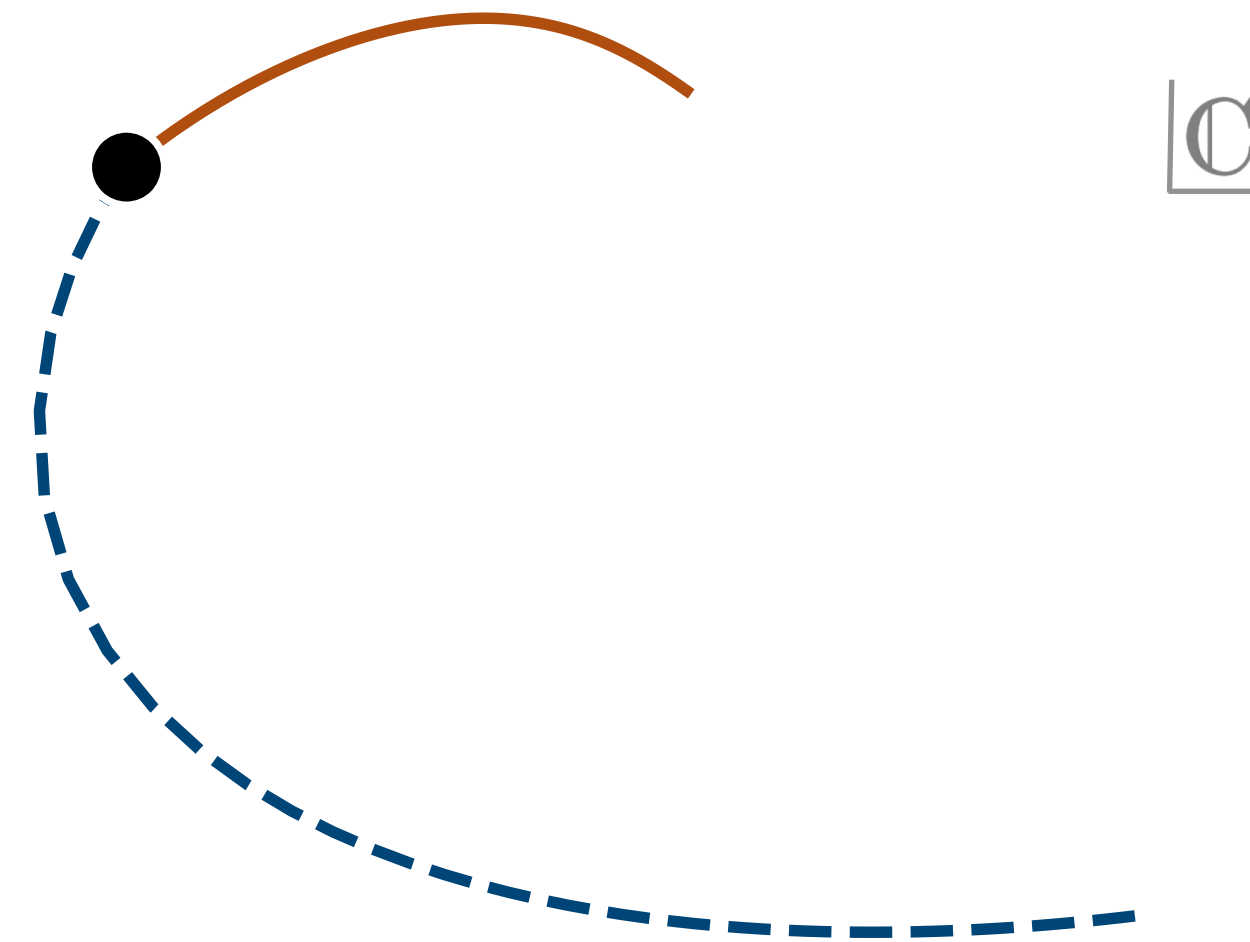


# Quasinormal modes and their excitation beyond general relativity

Hector O. Silva  
Max Planck Institute for Gravitational Physics  
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🏠 <https://www.phy.olemiss.edu/~hosilva>



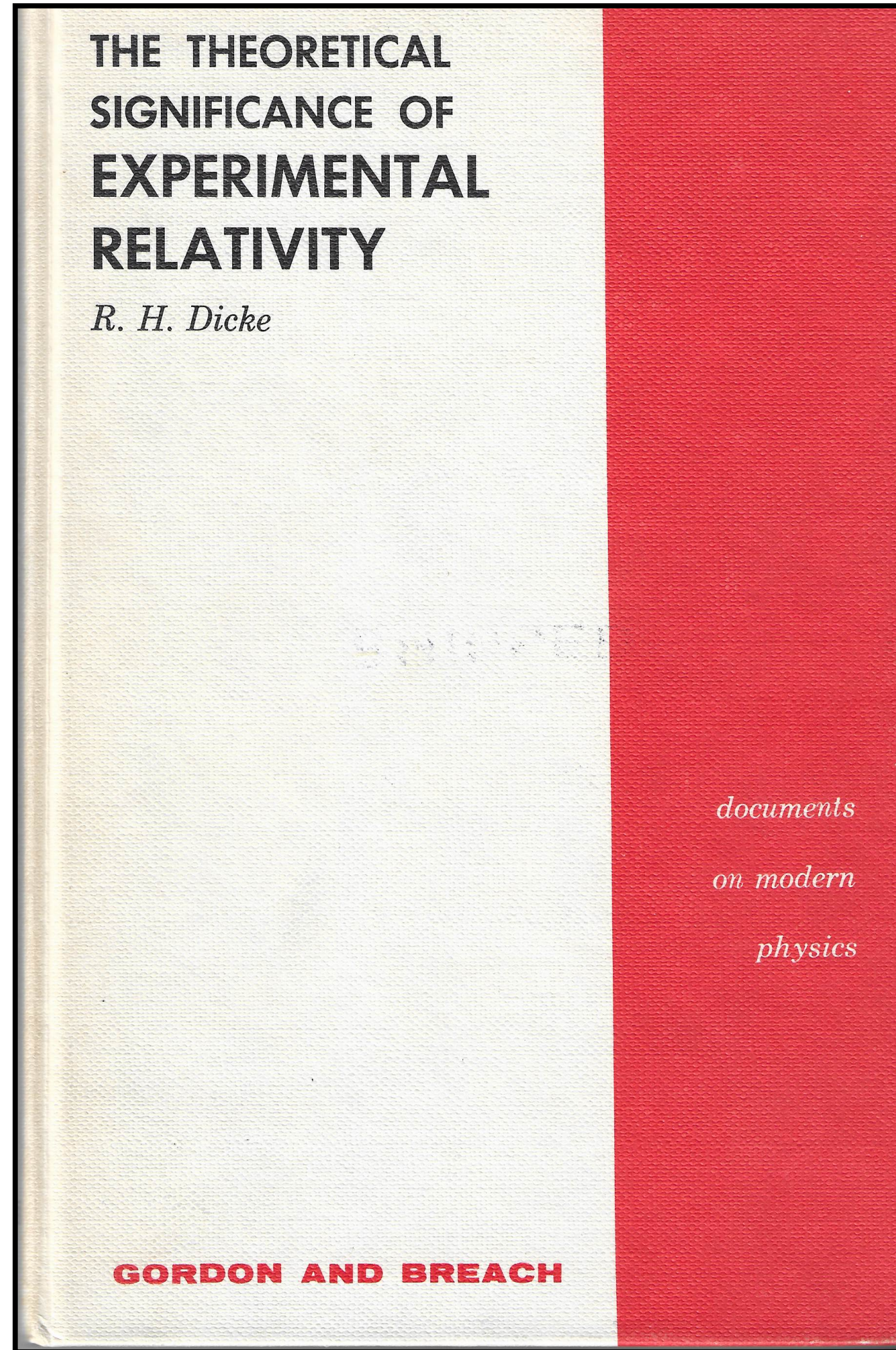
MAX-PLANCK-GESELLSCHAFT

New Frontiers in Strong Gravity  
Benasque, Spain  
12.07.2024



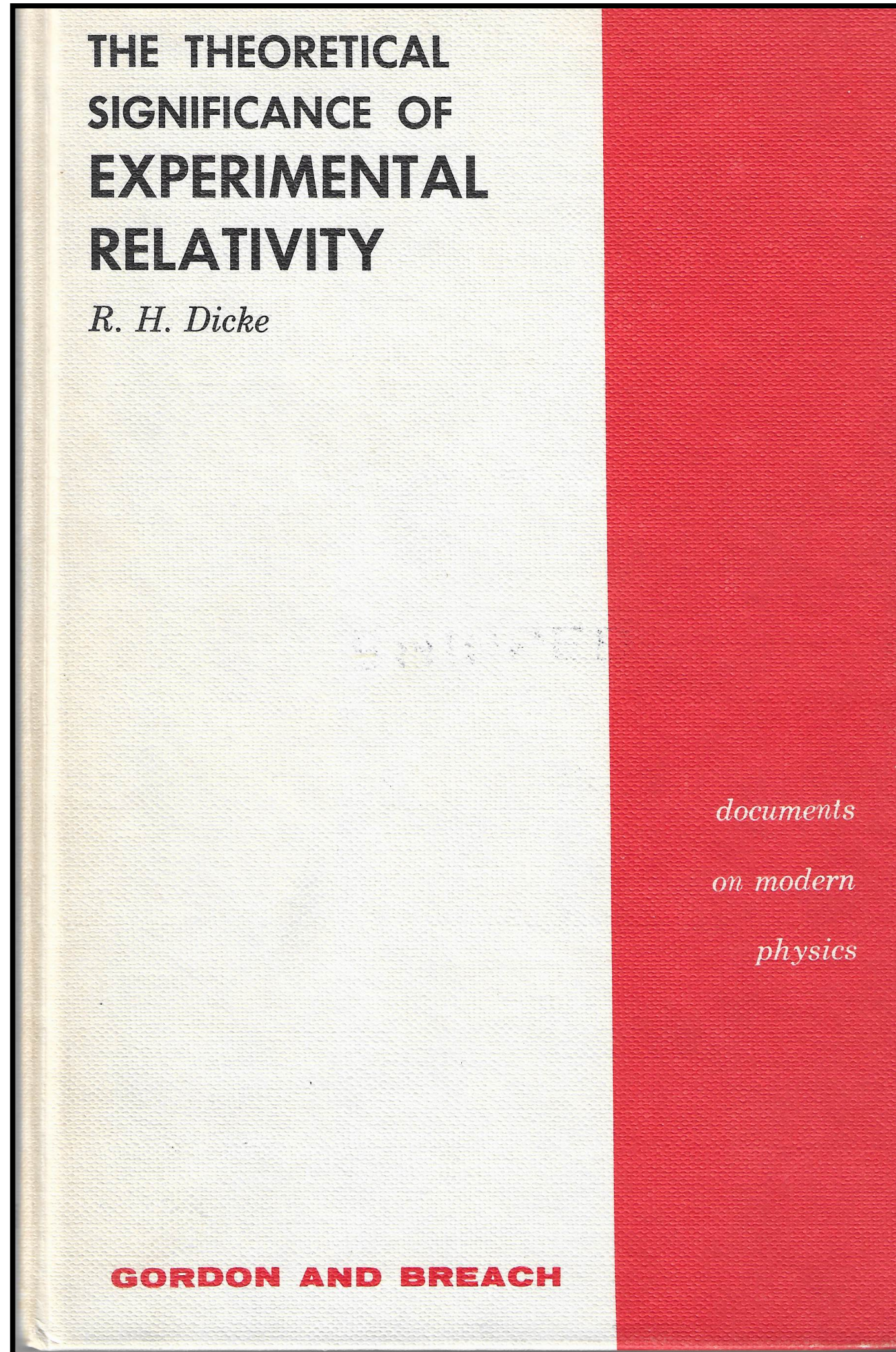
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Dicke (1964); reprinted as a GRG, "Golden Oldie" (2019)

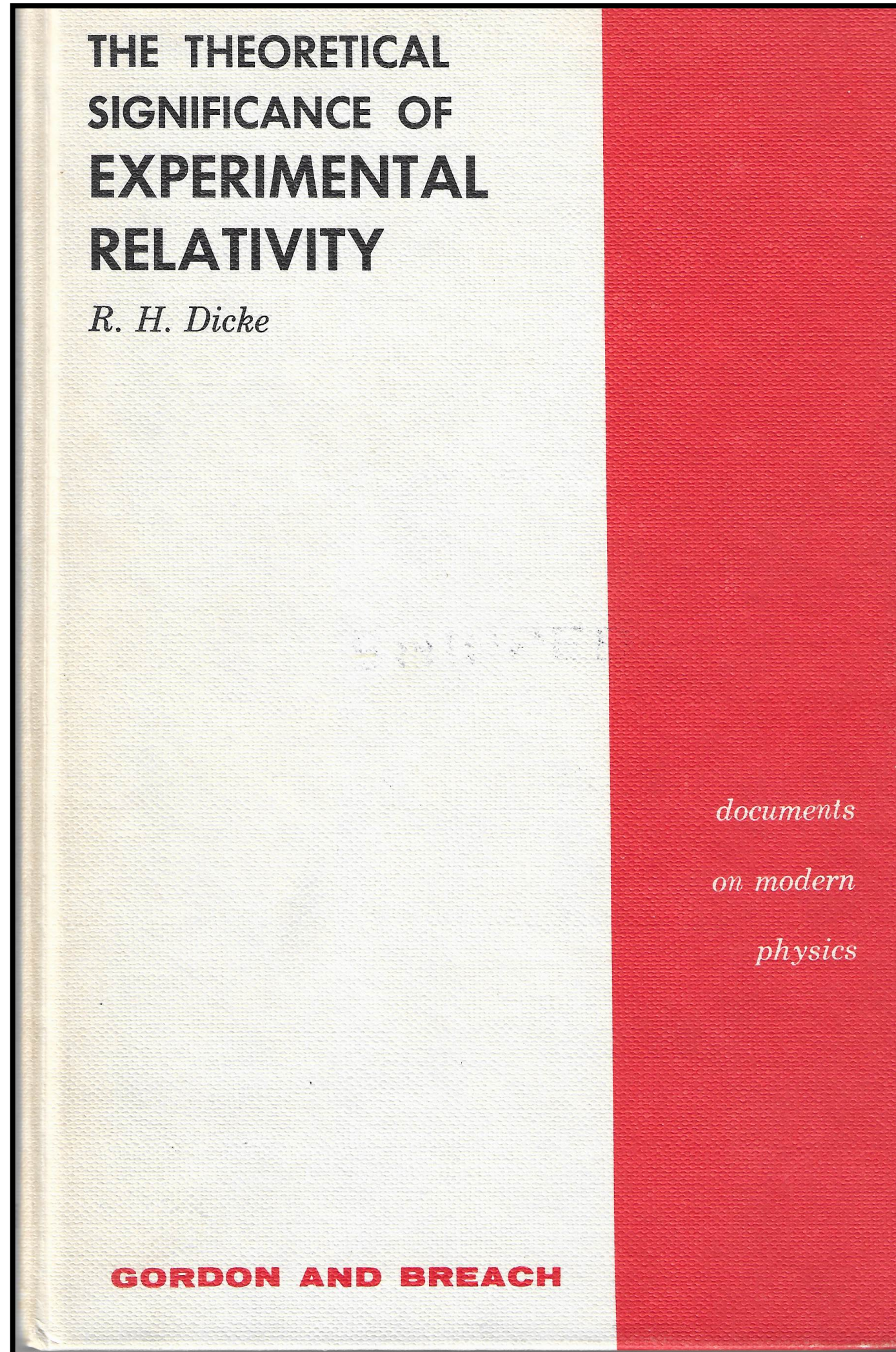
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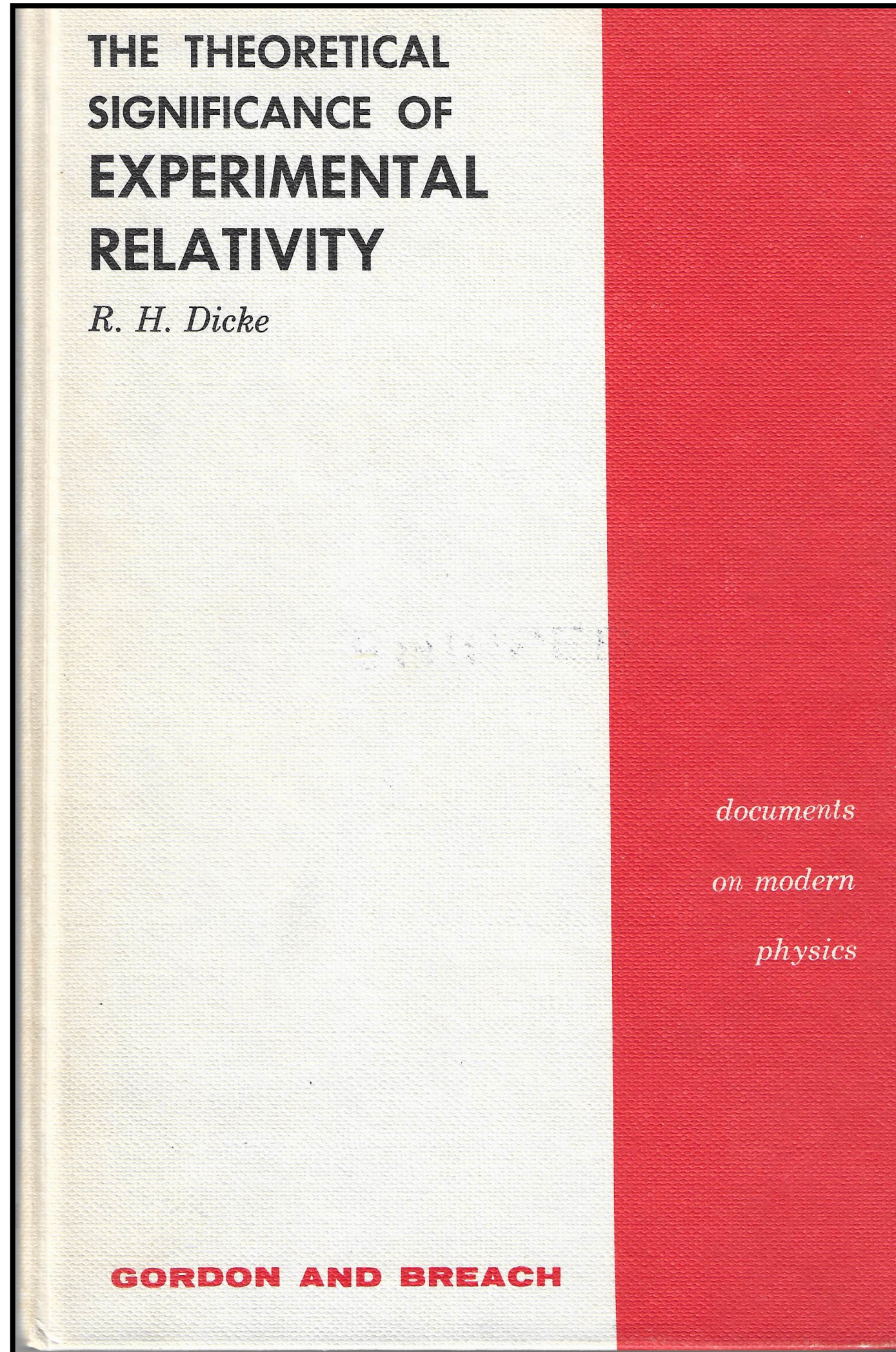
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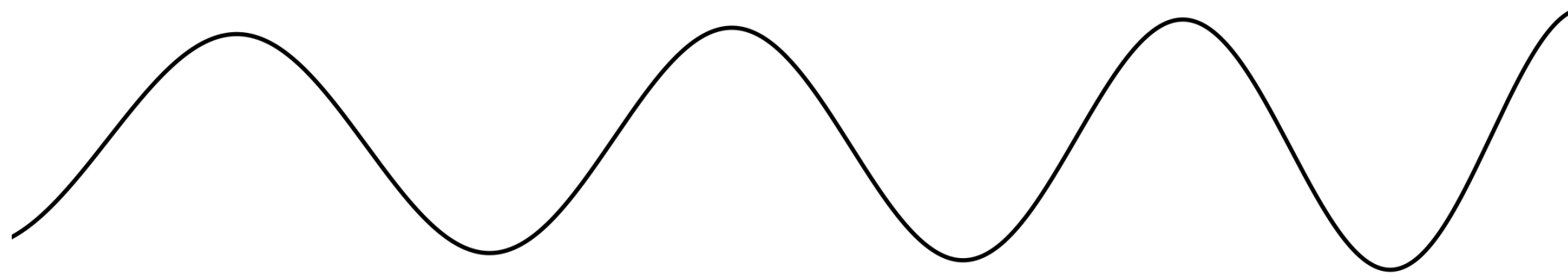
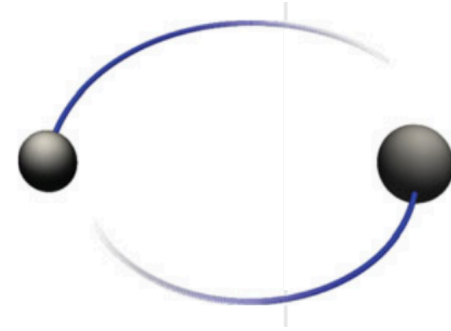
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- “**Experiment and observation can become intertwined in this field of physics** [gravitation], **and will probably become more so in the future** [...]. While his chances of finding exactly the experimental setup that he would have liked is rather small, there may be **many** of these ‘experiments’ that have some bearing on his **particular** problem.”

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# Coalescing binary black holes

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Inspiral

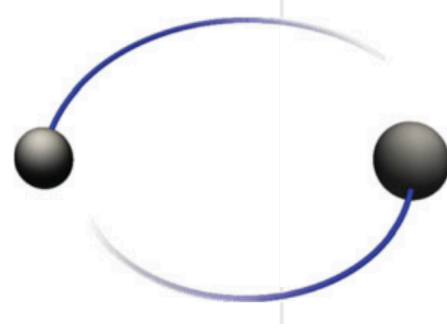


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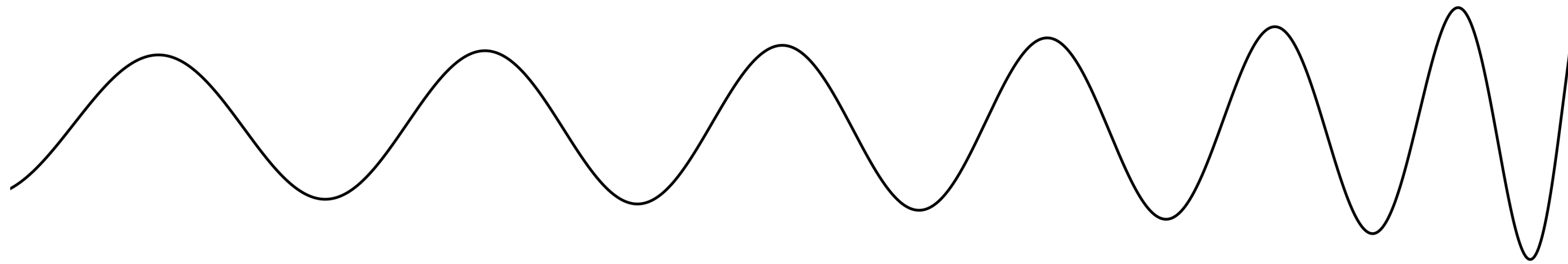
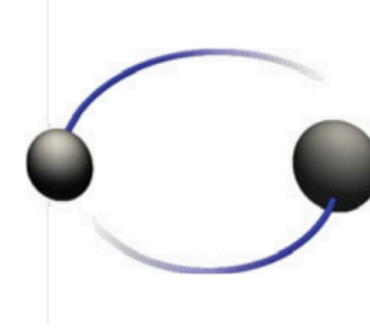


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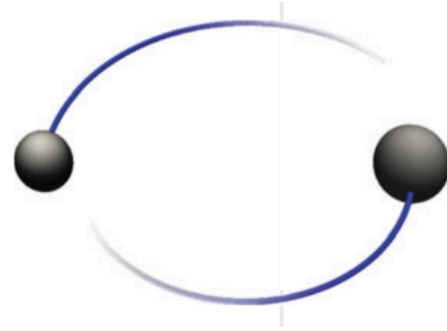
Plunge



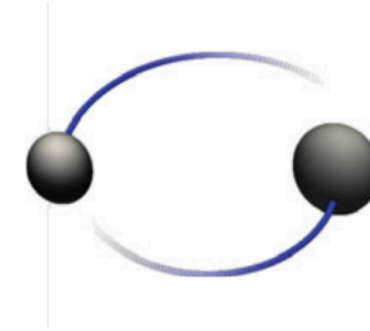
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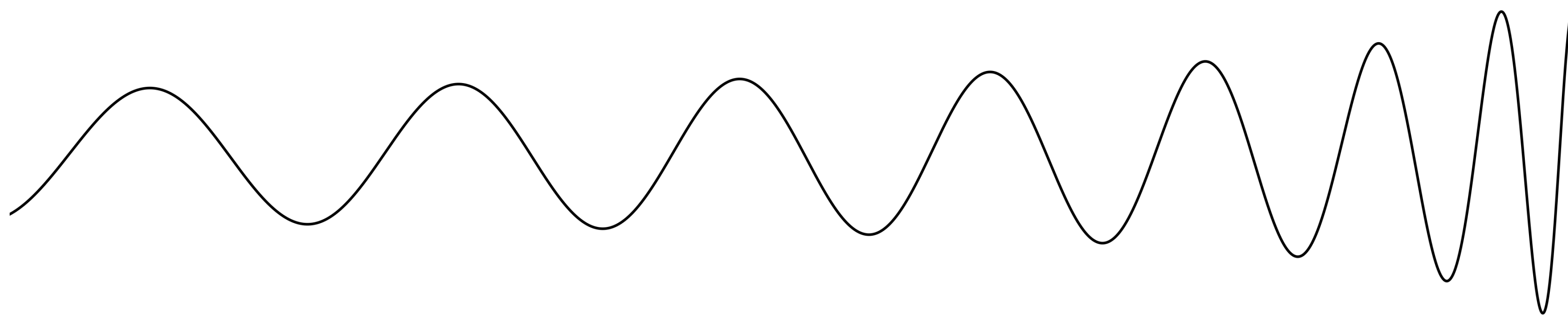
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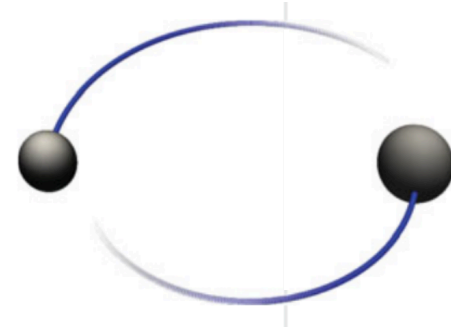
Merger



