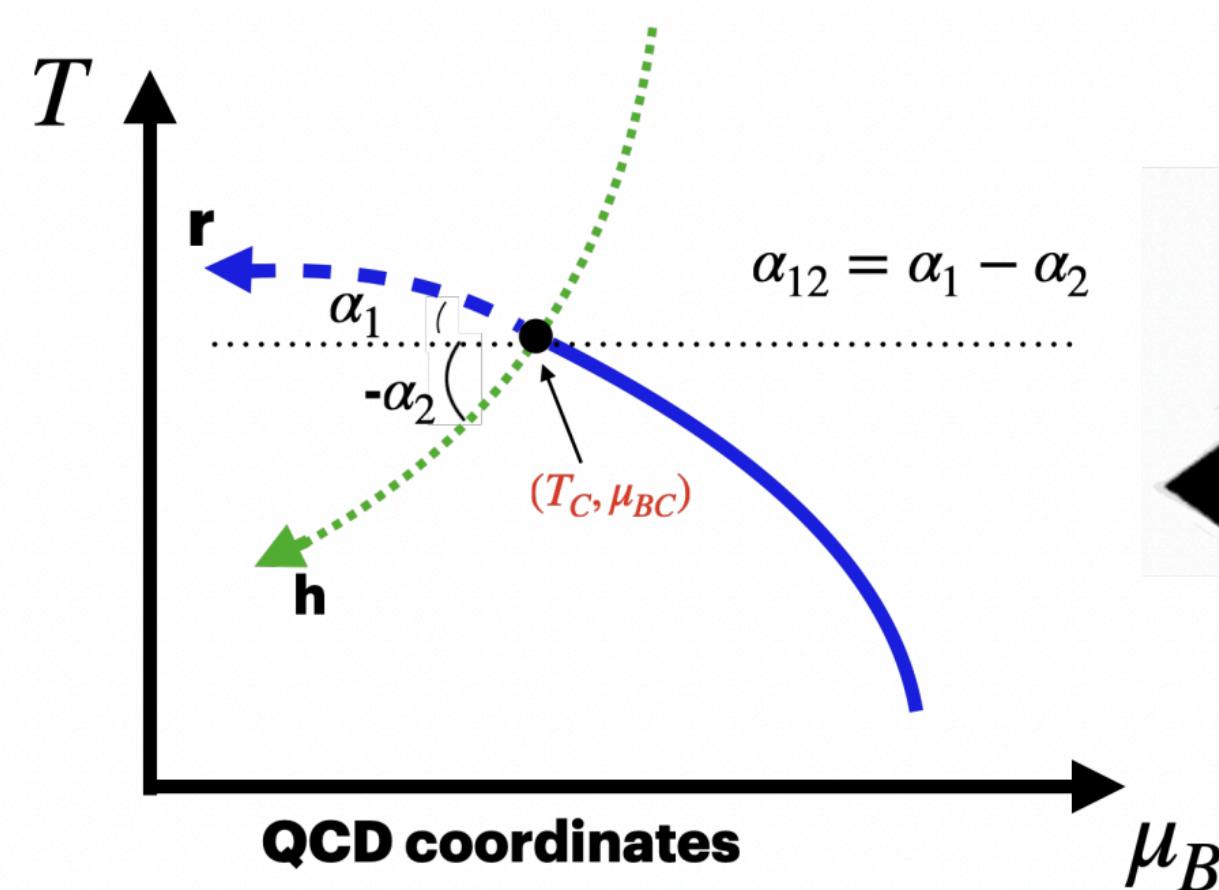
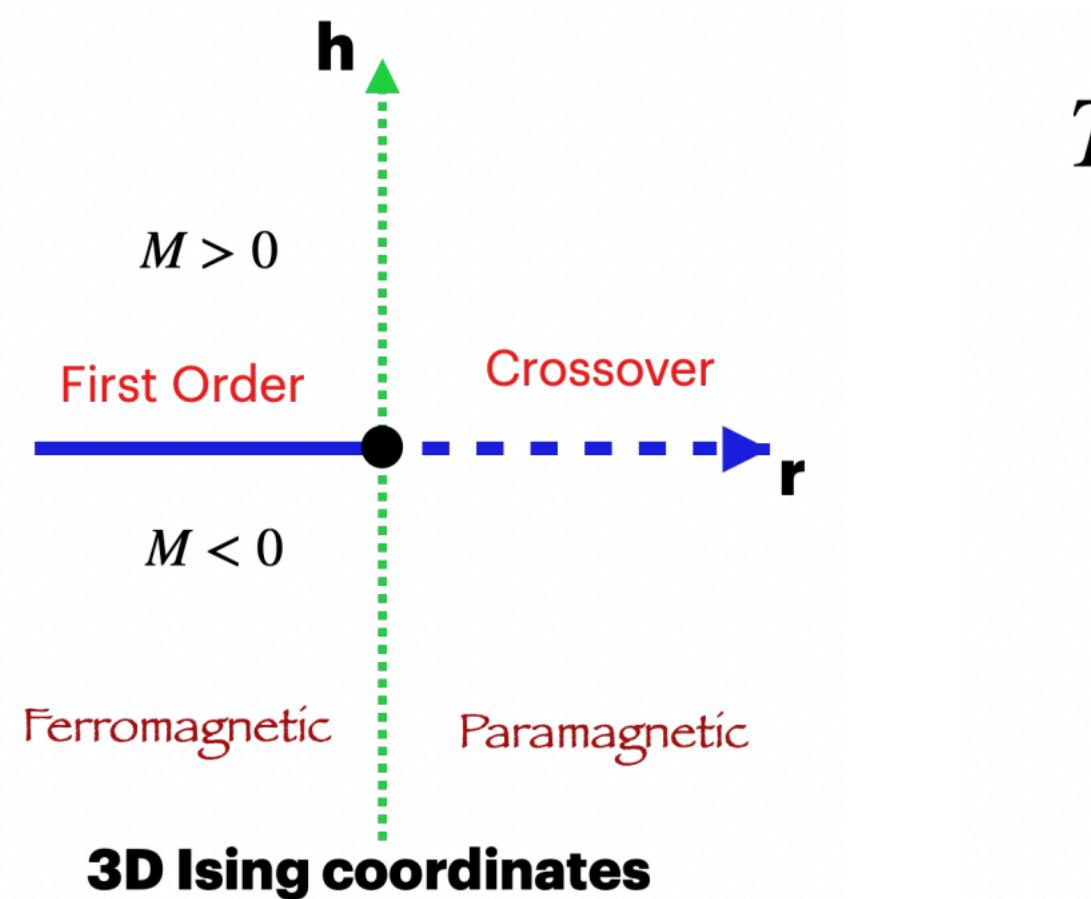
Summation scheme by WB collaboration
Borsanyi et al, 21

