

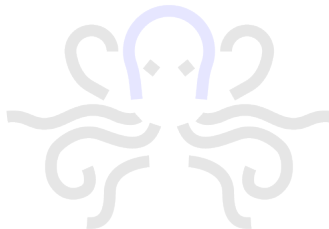
# Exponential convergence to steady-states for damped adhesive strings

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# Adhesive strings and reversible decohesion (“unzipping”)

## Aim

Modelling the (global) dynamics of an elastic string interacting with a rigid substrate through an adhesive layer and studying attachment–detachment regimes.

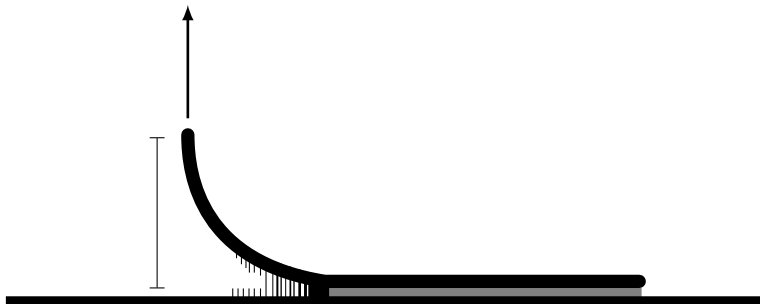


Figure: [Maddalena–Percivale–Puglisi–Truskinovsky, *Cont. Mech. Thermodyn.* 2009].

**Applications:** peeling of polymeric tapes, to rolling of cells, geckos’ fibrillar structures, denaturation of DNA...

# The model under consideration

*Dynamics of a one-dimensional linearly elastic body, whose reference configuration is  $(0, L)$ , interacting with a rigid substrate through an adhesive material, acting through the force  $\Phi'(u)$ :*

$$(W) \quad \begin{cases} \partial_{tt}^2 u - \partial_{xx}^2 u + \partial_t u + \Phi'(u) = 0, & (t, x) \in (0, +\infty) \times (0, L), \\ \partial_x u(t, 0) = \partial_x u(t, L) = 0, & t \in (0, +\infty), \\ u(0, x) = u_0(x), \quad \partial_t u(0, x) = v_0(x), & x \in (0, L). \end{cases}$$

- *wave operator* describing the dynamics of an adhesive string;
- a *damping term*,  $\partial_t u$ , accounts for the effect of friction;
- *forcing term*  $\Phi'$ : if the displacement  $u$  is small compared to  $u_*$ , the force is *purely elastic*; otherwise, if  $|u| \geq u_*$ , the adhesive material *ceases to act* on the elastic body;
- *Neumann boundary conditions*.

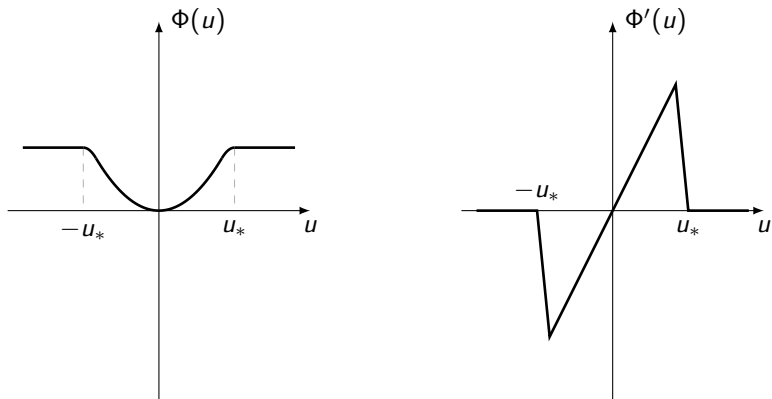


Figure: The potential  $\Phi$  and the force  $\Phi'$  as functions of  $u$ .

## Problem: long-time behavior

The attachment–detachment process ruled by the nonlinear force  $\Phi'(u)$  induces the natural question about the *long-time behavior* of such dynamics:

*Does the system converge towards a stationary state as  $t \rightarrow +\infty$ ?*

*Or does switching between the two states (attached/detached) persist?*

The problem contains two tasks:

- 1 study the global-in-time well-posedness of (W) *in a suitable function space* (more or less standard);
- 2 study the limit of  $u(t)$  as  $t \rightarrow +\infty$ .

