

p -Wasserstein metrics for gas networks

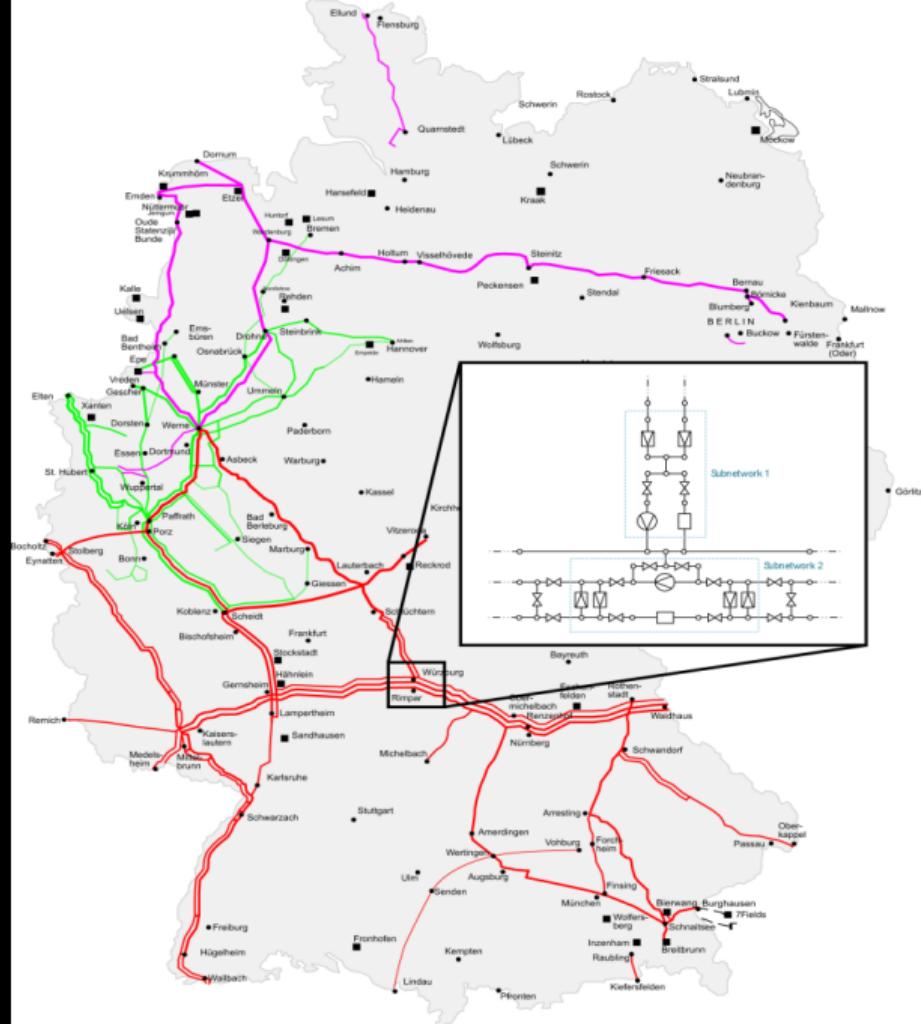
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p -Wasserstein metric



For two probability measures $\mu_0, \mu_T \in \mathbb{P}(\Omega)$ and $p \in [1, \infty)$:

$$\text{(STATIC)} \quad W_p(\mu_0, \mu_T) = \min_{\pi \in \Pi(\mu_0, \mu_T)} \left\{ \int_{\Omega \times \Omega} |x - y|^p \, d\pi(x, y) \right\}^{\frac{1}{p}}$$

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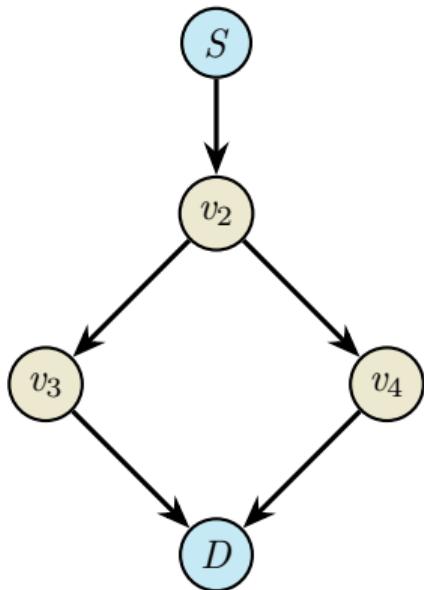
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Based on the work of Benamou and Brenier [BBOO]:

