

B decays and mixing

and consequences for the Standard Model

Preliminary remarks

Review + emphasis on B mixing

- Advertisement of other talks
 - 11:00 Monday: Alejandro Vaquero, $B \rightarrow D^* l \nu$ and related anomalies
 - 16:30 Monday: Enrico Lunghi, B and K physics, role of lattice QCD
 - 09:30 Tuesday: Chris Bouchard, $B \rightarrow K l l$ and $D \rightarrow K l \nu$ form factors and phenomenology
 - 11:00 Tuesday: Simon Kuberski, Heavy flavor physics with $O(a)$ improved Wilson quarks
- Survey of CKM status
- More in-depth summary of B meson mixing calculations

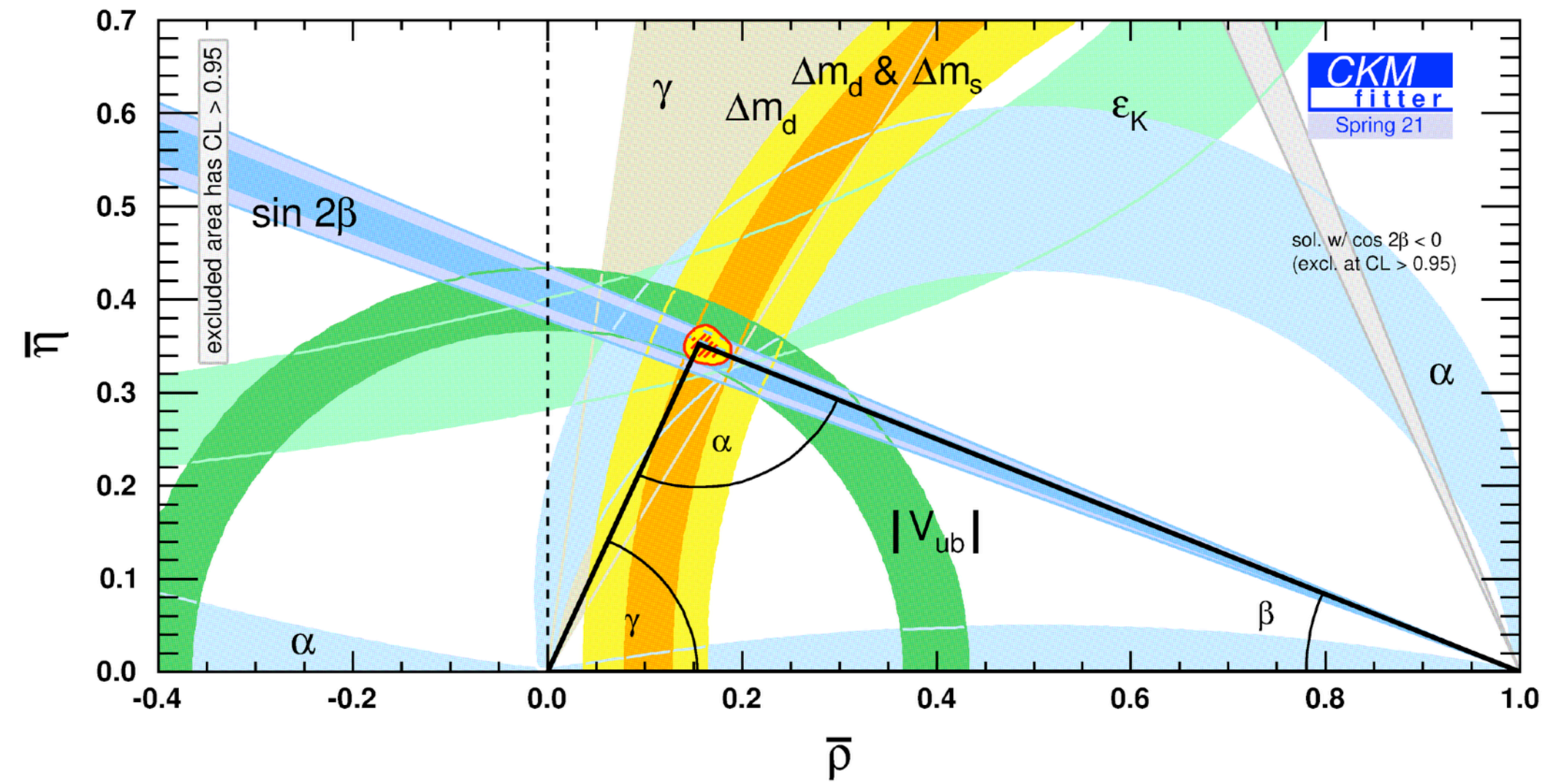
Introduction

Quark flavor in the Standard Model

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

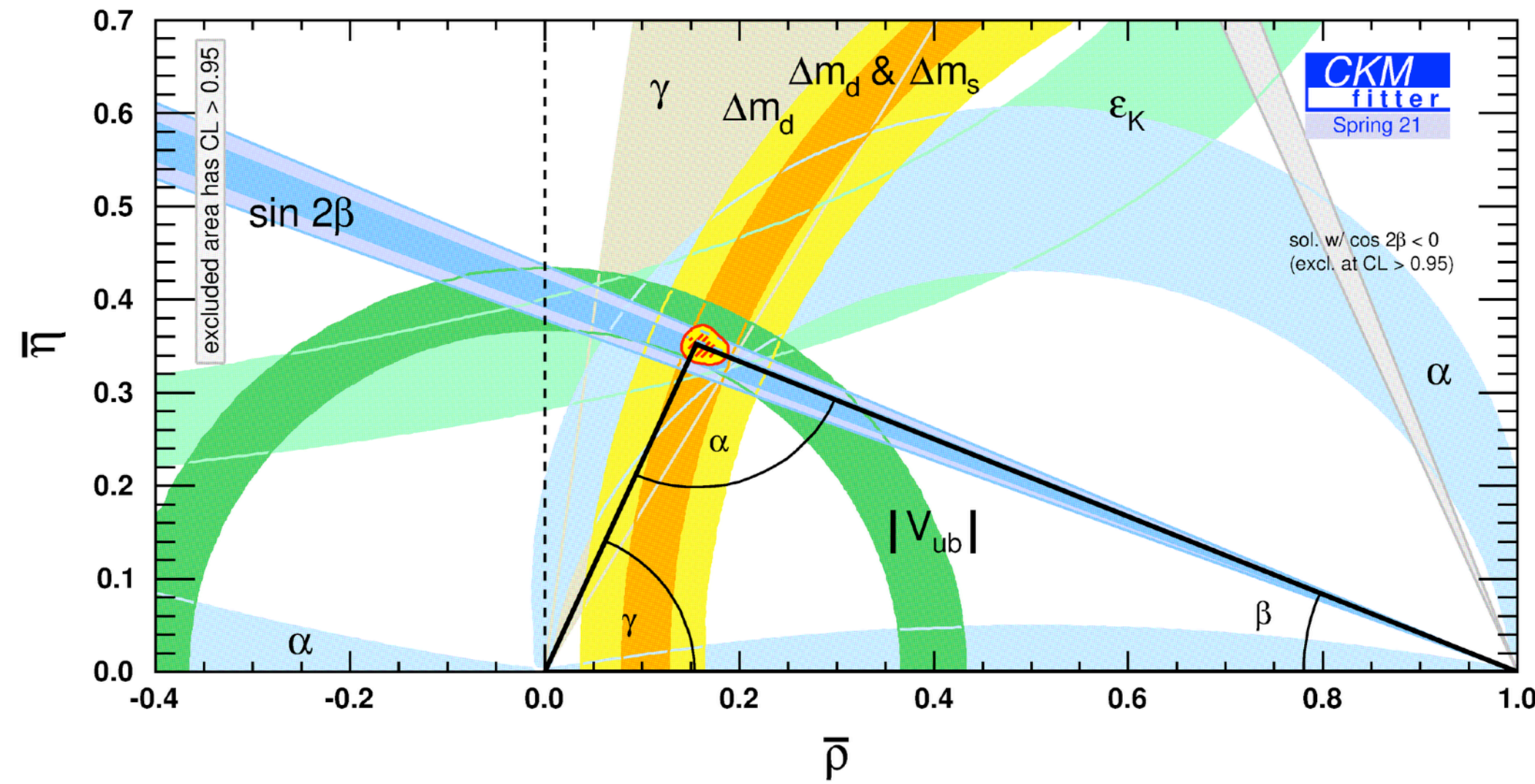
- Are the 4 free parameters in the CKM matrix sufficient to describe all **quark** flavor-changing interactions?
- High intensity (& high energy) era:
LOTS of experimental measurements of **hadronic** processes,
many very precise
- Require similarly precise SM predictions
- Lattice QCD connects hadronic observables to quark parameters