Non-universal map from QCD to Ising variables

$$P_{\text{QCD}}(\mu, T) = P_{\text{BG}}(\mu, T) + A G(r(\mu, T), h(\mu, T))$$

A general class of candidate EoSs

$$P_{\text{QCD}}(\mu, T) = P_{\text{BG}}(\mu, T) + A G(r(\mu, T), h(\mu, T))$$



Independent & non-universal parameters

μ_c , α_{12} , ρ , w

Weakly constrained in
the chiral limit

MP, Stephanov, 19

23

Kahangirwe et al., 24

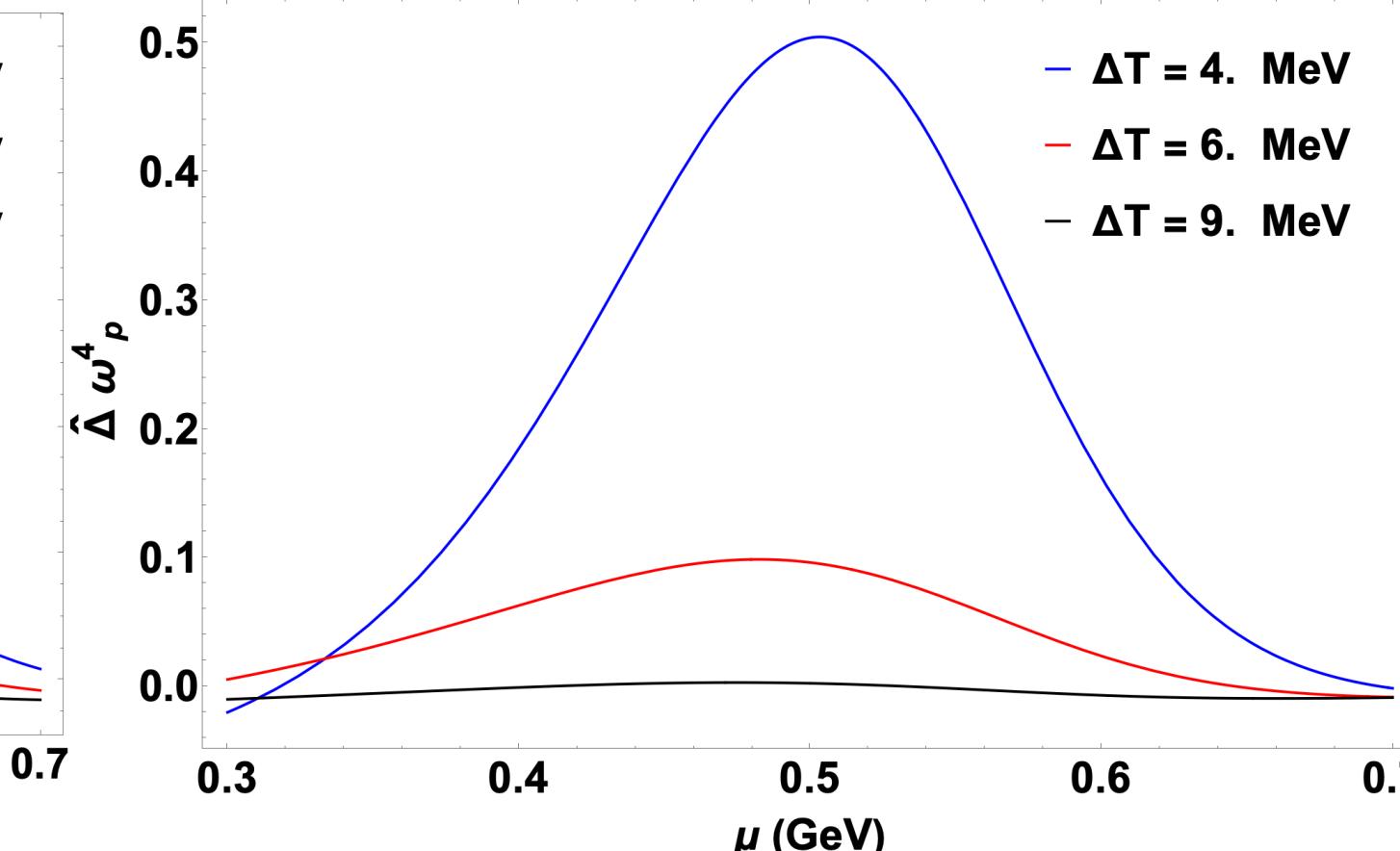
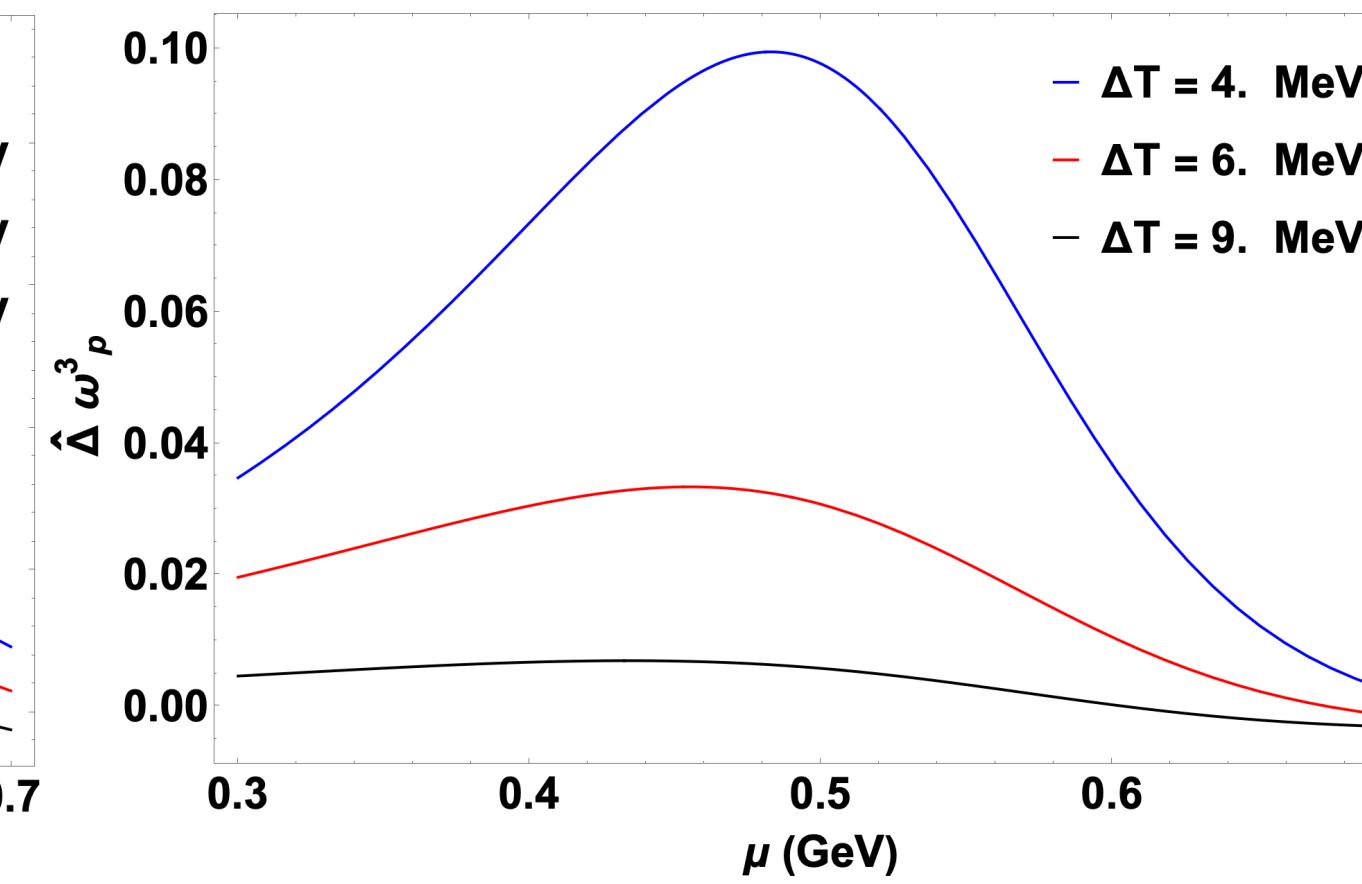
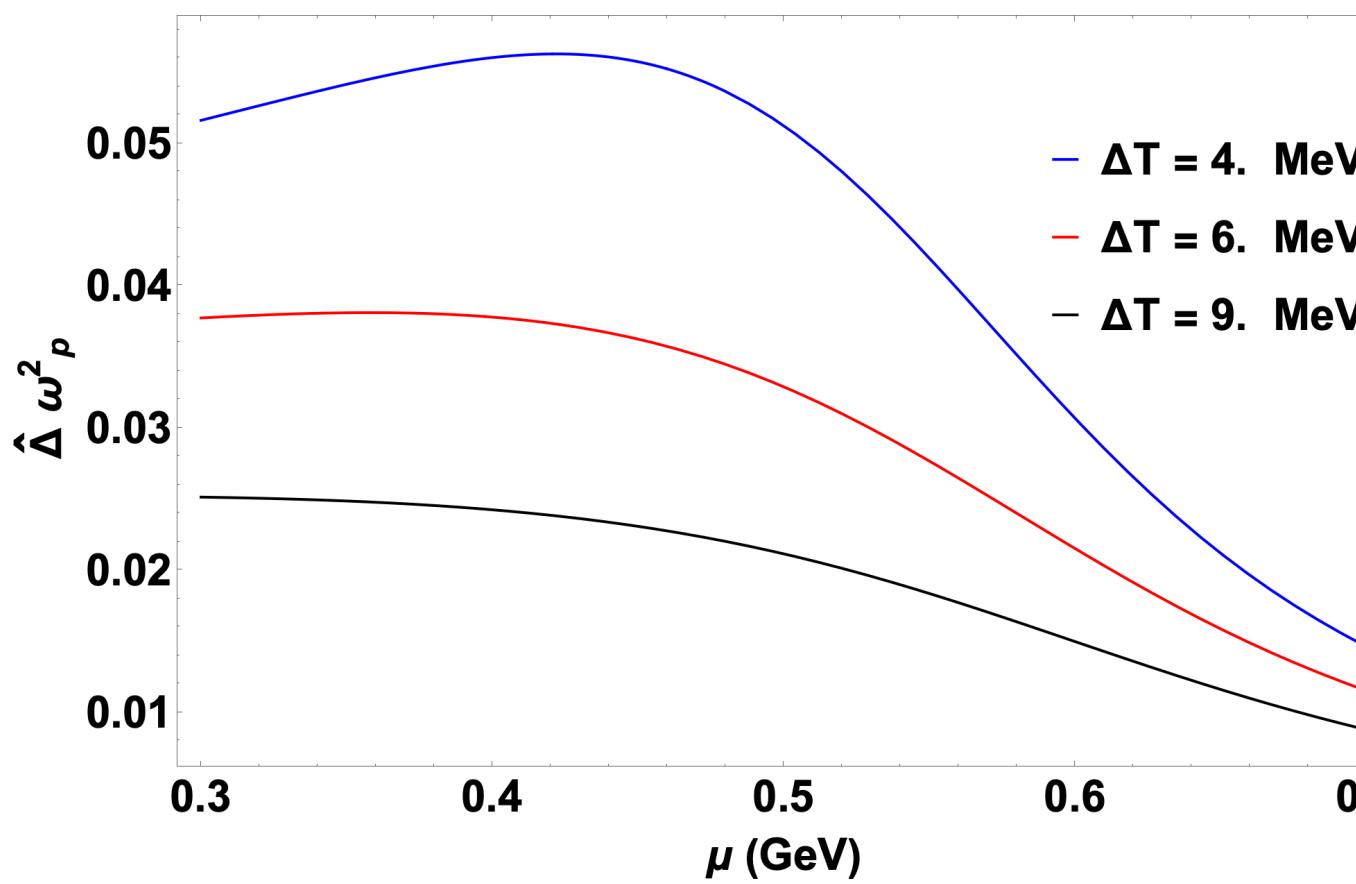
Range of Validity improved

$$0 \leq \mu_B \leq 700 \text{ MeV}, 25 \text{ MeV} \leq T \leq 800 \text{ MeV}$$