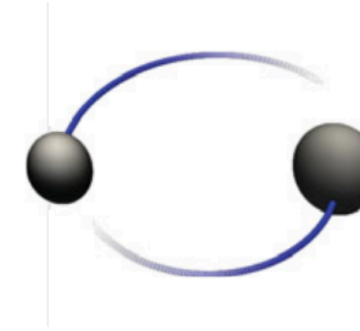
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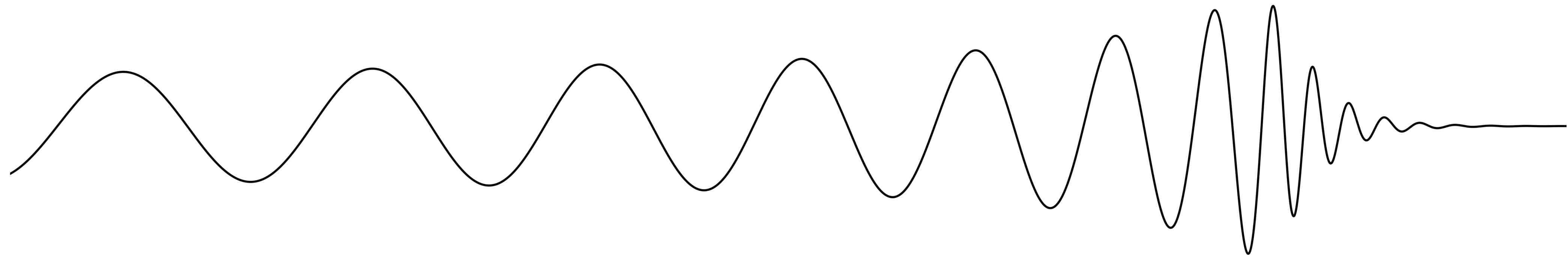
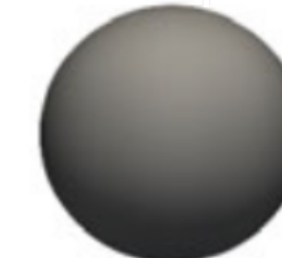
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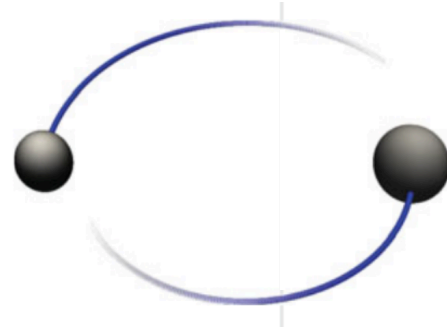
Ringdown



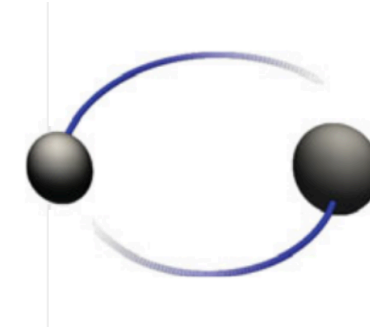
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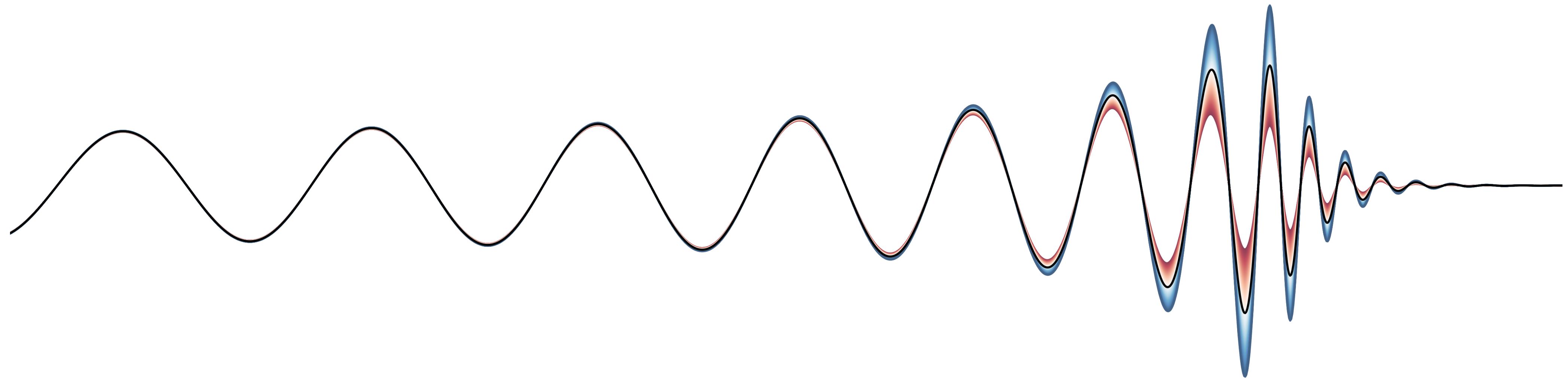
Plunge



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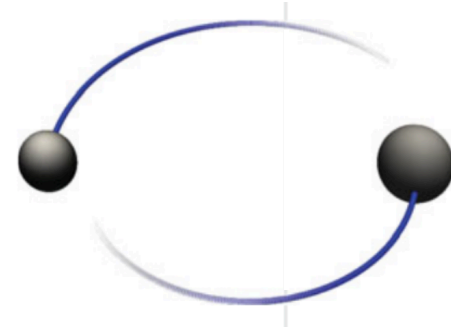
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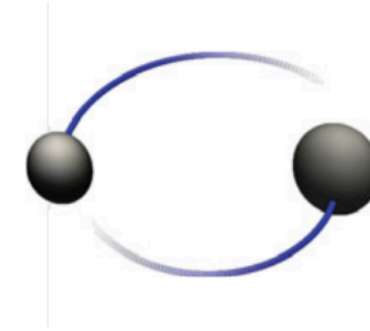
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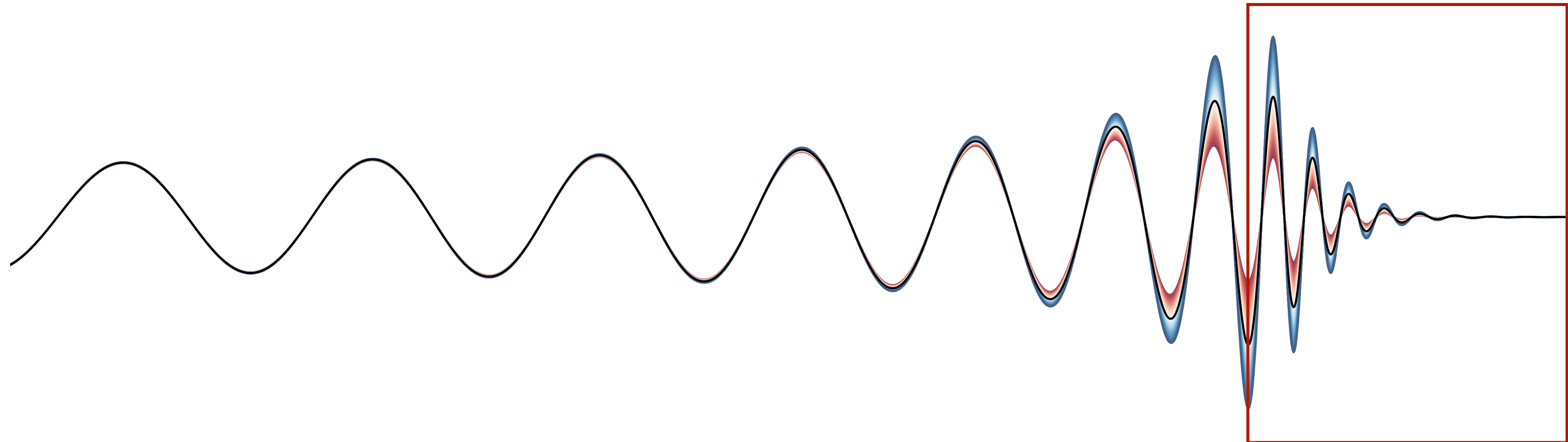
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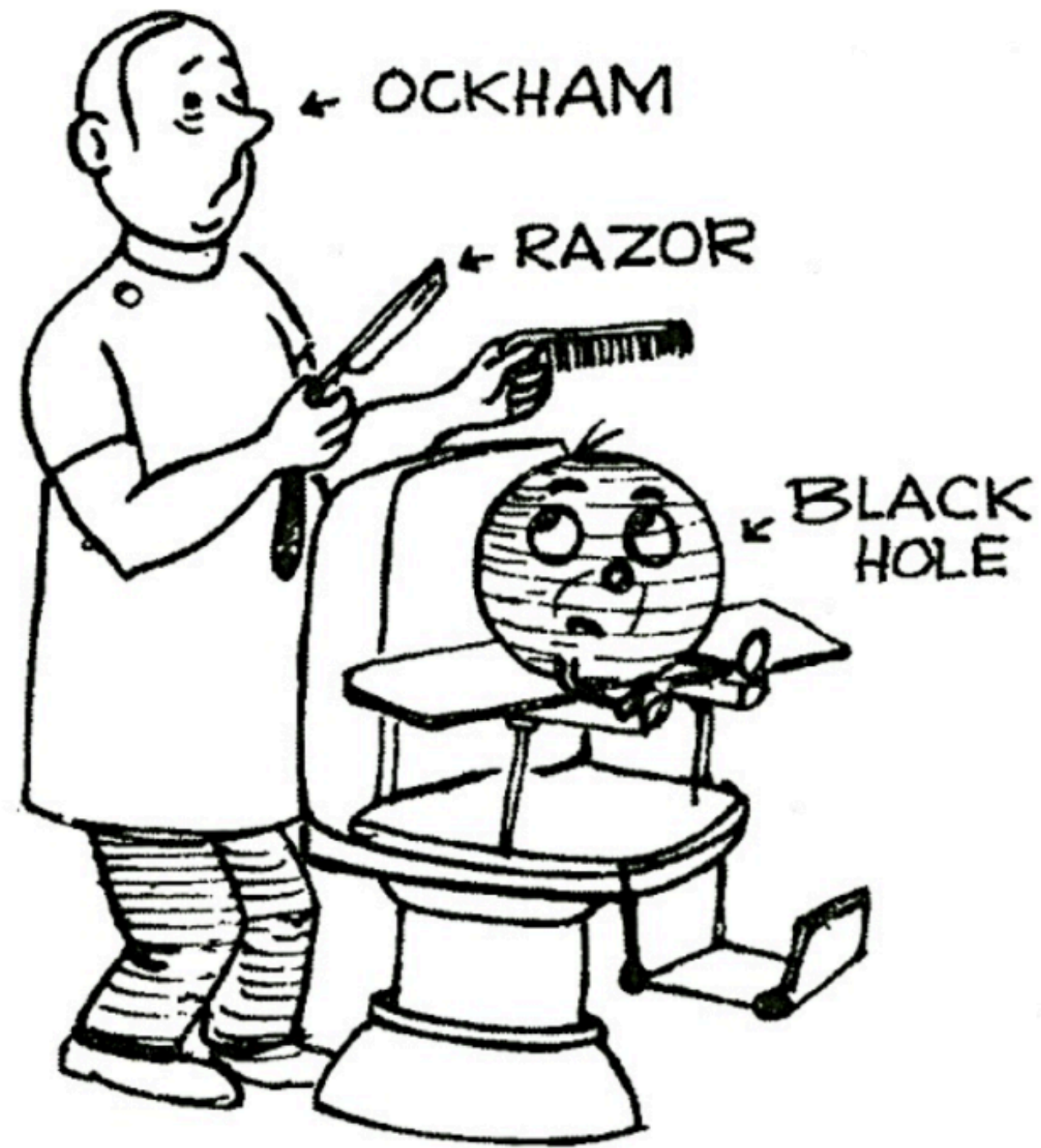
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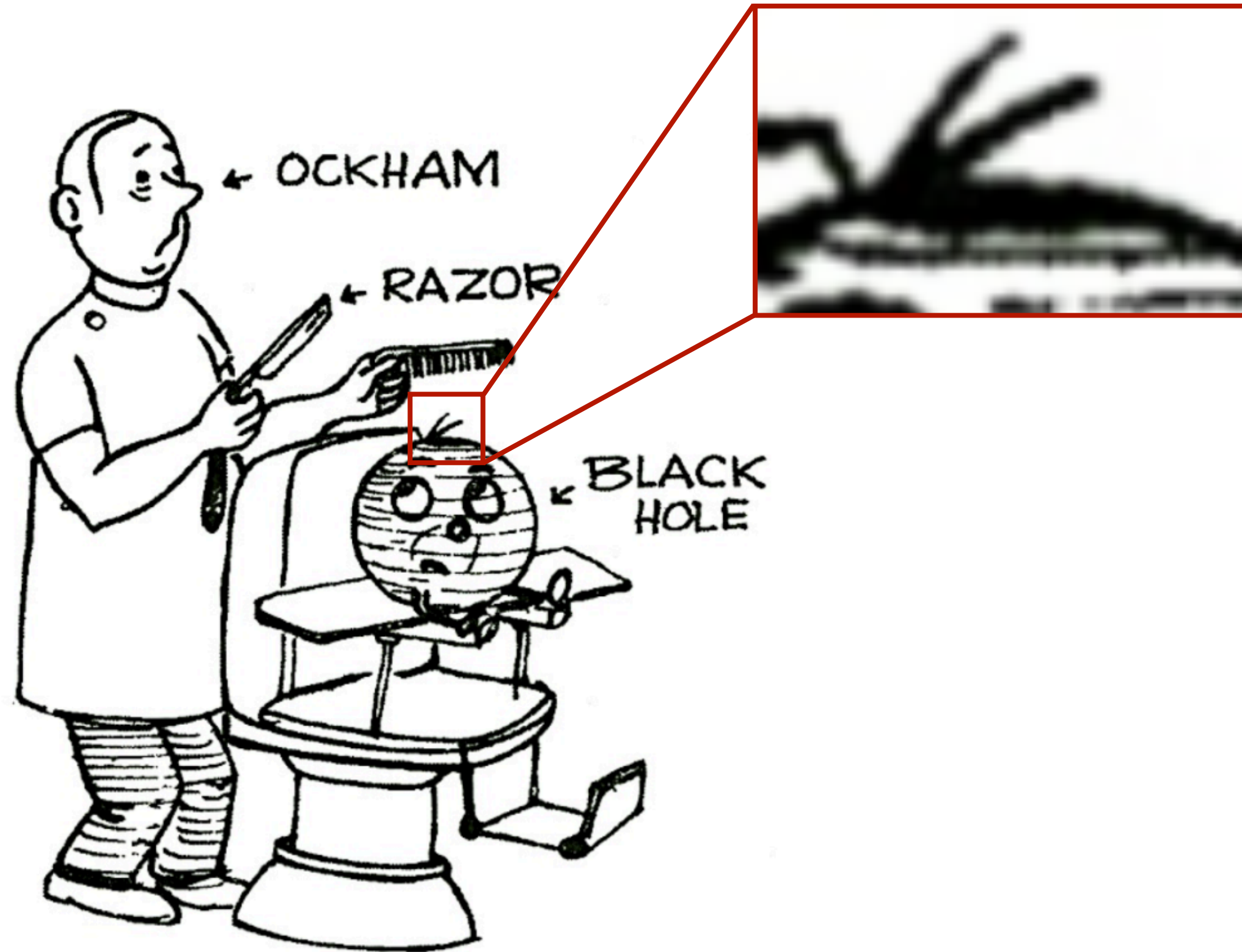
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Cartoon by C. V. Vishveshwara

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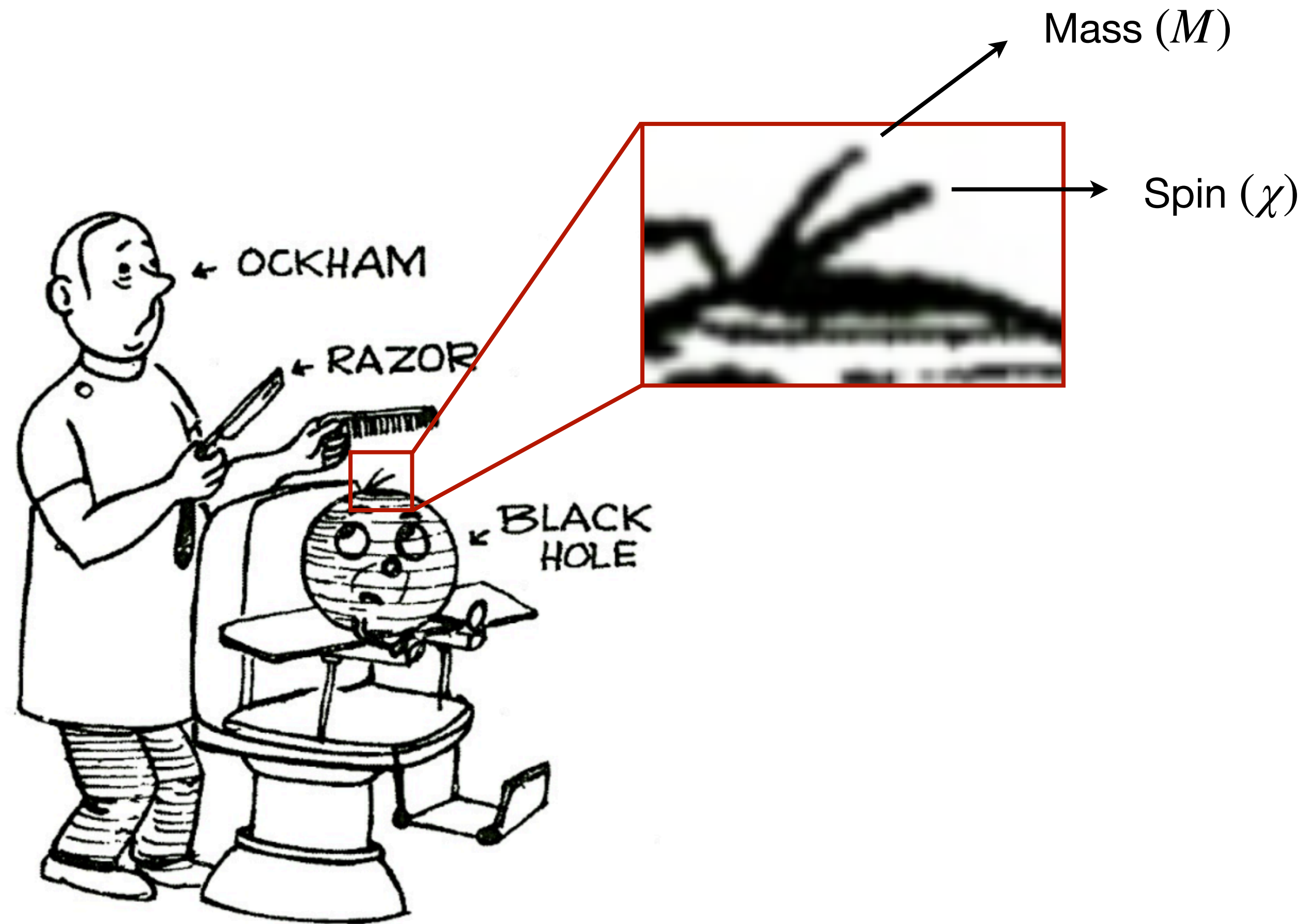