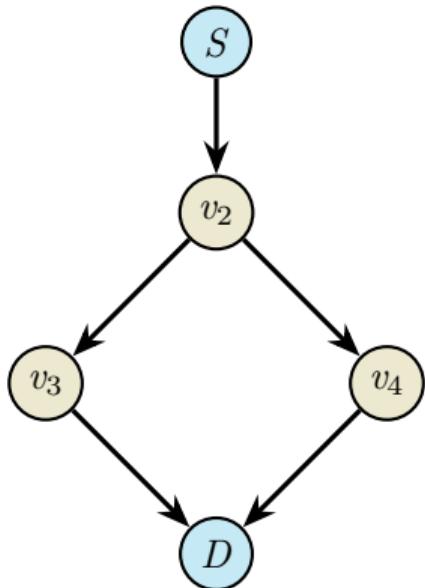
For a reference measure η , density μ and mass flux j :

$$\text{(DYNAMIC)} \quad W_p(\mu_0, \mu_T) = \min_{\substack{\mu \in \mathcal{M}_+(\Omega \times [0, T]), \\ j \in \mathcal{M}(\Omega \times [0, T])}} \left\{ T^{p-1} \int_0^T \int_{\Omega} \frac{|j(x, t)|^p}{\mu(x, t)^{p-1}} d\eta dt \middle| \begin{array}{l} \frac{\partial \mu}{\partial t} + \frac{\partial j}{\partial x} = 0, \\ \mu|_{t=0} = \mu_0, \\ \mu|_{t=T} = \mu_T \end{array} \right\}^{\frac{1}{p}}$$

Gas networks as metric graphs [FBP24]



- 1) Gas flow in individual pipes
- 2) Mass conservation at vertices
- 3) Storage of gas at interior vertices
- 4) Time-(in)dependent boundary conditions at boundary vertices

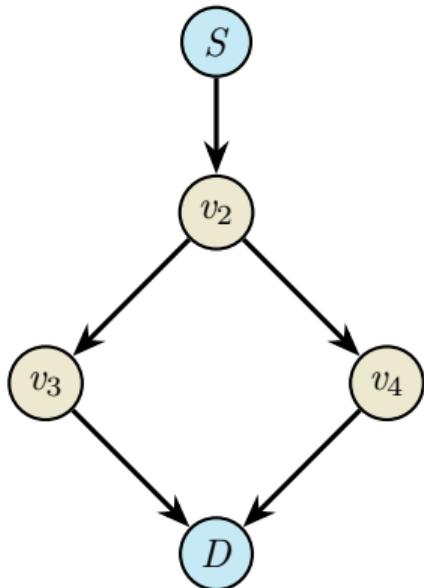


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For small flows $|v| \ll c$ and friction-dominance:

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} (\rho v) &= 0 \\ \text{(ISO3)} \quad \frac{\partial P}{\partial x} &= -\frac{\lambda}{2D} \rho v |v| - g \rho \sin(\alpha) \end{aligned}$$

Gas networks as metric graphs [FBP24]

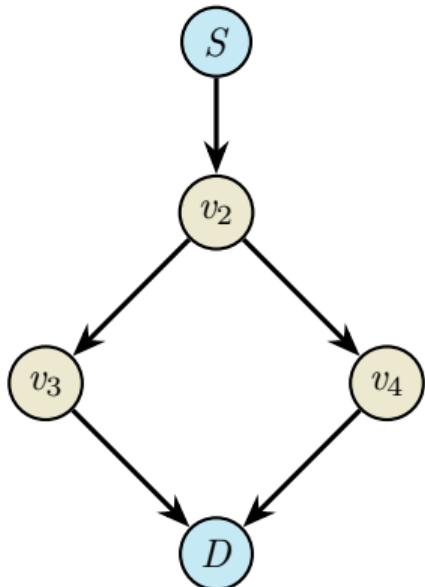


- 1) Gas flow in individual pipes
- 2) **Mass conservation at vertices**
- 3) Storage of gas at interior vertices
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Classical Kirchhoff's law at each vertex ν :

$$0 = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0}$$

Gas networks as metric graphs [FBP24]