The new construction is causal and stable for a larger range of ρ and w

Equilibrium estimates for critical contribution to factorial cumulants of proton multiplicities

Karthein, **MP**, Rajagopal, Stephanov, Yin (in preparation)

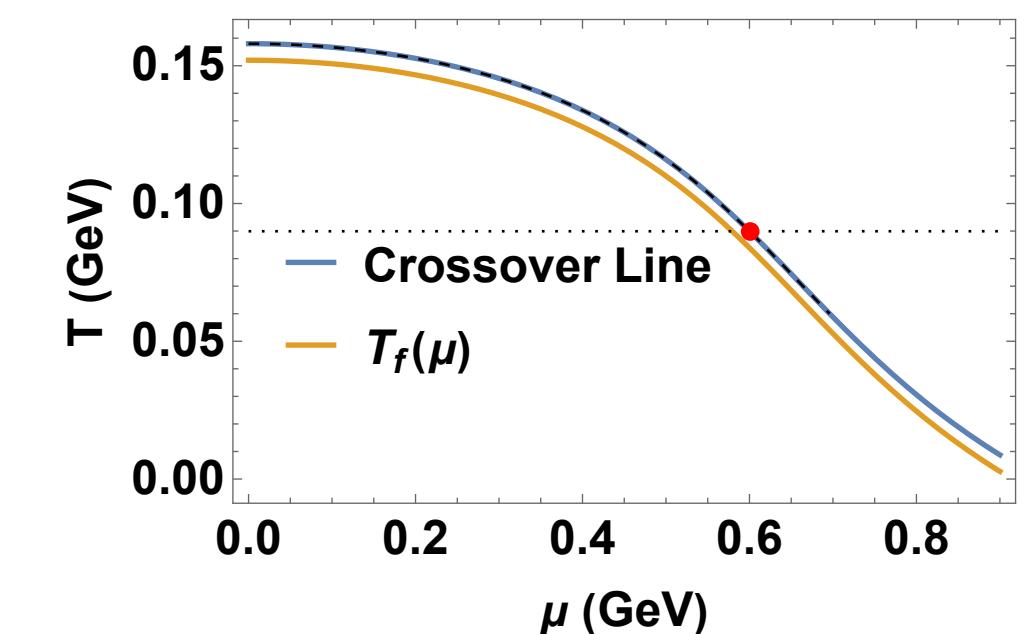


$\mu_c = 600$ MeV, $\alpha_2 = 0^\circ$, $\rho = 1$, $w = 20$

Example choice of Mapping Parameters

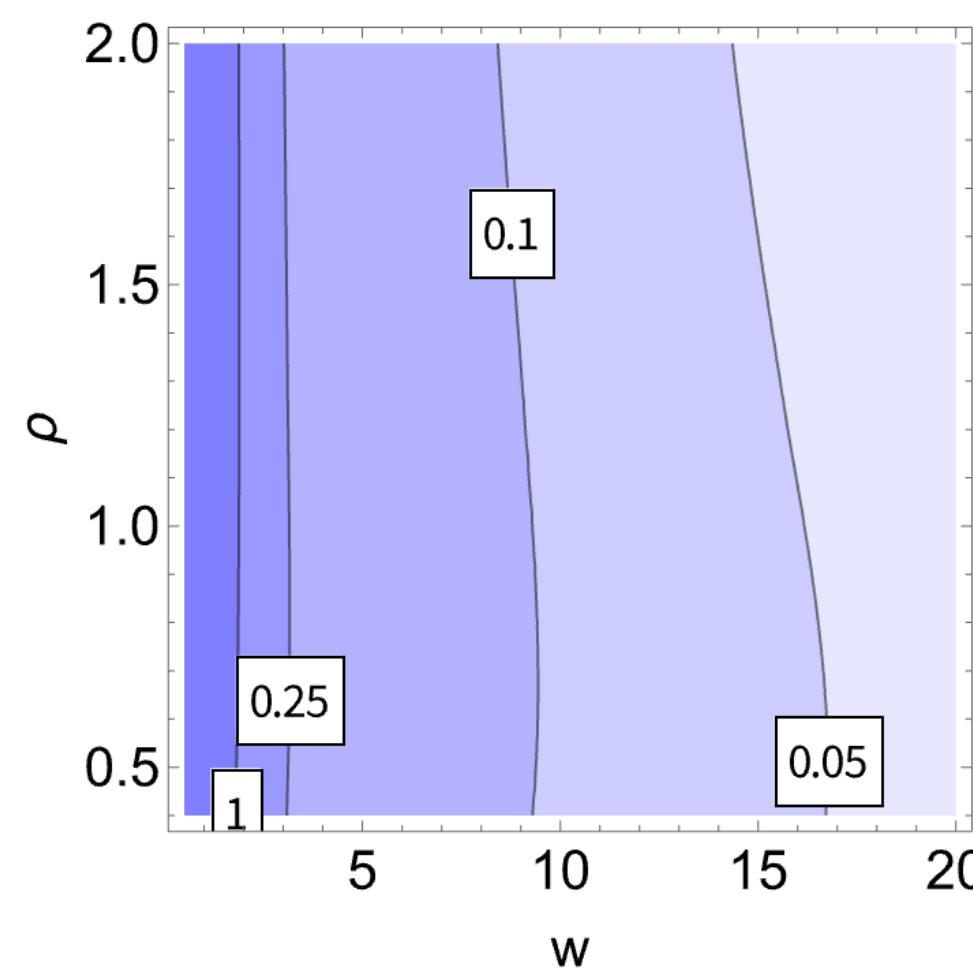
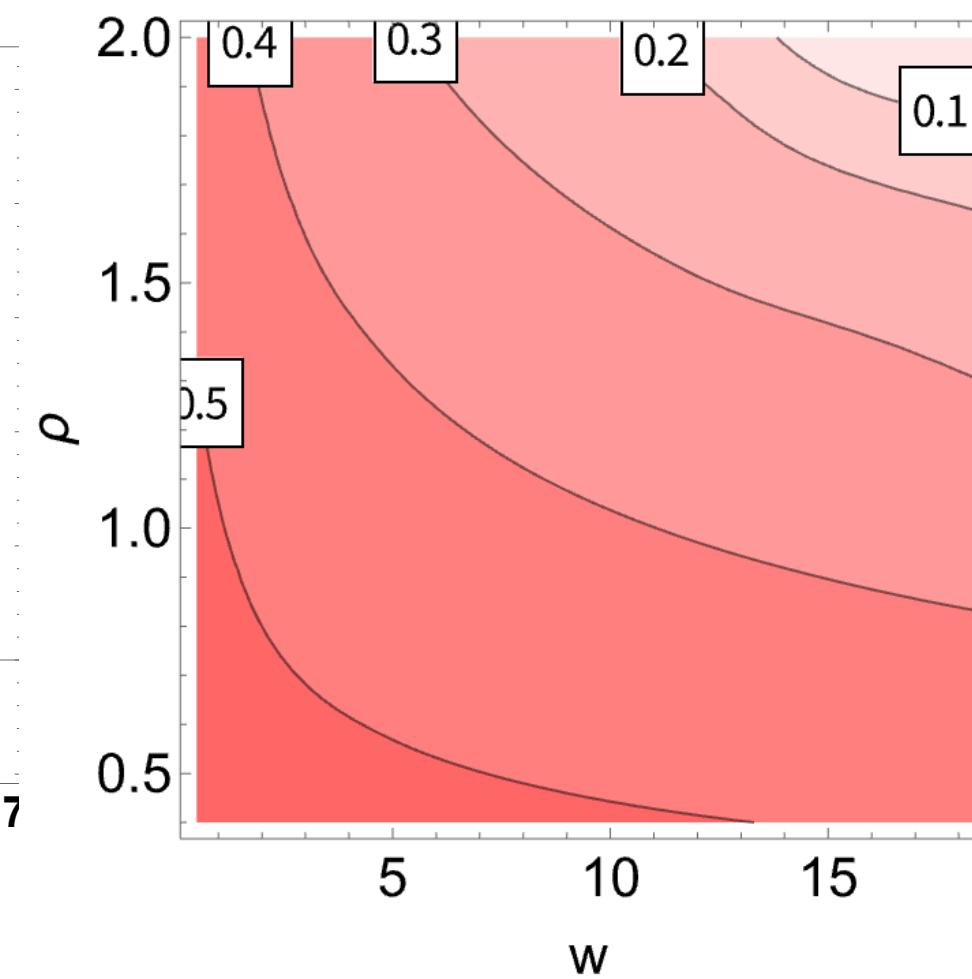
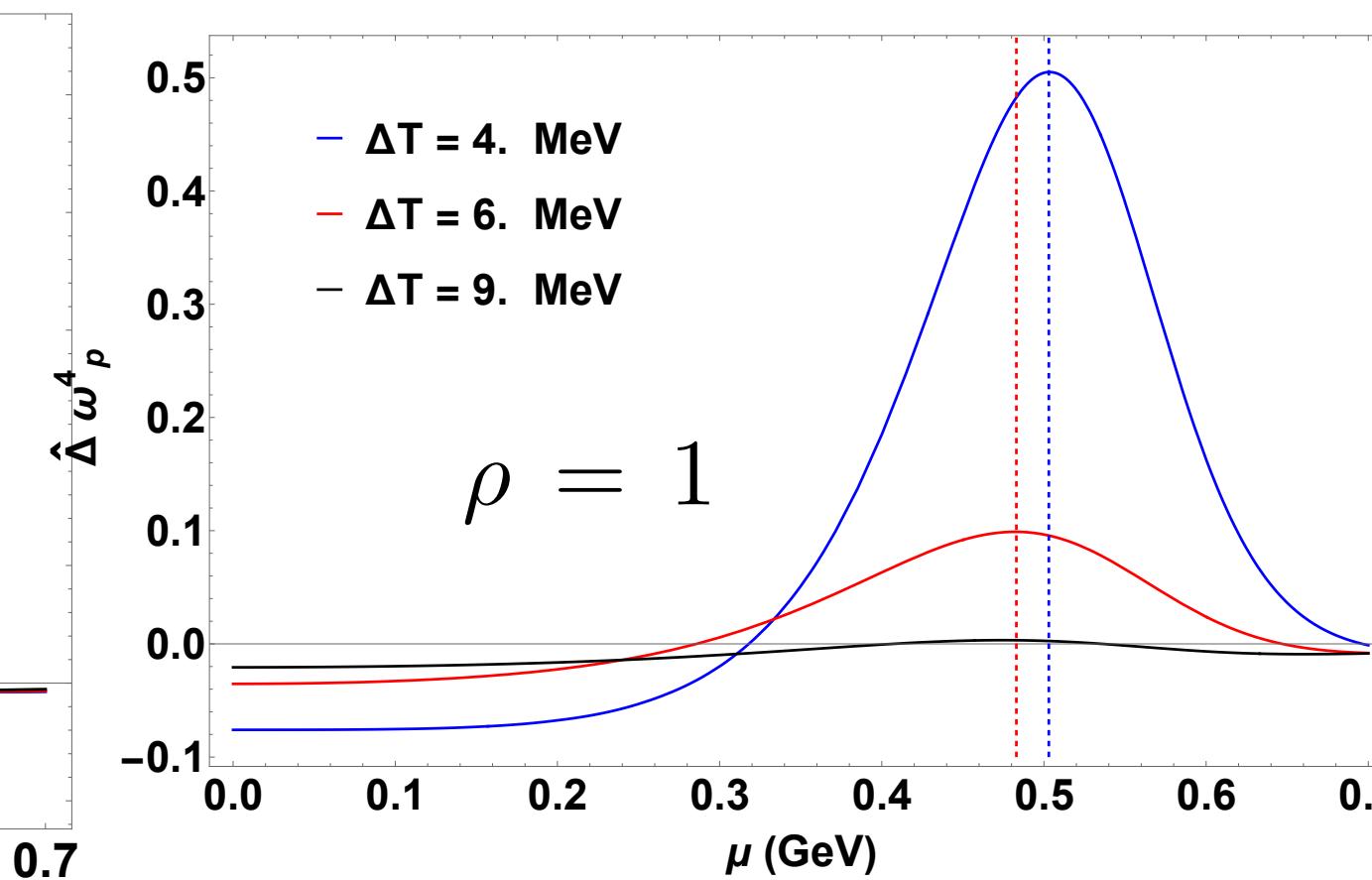