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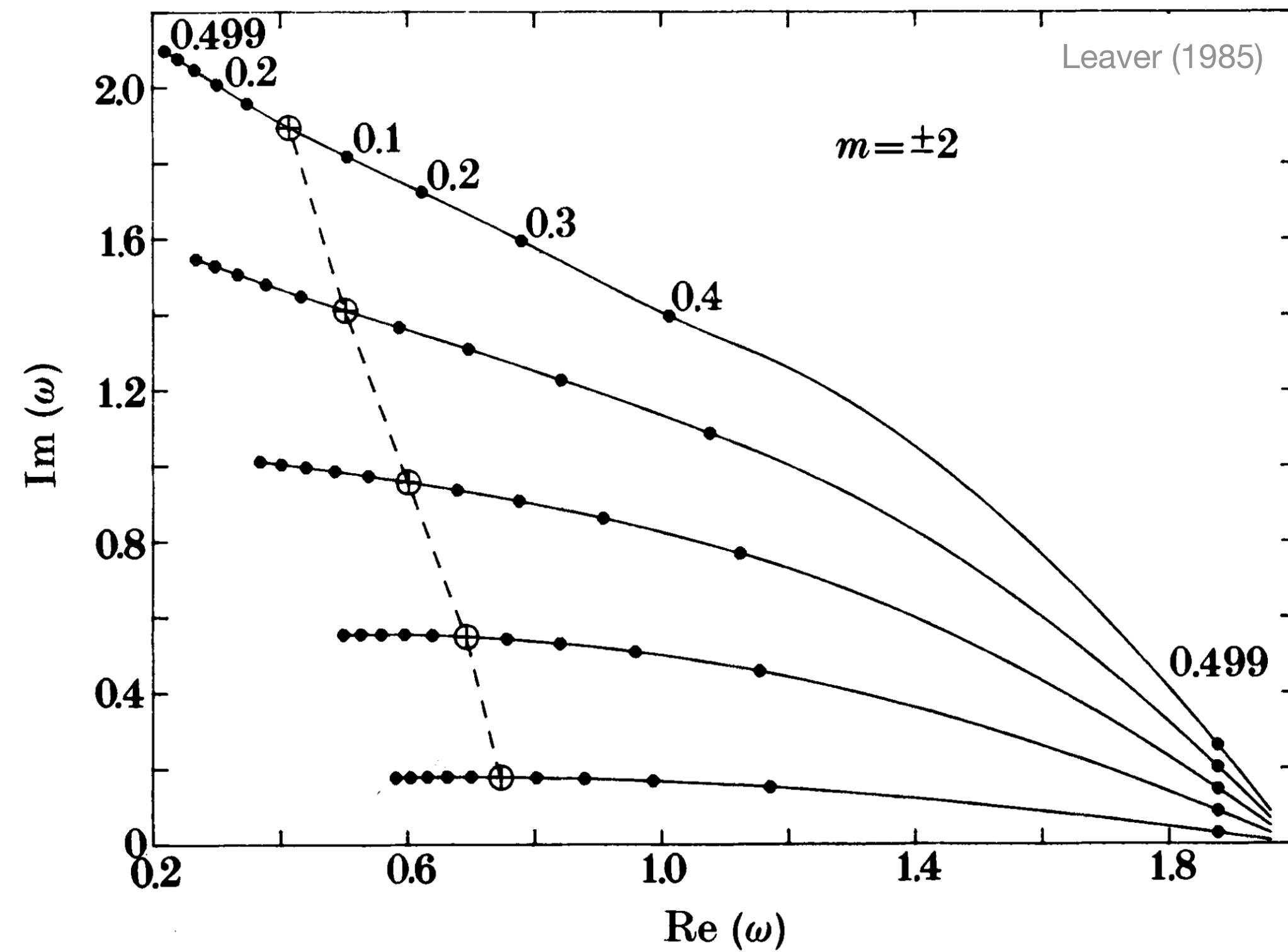
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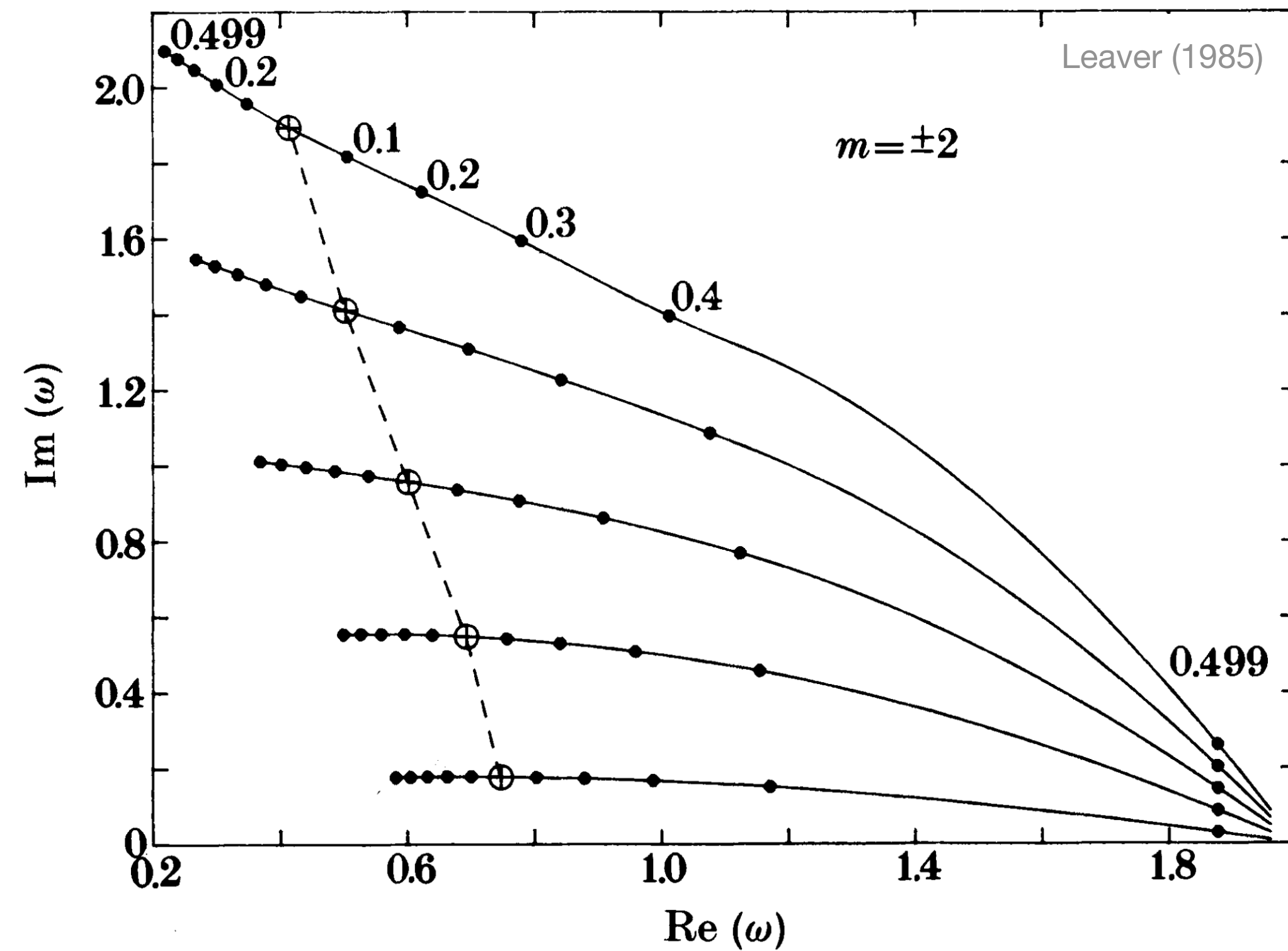
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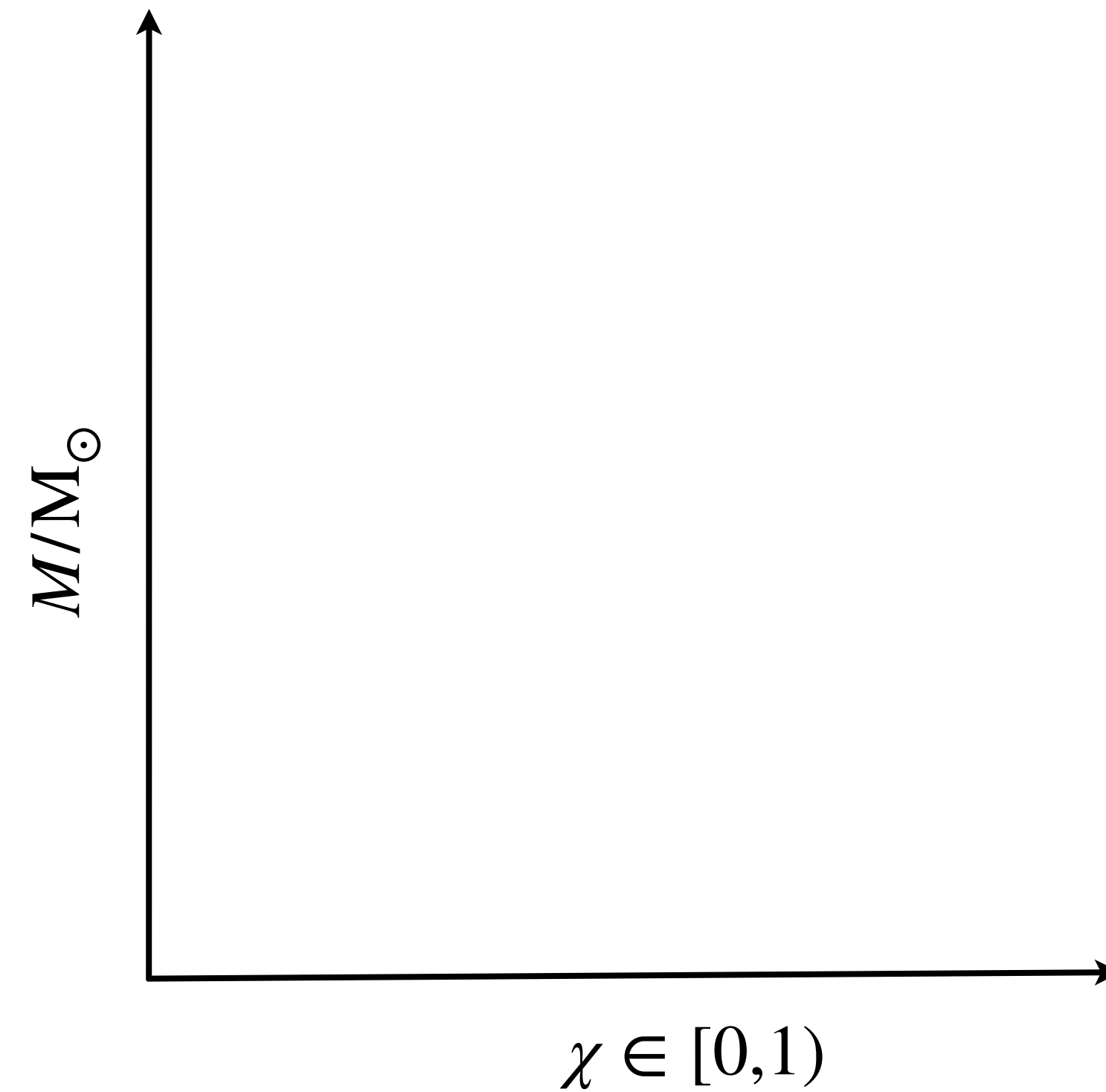
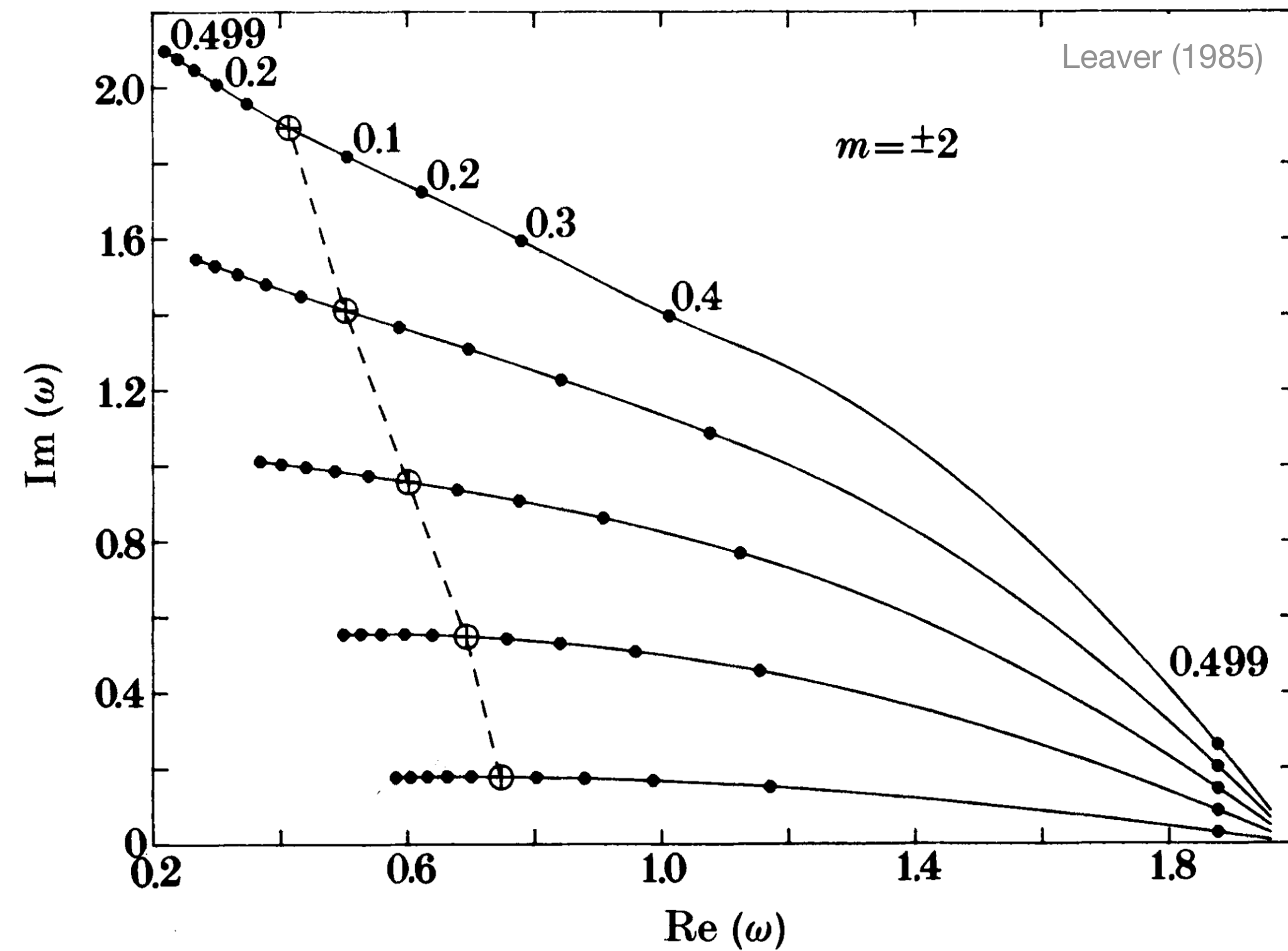
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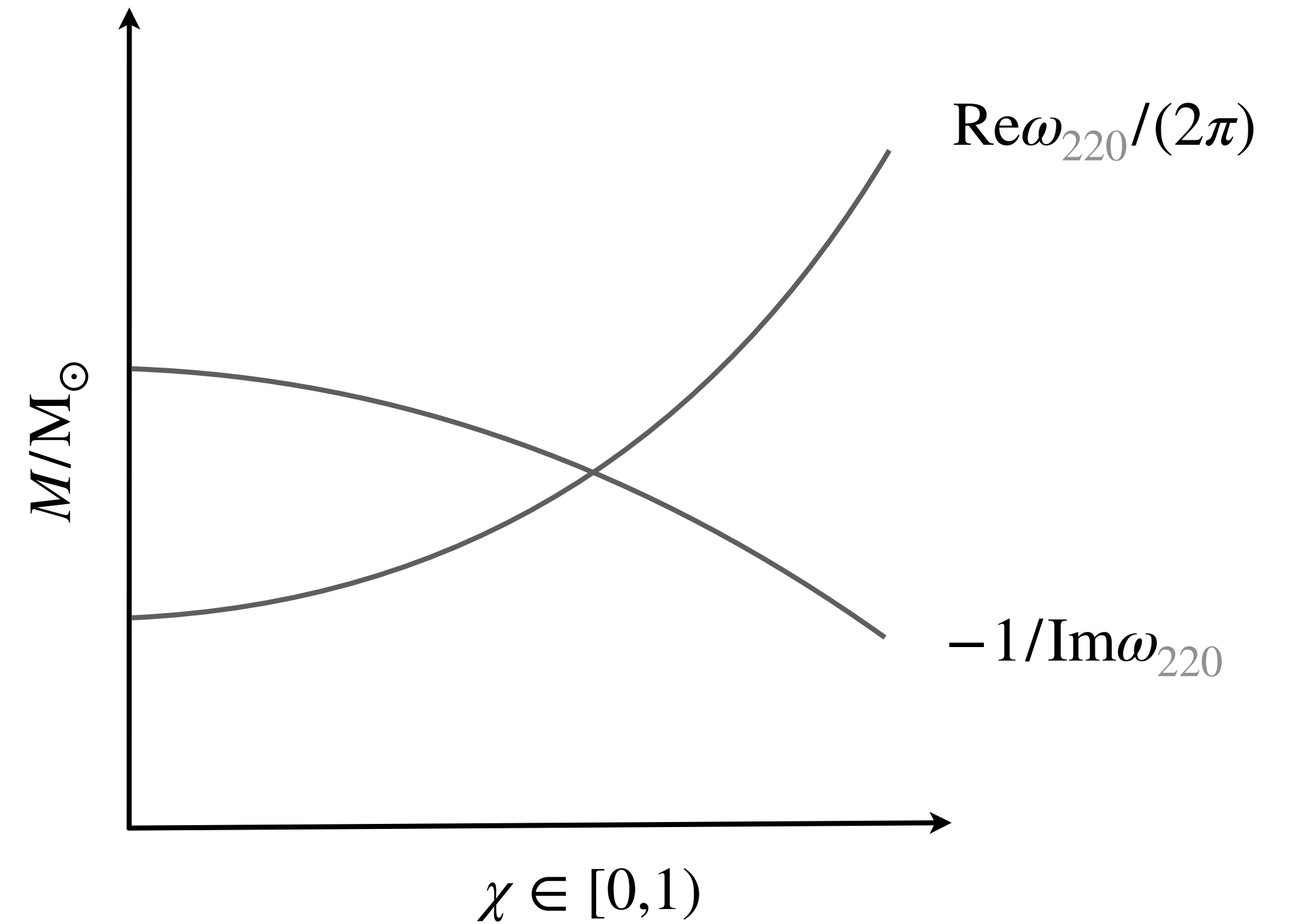
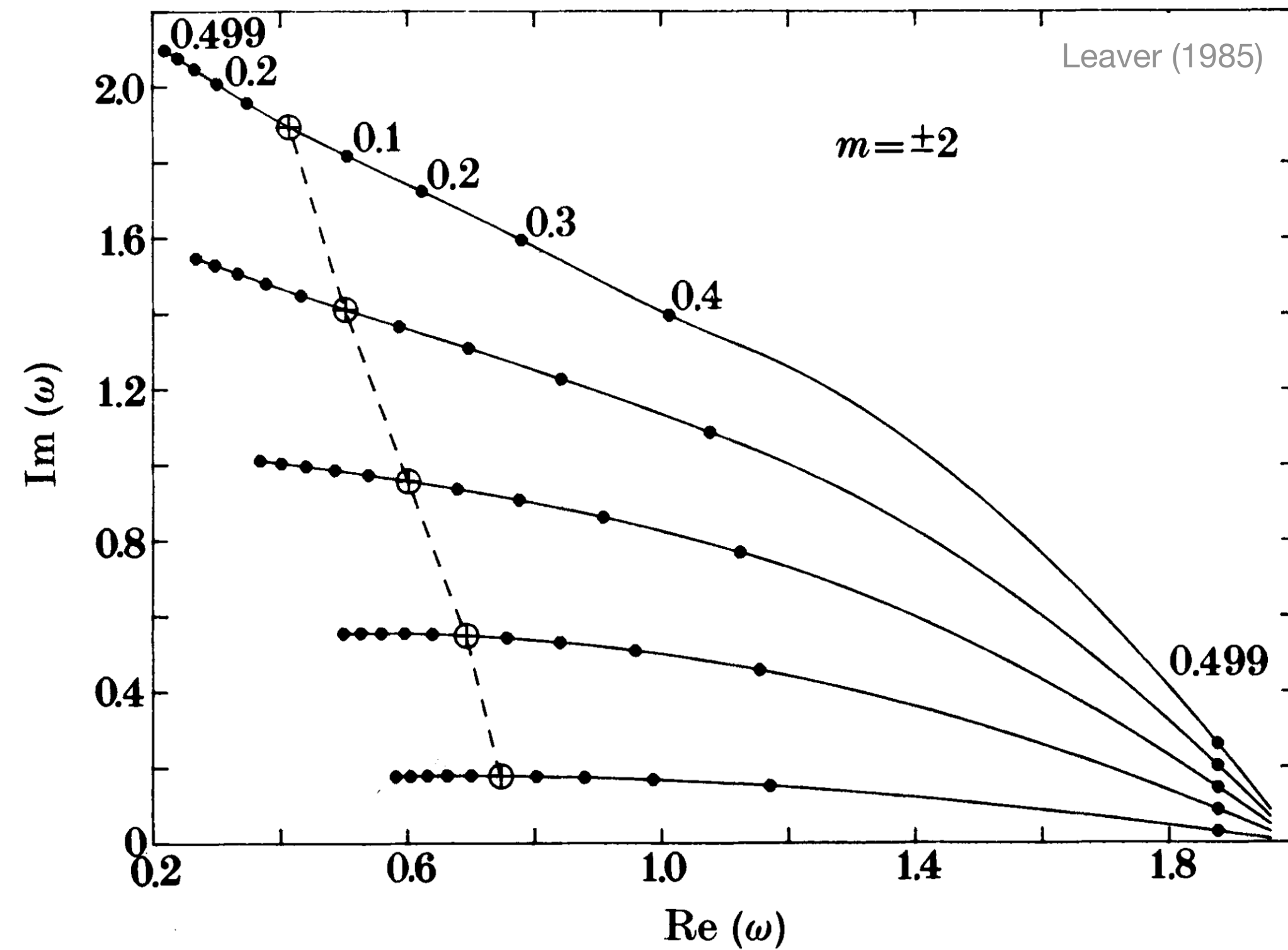
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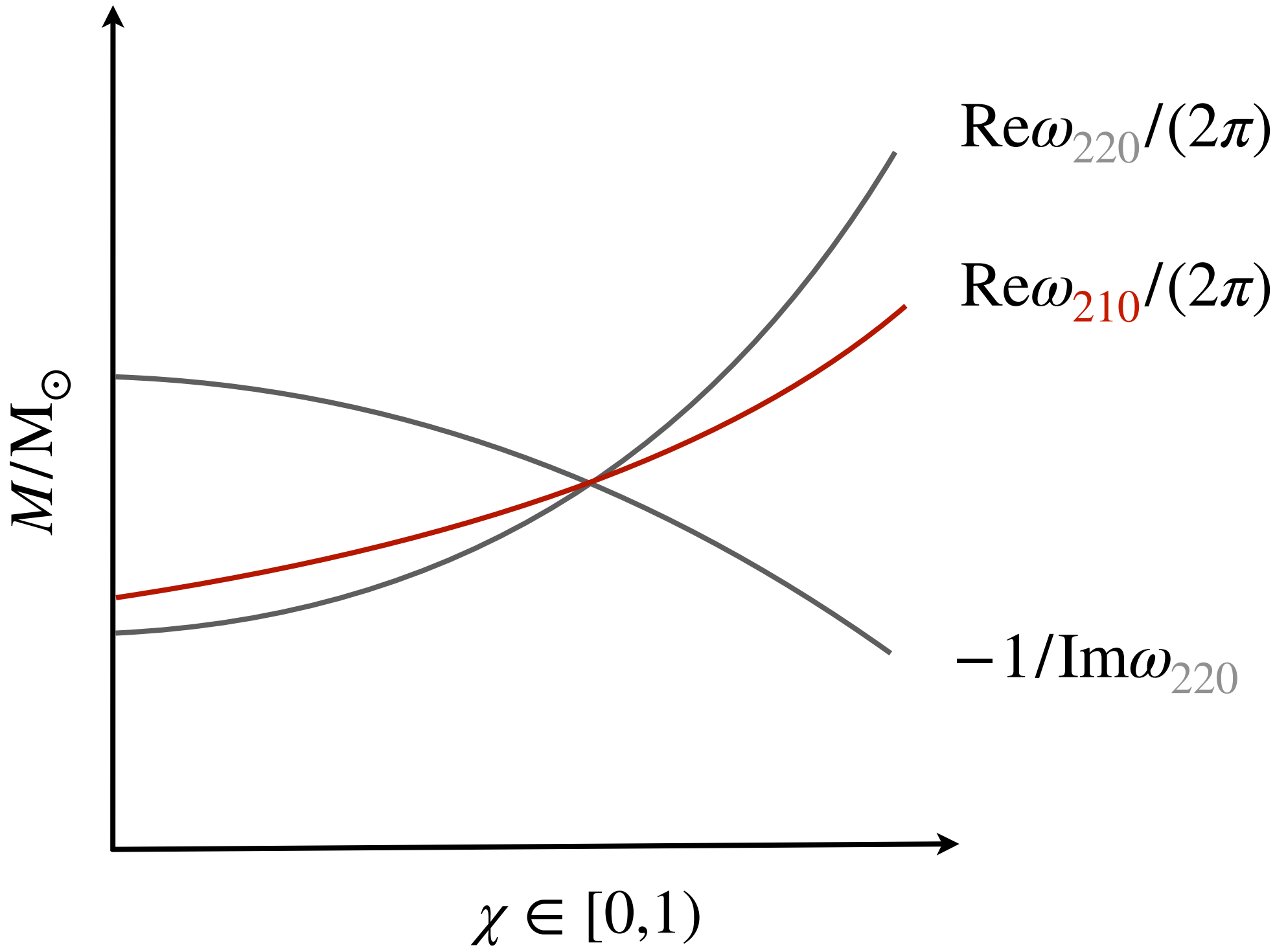
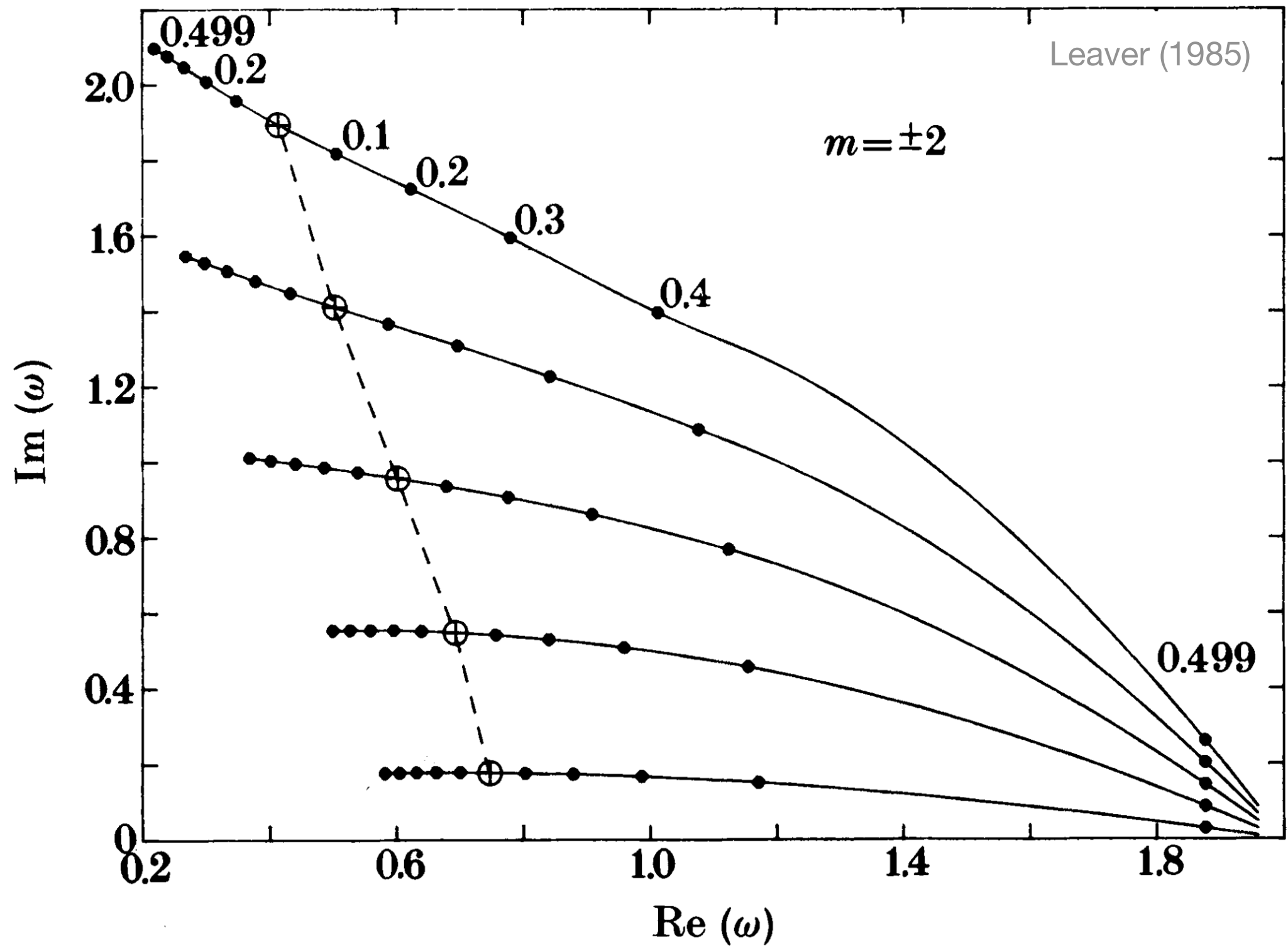
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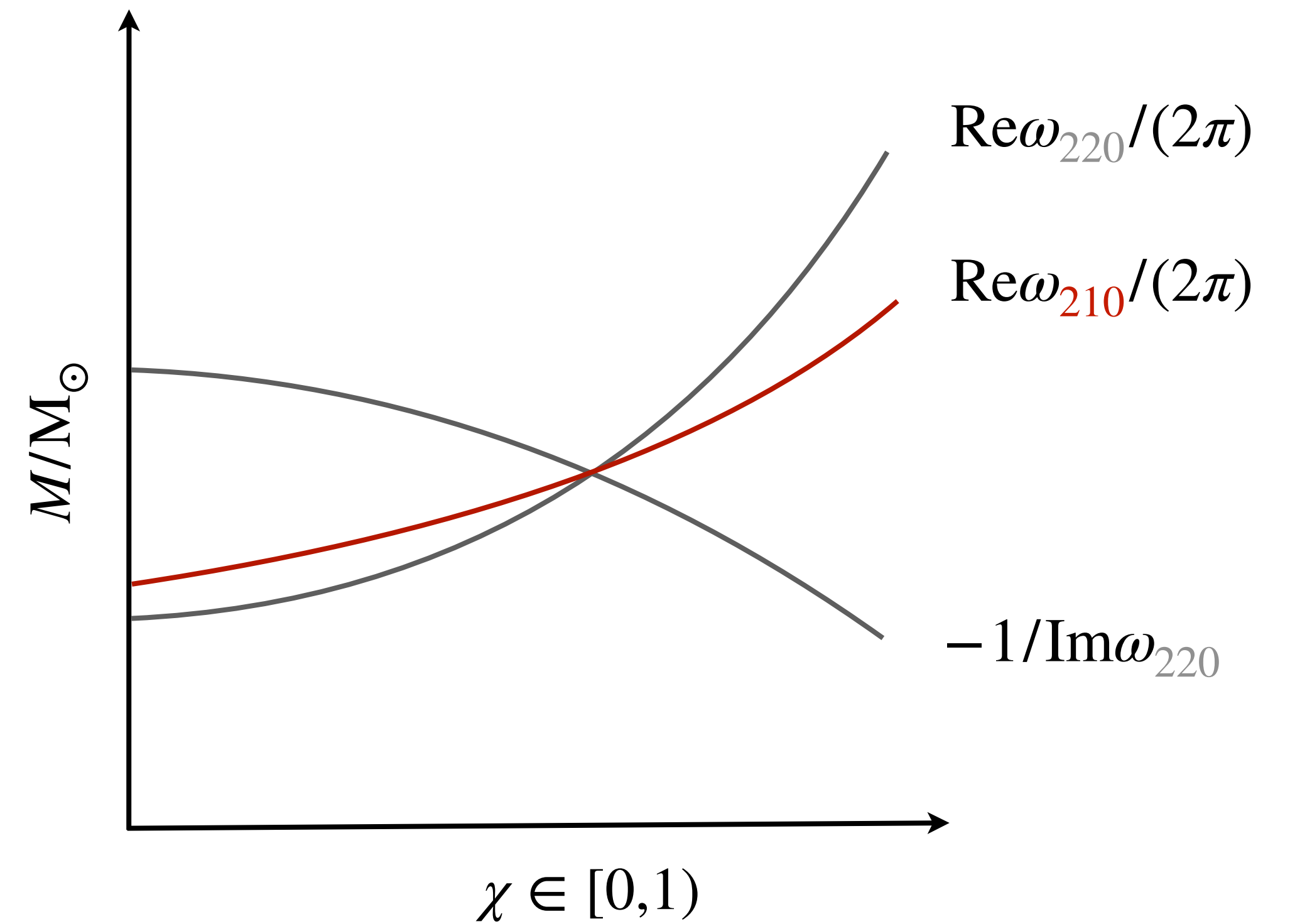
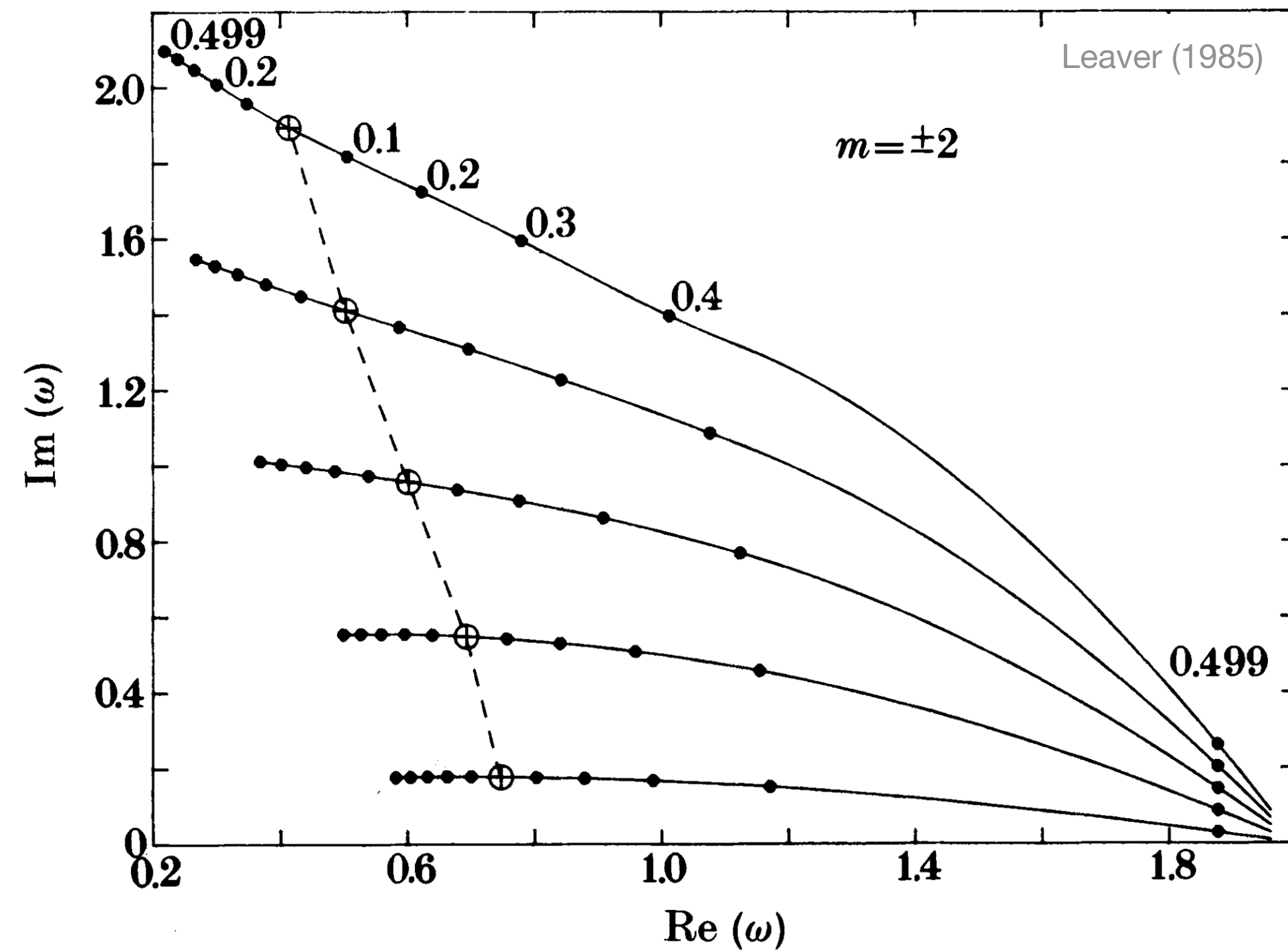




