# How much does it cost to control near a shock?

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Shock profiles for Burgers' equation.

2 The control problem.





2 The control problem.

3 Main result and sketch of proof

Consider the 1D inviscid Burgers' equation

$$\begin{cases} \partial_t u + u \partial_x u = 0, & t \in \mathbb{R}^+, x \in \mathbb{R}, \\ u(0, x) = u_0(x), & x \in \mathbb{R}. \end{cases}$$
(1)

By method of characteristics, jump discontinuities arise in finite time unless  $u_0$  is non-decreasing.

Rankine-Hugoniot condition : A jump from  $u^-$  to  $u^+ \; (u^- > u^+)$  is displaced at velocity

shock speed 
$$= \frac{u^- + u^+}{2}.$$

#### Family of stationary shocks

We restrict the study of (1) to a finite interval [-L, L] and add Dirichlet conditions :

$$\begin{cases} \partial_t u + u \partial_x u = 0, & t \in \mathbb{R}^+, x \in (-L, L), \\ u(t, -L) = 1, & t \in \mathbb{R}^+, \\ u(t, L) = -1, & t \in \mathbb{R}^+, \end{cases}$$
(2)

There is an infinite family of stationary shocks given by

$$U_{x_0}(x) = \begin{cases} 1 & \text{if } x < x_0 \\ -1 & \text{if } x > x_0 \end{cases}, \quad x_0 \in (-L, L).$$

Any solution to (2) with initial datum of bounded variation converges to some  $U_{x_0}$ .

We add a diffusive term  $\varepsilon\partial_{xx}$  to the Burgers' equation for a regularizing effect.

$$\begin{cases} \partial_t u + u \partial_x u = \varepsilon \partial_{xx} u, & t \in \mathbb{R}^+, x \in (-L, L), \\ u(t, -L) = 1, & t \in \mathbb{R}^+ \\ u(t, L) = -1, & t \in \mathbb{R}^+. \end{cases}$$
(3)

A unique stationary solution exists, given by

$$U^{\varepsilon}(x) = -c \tanh\left(\frac{cx}{2\varepsilon}\right),$$

where c > 0 is in order to verify the boundary conditions.

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# Uniform controllability problem.

We consider the problem of boundary null-controllability of Burgers equation linearized around the stationary shock.

$$\begin{cases} \partial_t u + \partial_x (U^{\varepsilon} u) = \varepsilon \partial_{xx} u, & t \in [0, T], x \in (-L, L), \\ u(t, -L) = h(t), & t \in [0, T], \\ u(t, L) = 0, & t \in [0, T], \\ u(0, x) = u_0(x), & x \in [-L, L], \end{cases}$$

with  $U^{\varepsilon}(x) = -\tanh\left(\frac{x}{2\varepsilon}\right)$ .

#### Uniform controllability

We want to find  $T_{\text{unif}}$  such that, for any  $T > T_{\text{unif}}$ ,

$$\begin{split} \exists \mathcal{C}, \forall \varepsilon > 0, \forall u_0 \in H^2 \cap H^1_0, \exists h \in L^2(0,T) : \\ u(T) \equiv 0 \text{ and } \|h\|_{L^2(0,T)} \leq \mathcal{C} \|u_0\|_{L^2(-L,L)} \end{split}$$

(4)

#### Dual observability problem.

By standard duality argument, this is equivalent to the uniform observability of the adjoint system

$$\begin{cases} \partial_t v - U^{\varepsilon} \partial_x v = \varepsilon \partial_{xx} v, & t \in [0,T], x \in (-L,L) \\ v(t,-L) = v(t,L) = 0, & t \in [0,T] \\ v(0,x) = v_0(x), & x \in [-L,L] \end{cases}$$
(5)

Namely we want to investigate the minimal time  $T_{\rm unif}$  such that, for any  $T>T_{\rm unif}$  ,

$$\exists \mathcal{C} = \mathcal{C}(T) > 0, \forall \varepsilon > 0, \forall v_0 \in H^2 \cap H^1_0,$$
$$\mathcal{C}^2 \int_0^T |\varepsilon \partial_x v(t, -L)|^2 \, \mathrm{d}t \ge \|v(T, \cdot)\|_{L^2(-L, L)}^2$$

# Related problems

#### Coron-Guerrero conjecture ('05) :

The uniform controllability of the transport equation

 $\partial_t y + M \partial_x y = \varepsilon \partial_{xx} y$ 

in the vanishing viscosity limit is possible whenever  $T>\frac{L}{M}$  if M>0, or  $T>\frac{2L}{|M|}$  if M<0.