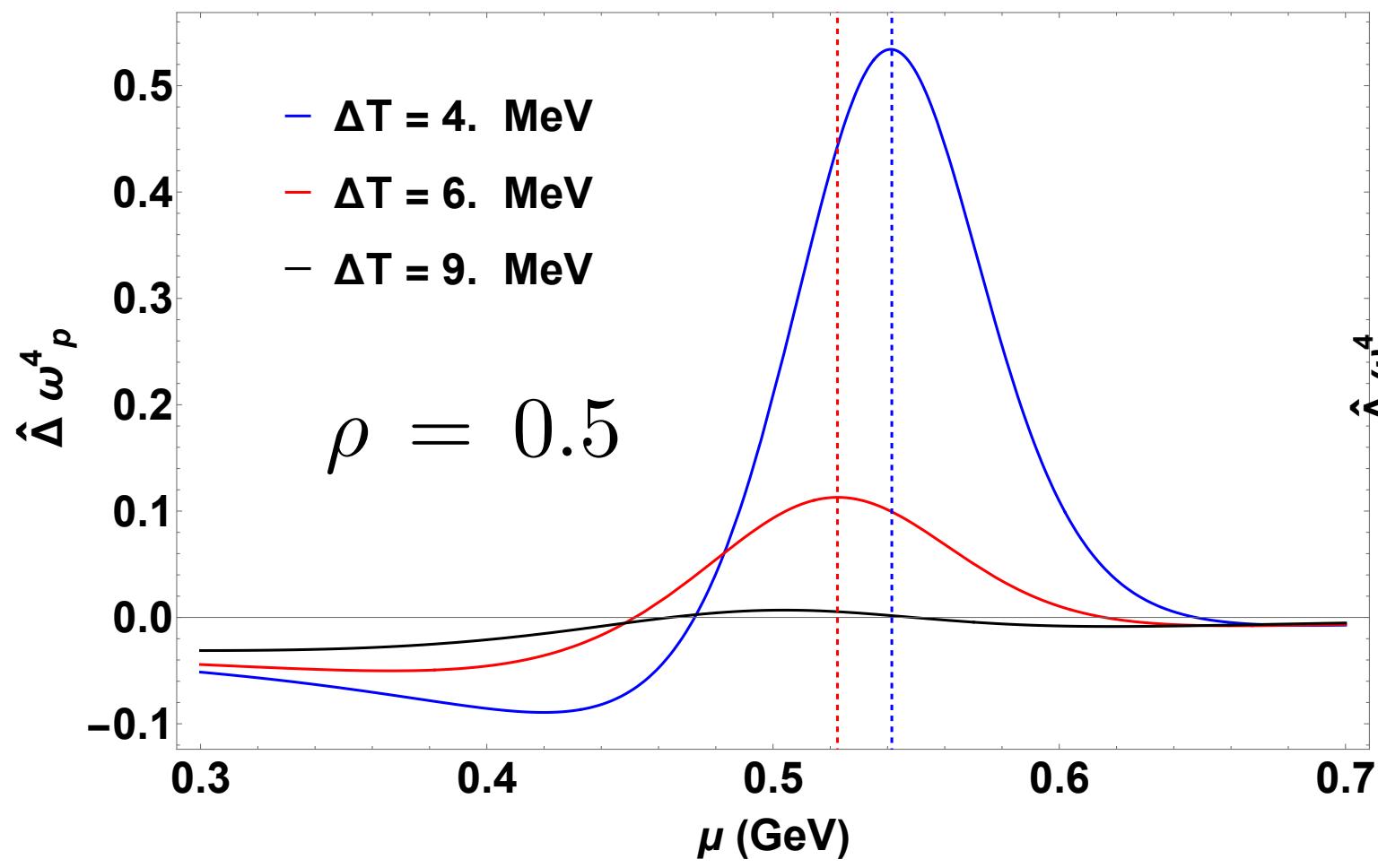
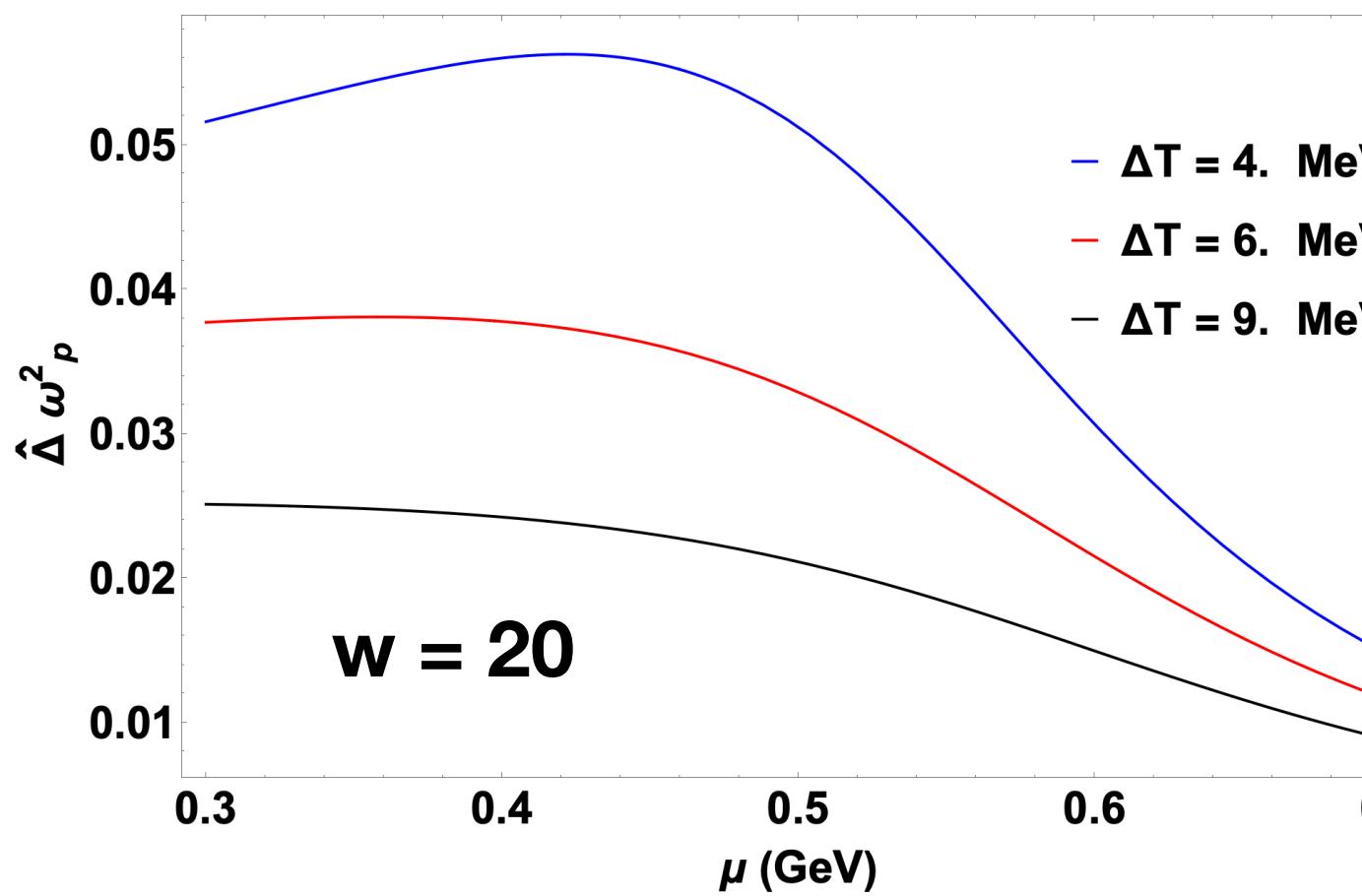
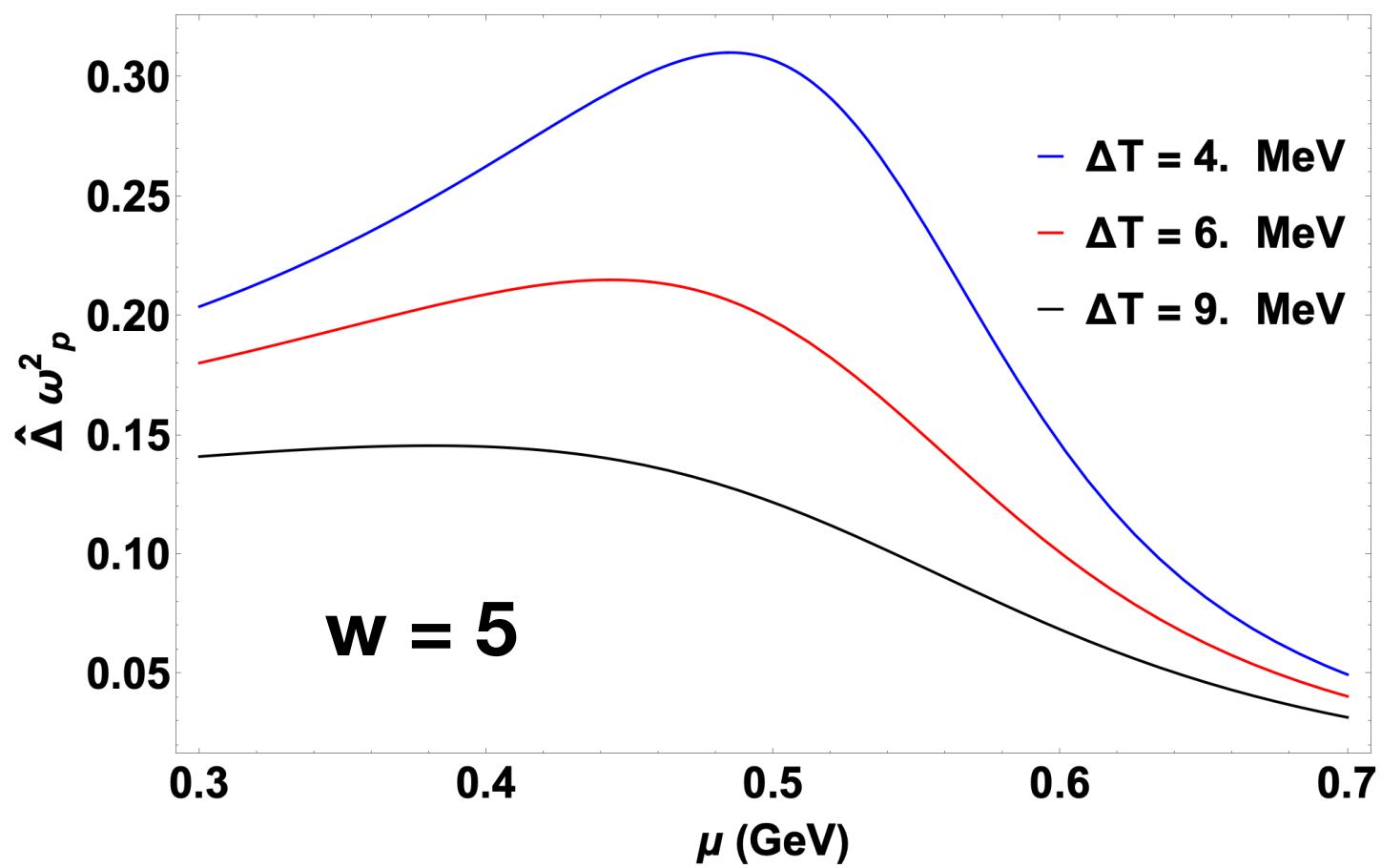
$$\hat{\Delta}\omega_p^k = \frac{\hat{\Delta} \langle \delta N_p^k \rangle}{\langle N_p \rangle}$$