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*"[...] direct evidence of black holes with the same certainty as, say, the 21 cm line identifies interstellar hydrogen."*  
 Detweiler (1980)

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Today: partial answer to some of these questions based on [2404.11110](#) (to appear in PRD), with [G. Tambalo](#), [K. Glampedakis](#), [K. Yagi](#), and [J. Steinhoff](#).



# Outline

1. Effective-field-theory of general relativity
2. Quasinormal modes: their frequencies
3. Quasinormal modes: their amplitudes
4. Conclusions and some conjectures

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- Analytical (rotating) black hole solutions.

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Mathematically, these assumptions translate into studying differential equations of the form:

$$\left[ \frac{d^2}{dx^2} + Q_{\ell m}(x, \omega, \vartheta) \right] \psi_{\ell m}(x) = s_{\ell m}(\omega, x)$$

# Unpacking

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- $\psi_{\ell m}$ : “master function”: a particular combination of metric perturbations  $h_{\mu\nu}$ , decomposed in tensor harmonics and in Fourier modes  $\psi_{\ell m}(t, x) \sim \exp(i\omega t) \psi_{\ell m}(x)$ .

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Quasinormal modes are solutions for which  $A_{\text{in}} = 0$ ; they are purely **ingoing** into the event horizon and **purely** outgoing at spatial infinity.

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Complex-valued **poles** of the scattering matrix: a **boundary-value problem**.

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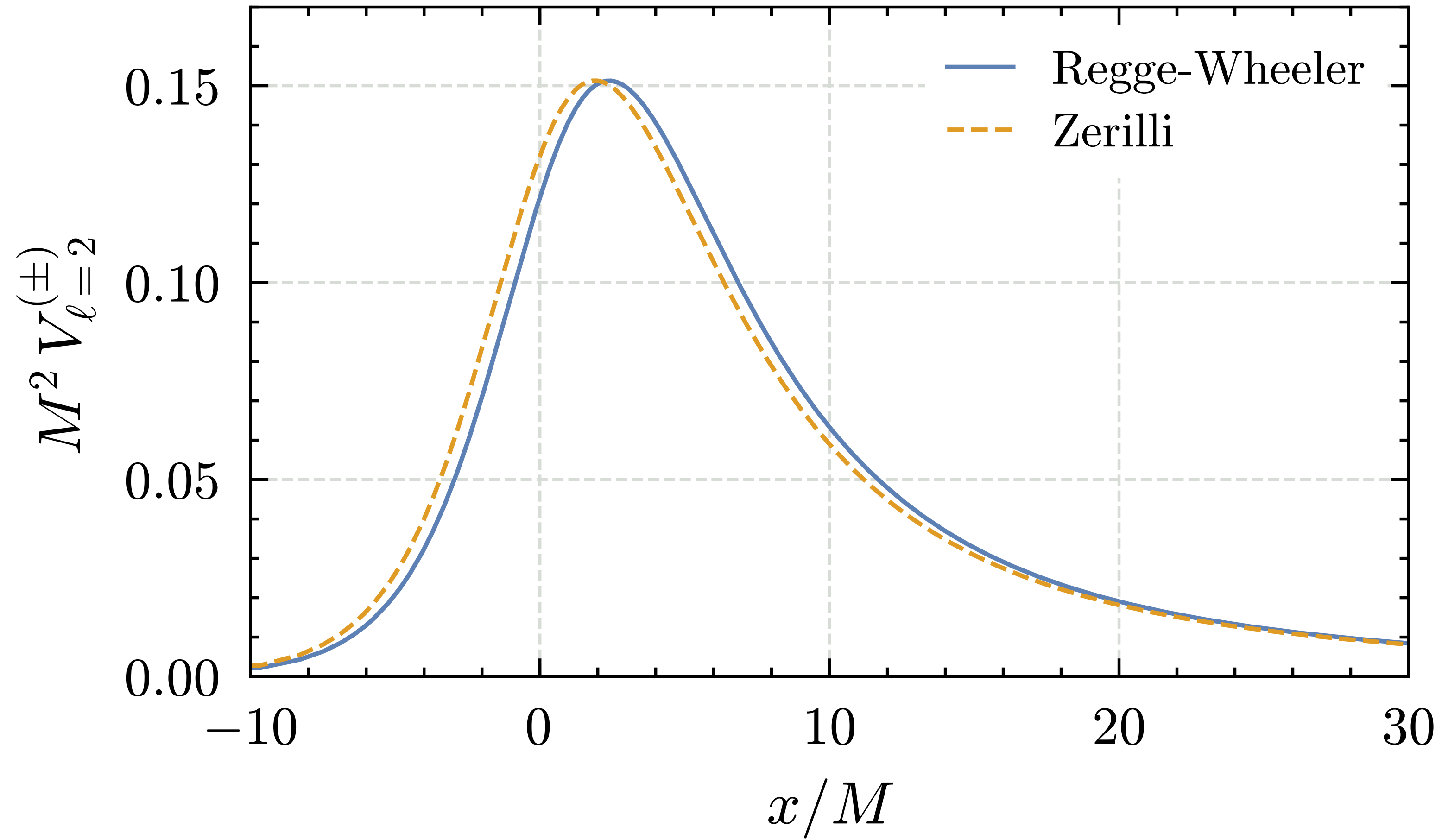
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where  $\lambda_{\ell} = (\ell + 2)(\ell - 1)/2$  and  $\Lambda_{\ell} = \lambda_{\ell} + 3M/r$ . Hereafter,  $V_{\ell m}^{(\pm)} = \omega^2 - Q_{\ell m}^{(\pm)}$  (the “effective potential”).

# Schwarzschild perturbations — II



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Chandrasekhar (1980), Lenzi and Sopuserta (2021).

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If two potentials are to have the same reflection and transmission coefficients they must satisfy an infinite hierarchy of integral equalities, in which the integrands are, formally, conserved quantities of the [Korteweg-de Vries](#) equation.

This implies that the quasinormal-frequency spectrum is the **same**,  $\omega_{\ell mn}^{(+)} = \omega_{\ell mn}^{(-)}$ .

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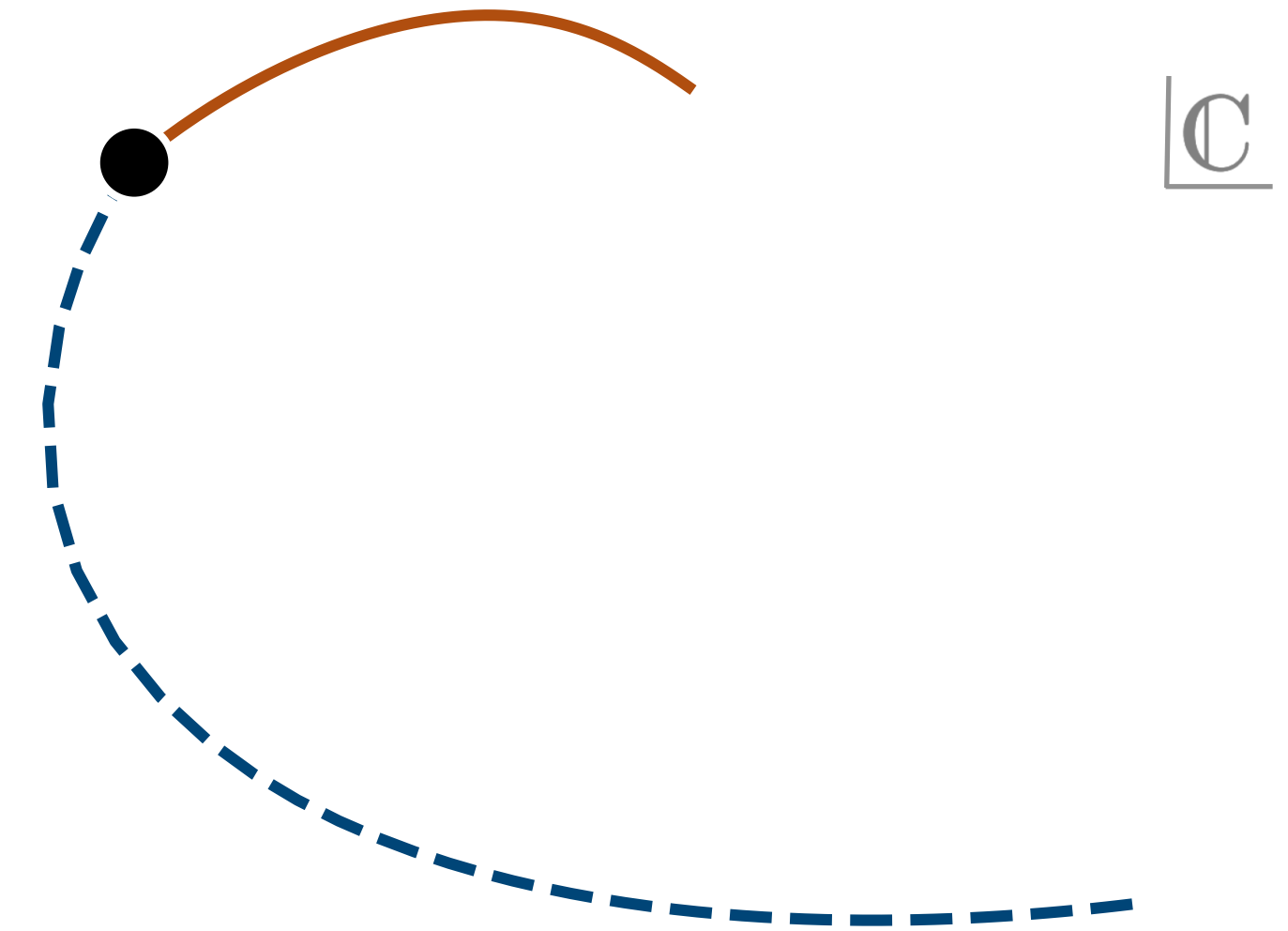
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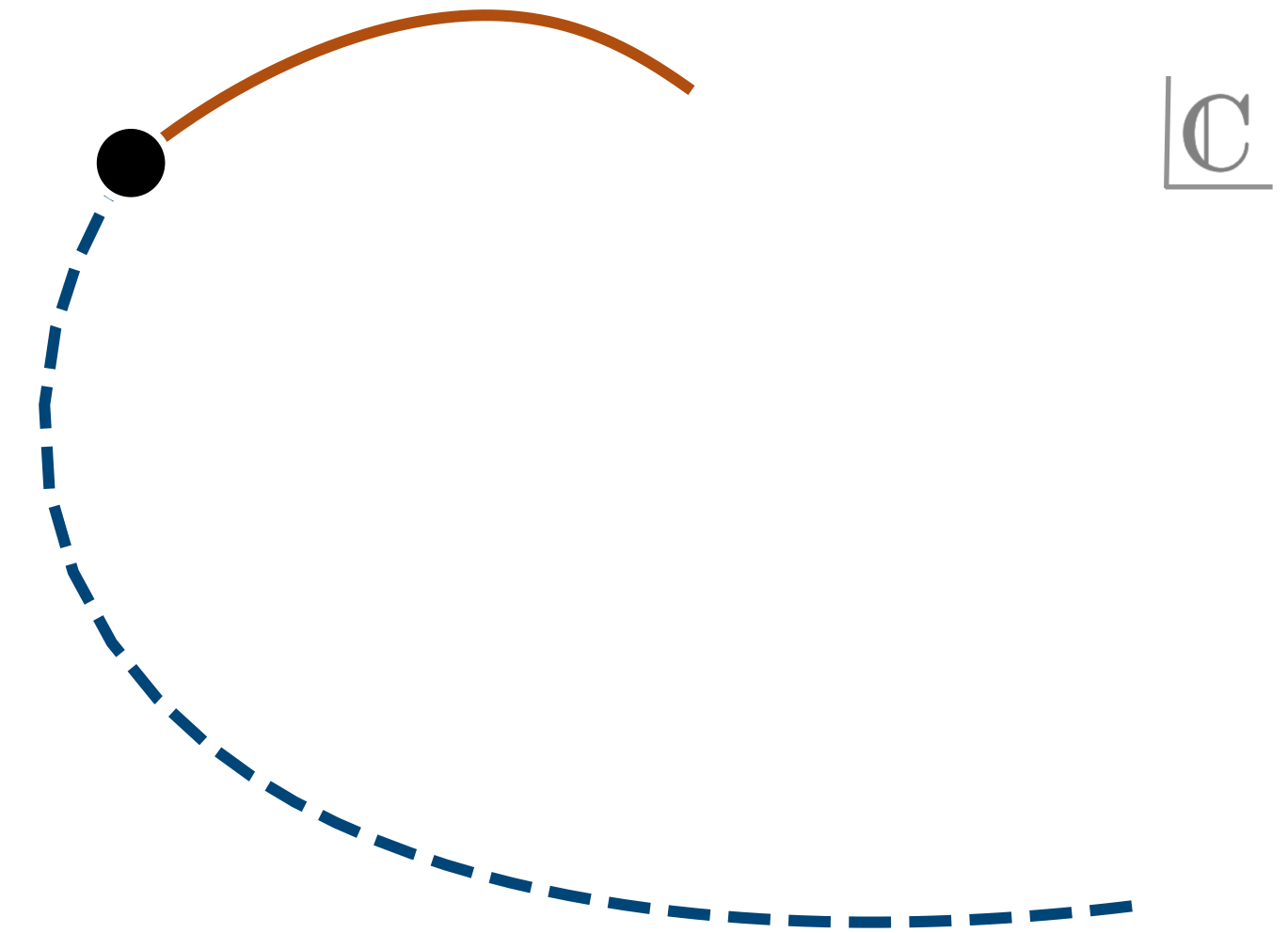


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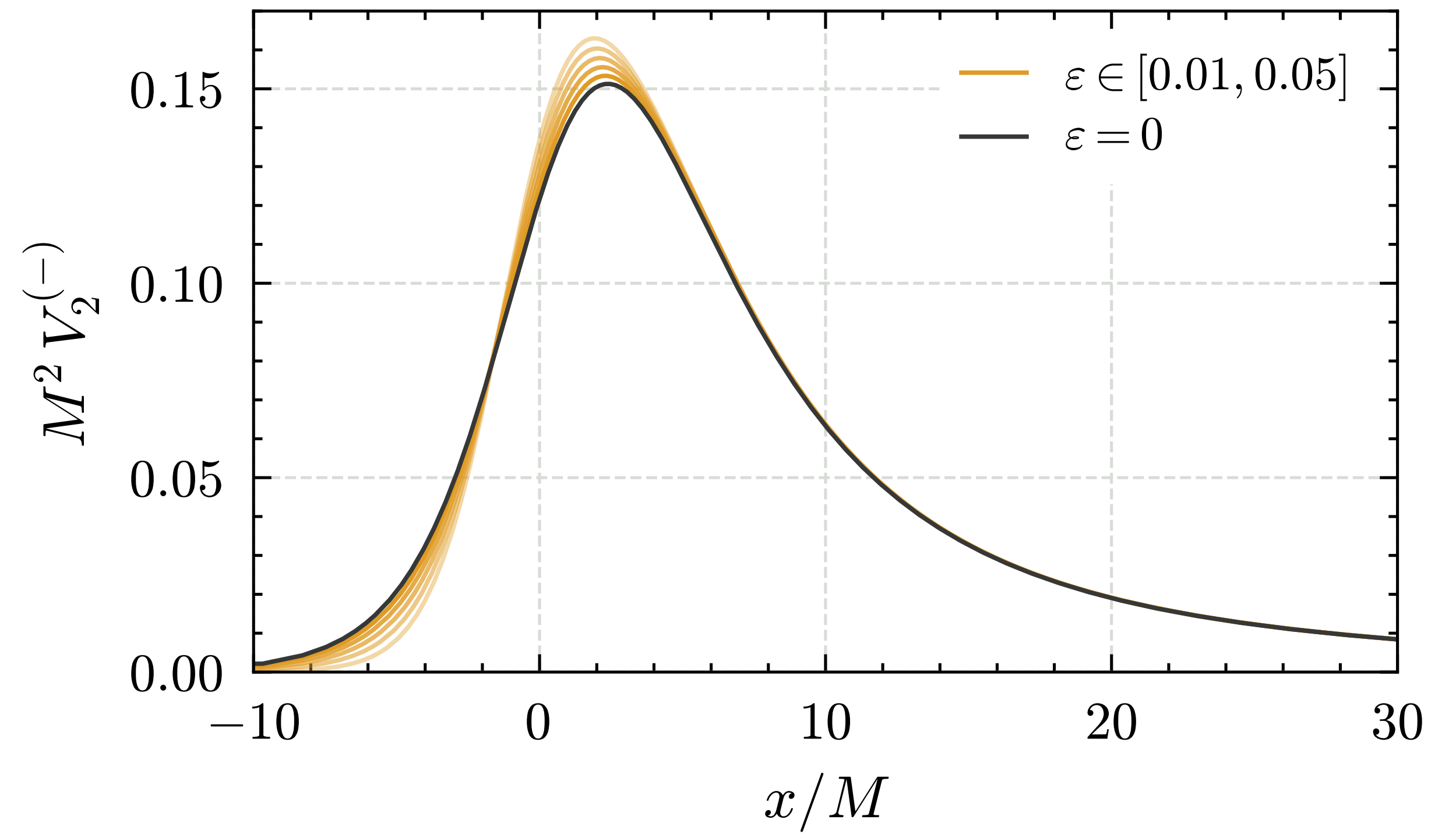
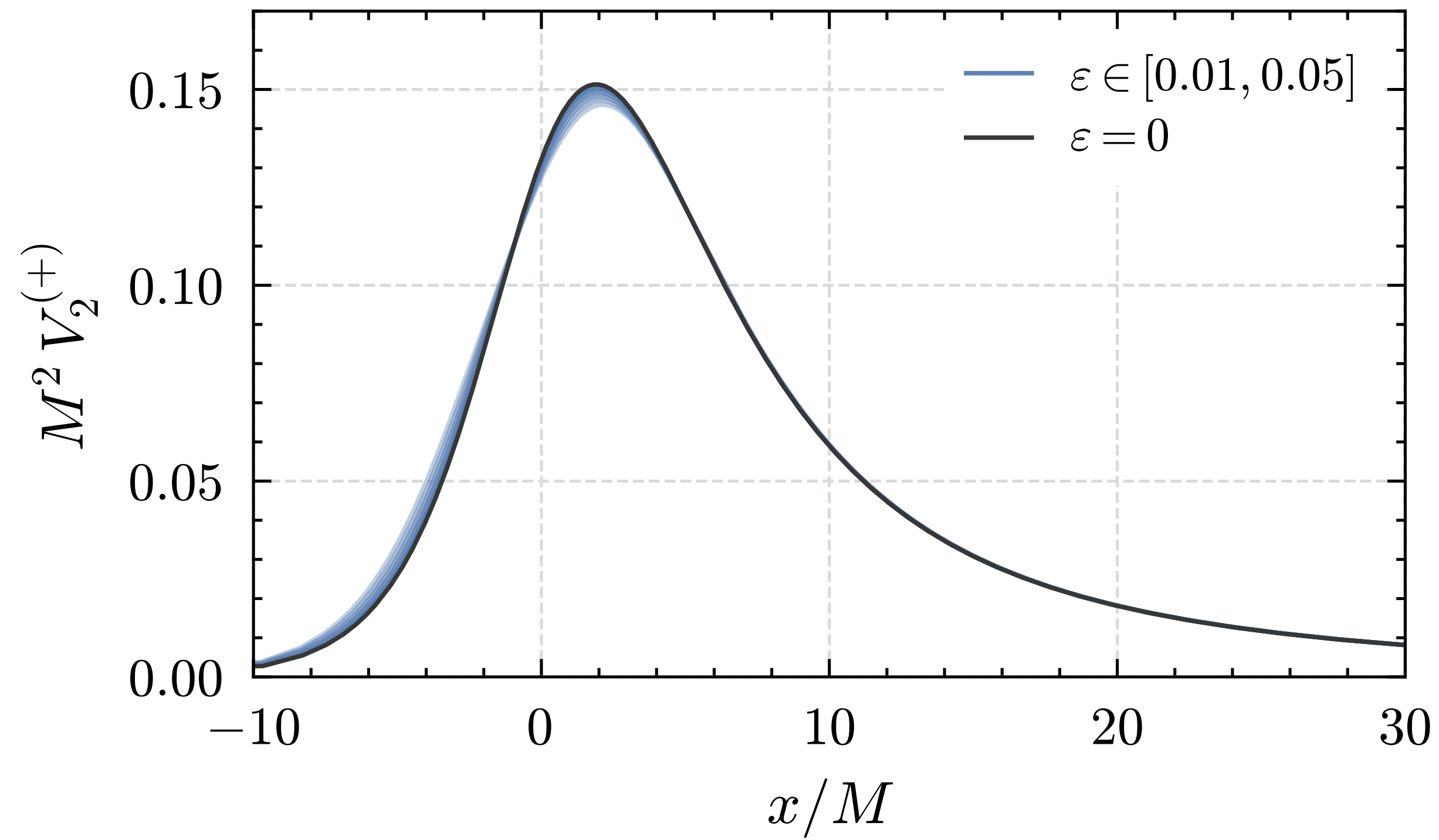
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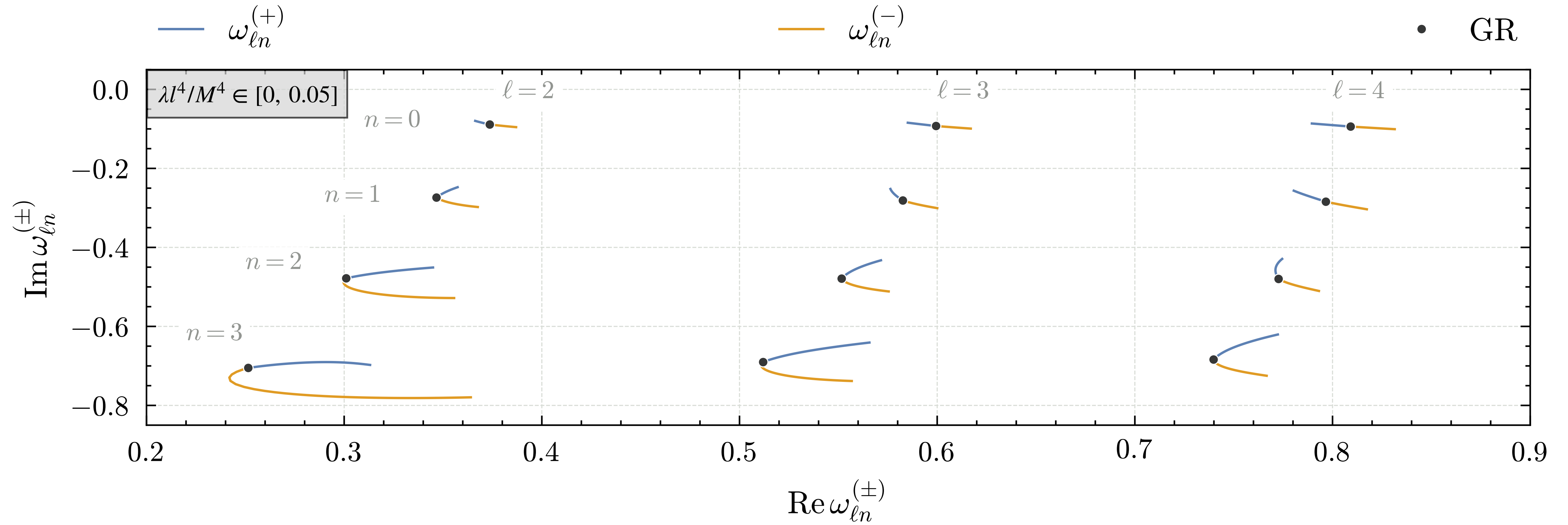