# Study of the long-time behavior of the dynamics

We expect that  $u(t) \rightarrow u_\infty$  as  $t \rightarrow +\infty$ , where  $u_\infty$  is a solution to the *stationary problem*:

$$(S) \quad \begin{cases} -\partial_{xx}^2 u_\infty(x) + \Phi'(u_\infty(x)) = 0, & x \in (0, L), \\ \partial_x u_\infty(0) = \partial_x u_\infty(L) = 0. \end{cases}$$

Continuous set of possible choices for the limit profiles:

$$\{\Phi' = 0\} = (-\infty, -u_*] \cup \{0\} \cup [u_*, +\infty).$$

**Part 1:** *characterization of the limit points.* We show that  $\{u(t)\}_{t \geq 0}$  converges, as  $t \rightarrow +\infty$  to a *uniquely determined* limit profile  $u_\infty$  satisfying (S).

**Part 2:** *rate of convergence.* We show that the convergence  $u(t) \rightarrow u_\infty$  as  $t \rightarrow +\infty$  occurs in an *exponential* fashion: *i.e.*,

$$\|u(t) - u_\infty\|_{H^1((0,L))} + \|\partial_t u(t)\|_{L^2((0,L))} \leq Me^{-\kappa t}.$$

# Main results

## Theorem 1 (Well-posedness)

Given initial data

$$u_0 \in H^1(\Omega) \quad \text{and} \quad v_0 \in L^2(\Omega),$$

there exists a unique solution  $(u, \partial_t u) \in C([0, +\infty); H^1(\Omega) \times L^2(\Omega))$  to (W).

## Theorem 2 (Long-time asymptotics)

Given initial data

$$u_0 \in H^1(\Omega) \quad \text{and} \quad v_0 \in L^2(\Omega),$$

let  $u$  be the unique solution to (W). Then there exists  $u_\infty$  constant a.e. in  $\Omega$  with  $u_\infty \in \{\Phi' = 0\}$  such that  $u(t) \rightarrow u_\infty$  in  $H^1(\Omega)$  and  $\partial_t u(t) \rightarrow 0$  in  $L^2(\Omega)$  for  $t \rightarrow +\infty$ .

**NB:** Both results hold for a bounded, open, and connected set  $\Omega \subset \mathbb{R}^d$ , with  $d \geq 1$  (assuming also that the boundary is  $C^2$  when  $d \geq 2$ ).

### Theorem 3 (Exponential decay rate)

Let  $\Omega = (0, L) \subset \mathbb{R}$ . Given initial data

$$u_0 \in H^1(\Omega) \quad \text{and} \quad v_0 \in L^2(\Omega),$$

let  $u$  be the unique solution to (W).

Let  $u_\infty$  be as in Theorem 2. Let us assume that  $u_\infty = 0$  or  $|u_\infty| > u_*$ .

Then

$$\|u(t) - u_\infty\|_{H^1(\Omega)} + \|\partial_t u(t)\|_{L^2(\Omega)} \leq M_\Phi e^{-\kappa t}$$

for some  $M_\Phi > 0$  (possibly depending on the Lipschitz constant of  $\Phi'$ ) and  $\kappa > 0$ .



# Main tool: energy balances

The *energy functional*

$$E(u, v) := \frac{1}{2} \|v\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\nabla u\|_{L^2(\Omega)}^2 + \|\Phi(u)\|_{L^1(\Omega)}$$

satisfies

$$E(u(t), \partial_t u(t)) + \int_0^t \|\partial_t u(s)\|_{L^2(\Omega)}^2 ds = E(u_0, v_0), \quad \text{for every } t \in [0, +\infty).$$

In particular,  $E$  is non-negative and decreasing on trajectories.

This suggests the appropriate function space (*the energy space*) to look for (global) solutions of (W).