$$T_f(\mu) = T_{\text{co}}(\mu) - \Delta T$$



Freeze-out curve

Role of scale factors



- Peak value of cumulants along the freeze-out curve is fixed by w
- Location of the maxima on the freeze-out curve fixed by ρw

Fluctuation dynamics

Stochastic Approach

- Stochastic differential equations for hydro fields
- Implementation using Metropolis algorithm developed and applied to various models, including Model H

Schaefer and Skokov, 22,
Chattopadhyay et al, 23, Florio
et al, 21, Basar et al,
24, Chattopadhyay et al, 24...

Deterministic Approach

- Deterministic evolution for hydro correlation functions
- Semi-realistic estimates for two-point correlation functions computed and connected to observables

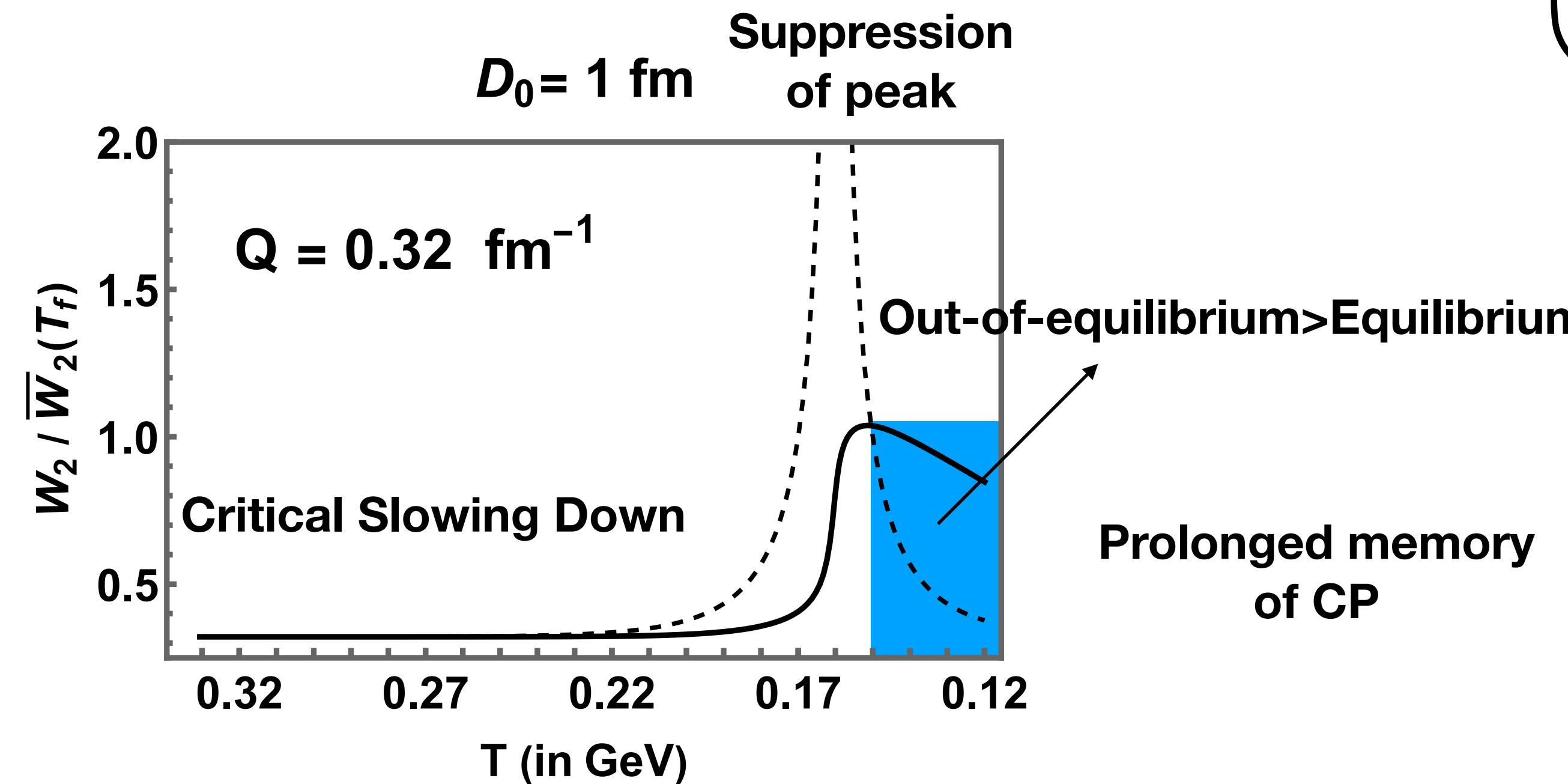
Teaney, Akamatsu, Mazeliauskas, 16 along with Yan, Yin , 19,
Stephanov, Yin, 17, An, Basar, Stephanov, Yee, 19, 20, 22, 24...

An, Basar, Stephanov, Yee, 22, 24
Rajagopal, Ridgway, Weller, Yin, 19, Du, Heinz, Rajagopal, Yin, 20
MP, Rajagopal, Stephanov, Yin, 22, Mukherjee, Venugpalan, Yin 15

Refer to Johannes Roth's talk

Out-of-equilibrium effects of fluctuations near the CP

$$\langle \delta\hat{s}(x_+) \delta\hat{s}(x_-) \rangle = \int e^{i\mathbf{Q} \cdot \Delta\mathbf{x}} W_2(\mathbf{Q}), \Delta\mathbf{x} = x_+ - x_-$$



Persistence of critical imprints in the fluctuation observables until freeze-out

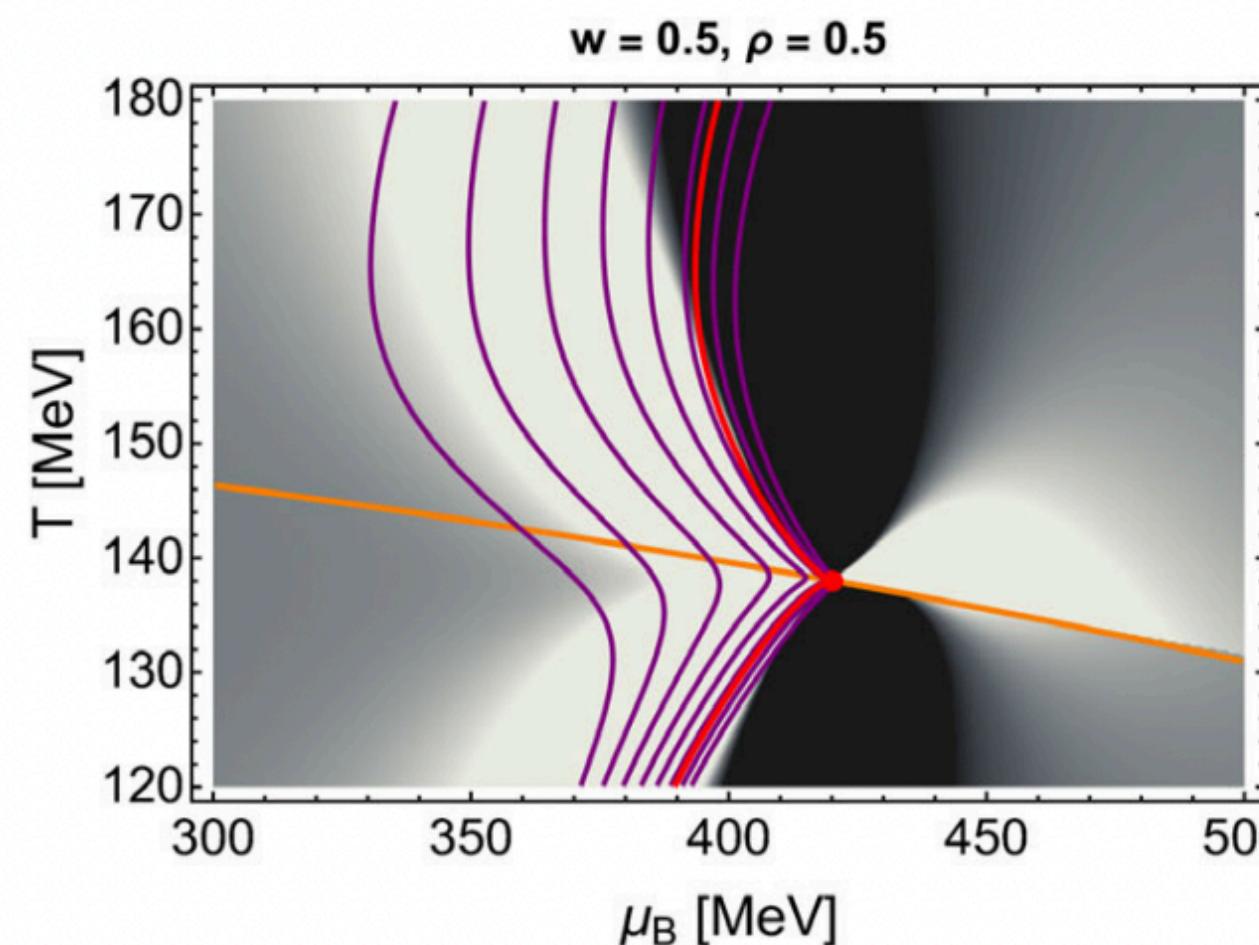
Rajagopal,Ridgway,Weller,Yin,19
Du,Heinz,Rajagopal, Yin,20
MP, Rajagopal, Stephanov, Yin, 22
Mukherjee, Venugopalan, Yin 15