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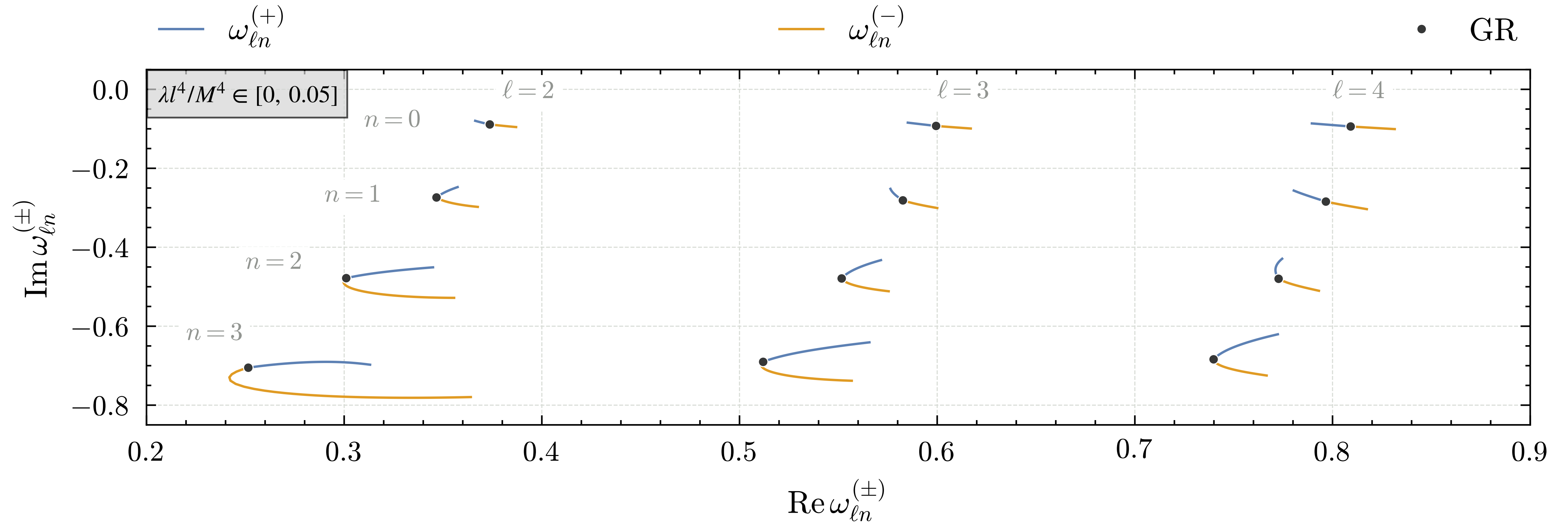
# EFT perturbations — II



# Quasinormal frequency spectra

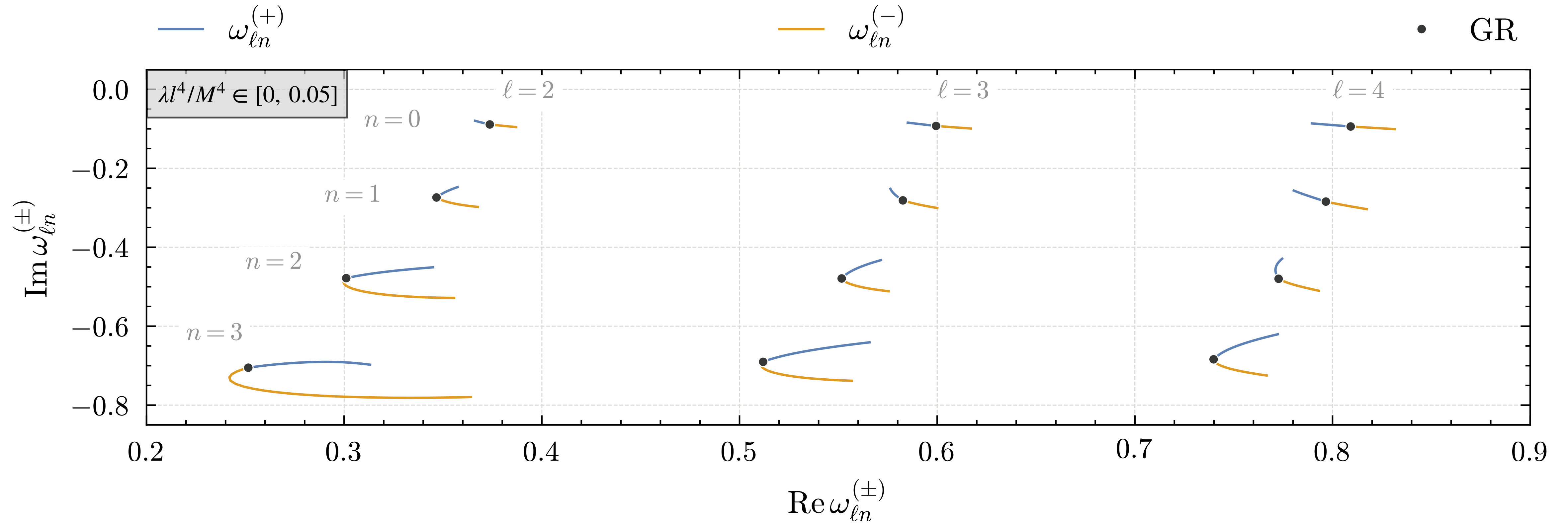


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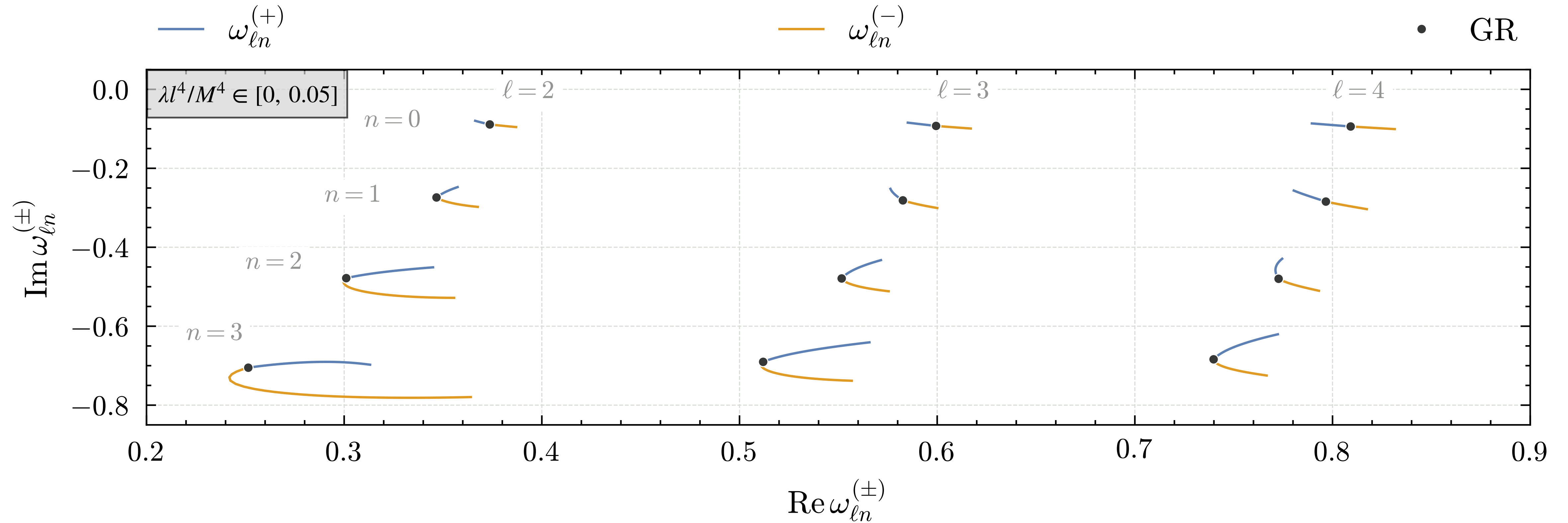
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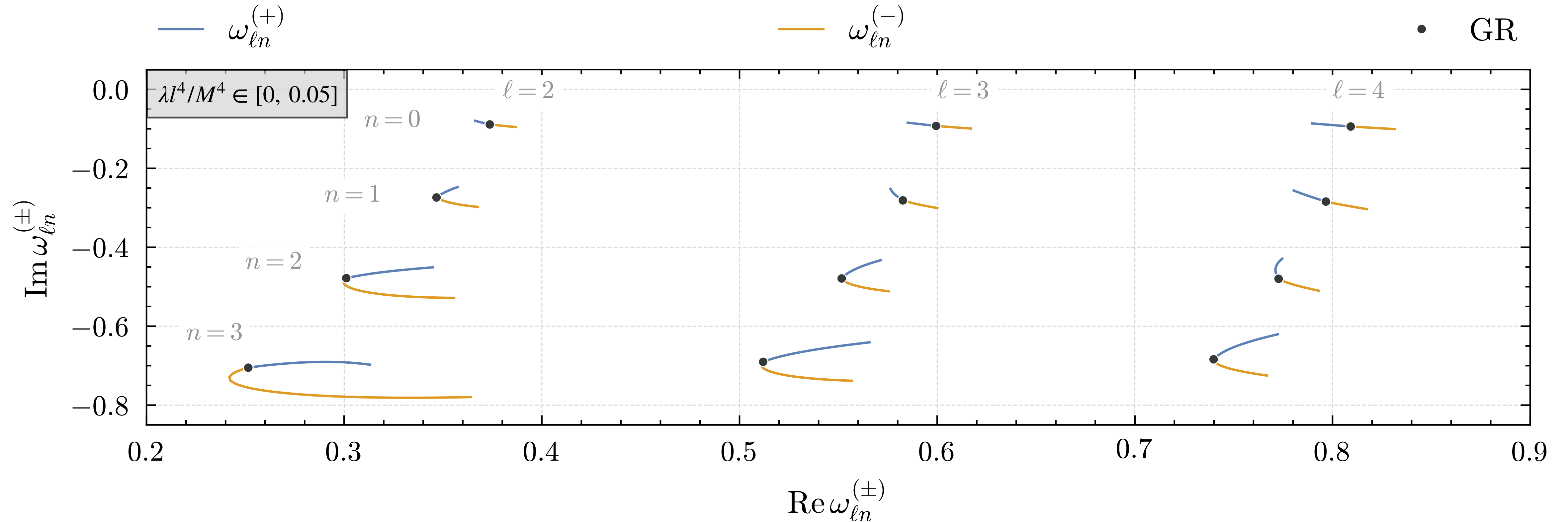
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- Scaling with  $\varepsilon = \lambda l^4 / M^4$  is nonlinear and depends on  $l$  and  $n$ .



# Overtones and the limits of the EFT description

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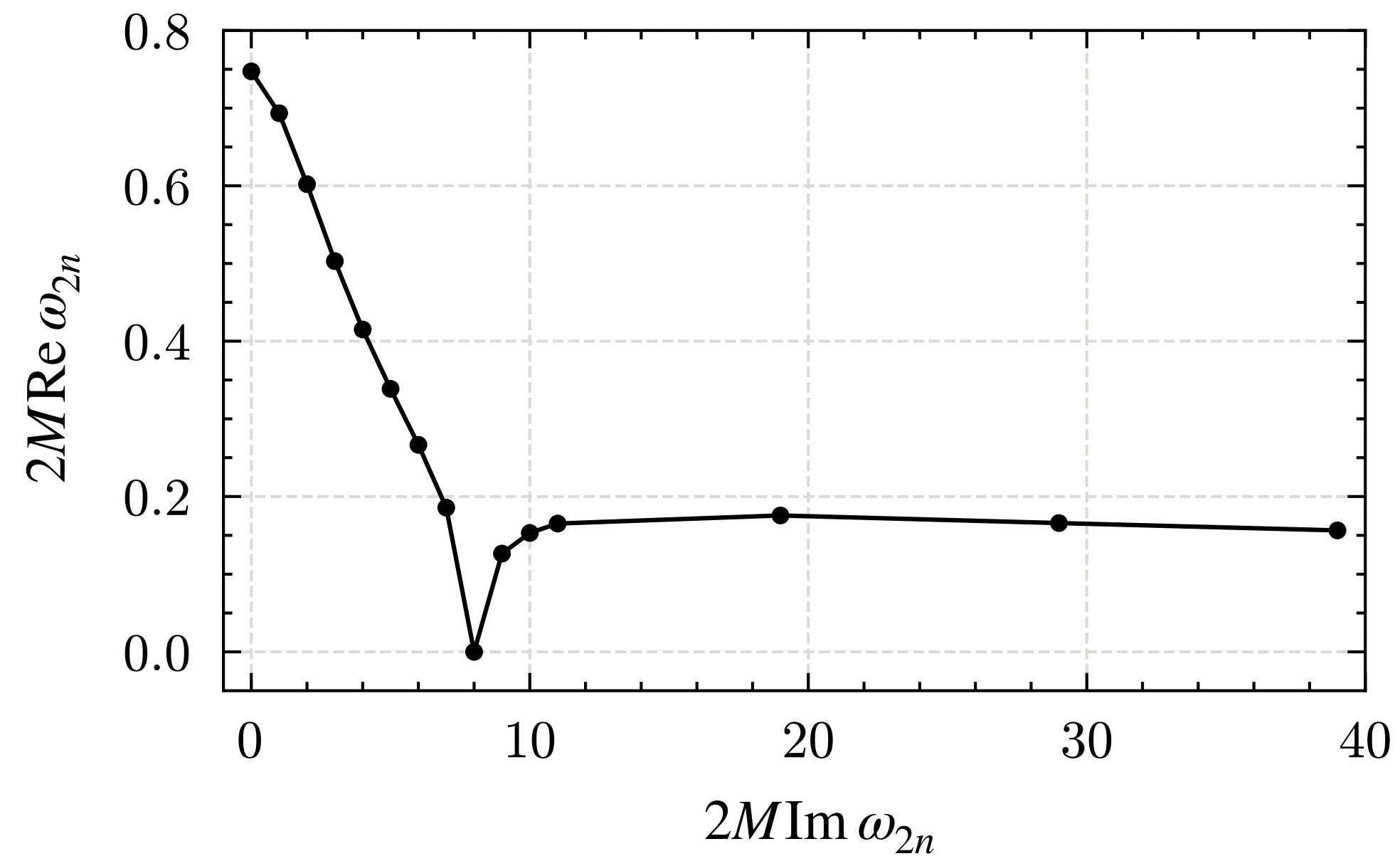
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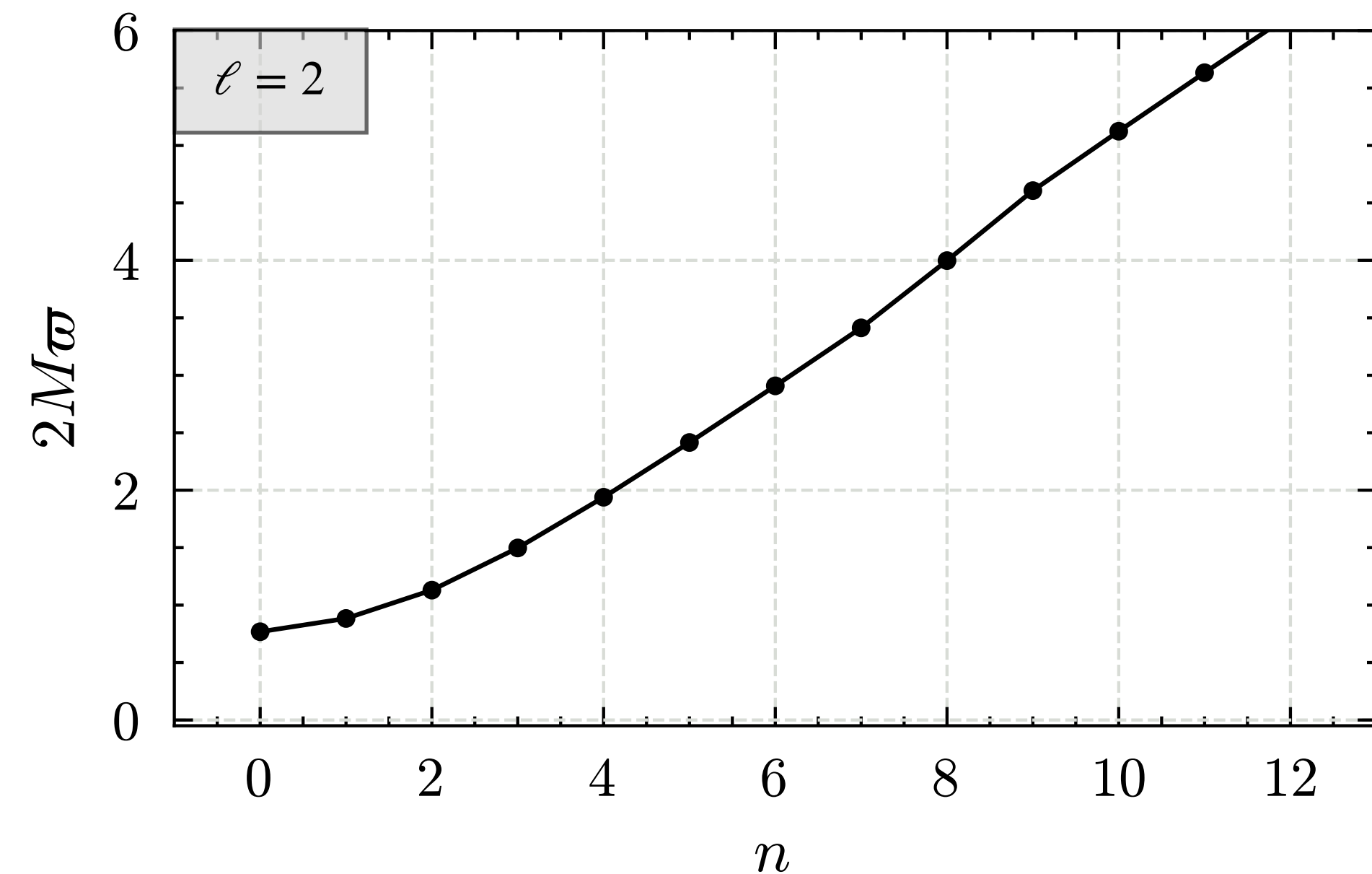
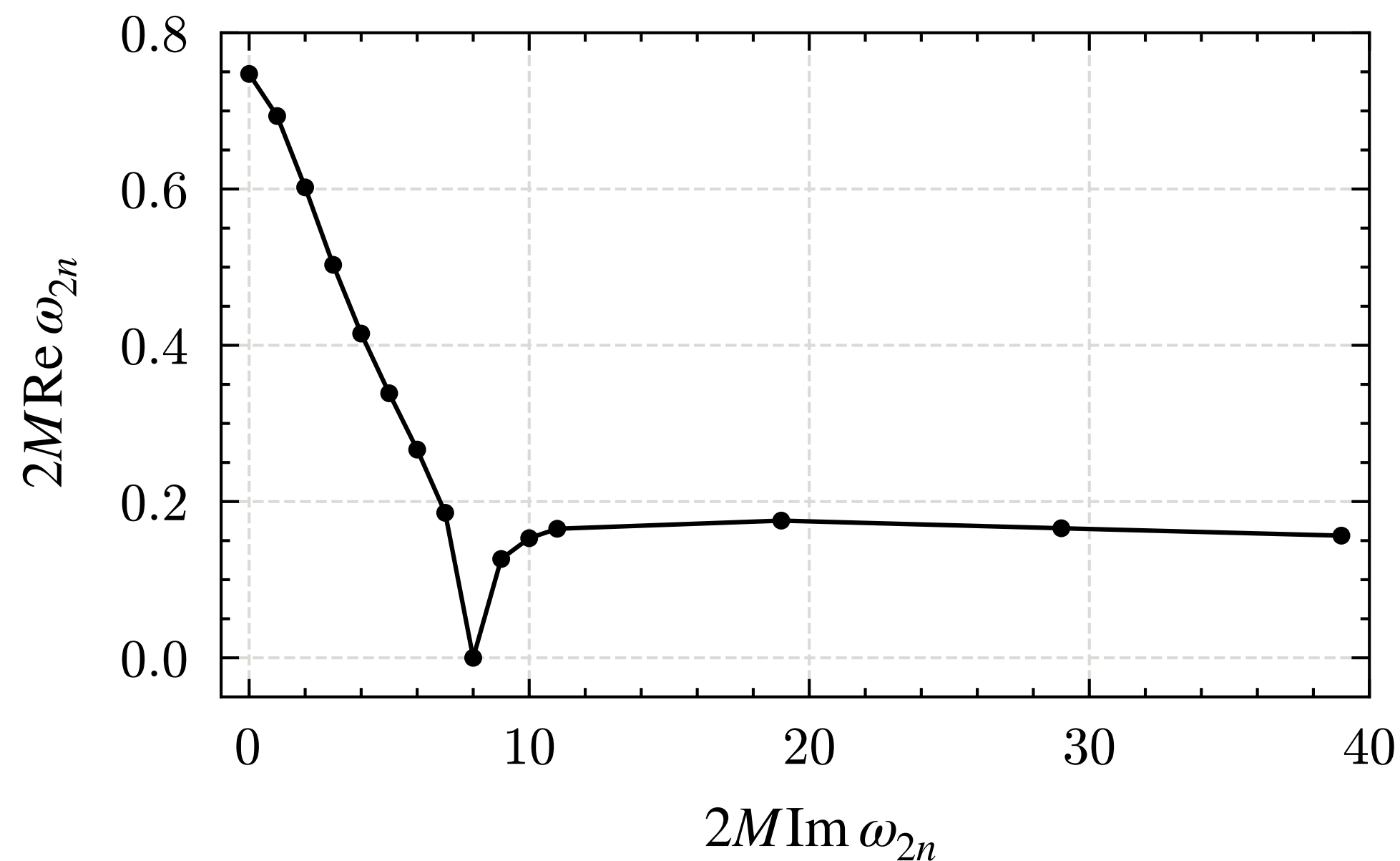