The *auxiliary functional*

$$J(u, v) := \frac{1}{2} \|u\|_{L^2(\Omega)}^2 + \langle u, v \rangle_{L^2(\Omega)}$$

(cf. [Haraux, *Bol. Soc. Bras. Mat.* 1986]) satisfies

$$\begin{aligned} J(u(t), \partial_t u(t)) + \int_0^t \left( \|\nabla u(s)\|_{L^2(\Omega)}^2 + \langle u(s), \Phi'(u(s)) \rangle_{L^2(\Omega)} \right) ds \\ = J(u_0, v_0) + \int_0^t \|\partial_t u(s)\|_{L^2(\Omega)}^2 ds, \quad \text{for every } t \in [0, +\infty). \end{aligned}$$

In particular,

$$J(u(t), \partial_t u(t)) + \int_t^{+\infty} \|\partial_t u(s)\|_{L^2(\Omega)}^2 ds$$

is non-negative and decreasing on trajectories.

# Outline of the proof of the qualitative convergence result

- 1 **Compactness argument** (via energy functional). Accumulation points of the trajectories  $\{u(t)\}_{t \geq 0}$  satisfy (S).

!! The compactness argument alone *does not suffice* to infer convergence of the whole trajectory as  $t \rightarrow +\infty$  to a uniquely determined limit profile.

*Why?* Because there is a *continuous set* of possible choices for the limit profiles: the set of solutions to (S) is given by *constant functions* valued in

$$\{\Phi' = 0\} = (-\infty, -u_*] \cup \{0\} \cup [u_*, +\infty).$$

- 2 **Selection of a unique limit profile.** The auxiliary functional  $J$  is a *Lyapunov functional* for the system and allows us to conclude the convergence  $u(t) \rightarrow u_\infty$  as  $t \rightarrow +\infty$ .

**NB:** the initial datum enforces the selection of a unique limit profile  $u_\infty \in \{\Phi' = 0\}$ .

# Outline of the proof of the exponential decay rate

## ① Compact embedding $H^1((0, L)) \subset\subset C([0, L])$ (true only in 1D).

This allows us to single out only one of the two possible attachment–detachment regimes (for time large enough) and study separately

$$\partial_{tt}^2 u + \partial_t u - \partial_{xx}^2 u = 0, \quad (t, x) \in (0, +\infty) \times (0, L),$$

or

$$\partial_{tt}^2 u + \partial_t u - \partial_{xx}^2 u + 2u = 0, \quad (t, x) \in (0, +\infty) \times (0, L).$$

## ② Grönwall-type inequalities for perturbed energy functionals.

[Haraux–Zuazua, *ARMA* 1988]:

$$G_\lambda(u, v) := \frac{1}{2} \|v\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\nabla u\|_{L^2(\Omega)}^2 + \int_{\Omega} (\Phi(u) - \Phi(u_\infty)) \, dx + \lambda \langle u - u_\infty, v \rangle_{L^2(\Omega)}$$

(perturbation, depending on  $\lambda \in (0, 1)$ , of  $E(u, v) - E(u_\infty, 0)$ ).

# Further questions

- **Abrupt attachment–detachment.** If the rate of the exponential decay  $Me^{-\kappa t}$  is *uniform* w.r.t. the slope of the decreasing part of the force  $\Phi'(u)$ , then we may consider a *discontinuous force*:

$$\Phi'(u) := \begin{cases} 2u & \text{for } |u| < u_* , \\ 0 & \text{for } |u| \geq u_* . \end{cases}$$

- We already have positive results for an ODE model!
- **System of thermoelasticity.** The temperature gradient acts as a force on the wave equation governing the elastic component and the pressure waves act as a heat-source on the diffusion equation governing the temperature:

$$(T) \quad \begin{cases} \partial_{tt}^2 u - \partial_{xx}^2 u - \partial_x \theta + \Phi'(u) = 0, & (t, x) \in (0, +\infty) \times (0, L), \\ \partial_t \theta - \partial_{xx}^2 \theta - \partial_{tx}^2 u = 0, & (t, x) \in (0, +\infty) \times (0, L). \end{cases}$$

Thank you for your attention!

## References



Giuseppe Maria Coclite, Nicola De Nitti, Francesco Maddalena, Gianluca Orlando, Enrique Zuazua. Exponential convergence to steady-states for trajectories of a damped dynamical system modelling adhesive strings. *Mathematical Models and Methods in Applied Sciences* 34, No. 08, 1445–1482 (2024). .