CKM constraints – (ρ, η) plane

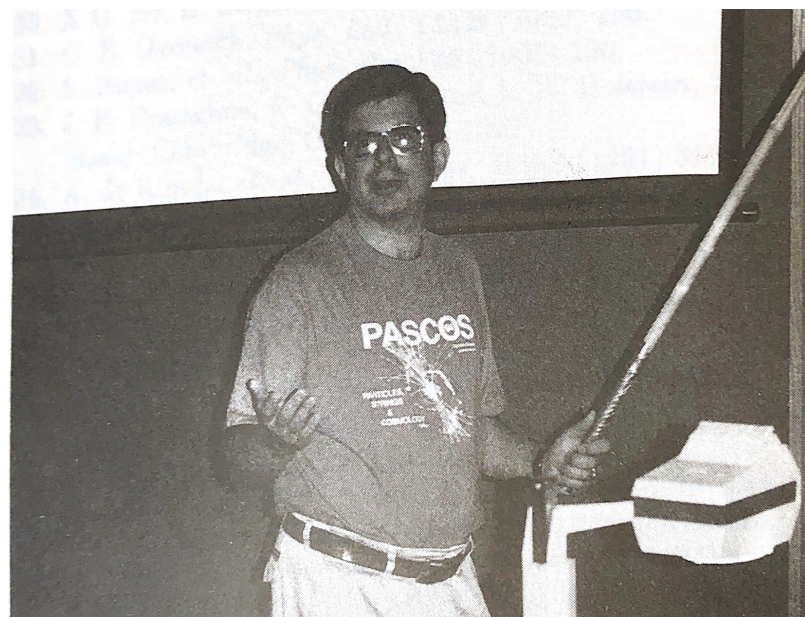


CKMfitter, Moriond 2021

CKM constraints – (ρ, η) plane



CKMfitter, Moriond 2021



Sheldon Stone, TASI 1994

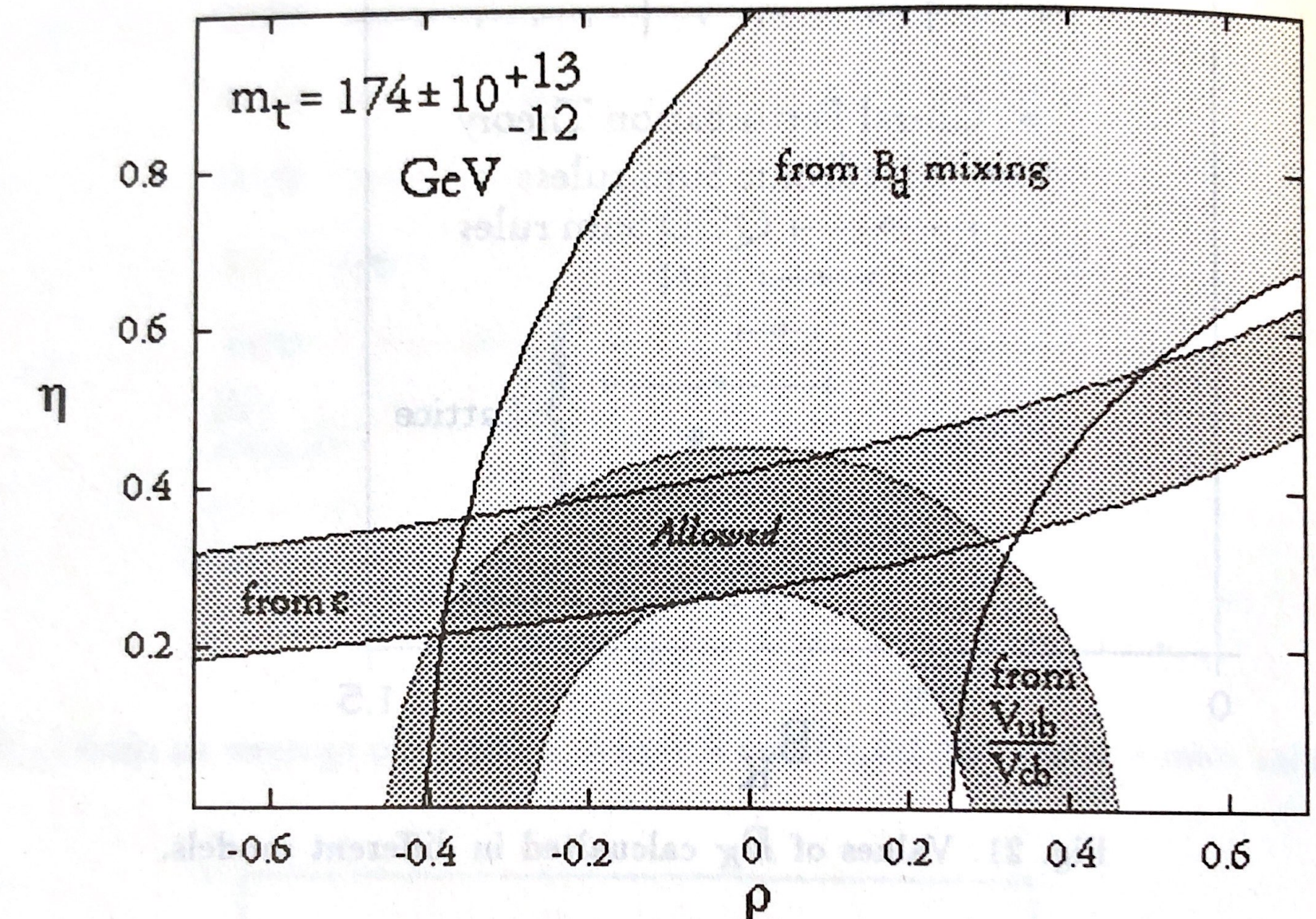
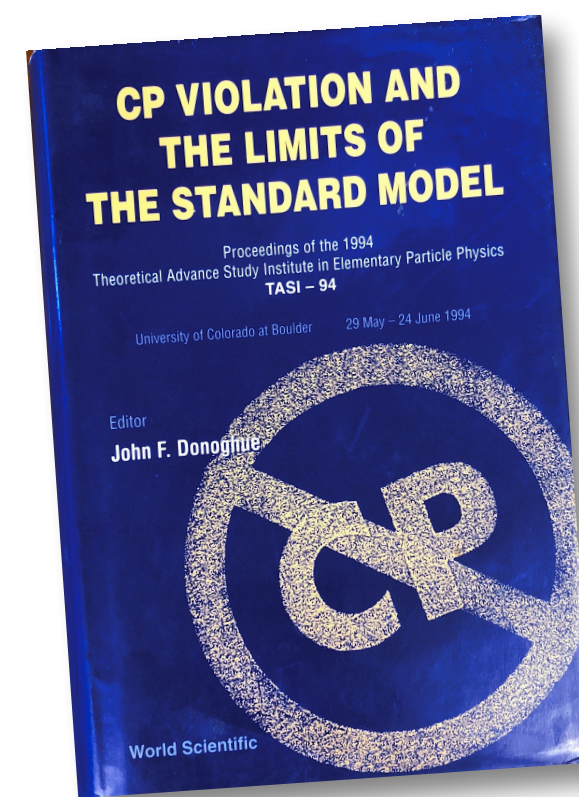
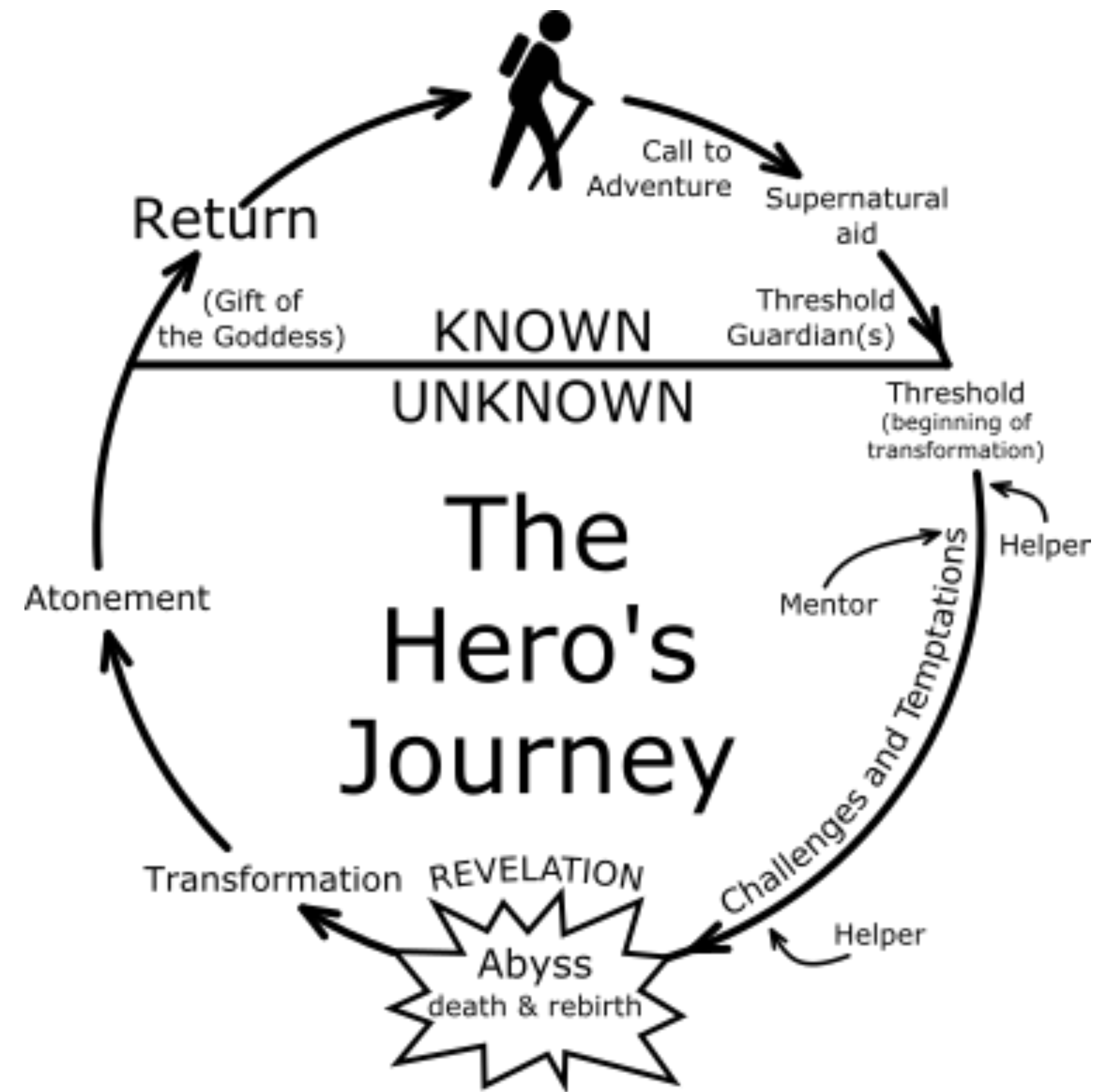
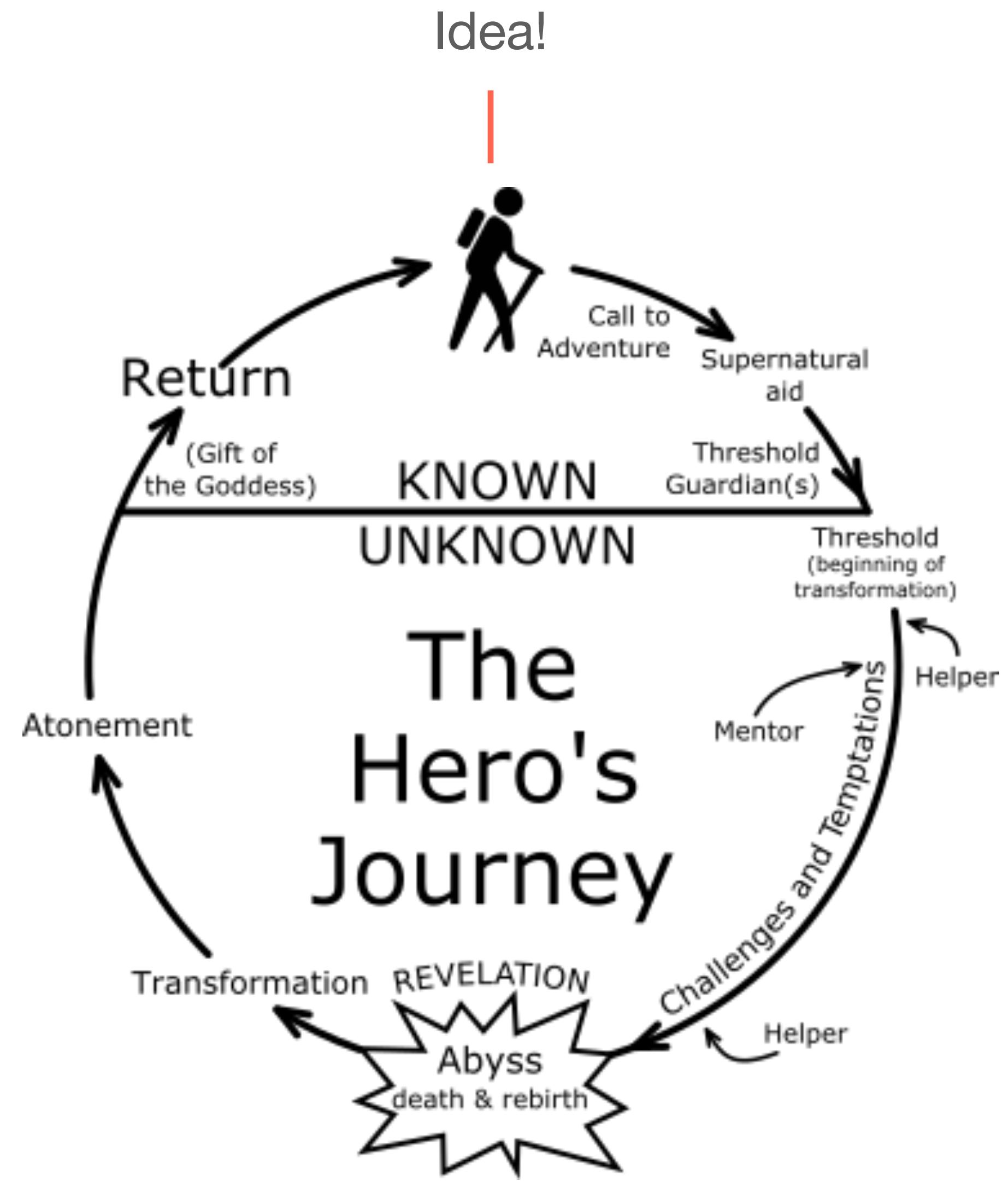
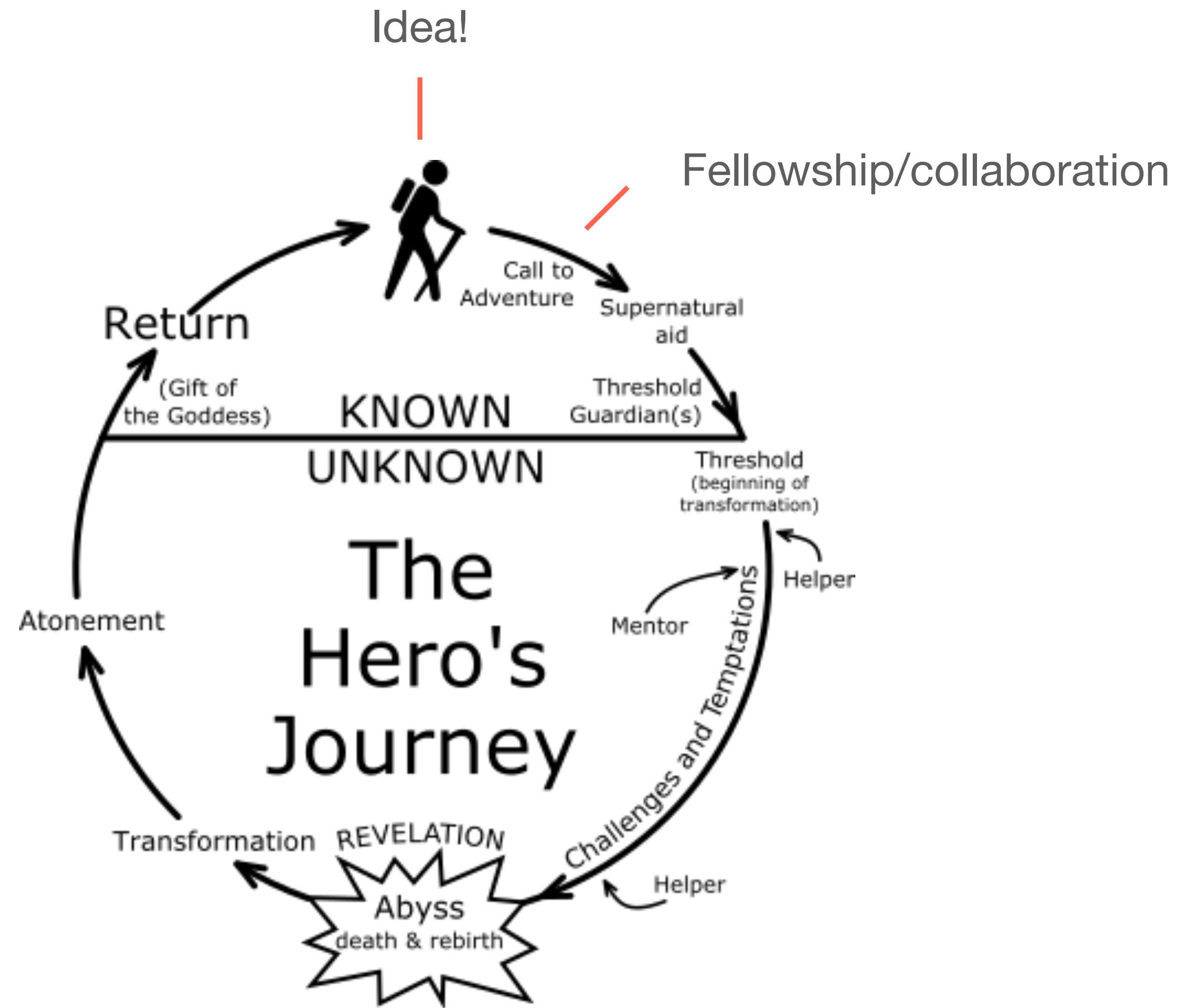
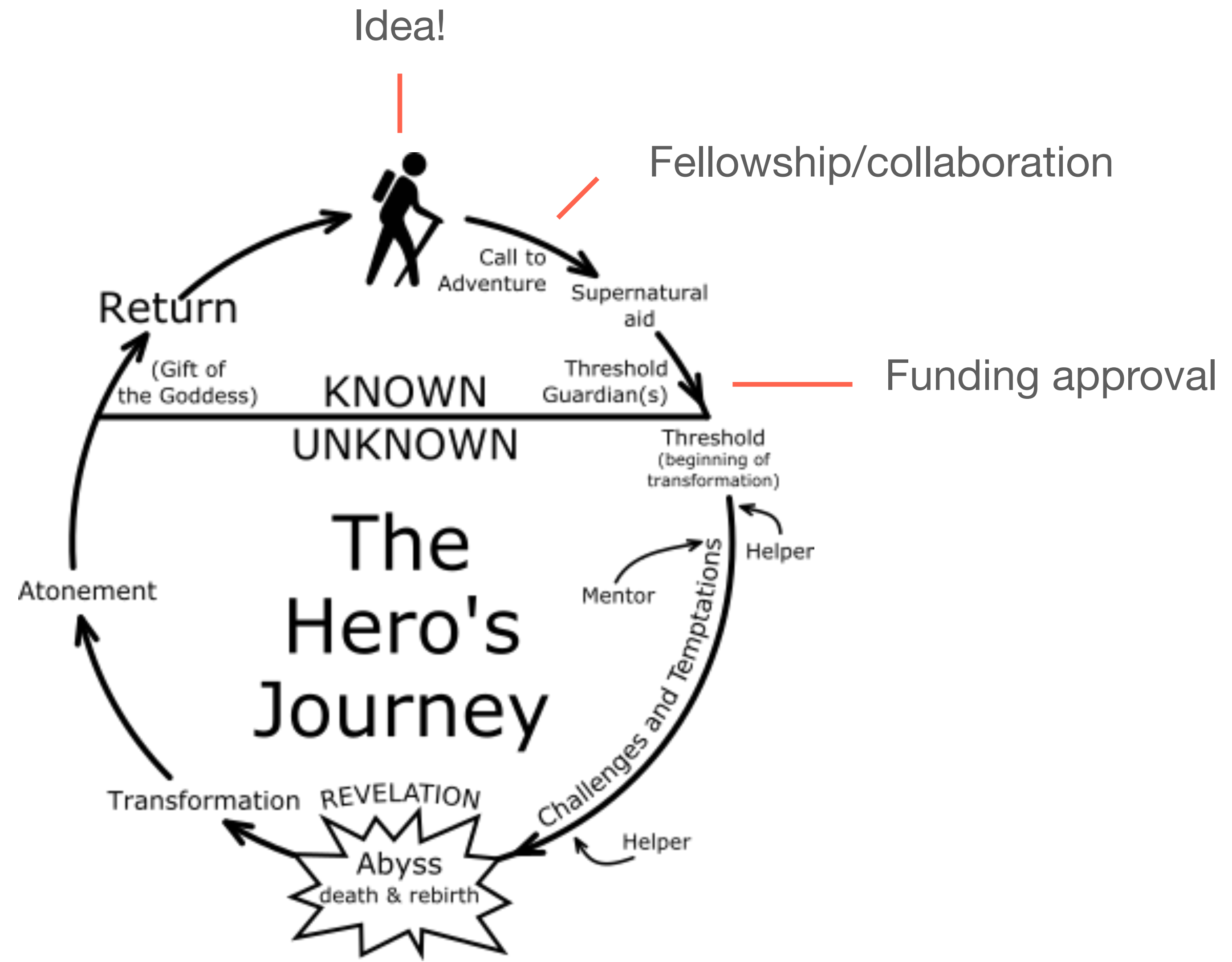