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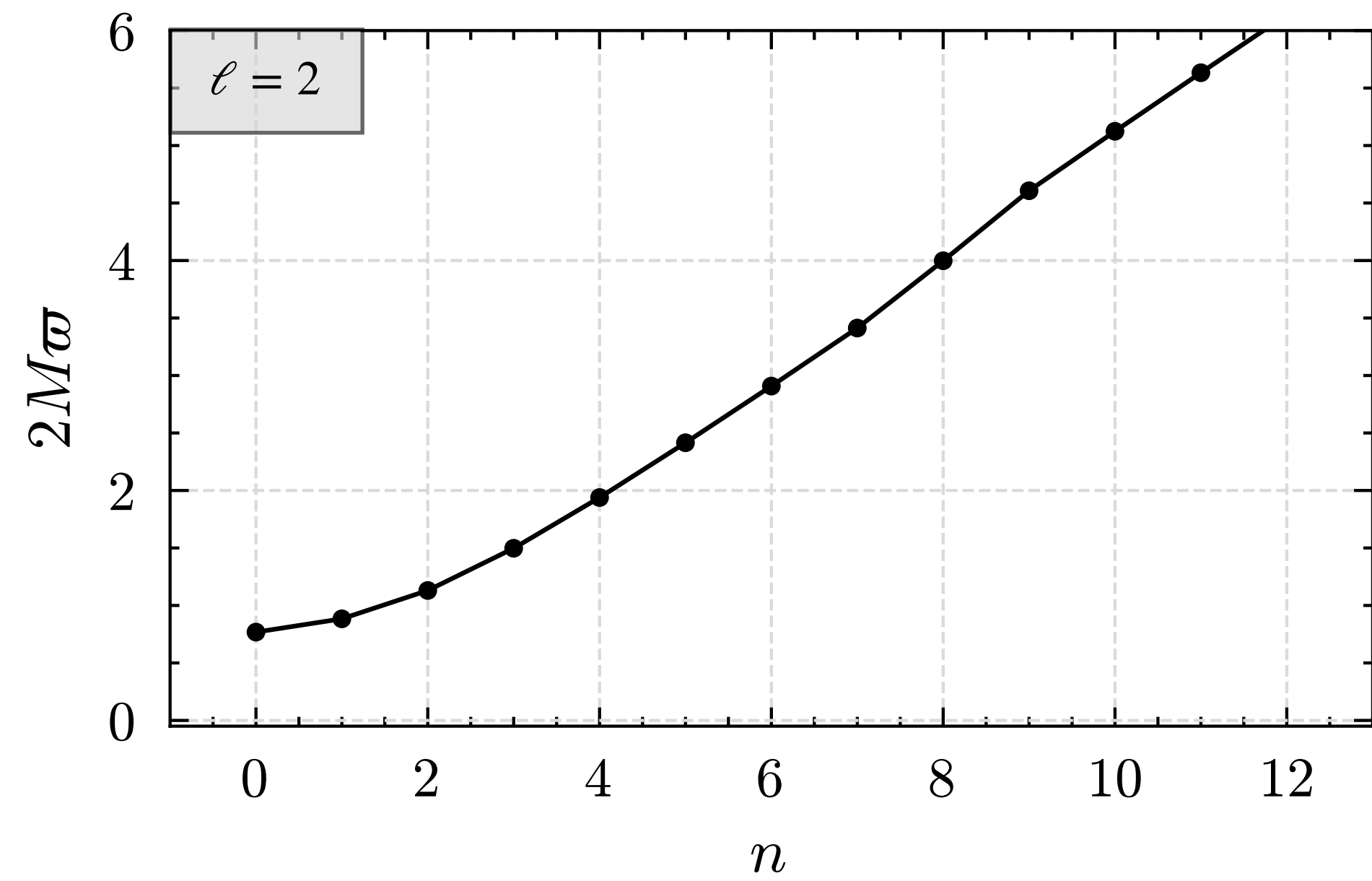
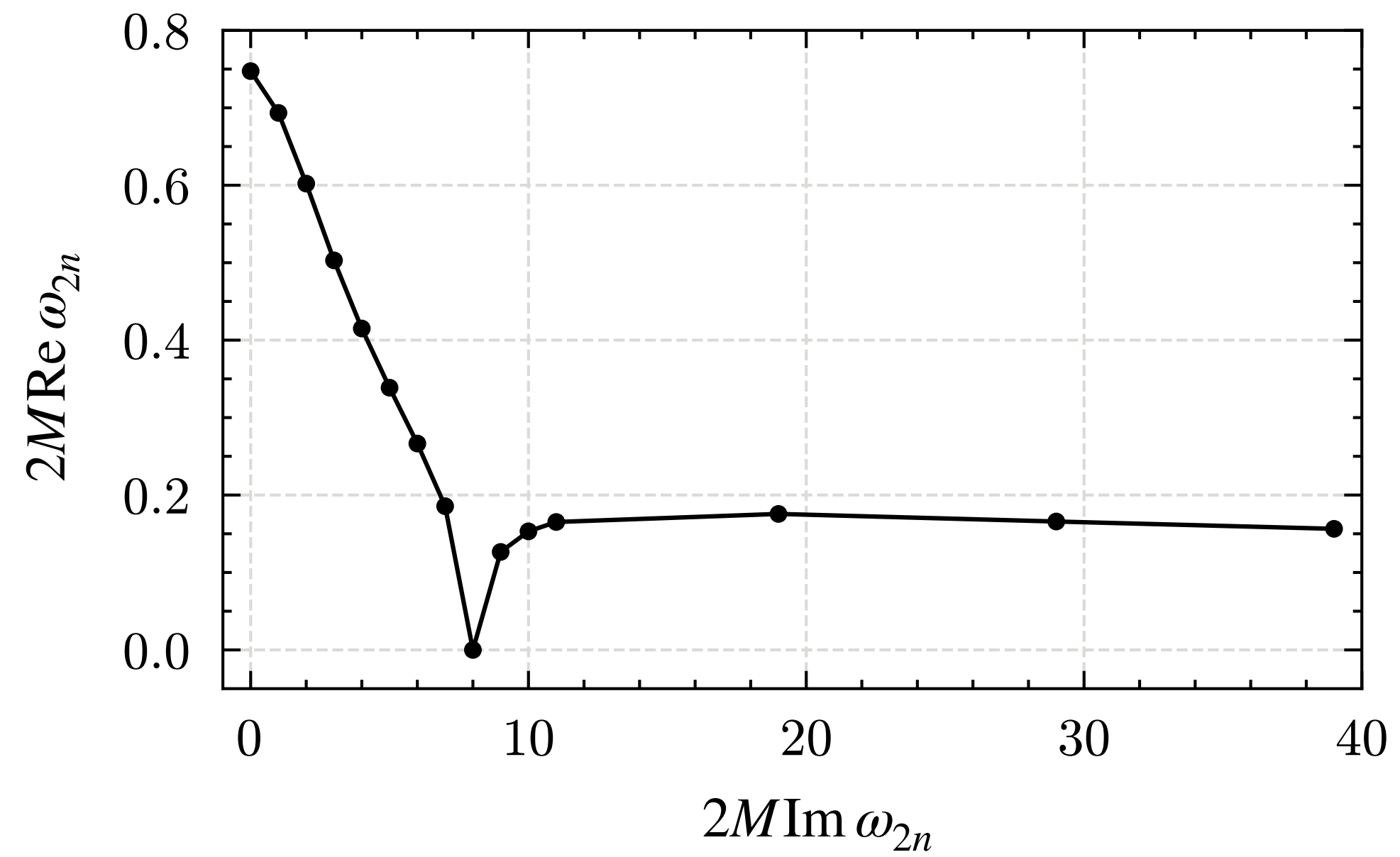
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- “Proper” mode frequency

$$f_{\ell n} = \varpi / (2\pi) = [\text{Re}(\omega_{\ell n})^2 + \text{Im}(\omega_{\ell n})^2]^{1/2}$$

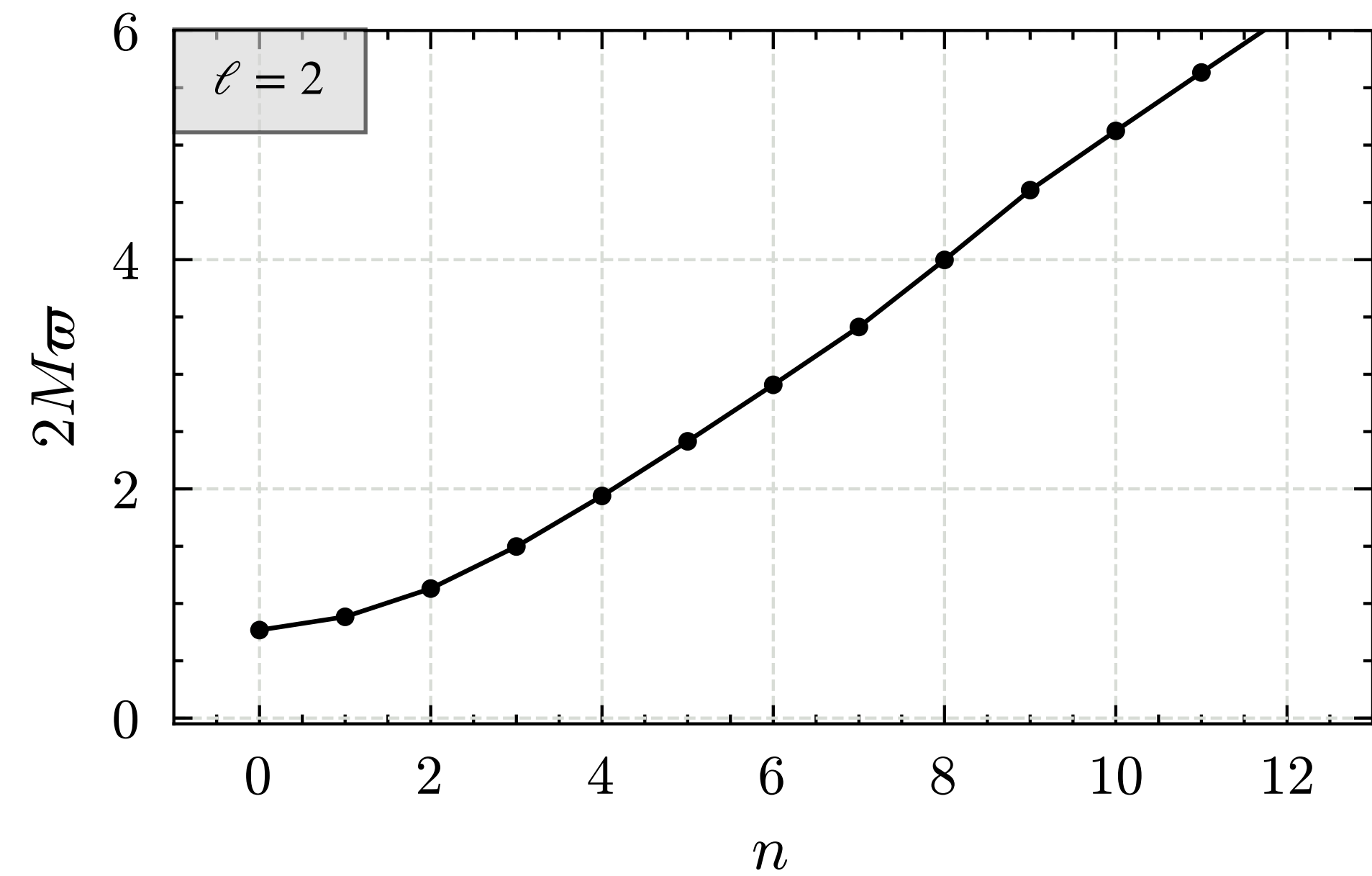
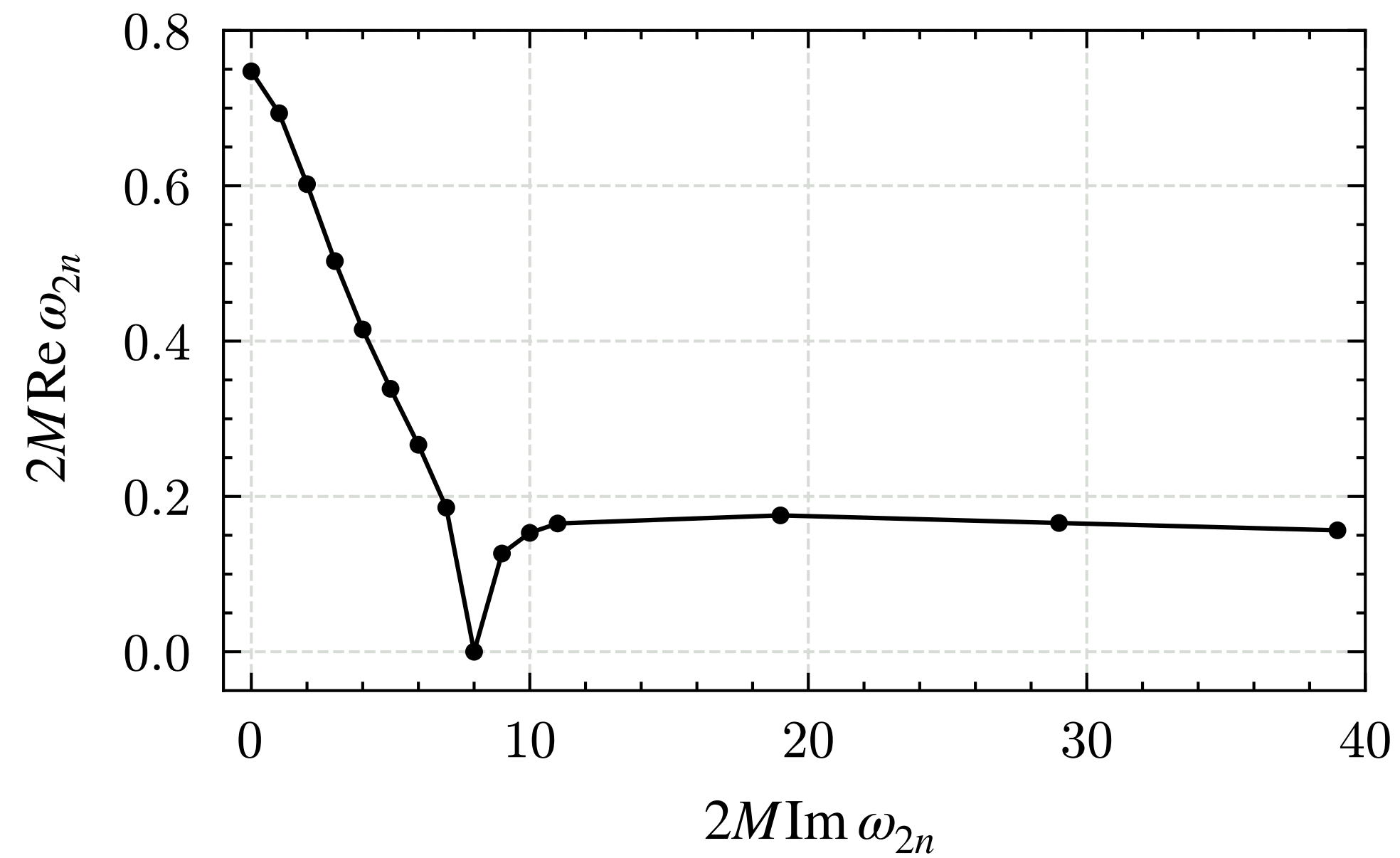


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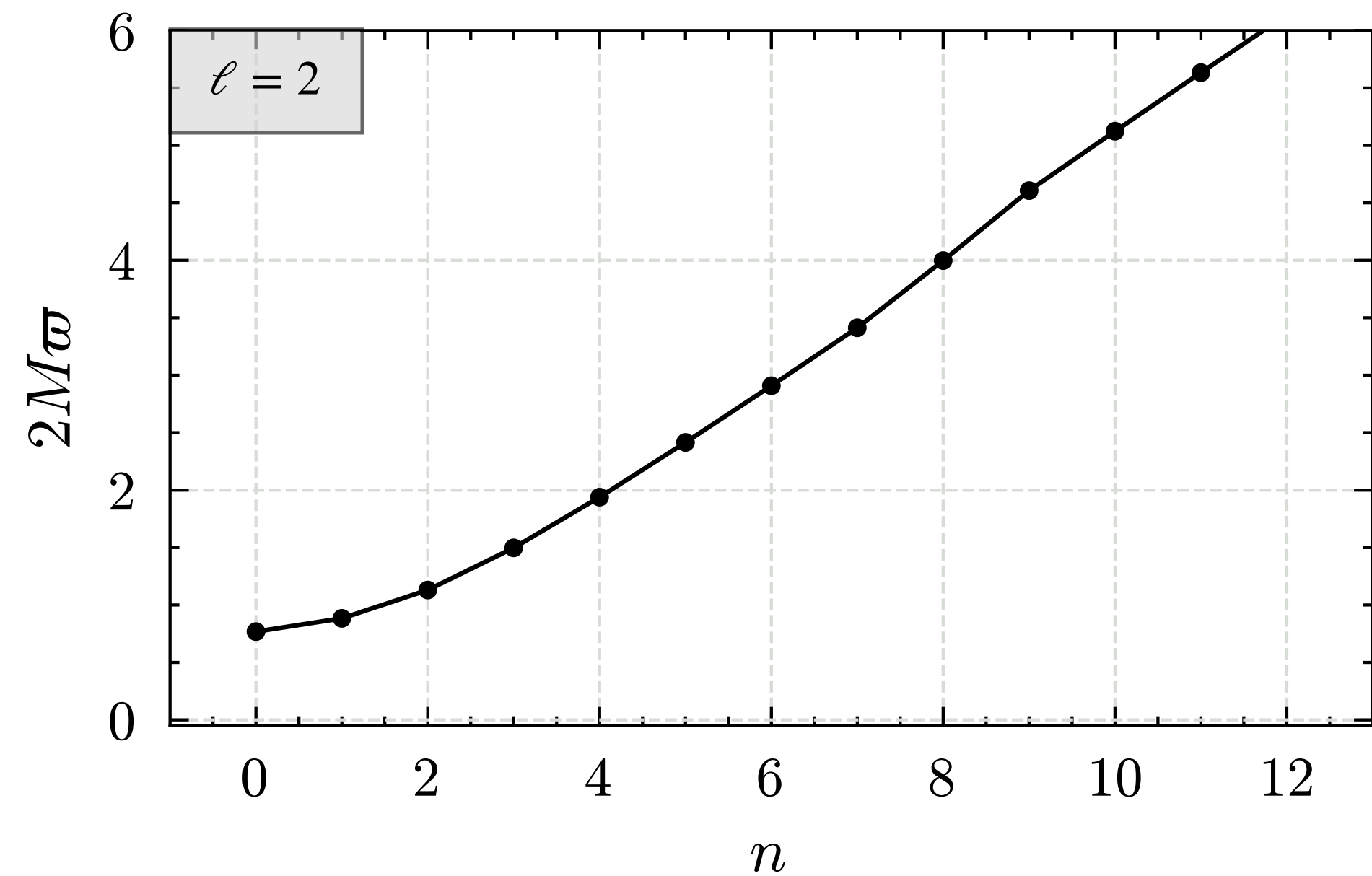
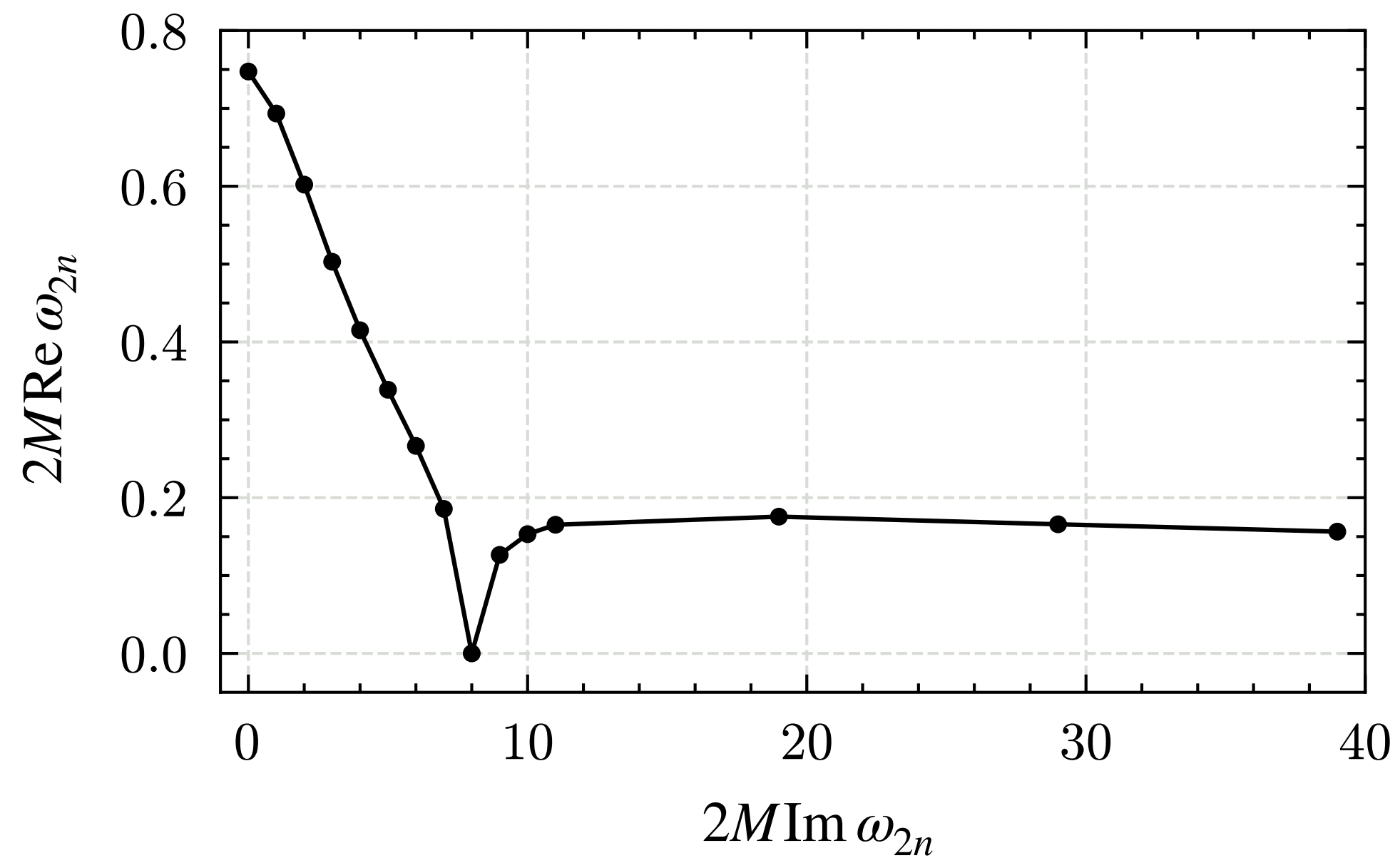
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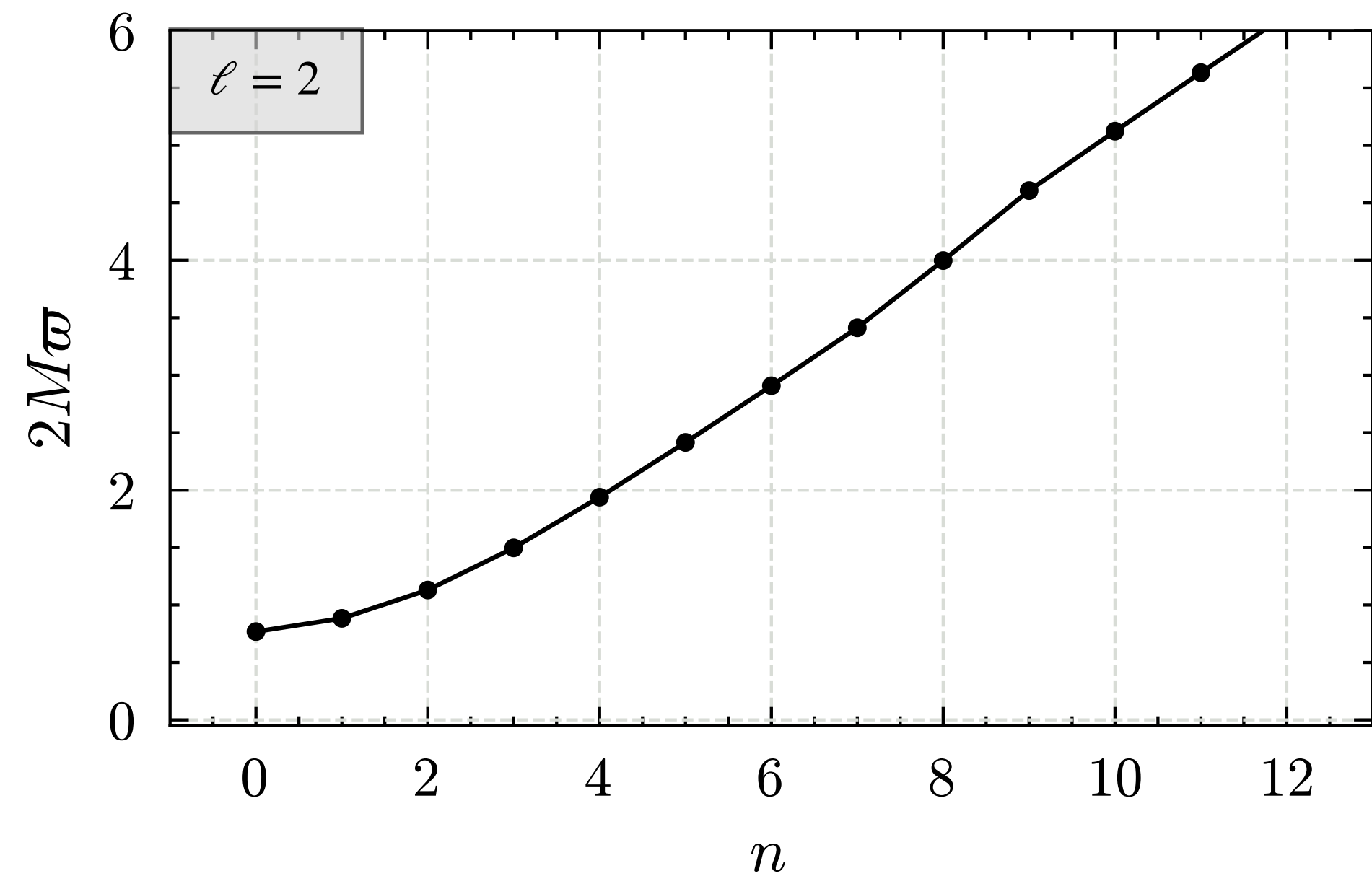
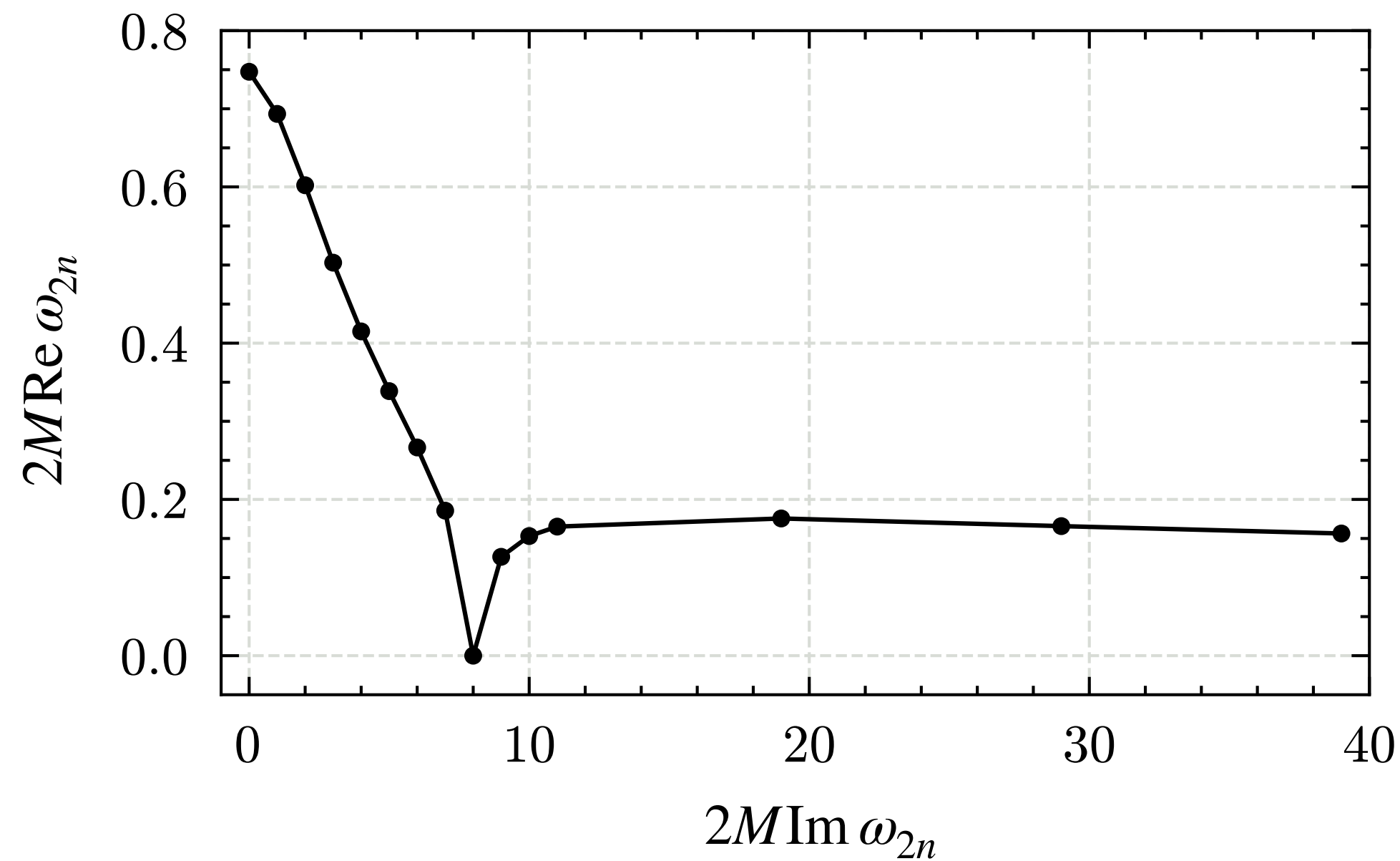