CP fluctuations in MD simulations : V. Kuznetsov et al

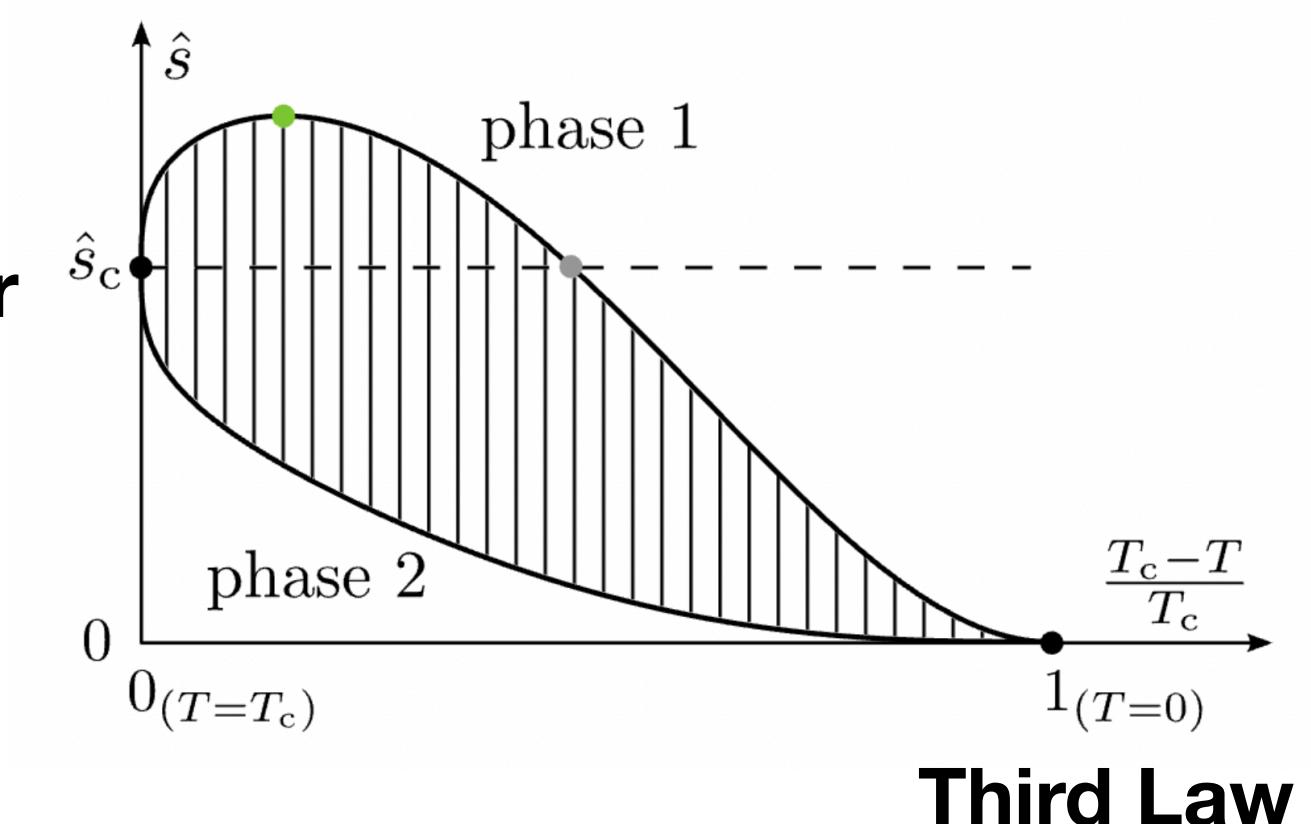
Deformation of hydrodynamic trajectories near CP

MP, Sogabe, Stephanov, Yee,24

Deformations can be broadly classified based on the value of the mapping parameter α_2



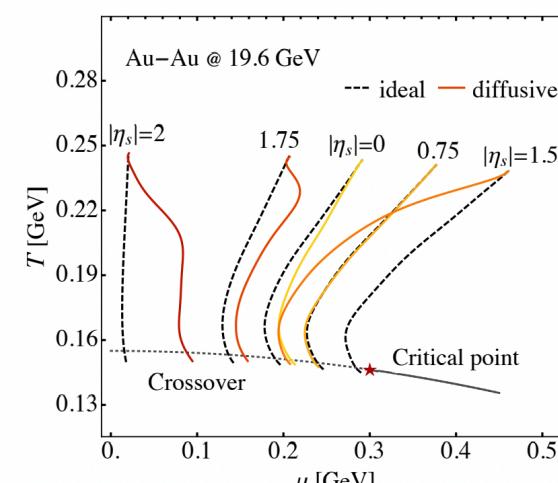
Universal scaling near CP



- Specific entropy is non-monotonic along one of the branches on the first-order curve

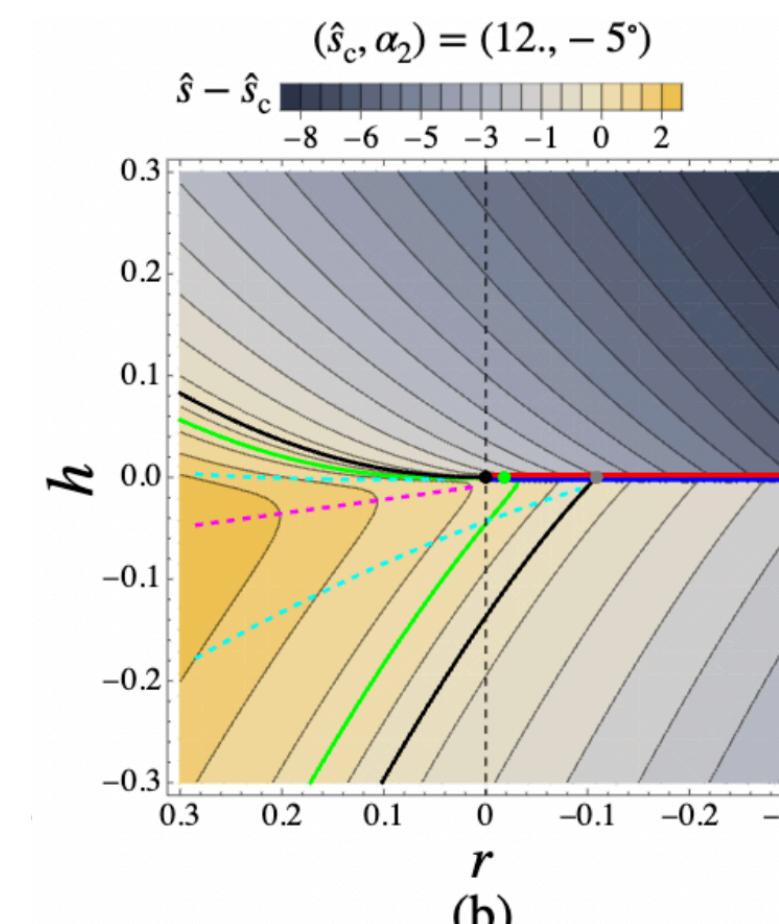
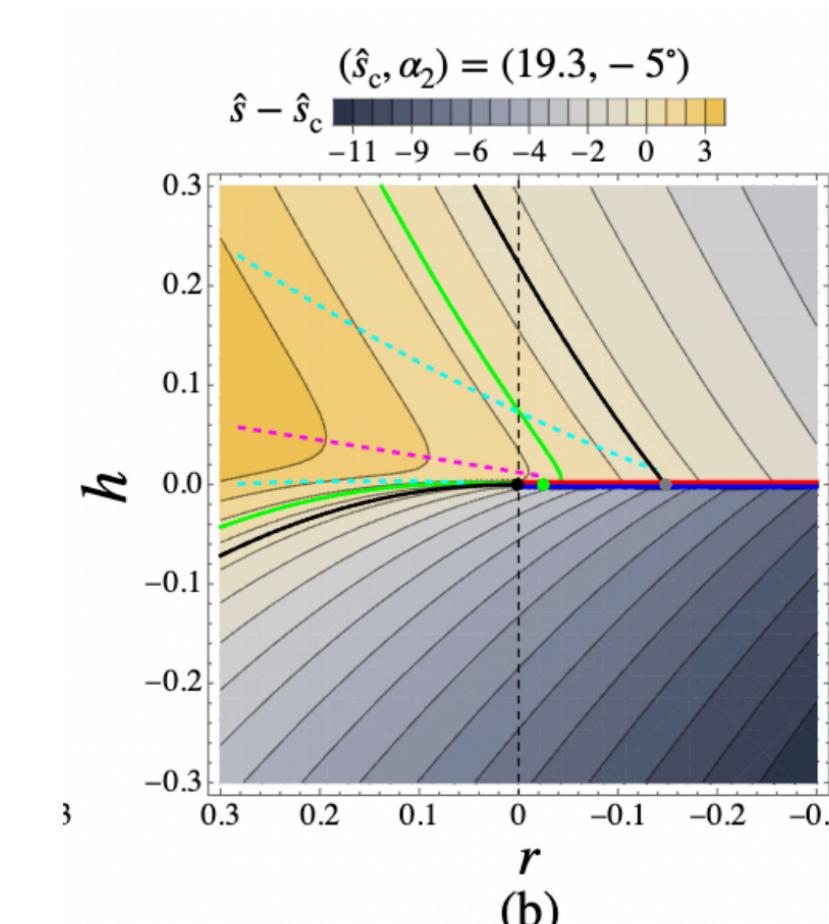
Critical lensing~Dore et al,22,
Nonaka&Asakawa, 05