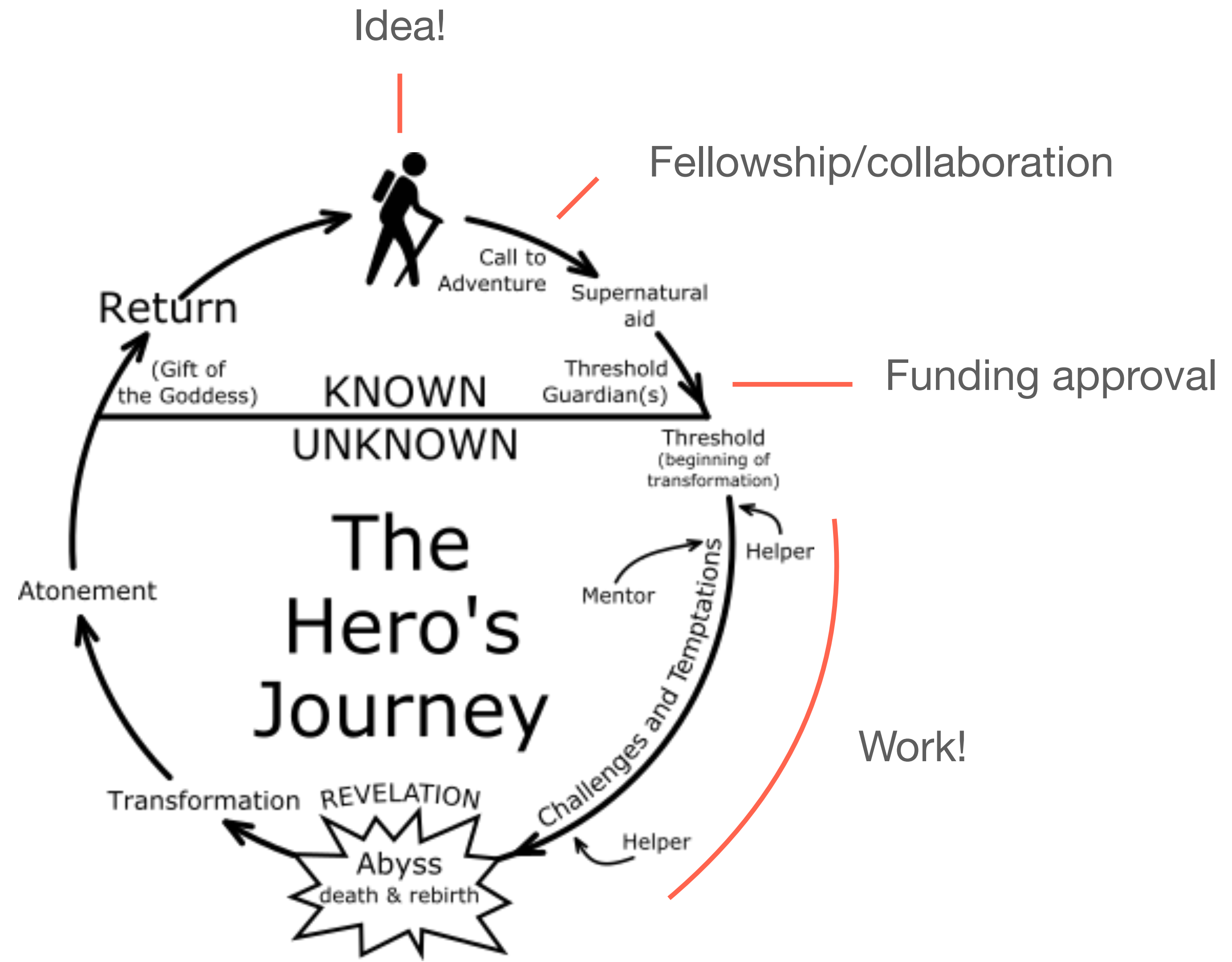
Fig. 22. Constraints in the $\rho - \eta$ plane from $|V_{ub}/V_{cb}|$, $B_d^0 - \bar{B}_d^0$ mixing and CP violation in the K_L^0 system (ϵ). The bands are $\pm 1\sigma$.