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- High-frequency waves probe deeper the potential barrier, hence the near-horizon region of the spacetime.





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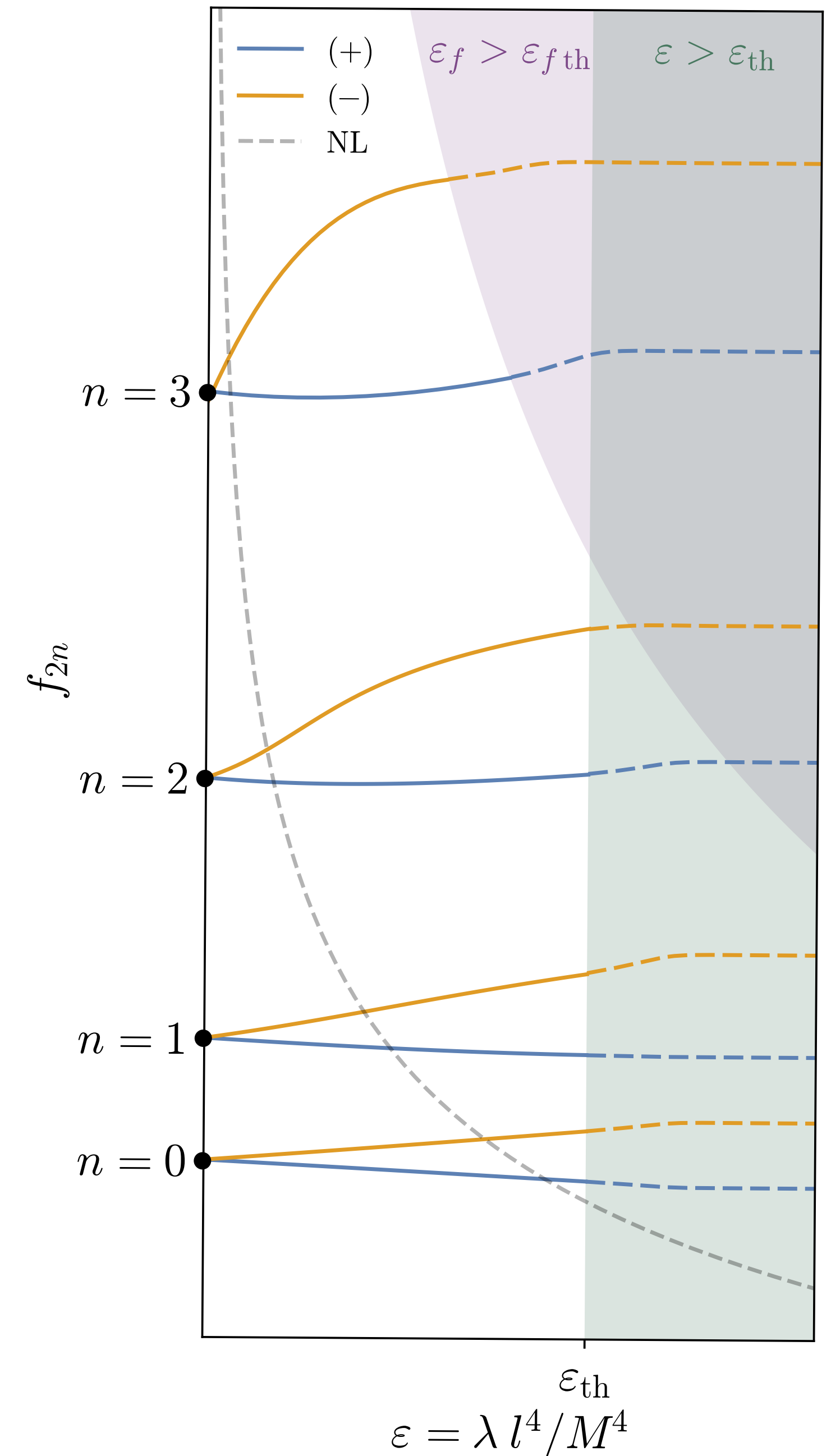
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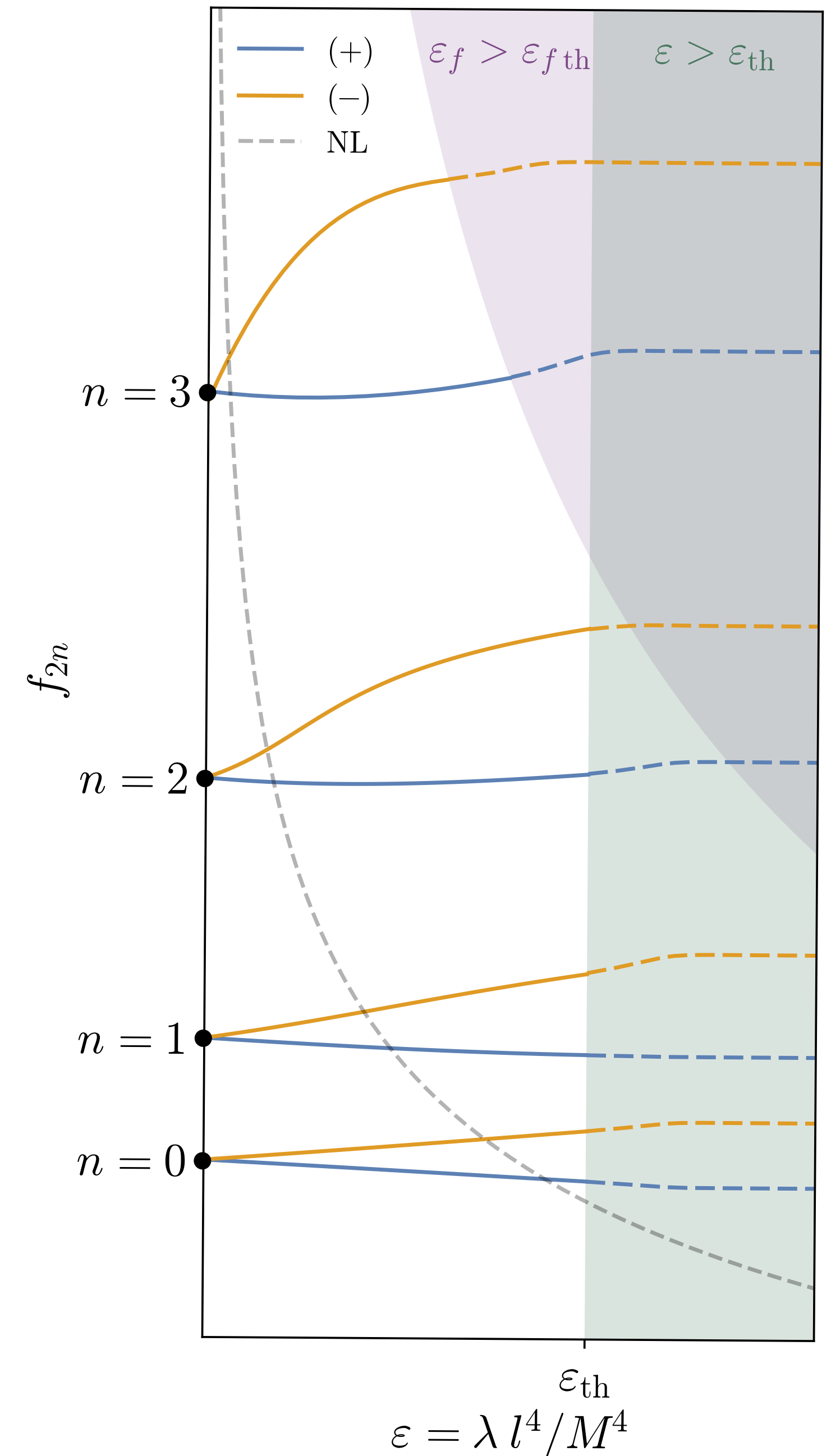
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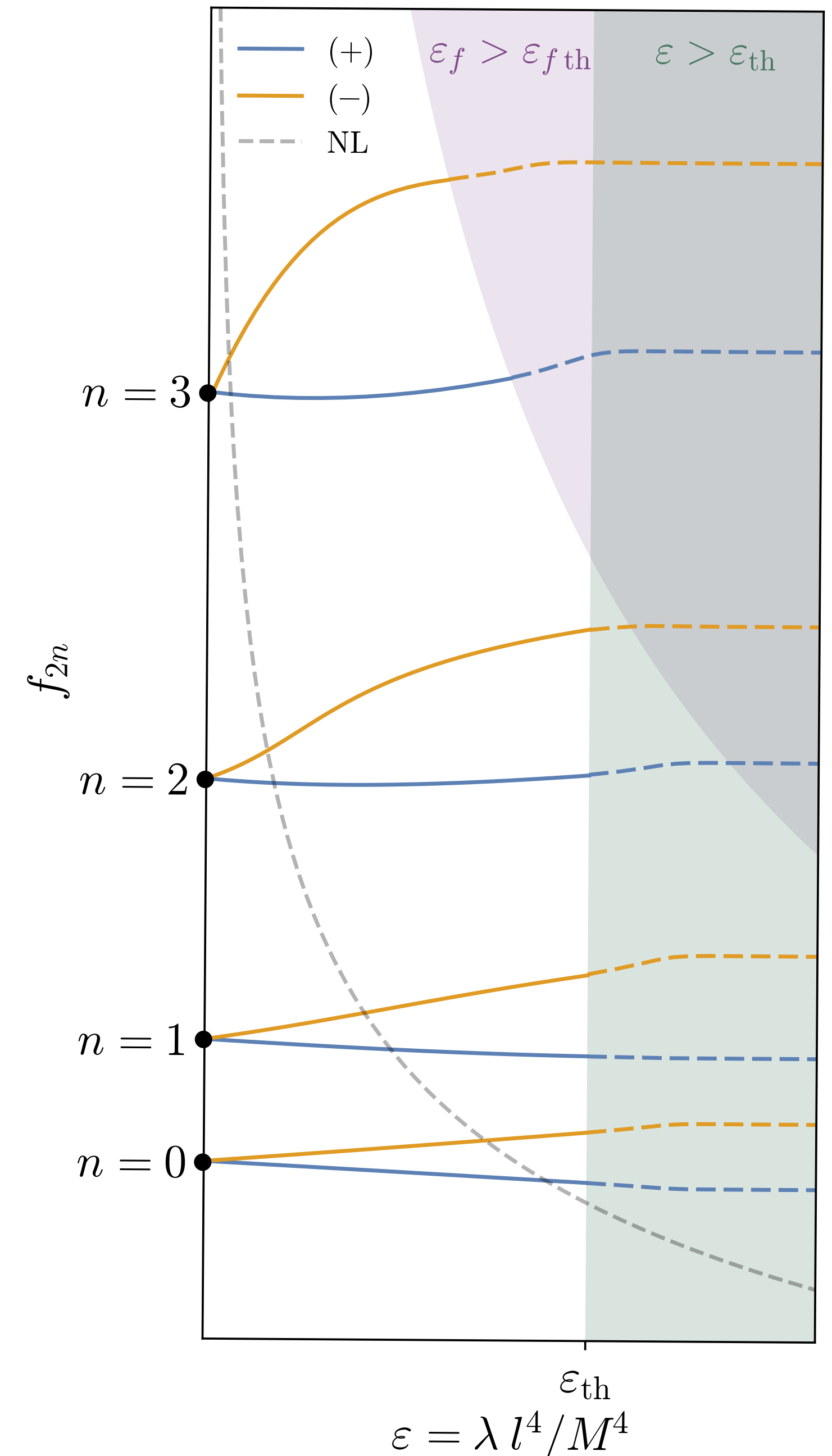
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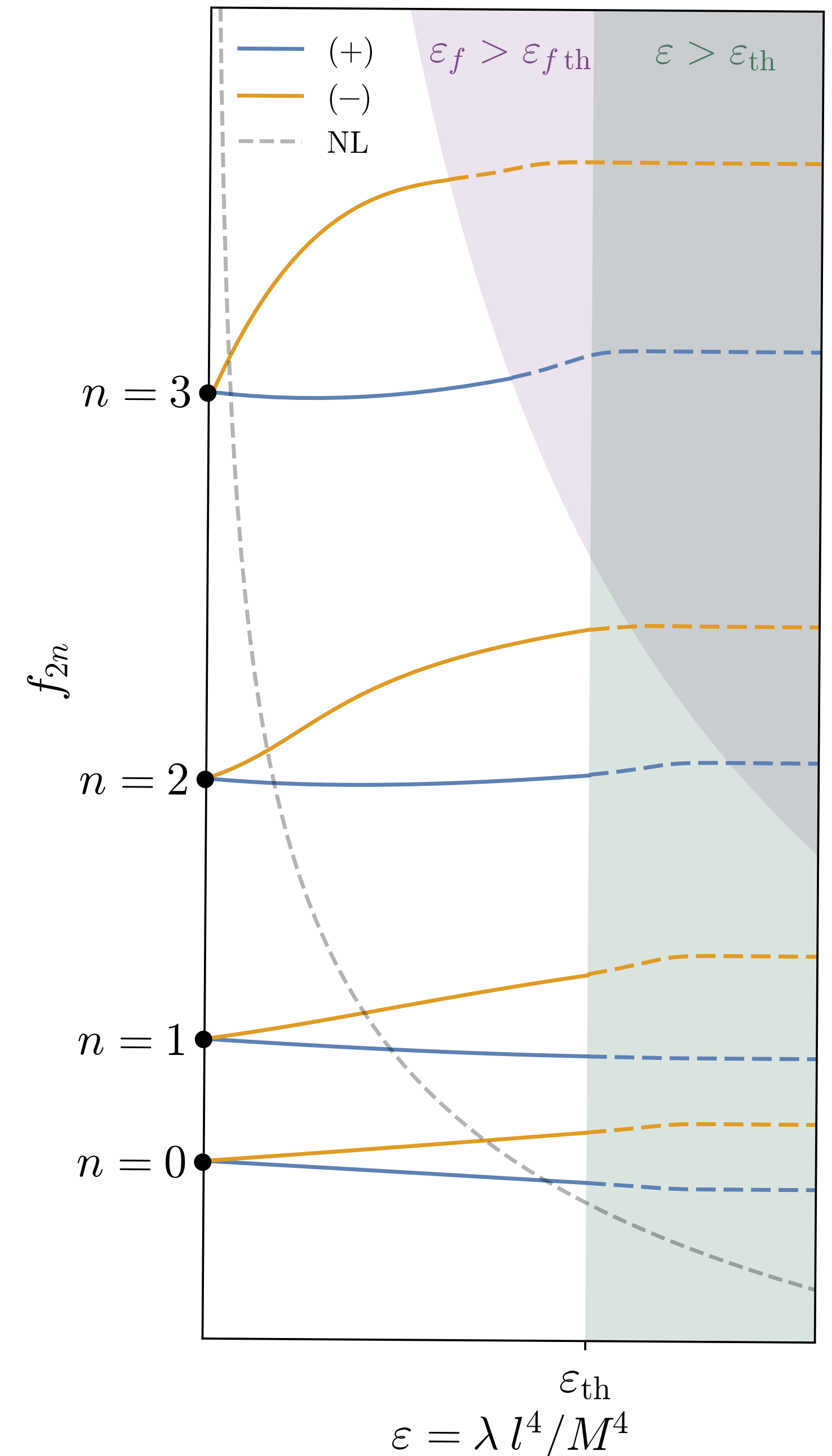
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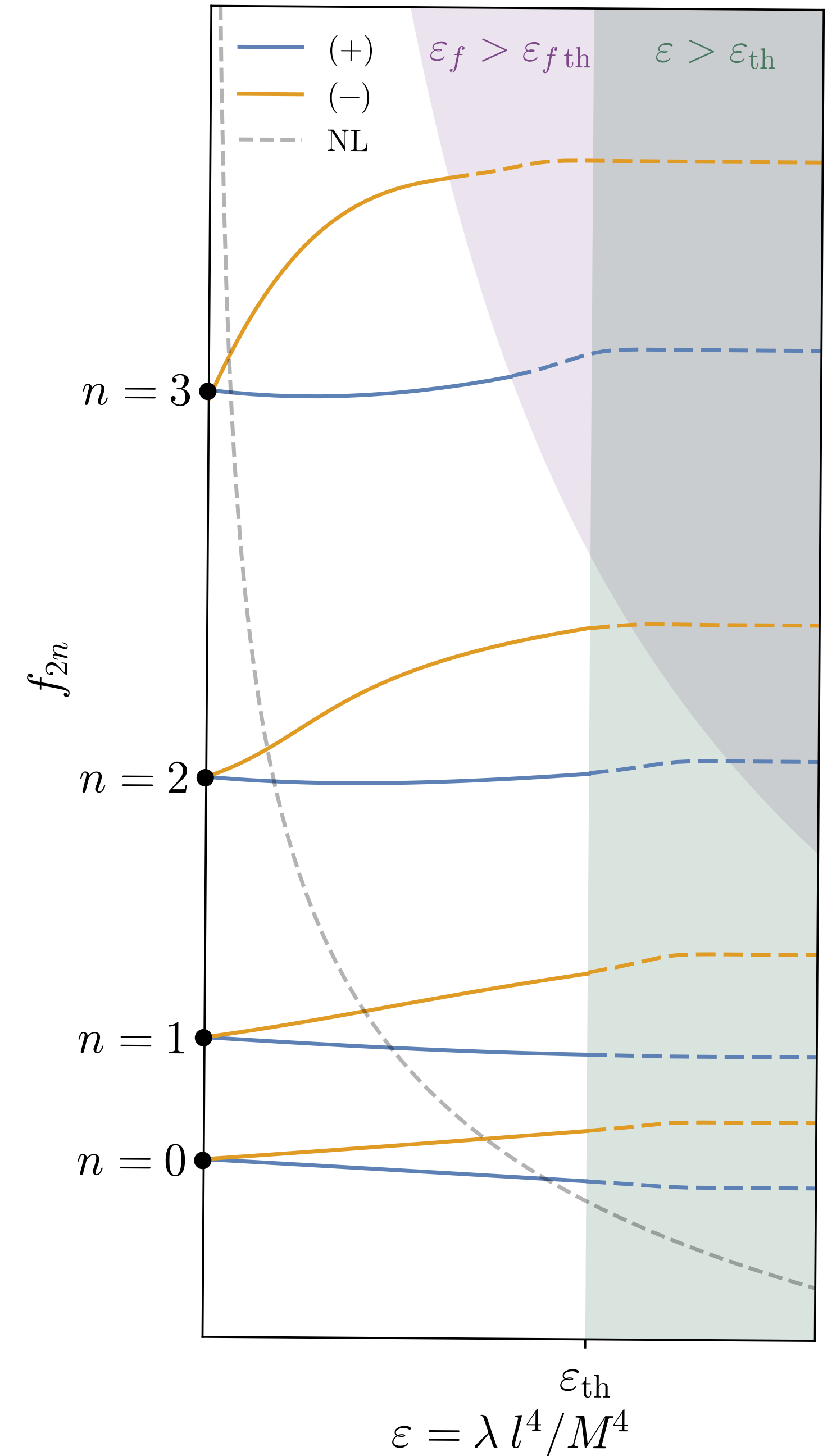
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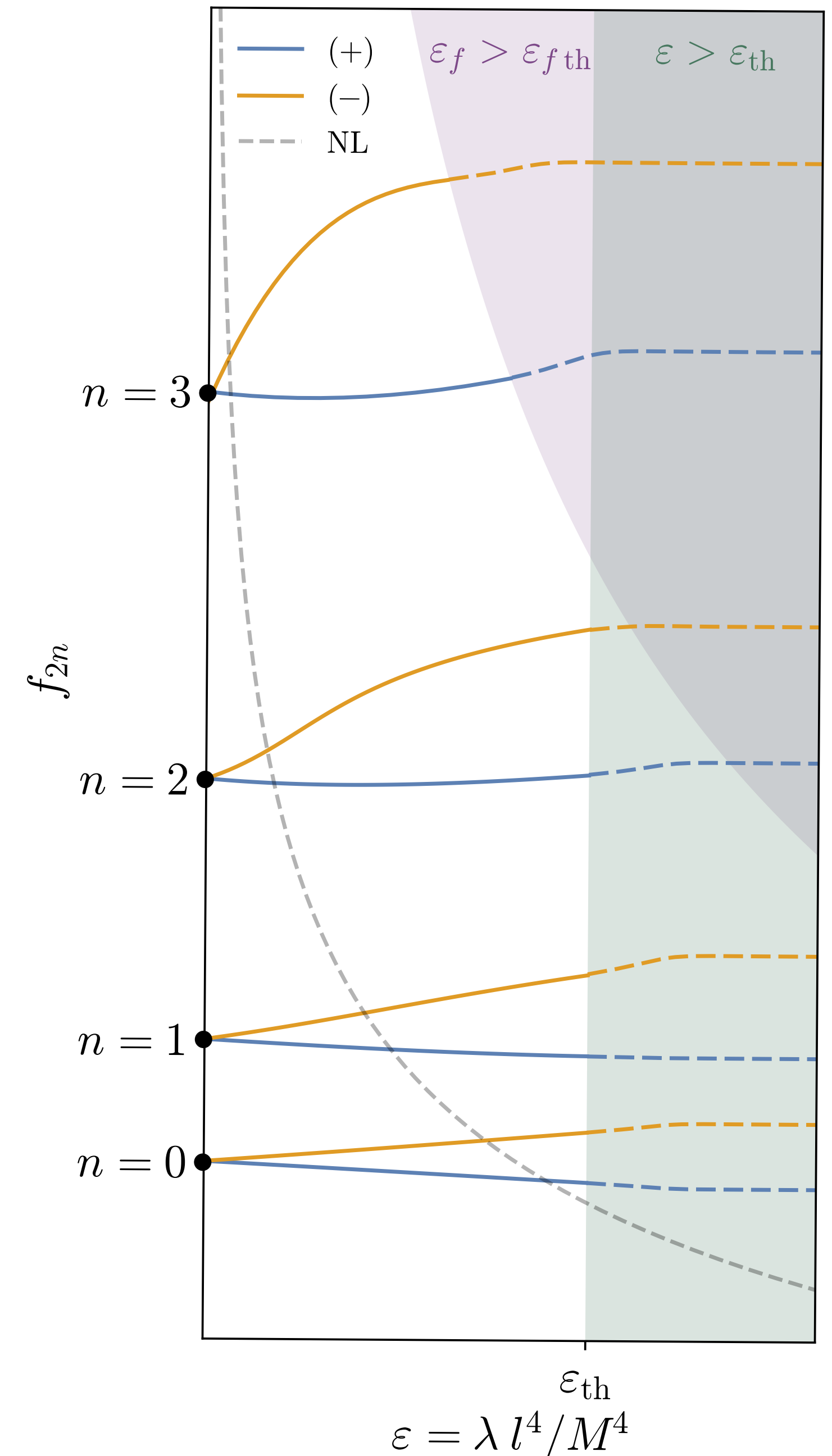
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The excitation coefficient  $C$  can be factorised into “background-” (**excitation factor  $B$** ) and “source-dependent” parts.