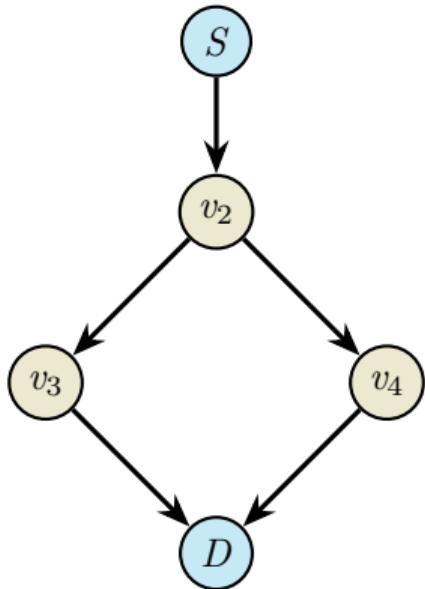
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Generalized Kirchhoff's law at each vertex ν :

$$f_\nu = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0}$$

with vertex flux f_ν and vertex density $\gamma_\nu \geq 0$ such that $\frac{\partial \gamma_\nu}{\partial t} = f_\nu$

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Time-dependent boundary condition at boundary vertex ν :

Given supply $s_\nu : [0, T] \longrightarrow \mathbb{R}_{\leq 0}$ for source vertex ν :

$$s_\nu = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0}$$

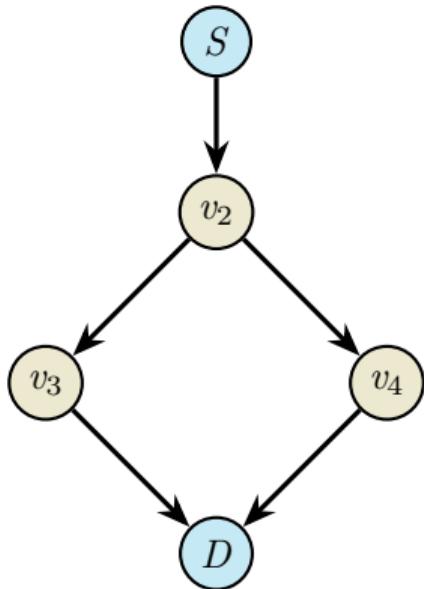
Given demand $d_\nu : [0, T] \longrightarrow \mathbb{R}_{\geq 0}$ for sink vertex ν :

$$d_\nu = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0}$$

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Time-independent boundary condition at boundary vertex ν :

Given total supply $S_\nu \leq 0$ for source vertex ν :

$$s_\nu = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0} \quad \text{such that} \quad \int_0^T s_\nu(t) dt = S_\nu$$

Given total demand $D_\nu \geq 0$ for sink vertex ν :

$$d_\nu = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0} \quad \text{such that} \quad \int_0^T d_\nu(t) dt = D_\nu$$

p -Wasserstein metric on gas networks [FBP24]



For a collection of edge and vertex measures $\mu_t = (\rho(t), \gamma(t), S(t), D(t))$:

$$W_p(\mu_0, \mu_T) = \inf_{\substack{\rho, j, \gamma, f, \\ S, s, D, d}} \left\{ T^{p-1} \int_0^T \left(\sum_e \int_0^{L_e} \frac{|j_e|^p}{\rho_e^{p-1}} dx + \sum_{\substack{\nu \text{ interior} \\ \text{vertex}}} \frac{|f_\nu|^p}{\gamma_\nu^{p-1}} + \sum_{\substack{\nu \text{ source} \\ \text{vertex}}} \frac{|s_\nu|^p}{S_\nu^{p-1}} + \sum_{\substack{\nu \text{ sink} \\ \text{vertex}}} \frac{|d_\nu|^p}{D_\nu^{p-1}} \right) dt \middle| \dots \right\}^{\frac{1}{p}}$$