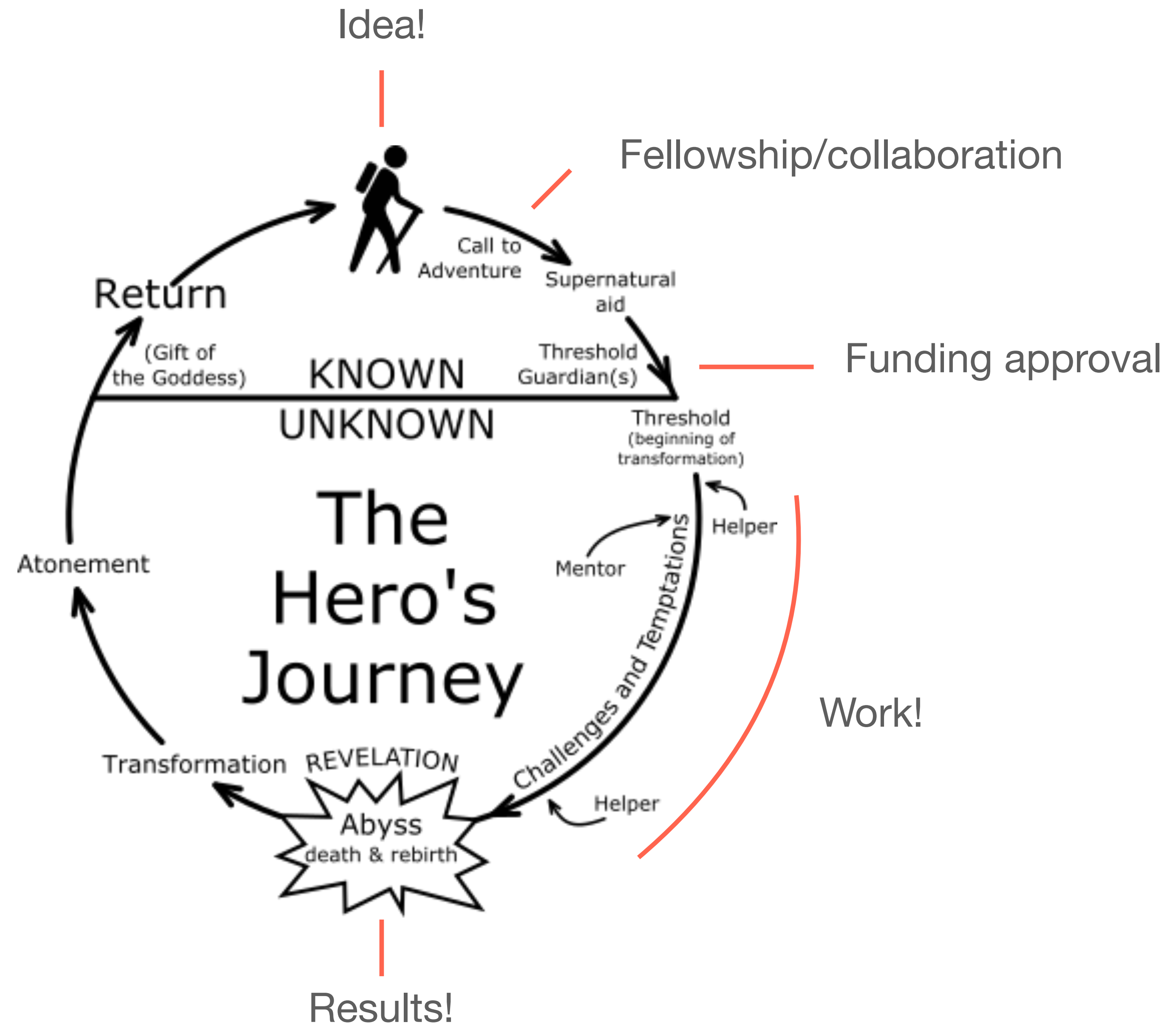


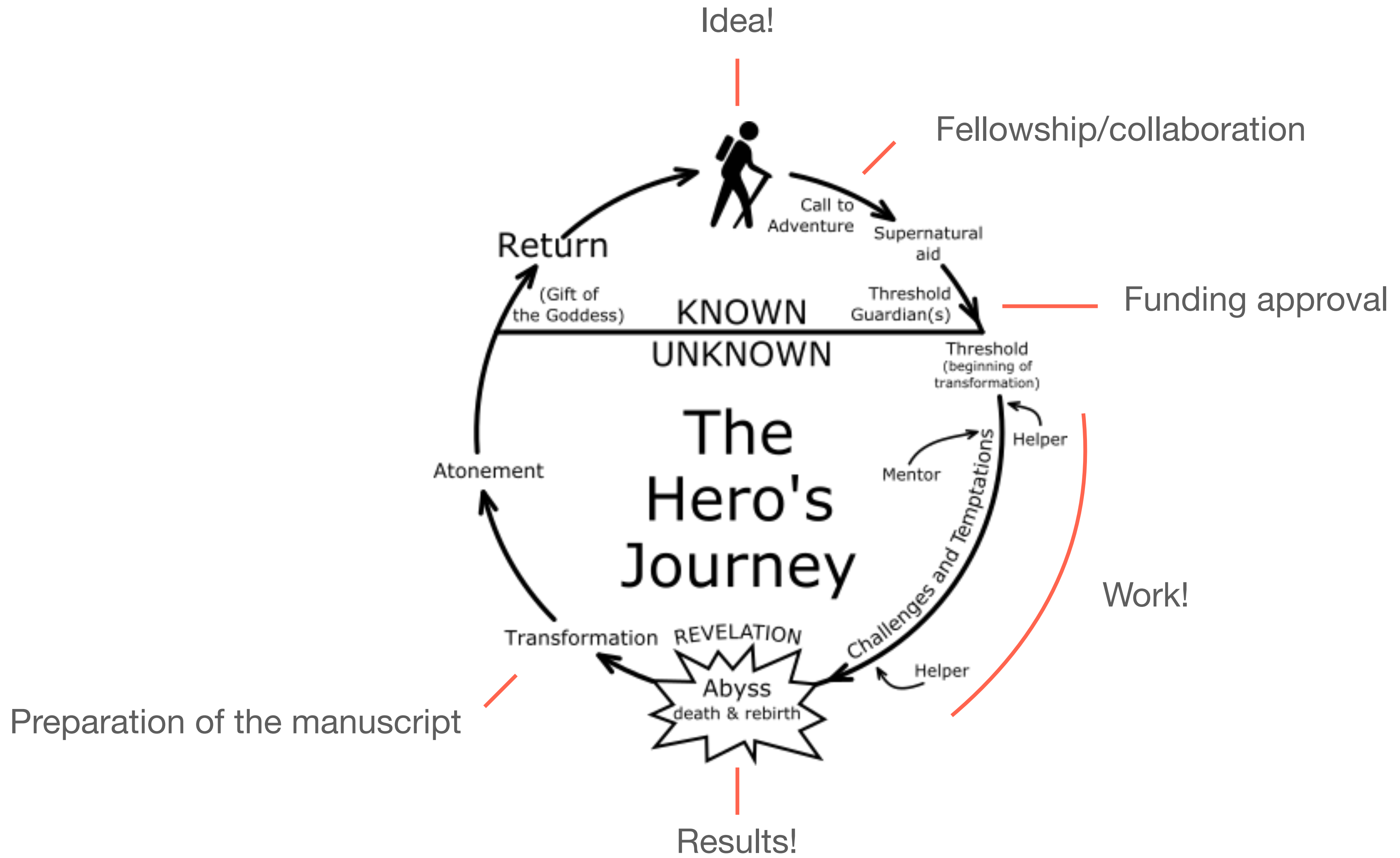


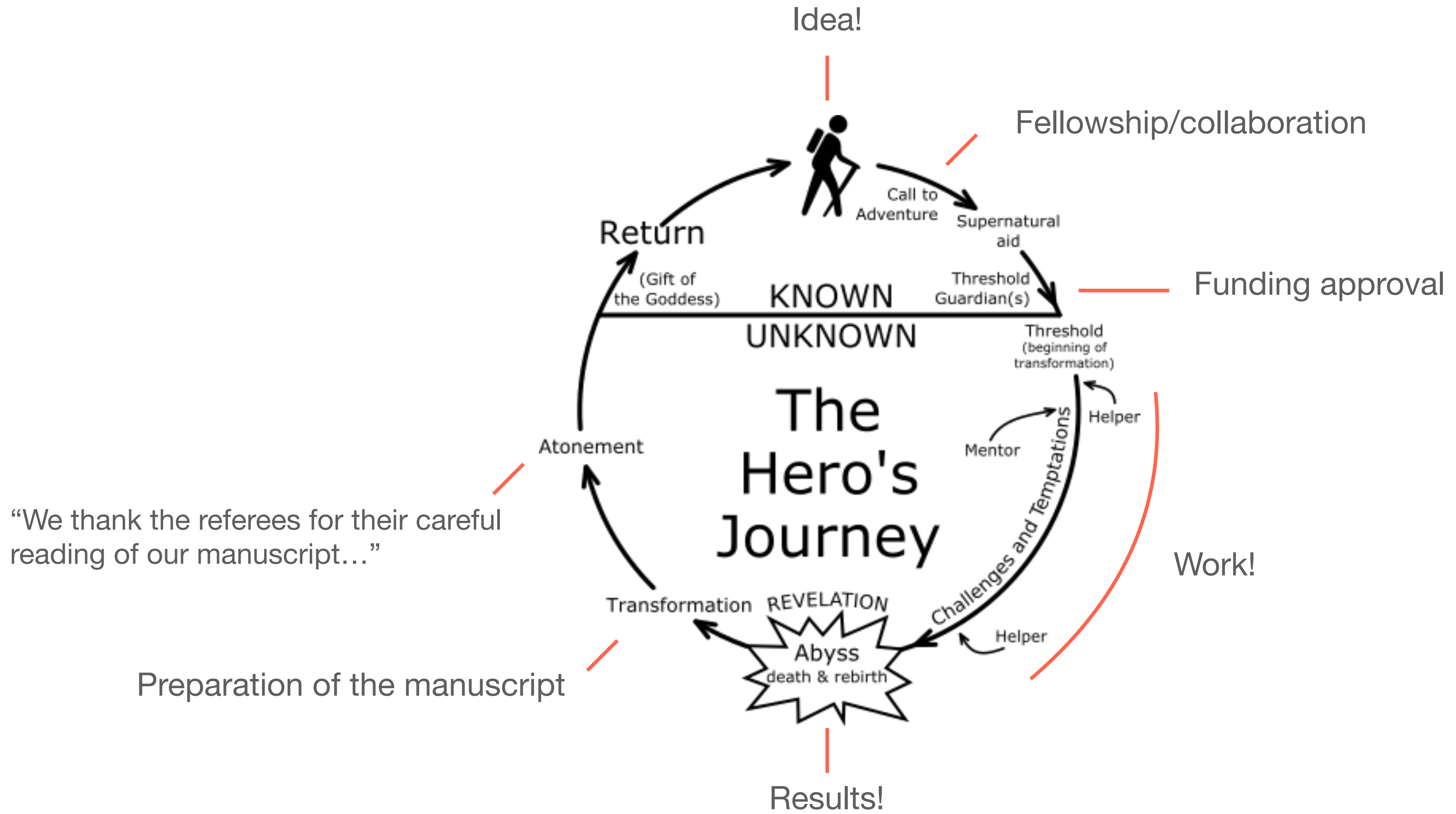


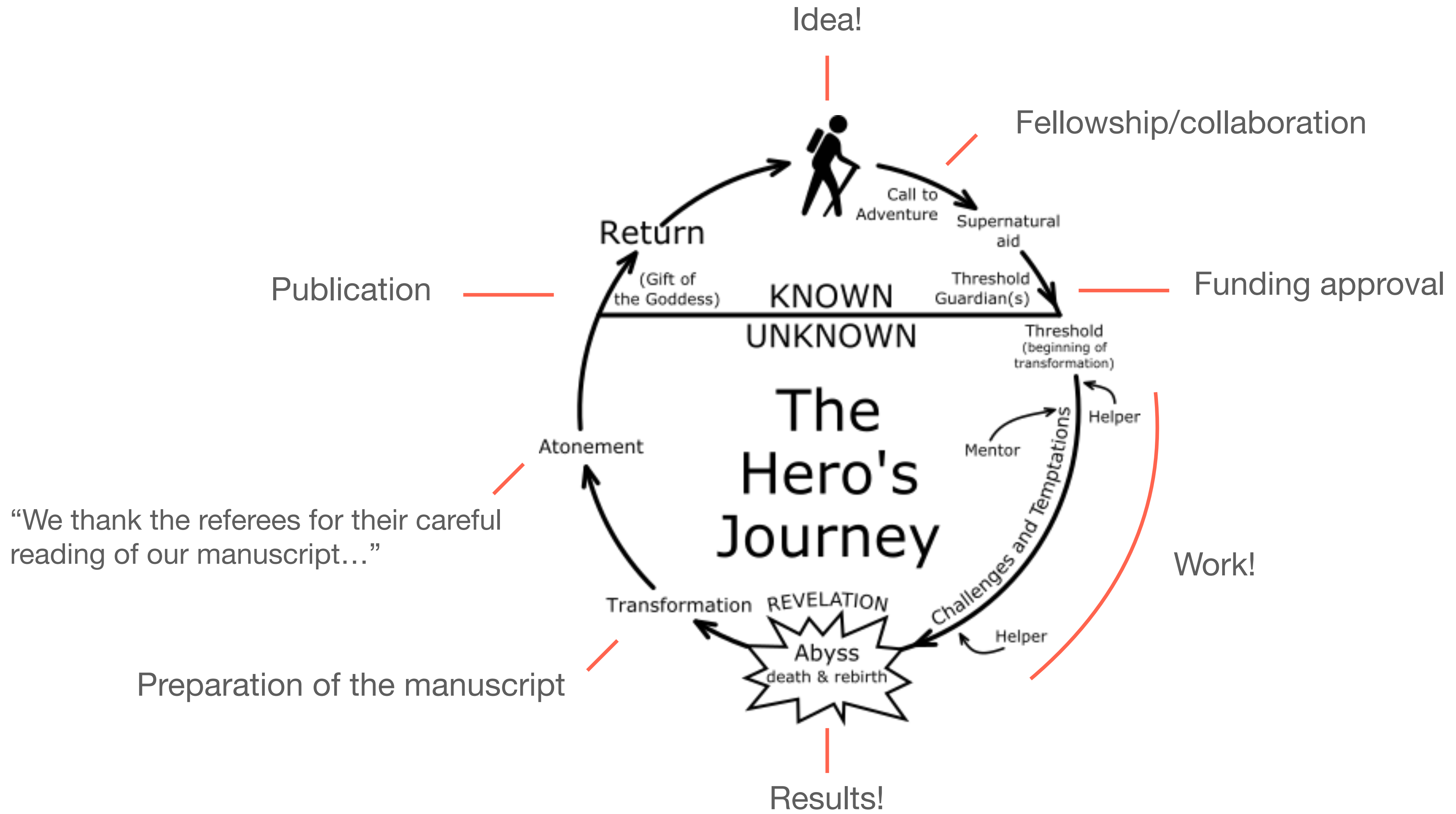


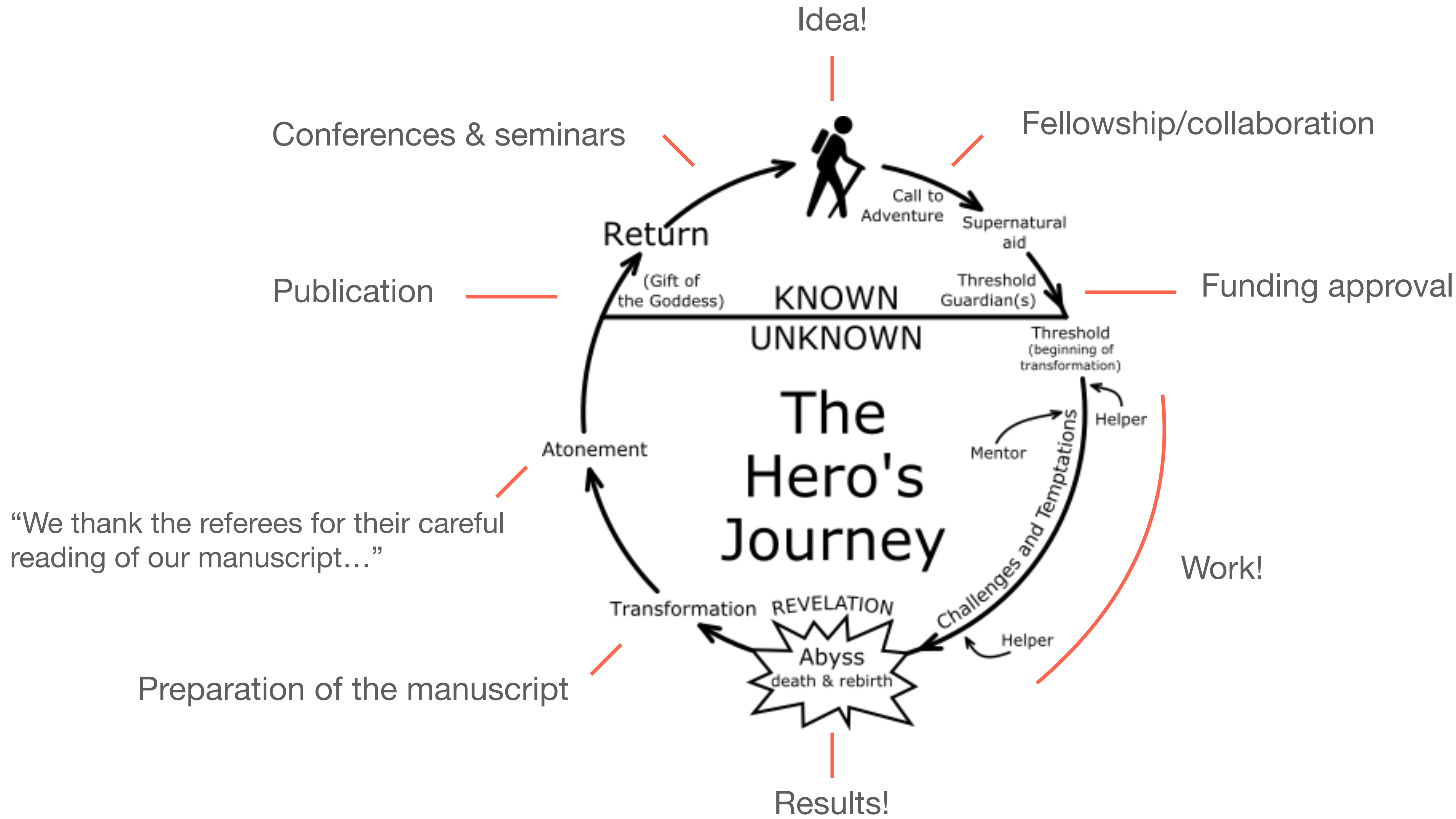