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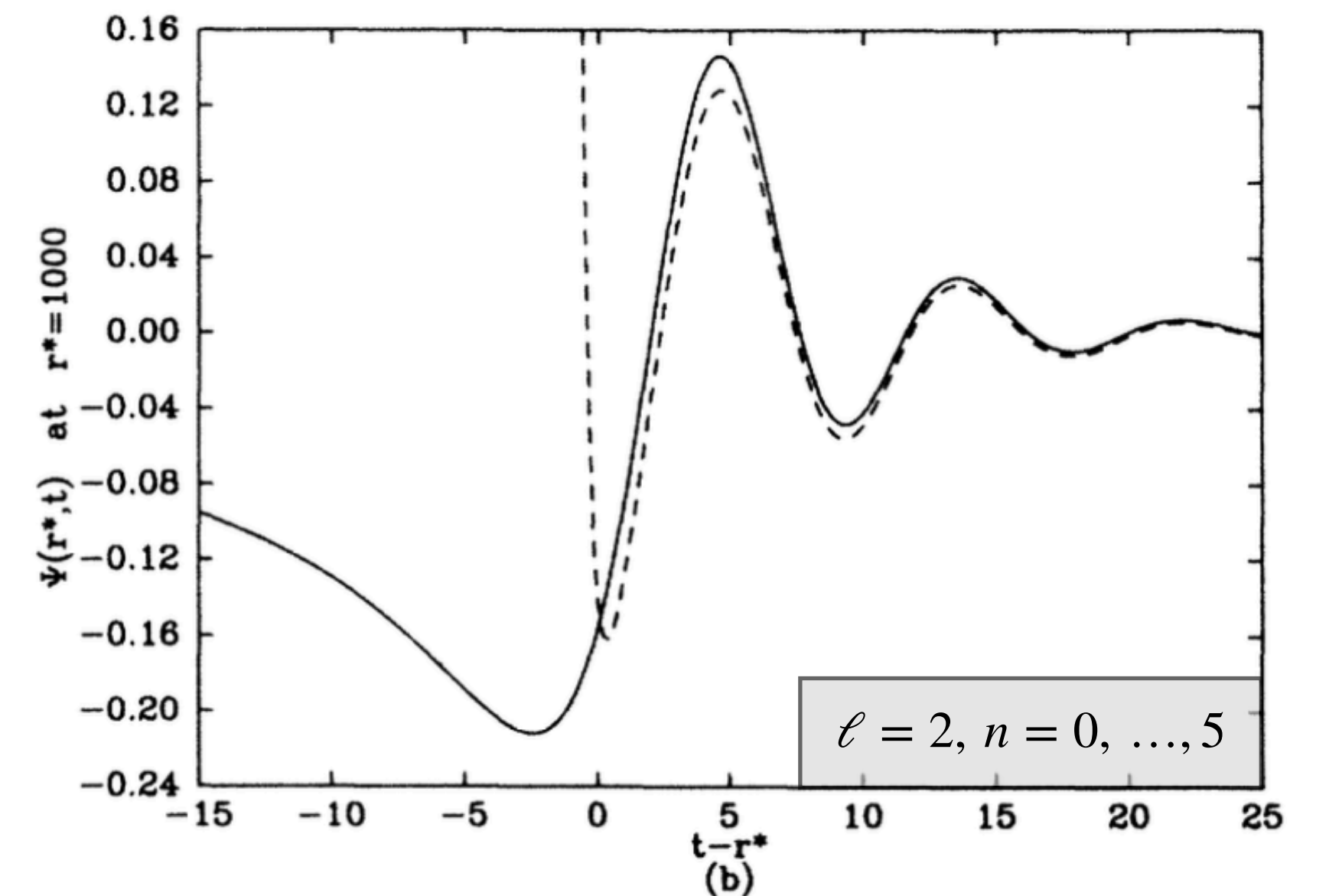
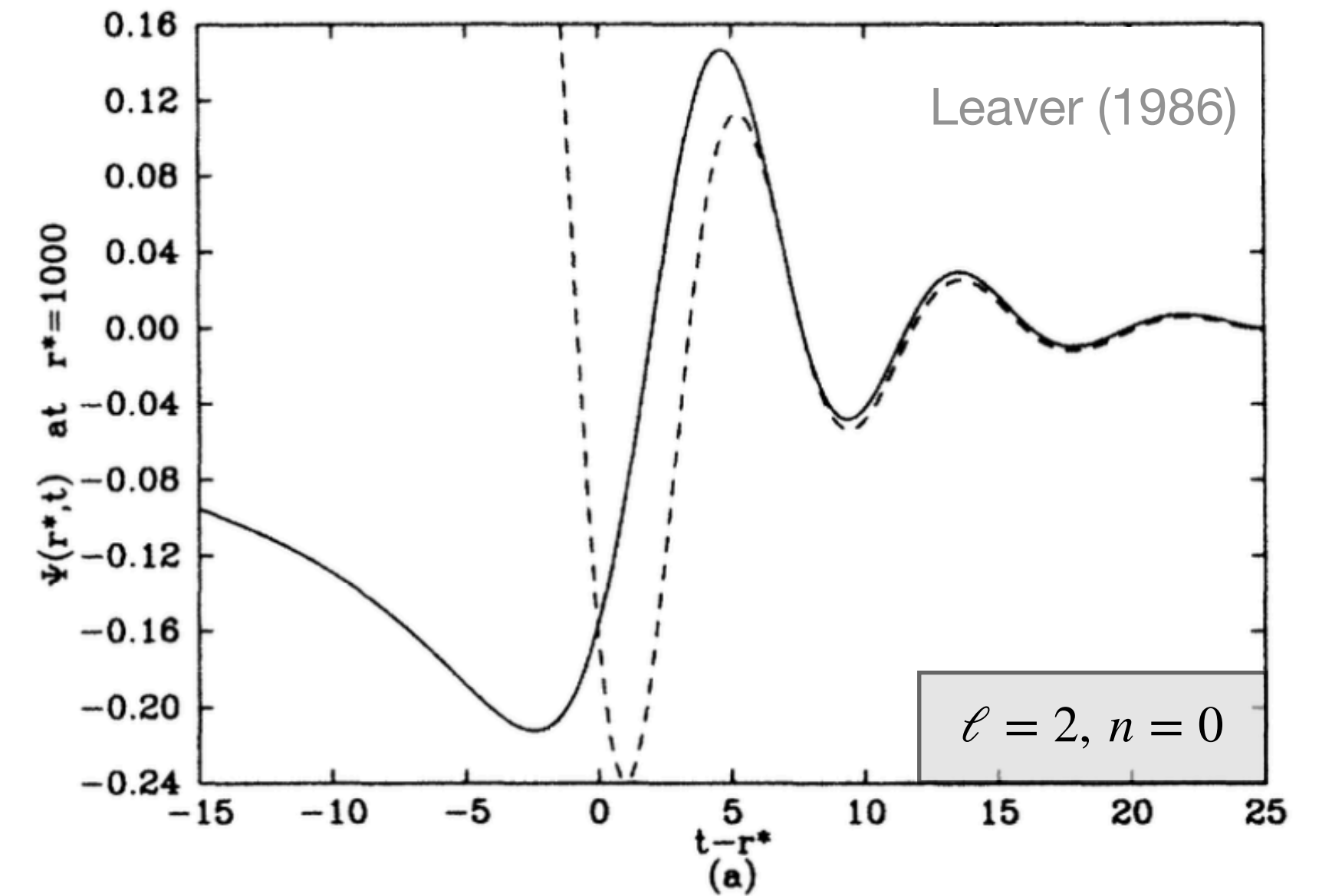
Leaver (1986), Andersson (1995)

$$\left[ \frac{d^2}{dx^2} + Q_{\ell m}(x, \omega, \vartheta) \right] \psi_{\ell m}(x) = s_{\ell m}(x)$$

Quasinormal mode contribution the **time-domain** signal:

$$\psi_{\ell m}(t, x) = -\operatorname{Re} \sum_n \left[ C_{\ell mn} e^{-i\omega_{\ell mn}(t-x)} \right]$$

The excitation coefficient  $C$  can be factorised into “background-” (**excitation factor  $B$** ) and “source-dependent” parts.



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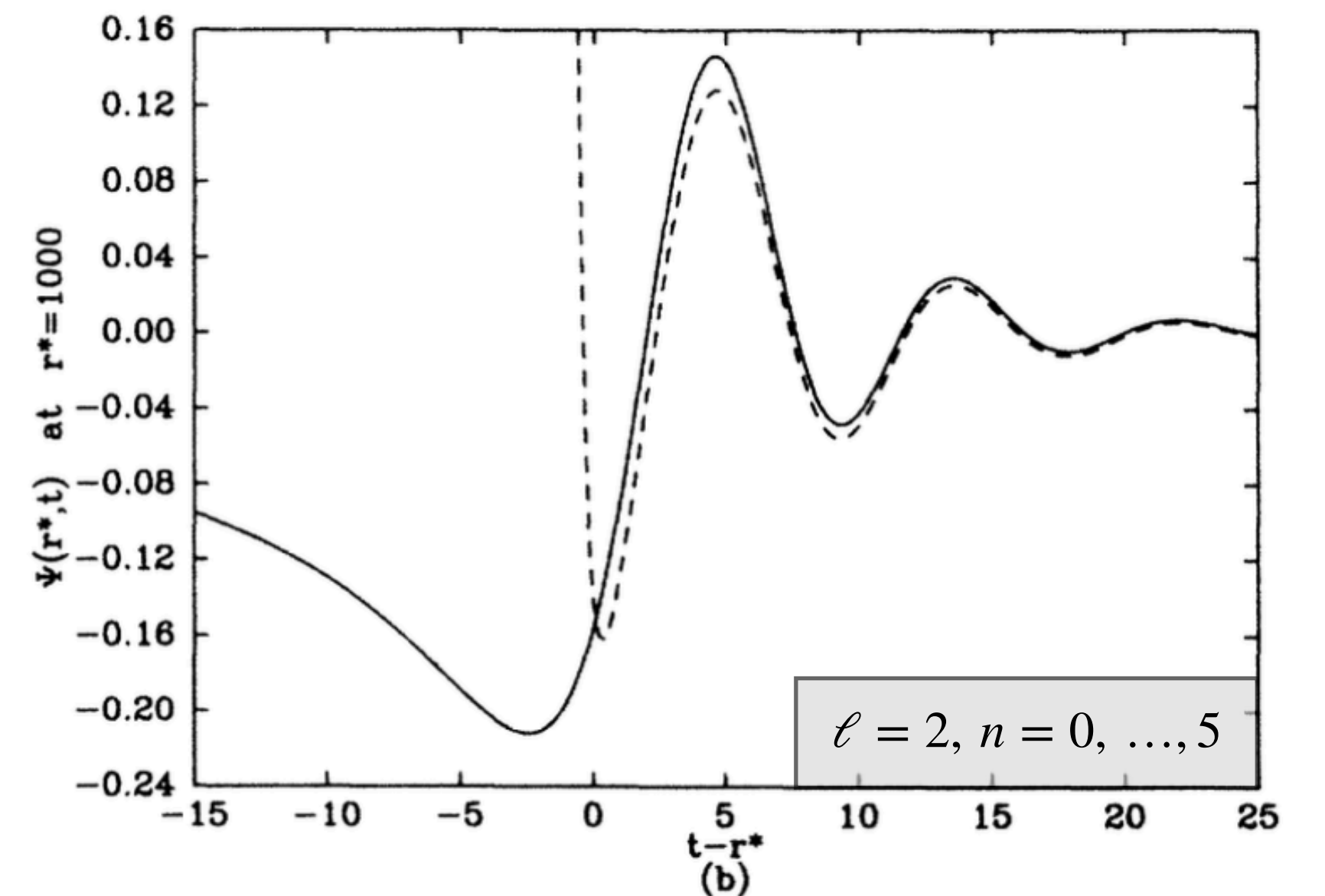
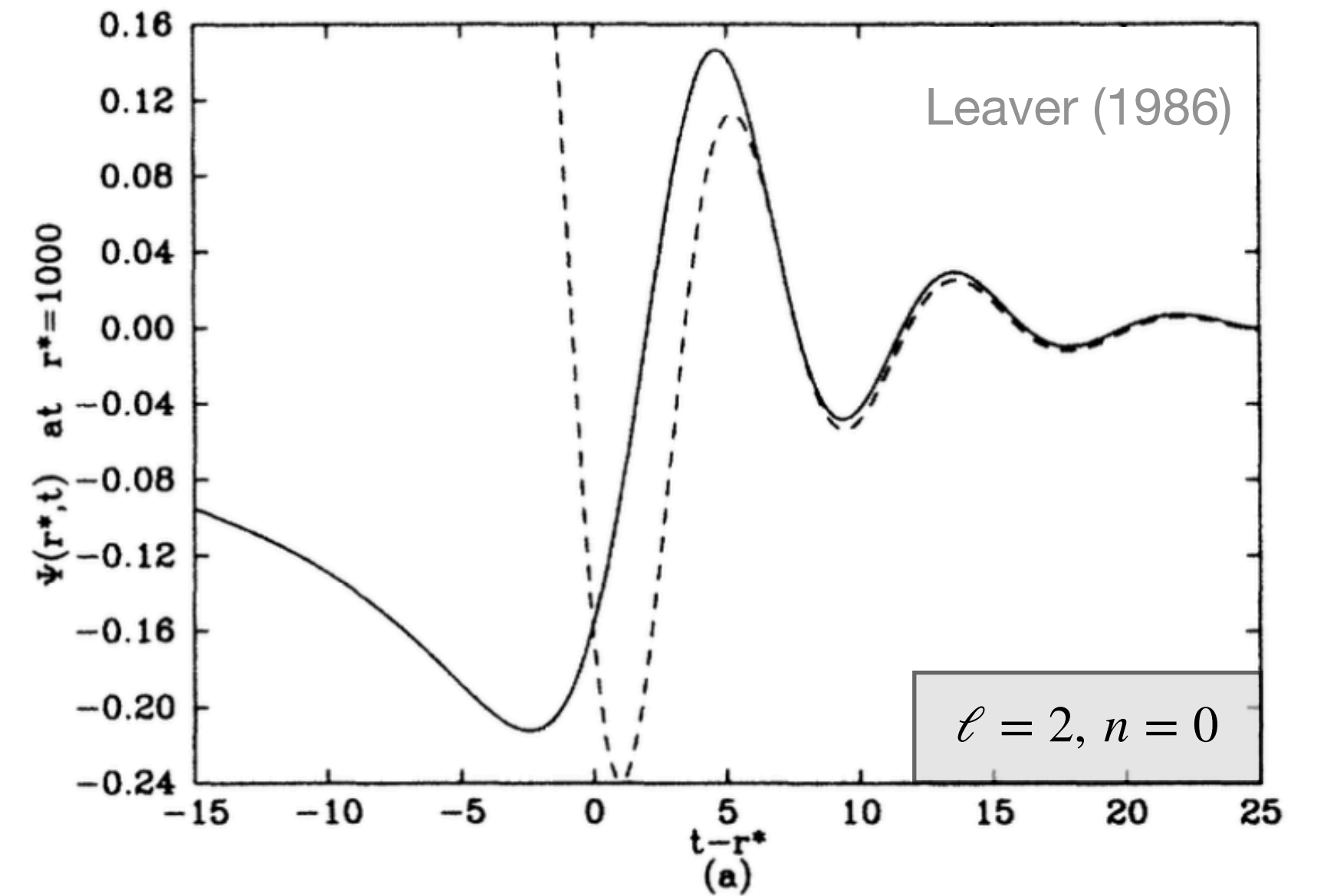
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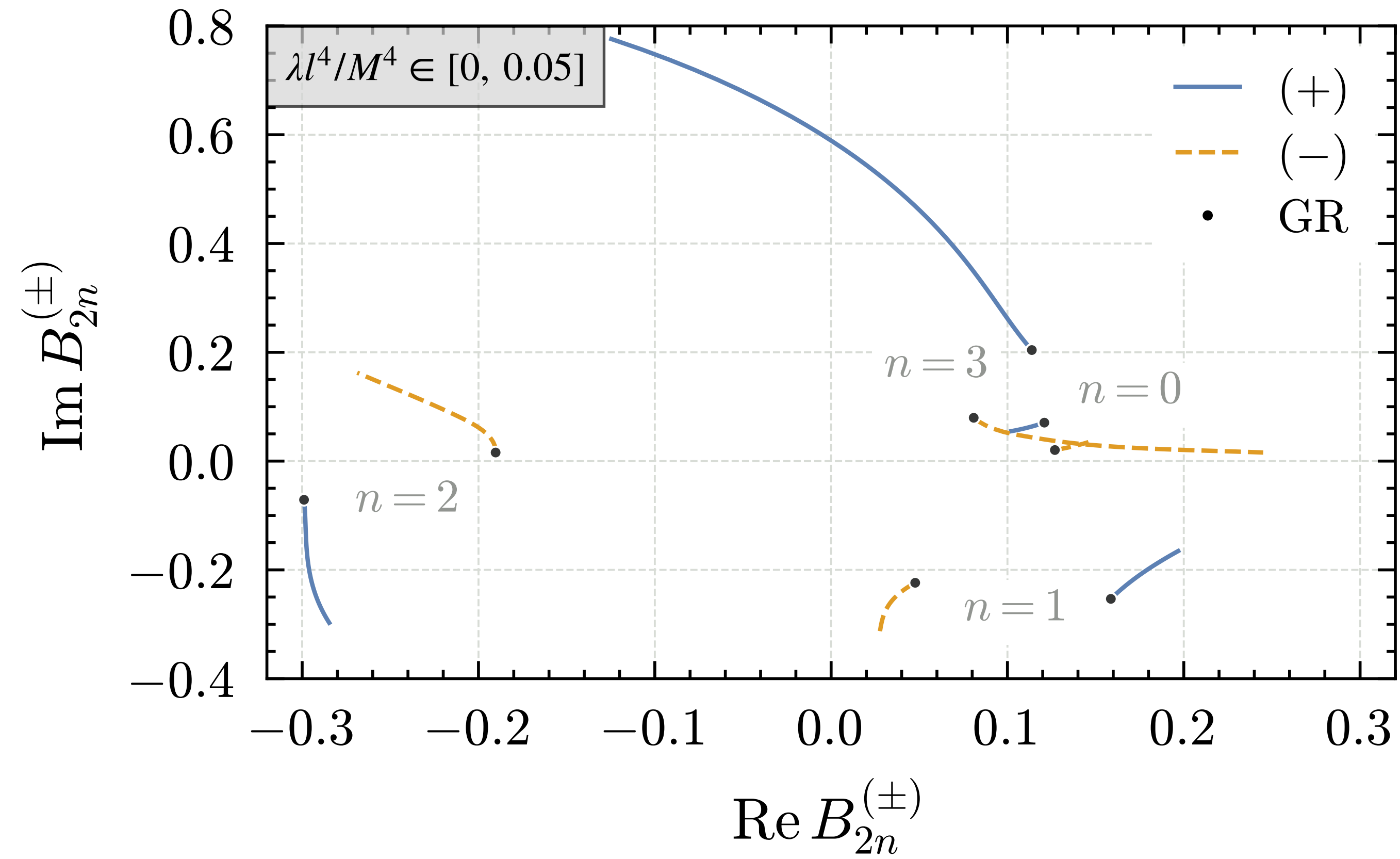
**Neither** have been computed in any beyond-general-relativity spacetime. I don't know why!





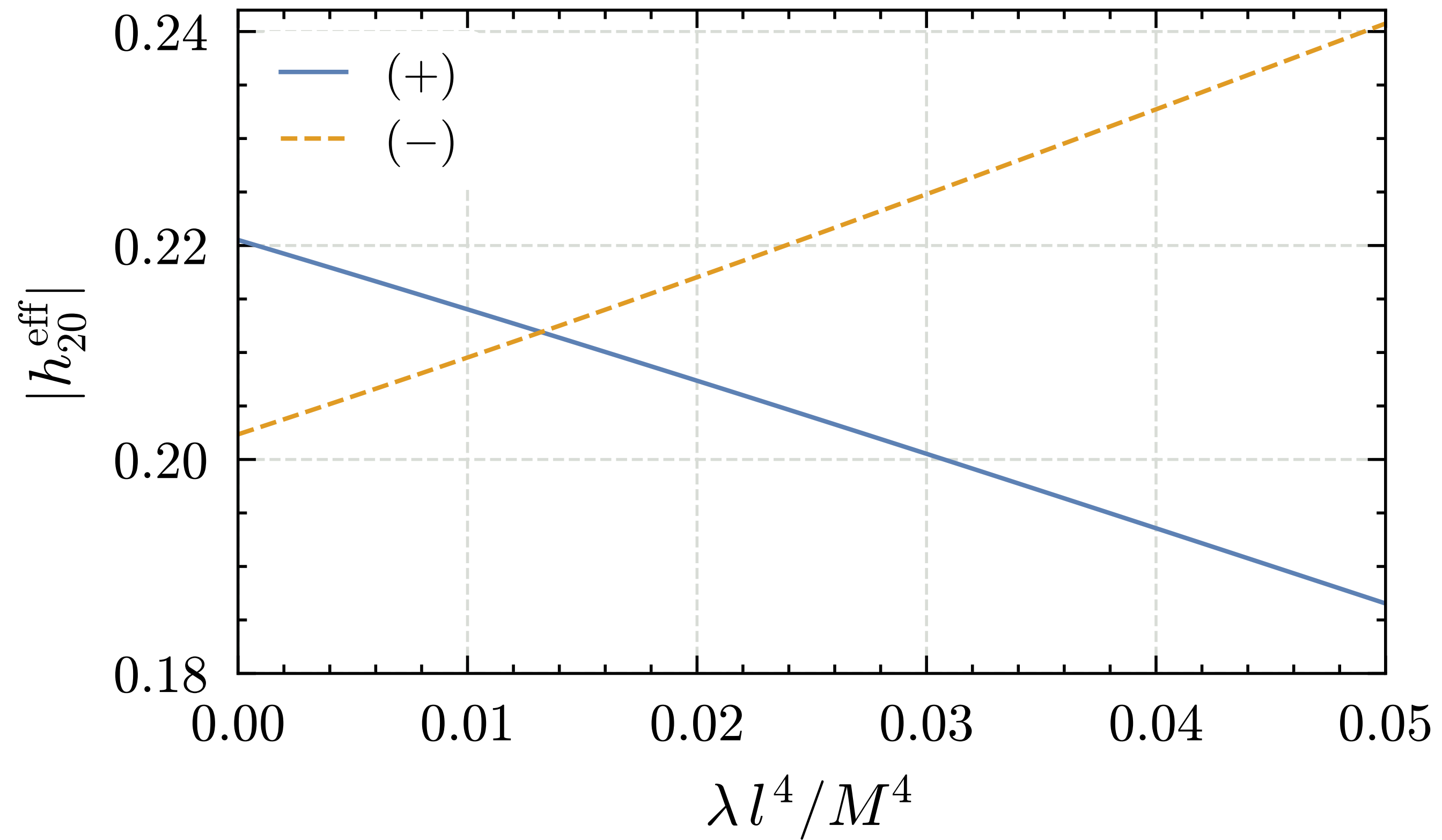
# Quadrupole excitation factors

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# Effective gravitational wave amplitudes

Glampedakis and Andersson (2001, 2003)



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- A lot to explore!