# Related problems

#### Coron-Guerrero conjecture ('05) :

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Results in this sense :

	M > 0	M < 0
Coron,Guerrero ('05)	$T > \frac{4.3L}{M}$	$T > \frac{57.2L}{ M }$
Glass ('09)	$T > \frac{4.2L}{M}$	$T > \frac{6.1L}{ M }$
Lissy ('12)	$T > \frac{2.35L}{M}$	$T > \frac{4.35L}{ M }$

Laurent, Léautaud ('23) : Similar results generalized to a wider class of problems of the form

$$\partial_t y + a \partial_x y + b y = \varepsilon \partial_{xx} y,$$

under some regularity assumptions on a, b.

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Spectral analysis of  $\mathcal{L}^{\varepsilon} = -U^{\varepsilon}\partial_x - \varepsilon\partial_{xx}$ 

#### Proposition (Mascia-Strani, 2012)

The spectrum of the operator  $\mathcal{L}^{\varepsilon}$  consists of simple eigenvalues  $\{\lambda_0 < \lambda_1 < \cdots\}$  such that :

$$\exists C > 0: 0 < \lambda_0 < Ce^{-\frac{C}{\varepsilon}}, \quad \lambda_1 \ge \frac{C}{\varepsilon}$$

Two main issues compared with the constant transport case :

- A small eigenvalue, caracteristic of a metastability phenomenon.
- Lack of precise knowledge regarding the higher eigenvalues.

Result based on the conjugation of the operator :

$$\operatorname{sech}\left(\frac{x}{2\varepsilon}\right)\mathcal{L}^{\varepsilon}\cosh\left(\frac{x}{2\varepsilon}\right) = \frac{1}{\varepsilon}\left(-\varepsilon^{2}\partial_{xx} - \frac{1}{4} + \frac{1}{2}\tanh^{2}\left(\frac{x}{2\varepsilon}\right)\right) =: \frac{1}{\varepsilon}P^{\varepsilon}$$

### Spectral analysis of $P^{\varepsilon}$

Schrödinger operator with a single-well symmetric potential  $\implies$  good estimates for the bottom of the spectrum, loss of precision far from the well.

We are saved by the factorisation

$$P^{\varepsilon} = a^* a,$$

where  $a = \varepsilon \partial_x + \frac{1}{2} \tanh\left(\frac{x}{2\varepsilon}\right)$ . We may compute

$$aa^* = -\varepsilon^2 \partial_{xx} + \frac{1}{4},$$

which is much simpler, and the eigenvalues of  $aa^*$  and  $a^*a$  are similar :

$$\begin{split} &aa^*\varphi = \lambda\varphi \iff a^*a(a^*\varphi) = \lambda(a^*\varphi),\\ &a^*a\psi = \mu\psi \iff aa^*(a\psi) = \mu(a\psi) \end{split}$$

### Spectral analysis of $P^{\varepsilon}$ .

Solving the eigenvalue problem

$$\begin{cases} P^{\varepsilon}\varphi = \lambda\varphi,\\ \varphi(\pm L) = 0, \end{cases}$$

is equivalent to solving

$$\begin{cases} (-\varepsilon^2 \partial_{xx} + \frac{1}{4})\psi = \lambda \psi \\ a^* \psi(\pm L) = 0. \end{cases}$$

#### Proposition

The spectrum of  $\mathcal{L}^{\varepsilon}$  consists of simple eigenvalues  $\{\lambda_0 < \lambda_1 < \cdots\}$  that verify

$$\begin{aligned} 0 &< \lambda_0 < Ce^{-\frac{C}{\varepsilon}}, \\ \frac{1}{4\varepsilon} + \varepsilon \frac{k^2 \pi^2}{4L^2} < \lambda_k < \frac{1}{4\varepsilon} + \varepsilon \frac{(k+1)^2 \pi^2}{4L^2}, \quad k \ge 1 \\ \sqrt{\lambda_{k+1} - \frac{1}{4\varepsilon}} - \sqrt{\lambda_k - \frac{1}{4\varepsilon}} > \sqrt{\varepsilon} \frac{\pi}{2L}, \quad k \ge 1 \end{aligned}$$