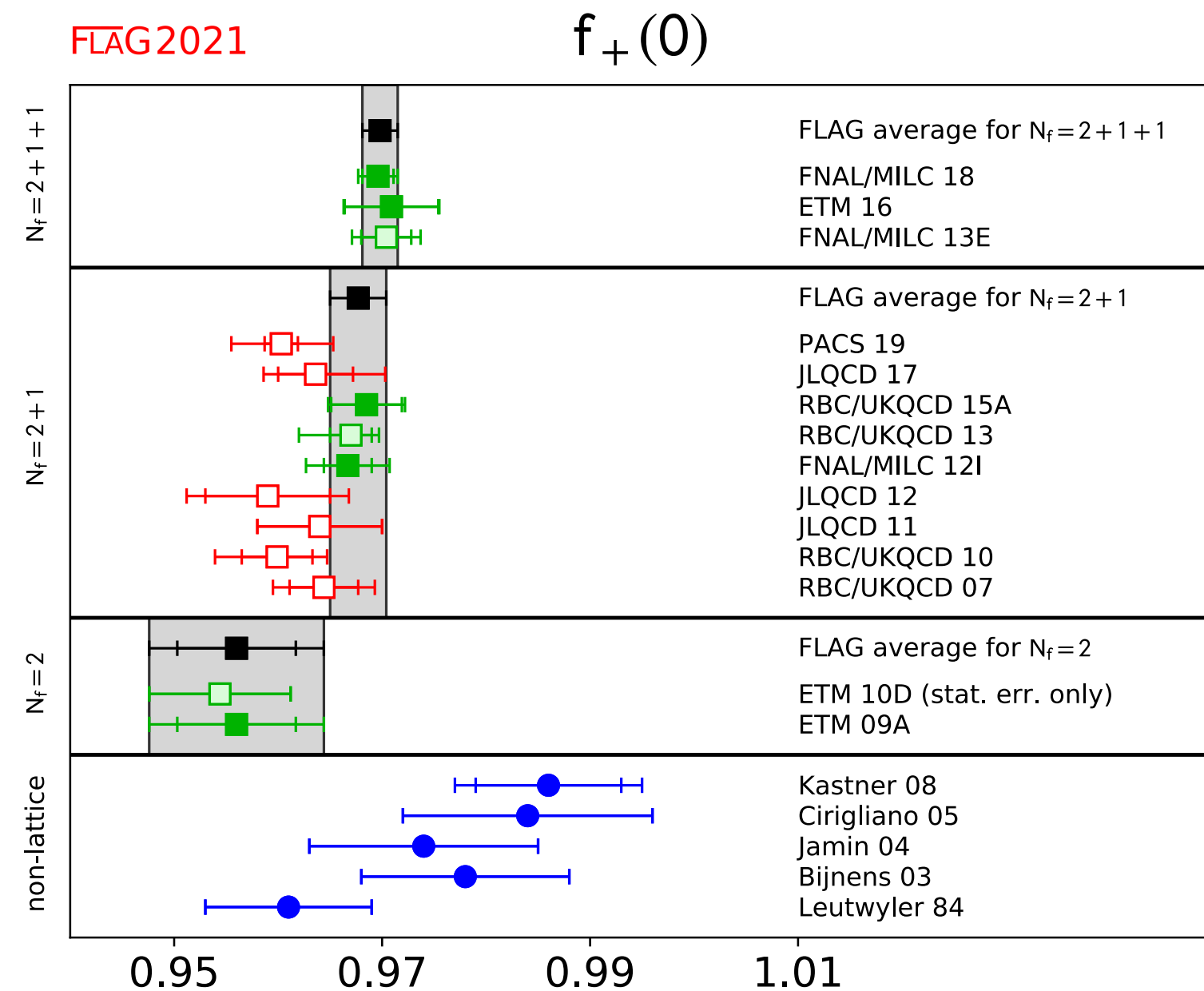
CKM: 1st row

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Precision test of actions

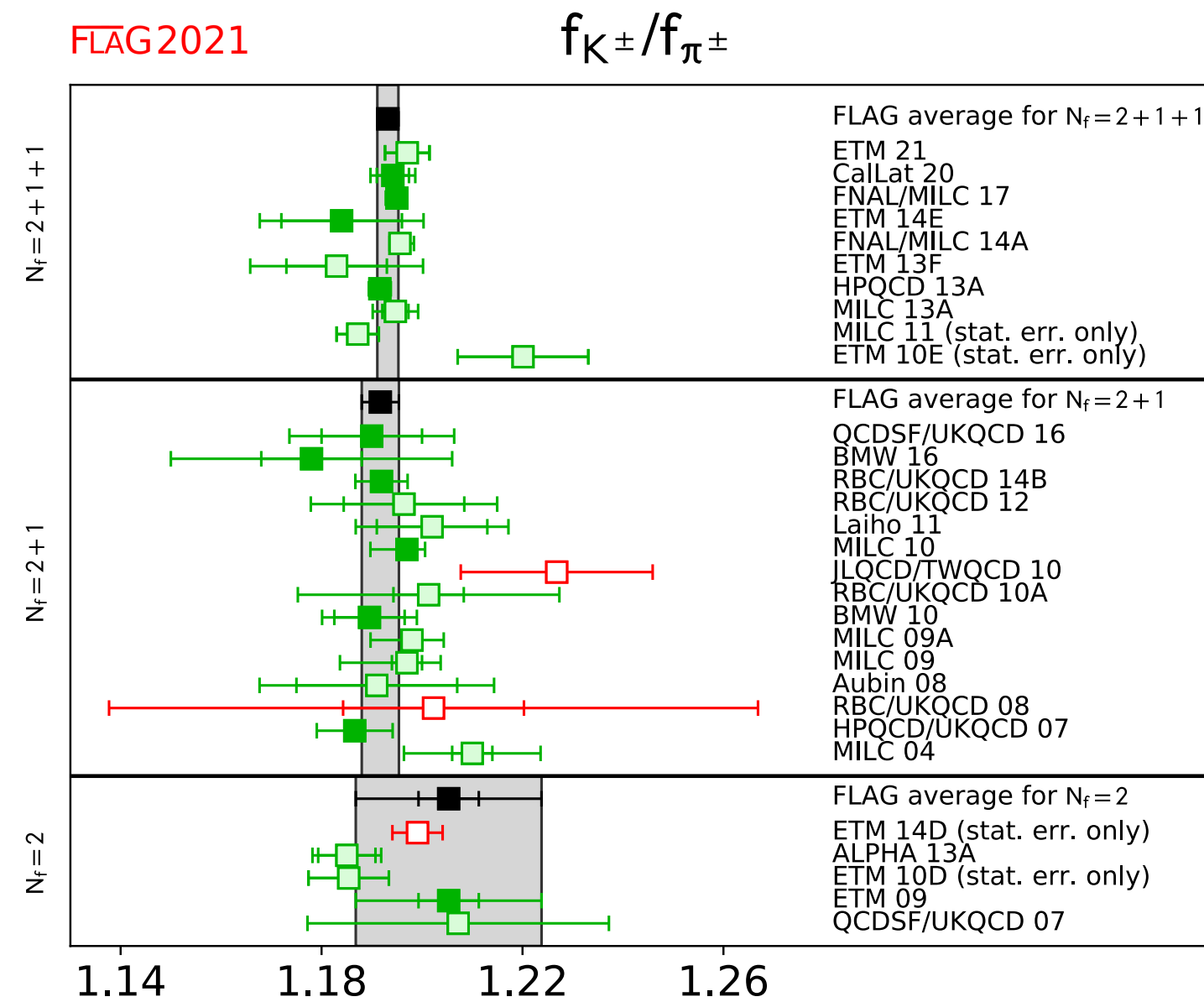
Kaon and pion decays

Semileptonic form factor
(normalization)