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Classical Kirchhoff's law and no boundary vertices, i.e. $\mu_t = (\rho(t))$:

$$W_p(\mu_0, \mu_T) = \inf_{\rho, j} \left\{ T^{p-1} \int_0^T \sum_e \int_0^{L_e} \frac{|j_e|^p}{\rho_e^{p-1}} dx dt \middle| \begin{array}{l} \forall e : \frac{\partial \rho_e}{\partial t} + \frac{\partial j_e}{\partial x} = 0, \\ \rho_e|_{t=0} = (\rho_0)_e, \quad \rho_e|_{t=T} = (\rho_T)_e, \\ \forall \nu : 0 = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0} \end{array} \right\}^{\frac{1}{p}}$$

p -Wasserstein gradient flows [FBP24]



p -Wasserstein gradient flow derived in the limit $\tau \rightarrow 0$ for $\tau > 0$, from the minimizing movement scheme:

$$\rho_{(k+1)\tau}^\tau = \arg \min_\rho \left\{ \frac{1}{p\tau^{p-1}} W_p(\rho, \rho_{k\tau}^\tau)^p + \sum_e E(\rho_e) \right\}$$

with energy E , entropy F and integration constants c_e such that $\sum_e c_e = 0$:

$$E(\rho_e) = \int_0^{L_e} \left(F_e(\rho_e) + \frac{2\mathcal{D}_e}{\lambda_e} g \sin(\alpha_e) x \rho_e + c_e \rho_e \right) dx \quad F_e''(\rho_e) = \frac{2\mathcal{D}_e}{\lambda_e} \frac{P'_e(\rho_e)}{\rho_e}$$

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Result:

$$\begin{aligned} \forall e : \quad & \frac{\partial \rho_e}{\partial t} + \frac{\partial j_e}{\partial x} = 0, \quad j_e |j_e|^{p-2} = -\rho_e^{p-1} \frac{\partial}{\partial x} (E'(\rho_e)) \\ \forall \nu : \quad & 0 = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0} \end{aligned}$$

p -Wasserstein gradient flows [FBP24]



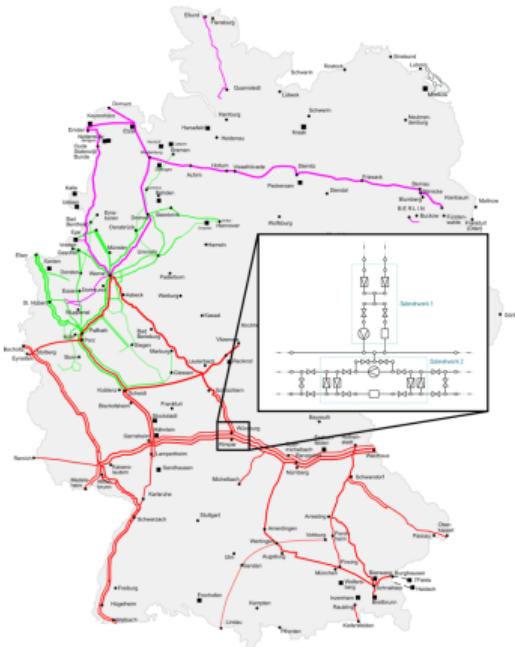
Corresponds to the (ISO3) formulation with classical Kirchhoff's law for $p = 3$:

$$\forall e : \frac{\partial \rho_e}{\partial t} + \frac{\partial j_e}{\partial x} = 0$$

$$\forall e : j_e |j_e| = -\rho_e^2 \left(\frac{2\mathcal{D}_e}{\lambda_e \rho_e} \frac{\partial P_e}{\partial x} - \frac{2D_e}{\lambda_e} g \sin(\alpha_e) \right) \iff \frac{\lambda_e}{2D_e} \rho_e v_e |v_e| = -\frac{\partial P_e}{\partial x} - g \rho_e \sin(\alpha_e) \iff \frac{\partial P_e}{\partial x} = -\frac{\lambda_e}{2\mathcal{D}_e} \rho_e v_e |v_e| - g \rho_e \sin(\alpha_e)$$

$$\forall \nu : 0 = \sum_{e=(\omega, \nu)} j_e|_{x=L_e} - \sum_{e=(\nu, \omega)} j_e|_{x=0}$$

Current work and next steps...



- Properties of the generalized p -Wasserstein metric
- Port-Hamiltonian formulation of (ISO2) in Wasserstein spaces
- Gas mixture model in (ISO3) formulation

Thank you for your attention!



References:

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