#### The main result

#### Theorem (L. 2024) [WIP].

There exists a time

$$T_{\text{unif}} = 8.4L$$

such that, whenever  $T > T_{unif}$ , the system

$$\begin{cases} \partial_t u + (\mathcal{L}^{\varepsilon})^* u = 0\\ u(t, -L) = h(t), \quad u(t, L) = 0,\\ u(0, x) = u_0(x) \end{cases}$$

is uniformly controllable to the first mode in the vanishing viscosity limit. Namely, for  $T > T_{unif}$ , there exists a control h uniformly bounded with respect to  $\varepsilon$  such that

$$\left\langle \operatorname{sech}\left(\frac{x}{2\varepsilon}\right)u(T),\varphi_{k}\right
angle =0,\quad k\geq1.$$

Moreover the controllability cost verifies

$$\mathcal{C}(T,\varepsilon) \le ce^{-\frac{c}{\varepsilon}}.$$

#### Construction of a biorthogonal family

The first step is to construct a bi-orthogonal family to the  $(e^{-\lambda_k t})_{k\geq 1}$  with suitable estimates, namely a family  $(q_j)_{j\geq 1} \subset L^2(0,T)$  such that

$$\int_0^T q_j(t) e^{-\lambda_k t} \, \mathrm{d}t = \delta_{j,k}$$

Estimates on  $\lambda_k$  allow us to build such a family with the bound

$$||q_k|| \le C \exp\left(f(L,\varepsilon,\lambda_k,T)\right)$$

where f is such that

$$-\lambda_k T + f(L,\varepsilon,\lambda_k,T) \le -\alpha\lambda_k, \quad \alpha > 0$$
(6)

for T > 8.4L.

#### Conclusion by method of moments.

Then, take any initial datum  $v_0$  of the form

$$v_0(x) = \sum_{k \ge 1} c_k \psi_k,$$

where  $\psi_k := \cosh\left(\frac{x}{2\varepsilon}\right)\varphi_k$  are the eigenvectors of  $\mathcal{L}^{\varepsilon}$ . The solution at time t is immediately given by

$$v(t,x) = \sum_{k \ge 1} c_k \psi_k e^{-\lambda_k t}$$

By taking the space derivative in -L and integrating in time against  $q_j$ , we obtain

$$\int_0^T \varepsilon \partial_x v(t, -L) q_j(t) \, \mathrm{d}t = c_j \varepsilon \psi'_j(-L).$$

### Conclusion by method of moments.

It follows, for  $j \ge 1$ ,

$$|c_j| \le \frac{\|\varepsilon \partial_x v(\cdot, -L)\| \|q_j\|}{\varepsilon \psi'_k(-L)}$$

Finally, we plug the latter estimate into  $\boldsymbol{v}(T)$  :

$$\|v(T,\cdot)\| = \left\| \sum_{k\geq 1} c_k \psi_k e^{\lambda_k T} \right\|$$
  
$$\leq \|\varepsilon \partial_x v(\cdot, -L)\| \sum_{k\geq 1} e^{-\lambda_k T} \frac{\|q_k\| \cdot \|\psi_k\|}{\varepsilon \psi'_k(-L)}$$

Using (6) finally gives

$$\|v(T,\cdot)\| \le \|\varepsilon \partial_x v(\cdot,-L)\| \sum_{k\ge 1} e^{-\alpha\lambda_k} \cdot \frac{\|\psi_k\|}{\varepsilon \psi'_k(-L)},$$

and we conclude by having a precise description of the eigenfunctions  $\psi_k$ .

# Work in progress and possible leads.

- Refine the constant 8.4 by improving the estimates on the bi-orthogonal family.
- Possibly obtain a time  $T_{\min}$  for which we can prove the cost of controllability to the first mode explodes as  $\varepsilon \to 0$  whenever  $T < T_{\min}$ .
- Prove non-uniform controllability for the full system for any finite time T.

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#### Thank you for your attention !