FLAG, arXiv:2111.09849

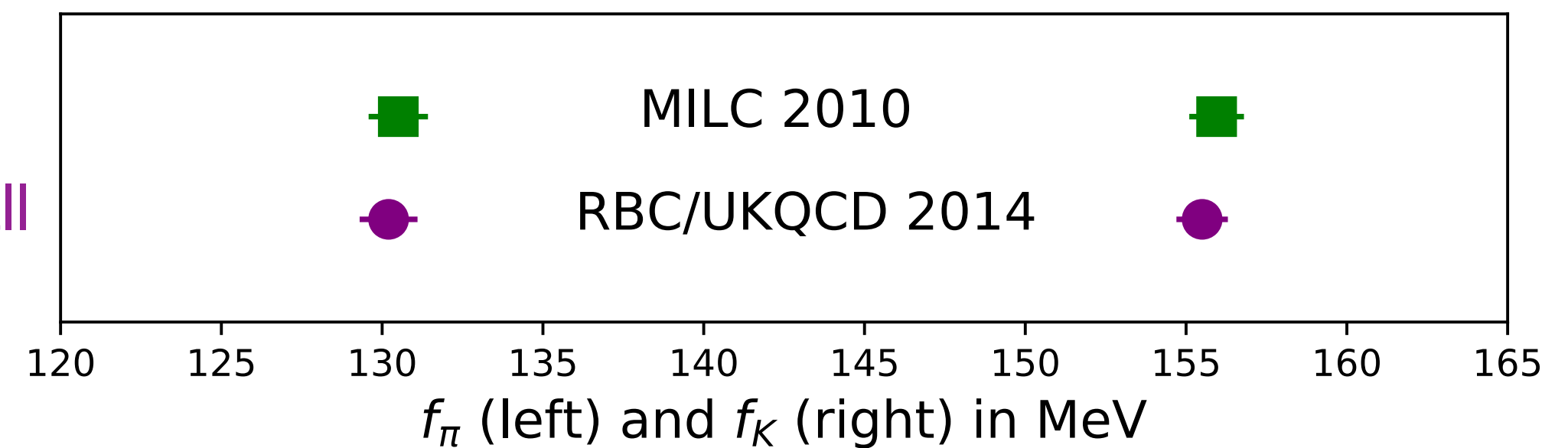
Leptonic decay constants



Agreement @ level < 1%

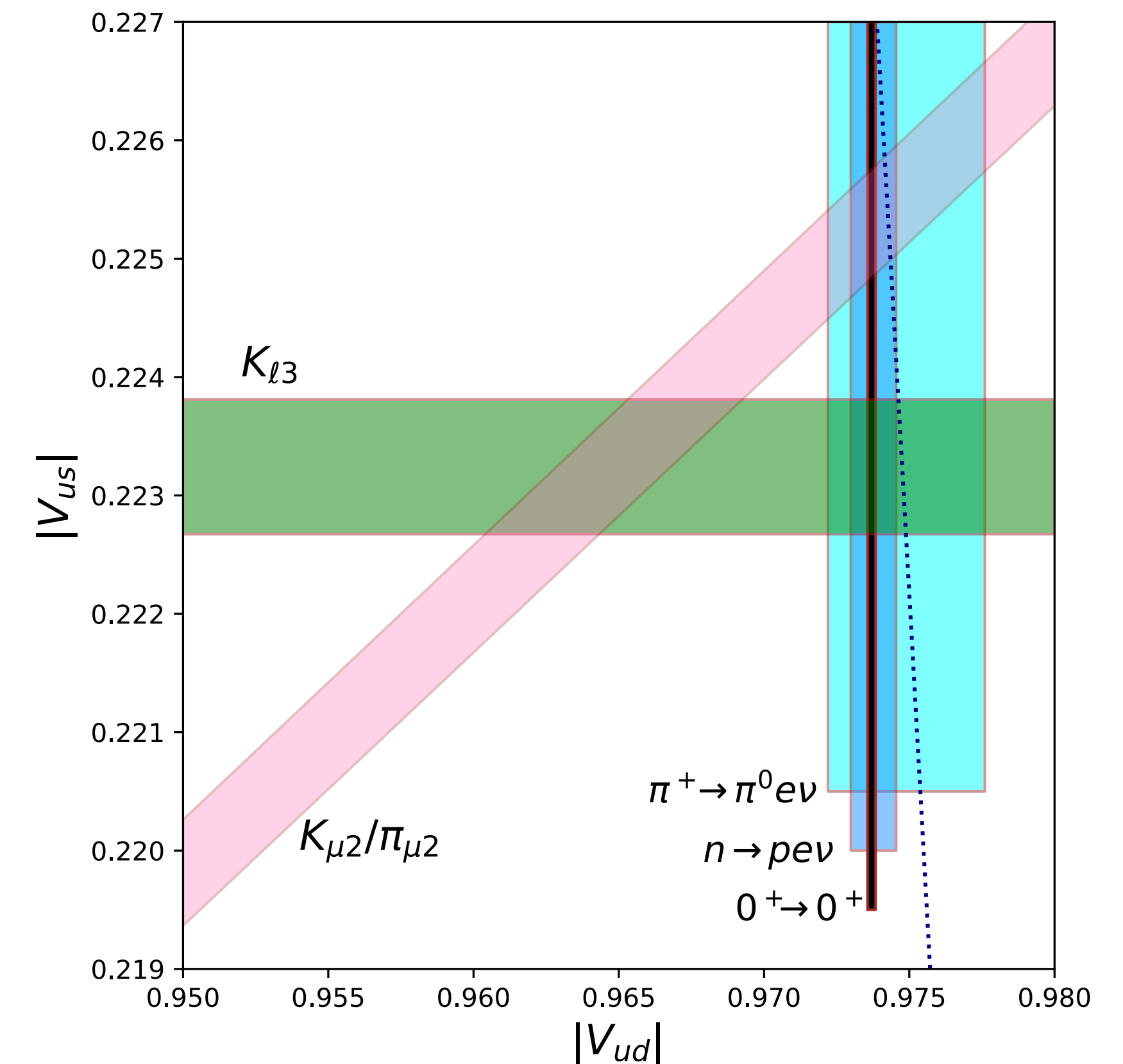
$n_f = 2 + 1$

Staggered
Domain wall



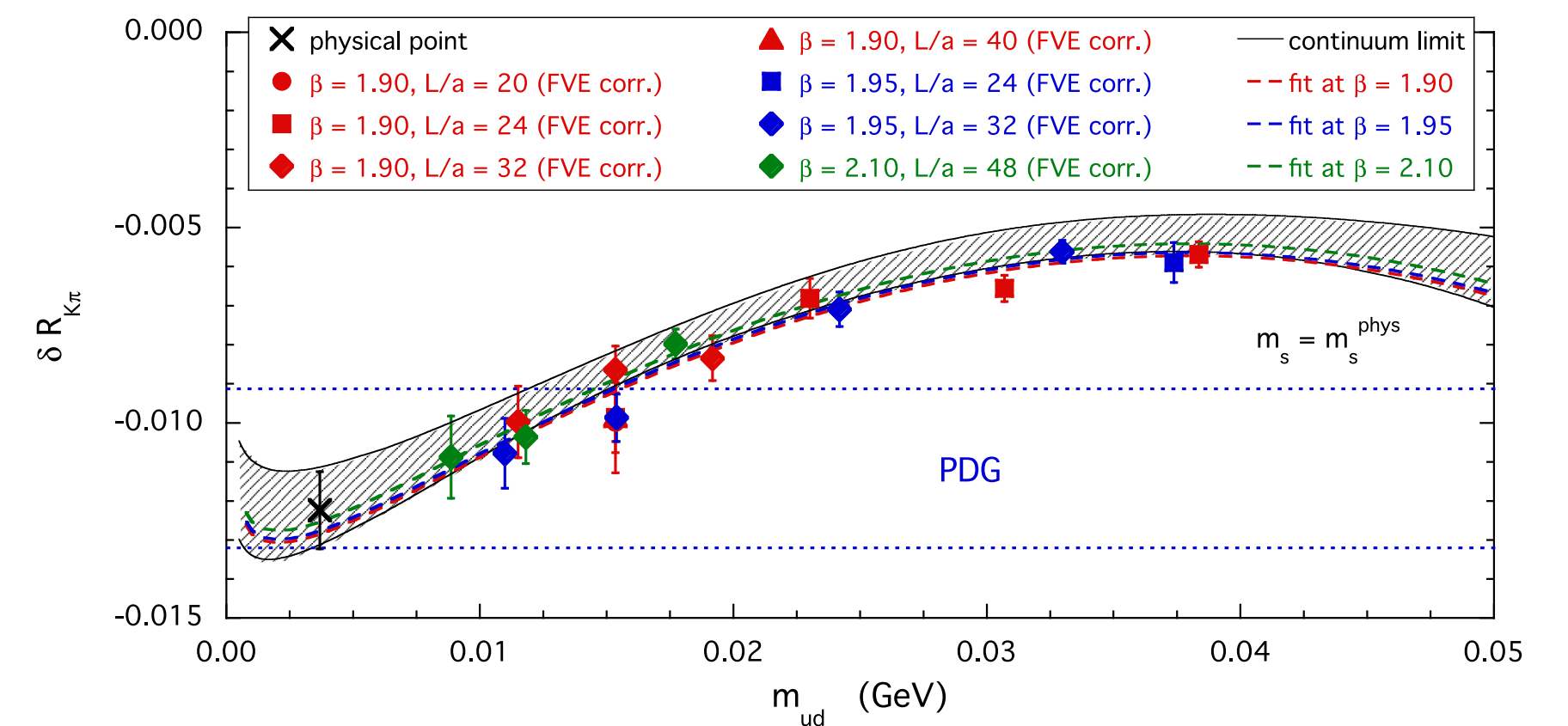
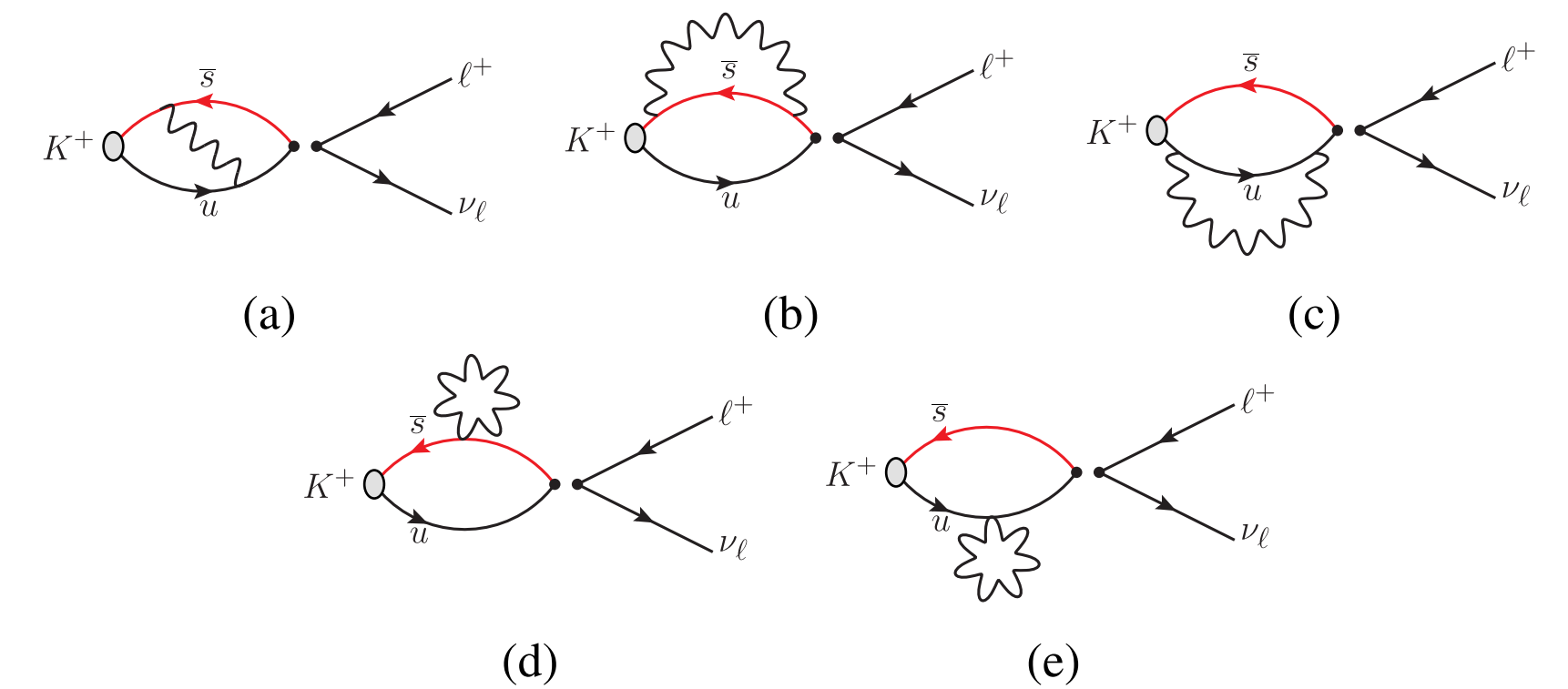
First row, unitarity, “Cabibbo anomaly”

- Note the precision on the axes! Lots of effort to shrink the constraints!
- Dotted line corresponds to $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$
- Two puzzles:
 - Are 2 parameters enough to explain the 3 classes of constraints?
 - Are the parameters consistent with CKM unitarity?
- Tension with unitarity due to new result from nuclear decays
- Precision of lattice matrix elements now means unequal u/d mass and QED effects are required



Including QED effects

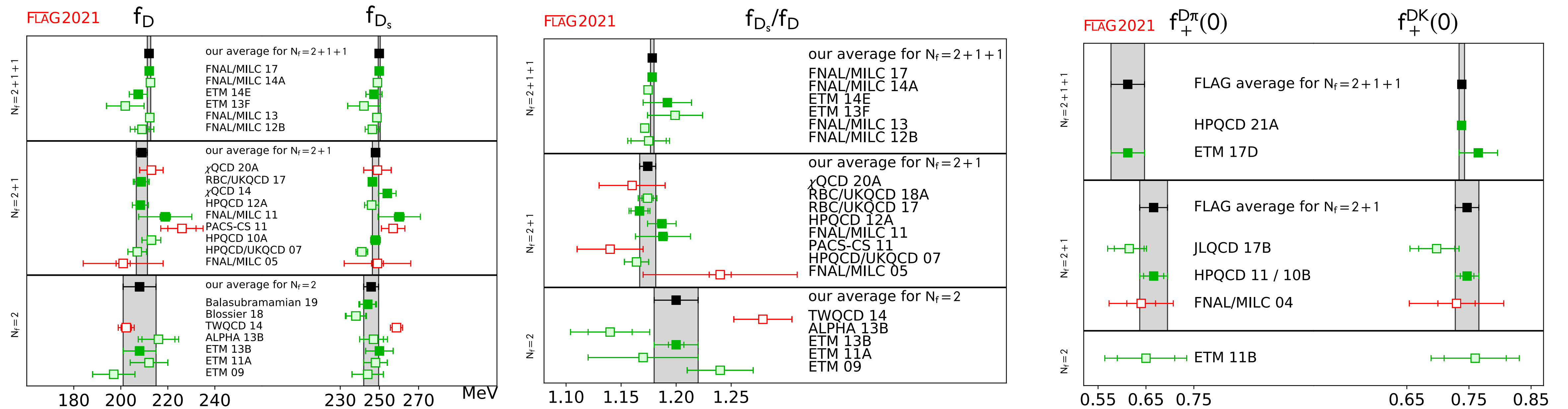
- RM123 method to include isospin breaking effects in LQCD determinations of leptonic decays of mesons (must include soft photon in final state)
- Several technical issues to face: finite volume effects, cancellation of IR divergences
- First calculation found good agreement with chiral P.T. predictions
- Continuing progress to extend applicability to larger photon momentum, requires electromagnetic form factors



CKM: 2nd row

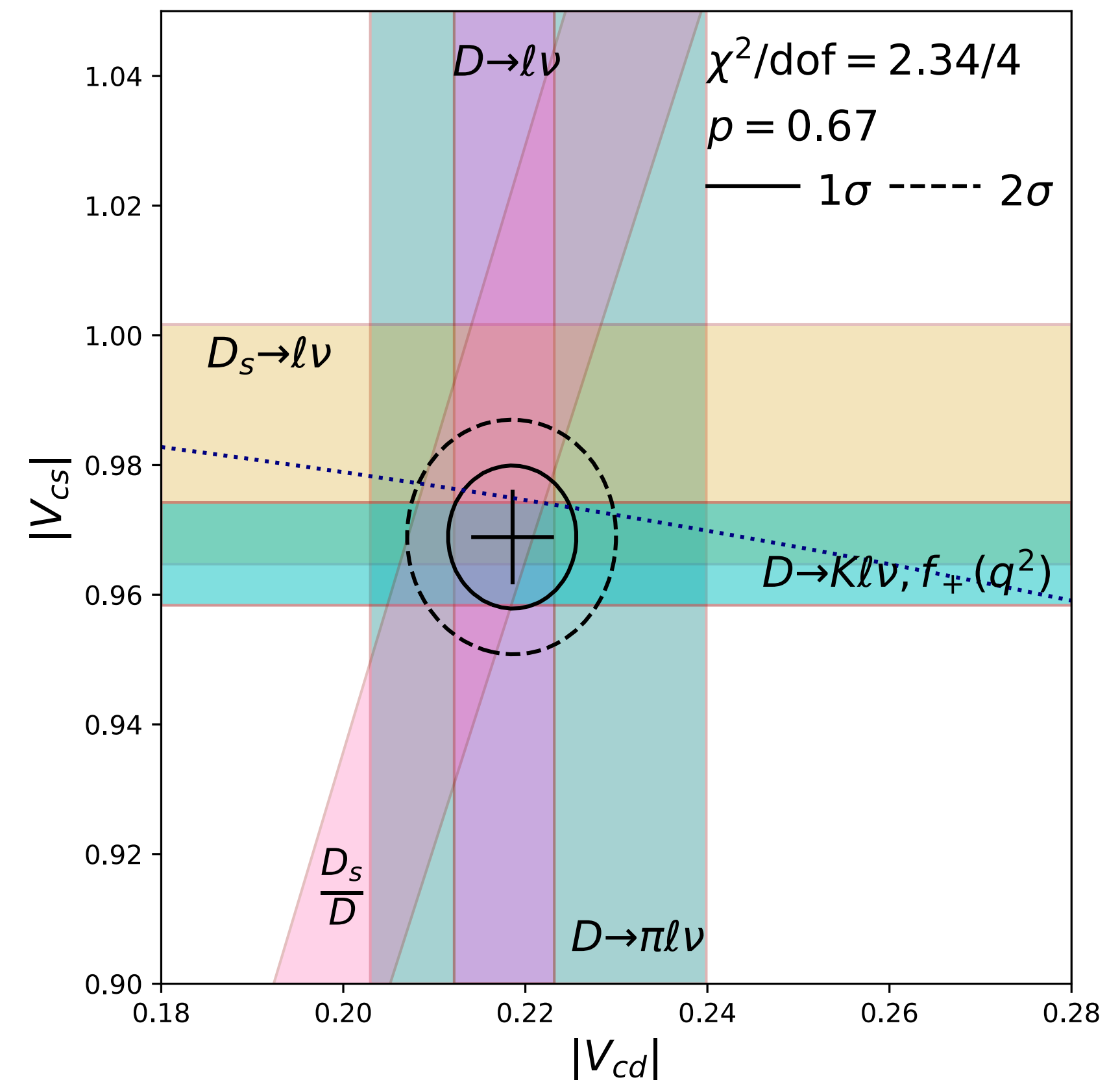
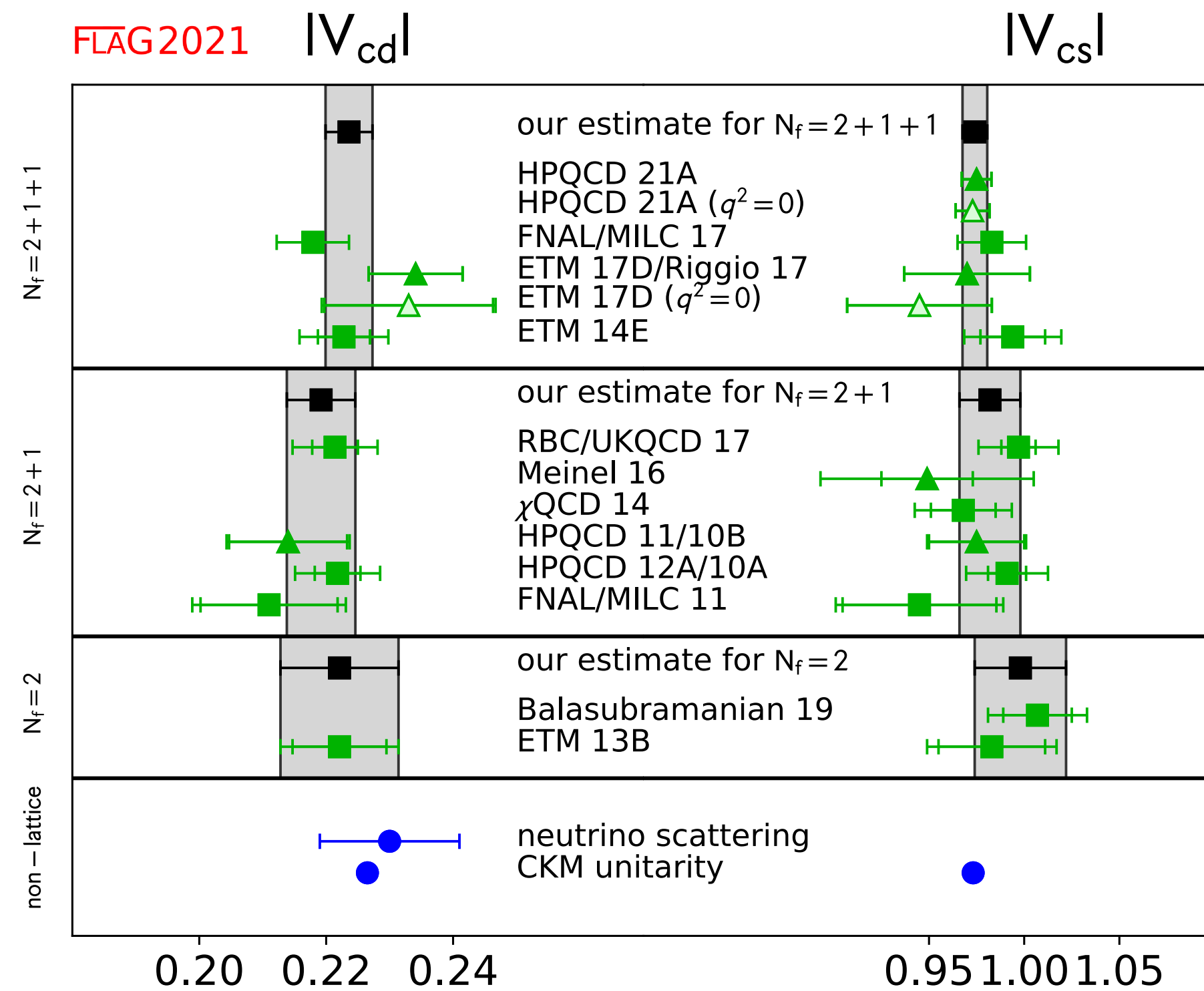
$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