- Consequence of universal ridge-like structure of the isentropes near CP



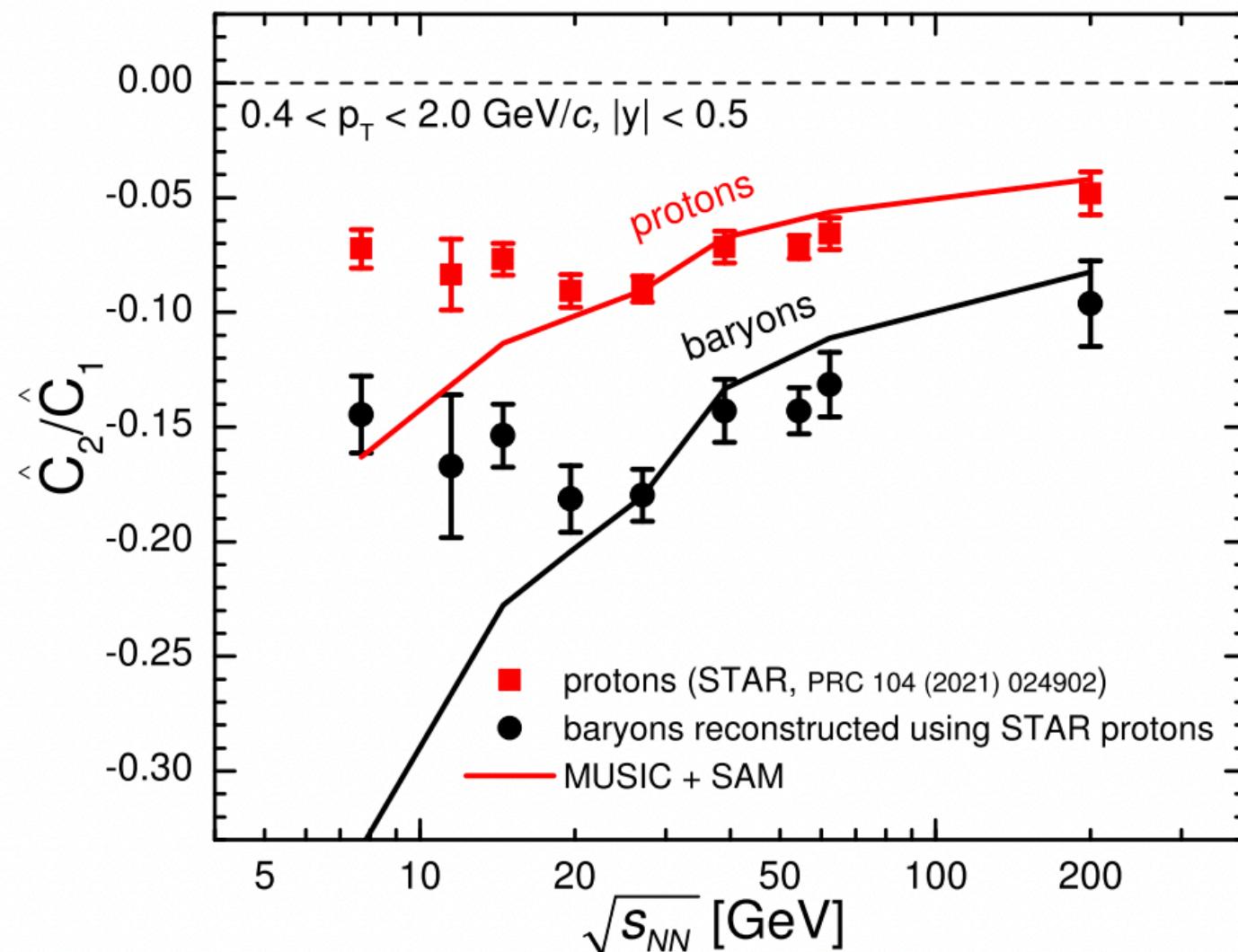
Smearing effect - Du et al, 22

- Phenomenological implications



- Critical point
- Maximum
- Critical double
- $h = +0$
- $h = -0$
- $r = 0$
- $\hat{s} = \hat{s}_c$
- $\hat{s} = \hat{s}_{\max}$
- Ridge curve
- Valley curve

Summarizing & Looking forward



$$\chi_j(\mu, T; \mu_c, \alpha_{12}, w, \rho) \xrightarrow{\text{ME}} C_A^k(\mu_F(T_F); \mu_c, \alpha_{12}, w, \rho, \Gamma)$$

Dynamics

A *family of candidate EoSs with a CP* that match with the lattice have been developed

Quantitative estimates for *out-of equilibrium* corrections to higher order cumulants - needed

Vovchenko, Koch, Shen, 22

EoS \longleftrightarrow Particle multiplicity distributions

Bayesian Analysis of experimental data pertaining to *multiple observables* with the theoretical framework may possibly help us learn about QCD EoS near CP, if it exists in the regime scanned by HICs

Thank you!

Back up

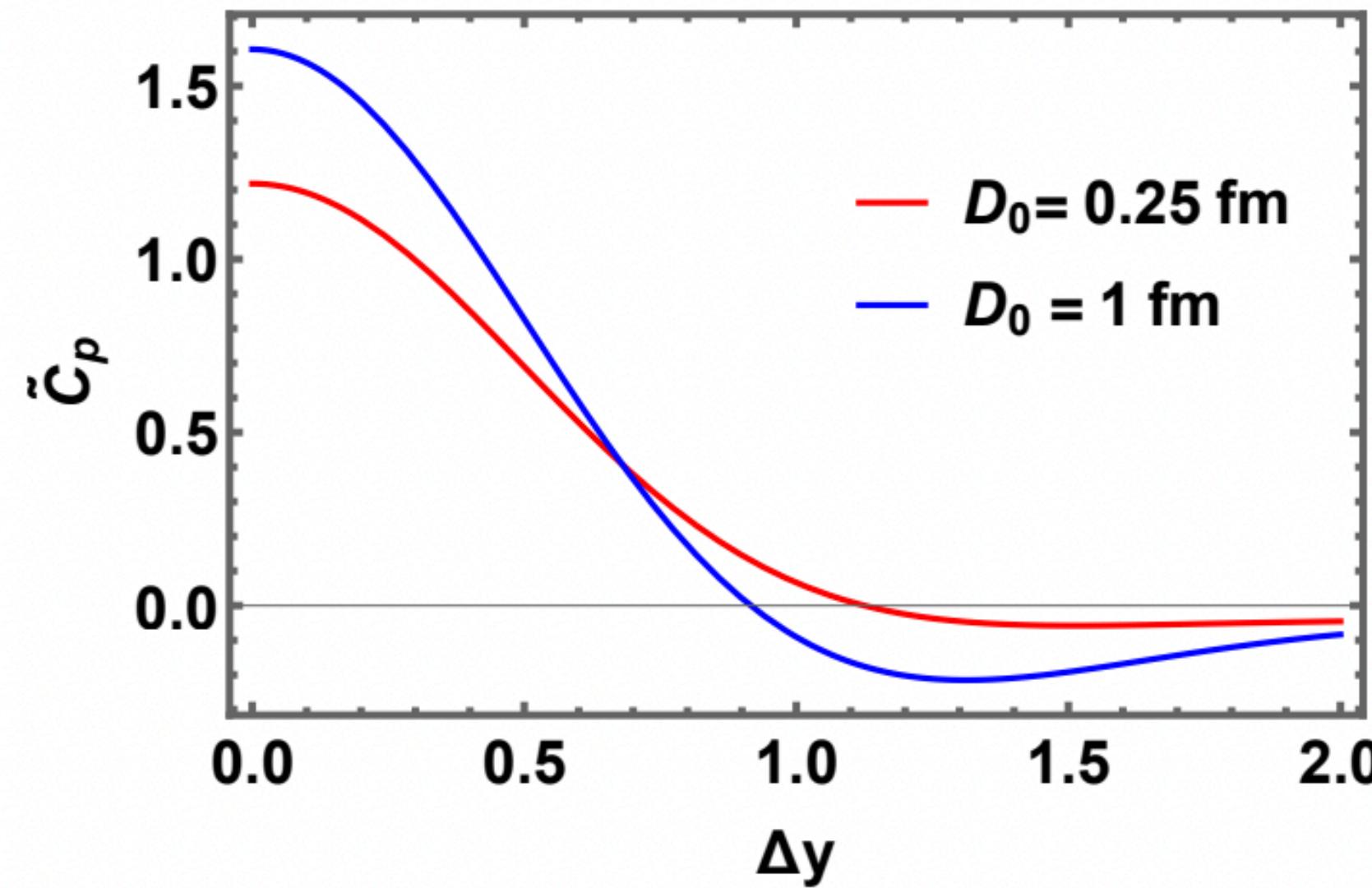


FIG. 6: Normalized proton multiplicity correlator $\tilde{C}(\Delta y)$ for protons from Eq. (60) as a function of the rapidity gap Δy in the Bjorken scenario for two choices of the diffusion parameter D_0 . **MP, Rajagopal, Stephanov, Yin, 22**