D leptonic and semileptonic decays

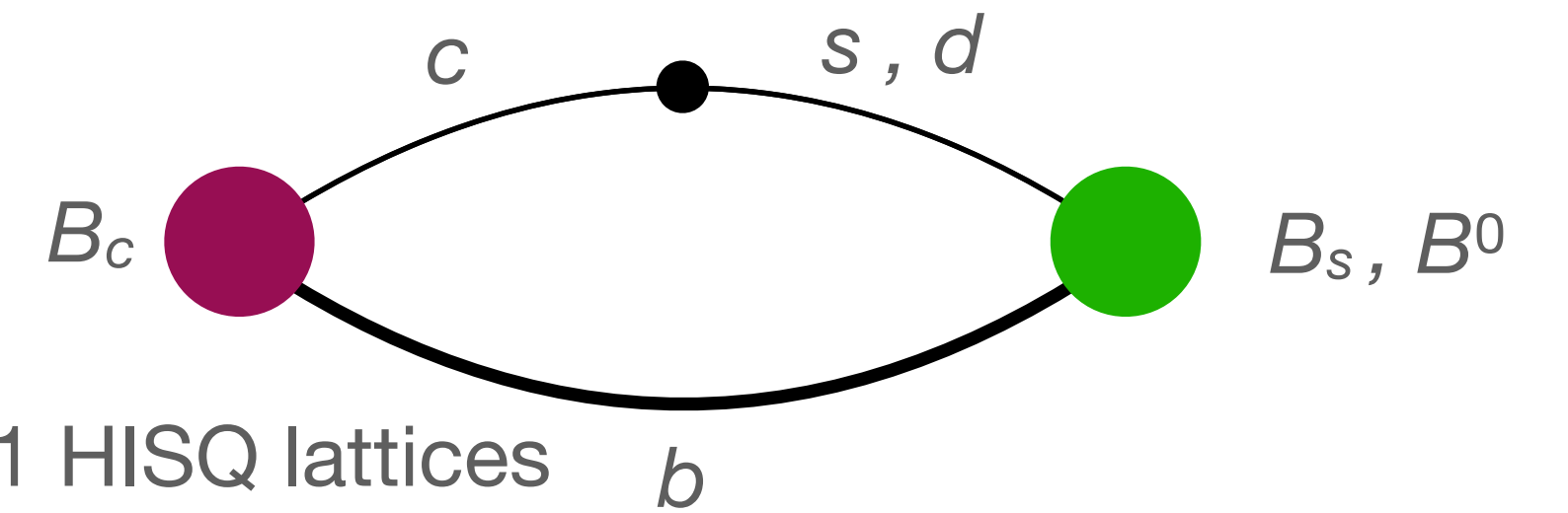


Much more on $D \rightarrow K\ell\nu$ from Chris Bouchard, Tues 09:30

Second row and unitarity

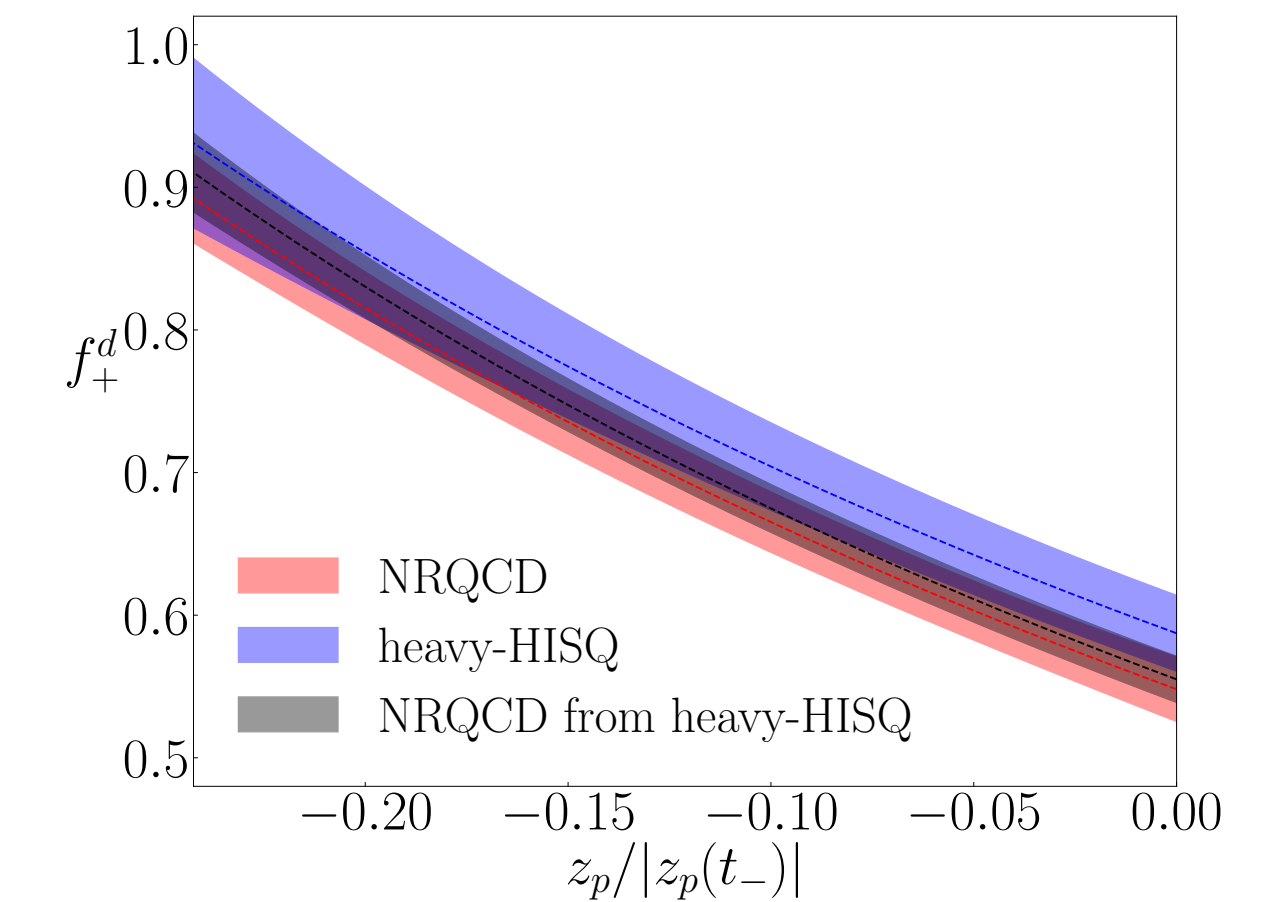
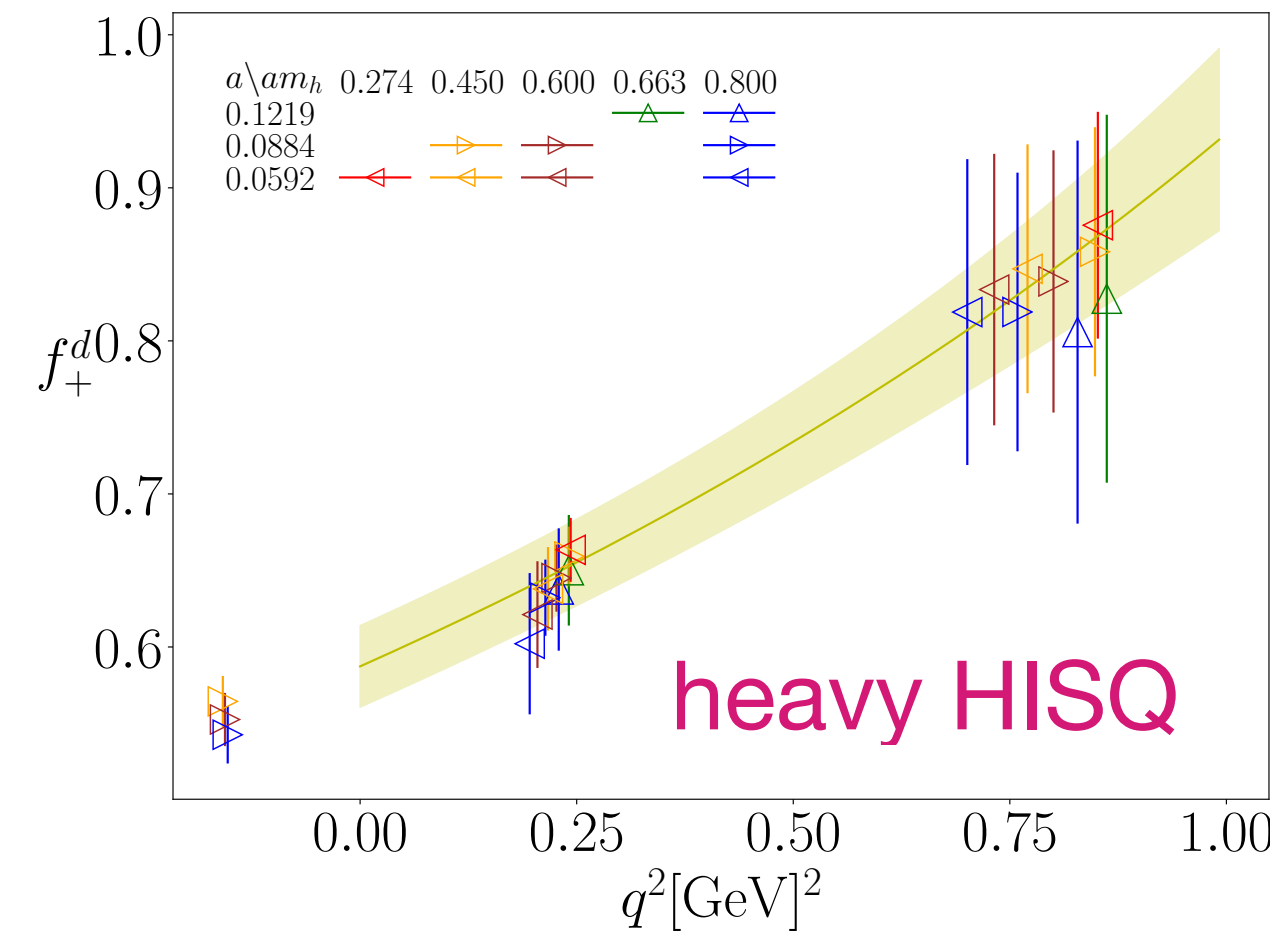
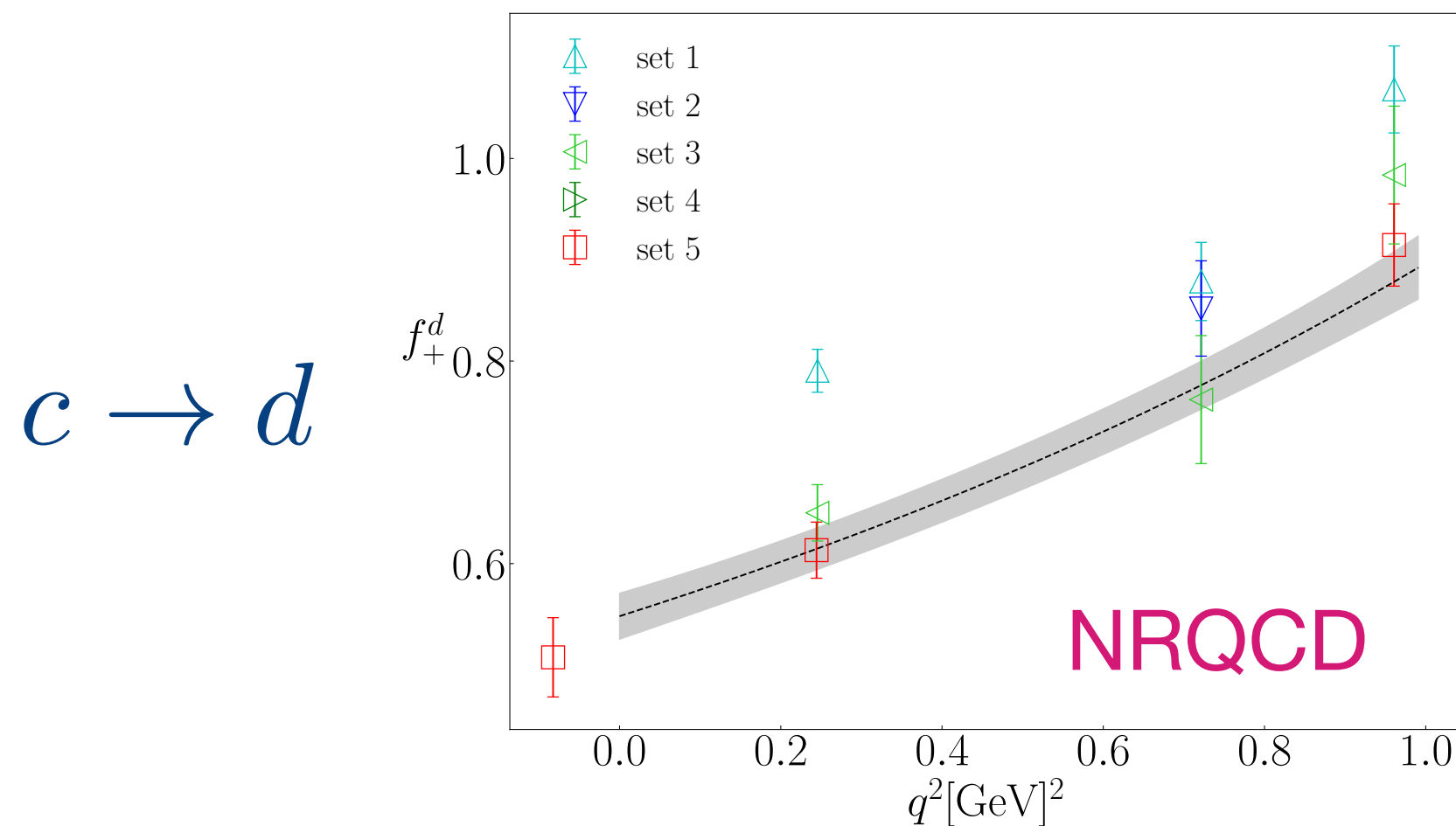
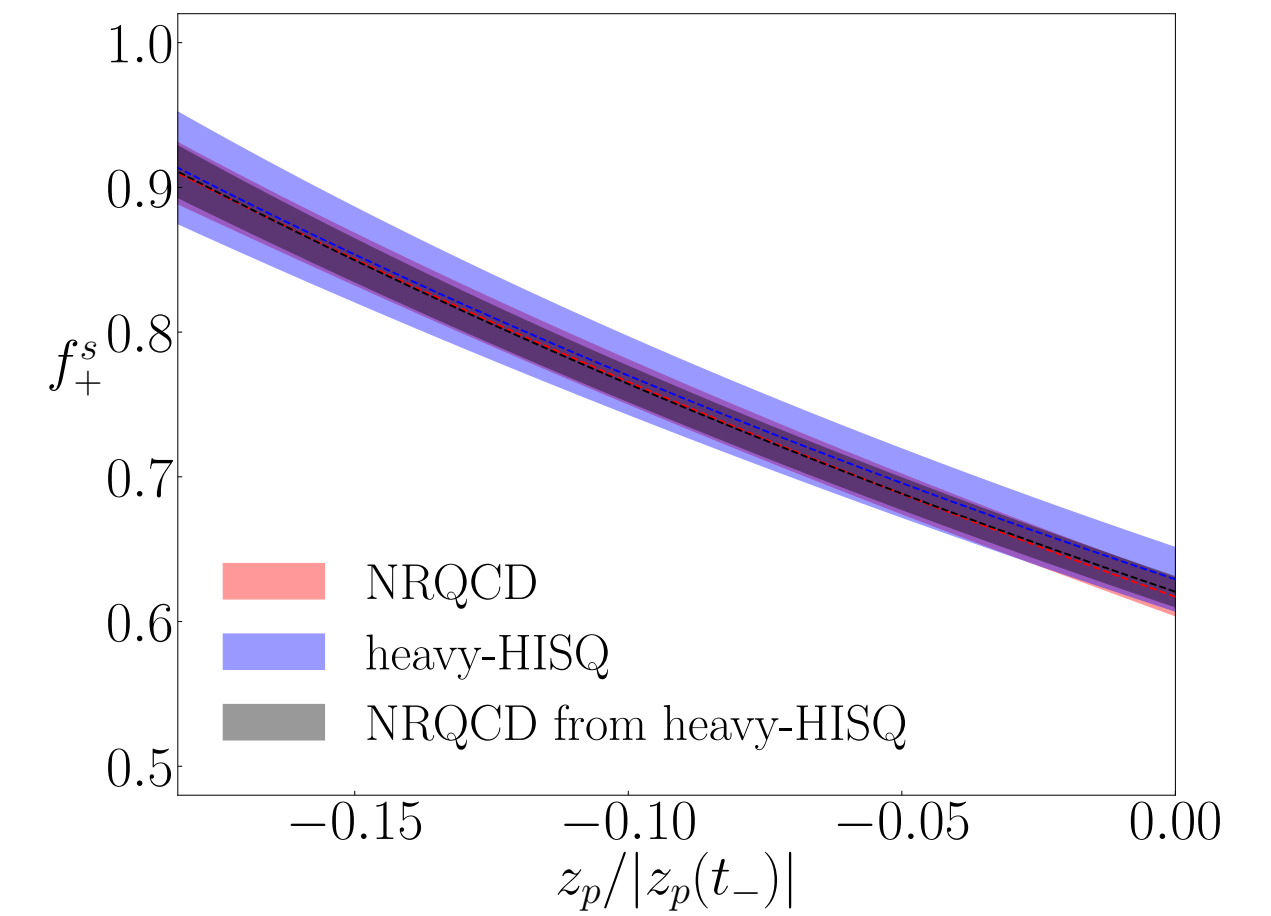
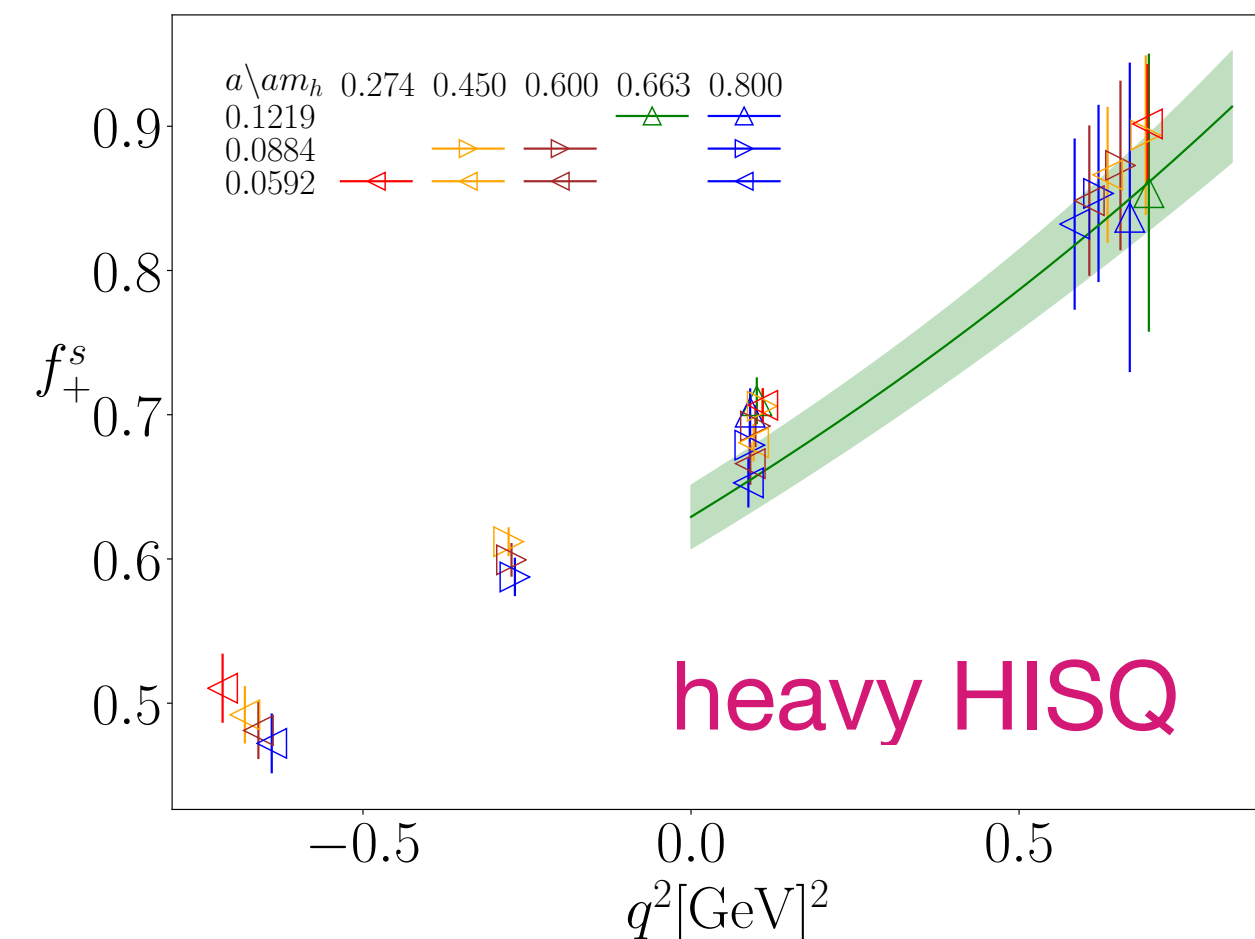
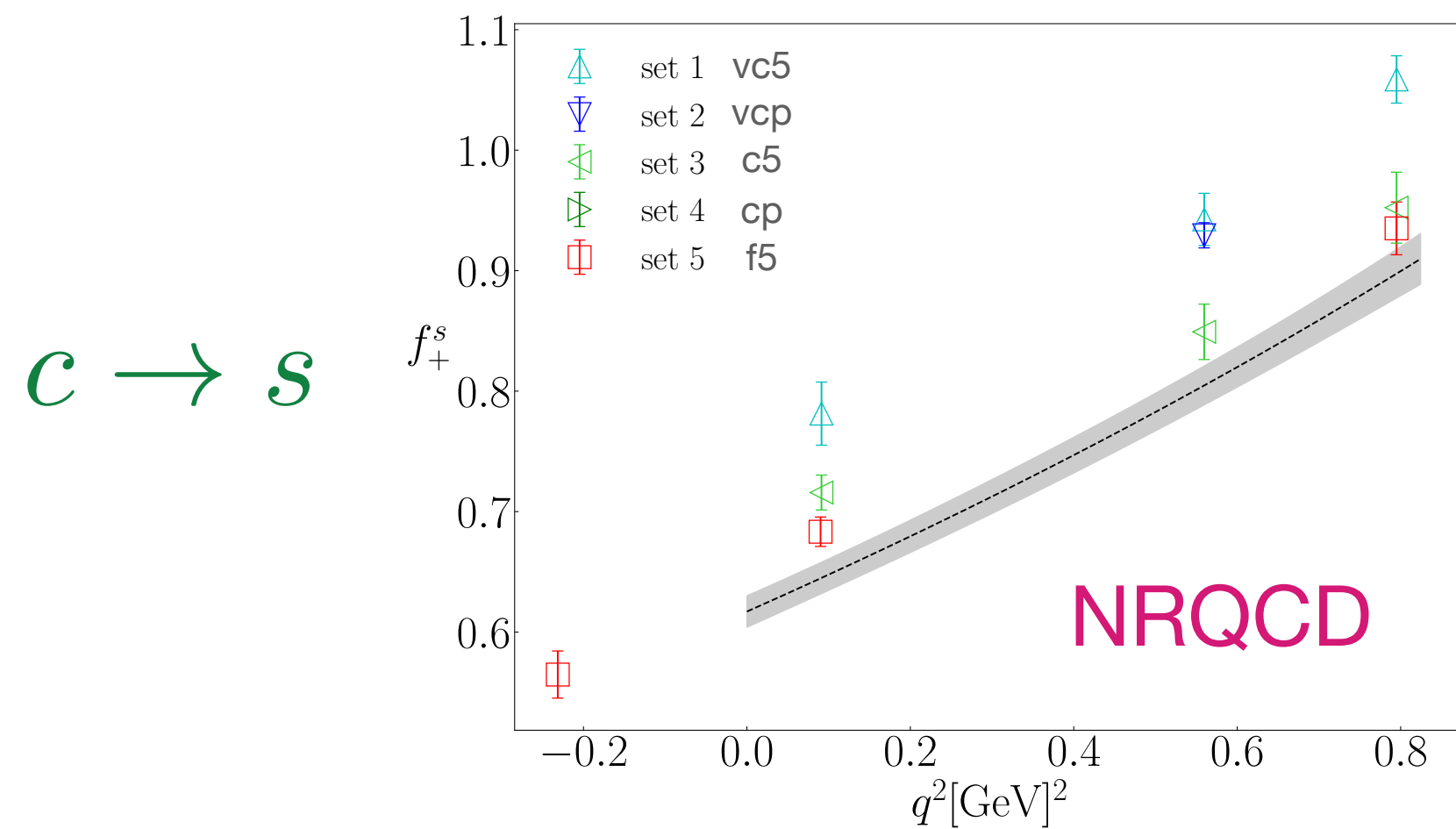


Charm decays in B_c

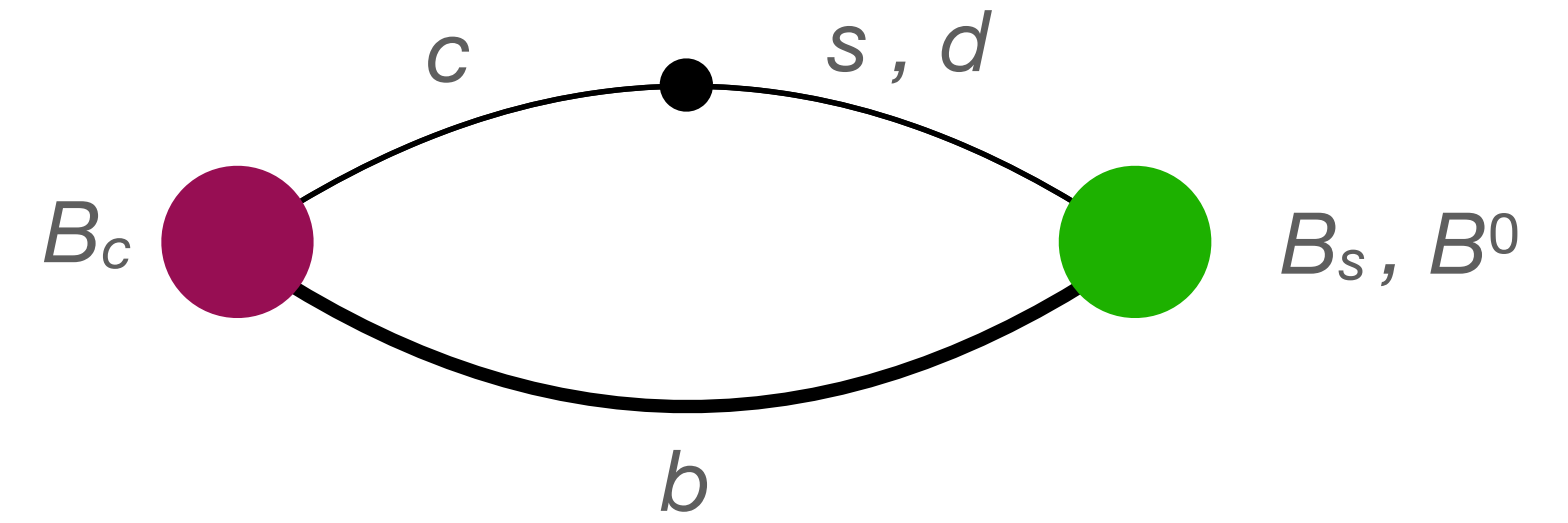


HPQCD: Combined analysis of NRQCD and heavy-HISQ results on MILC 2+1+1 HISQ lattices

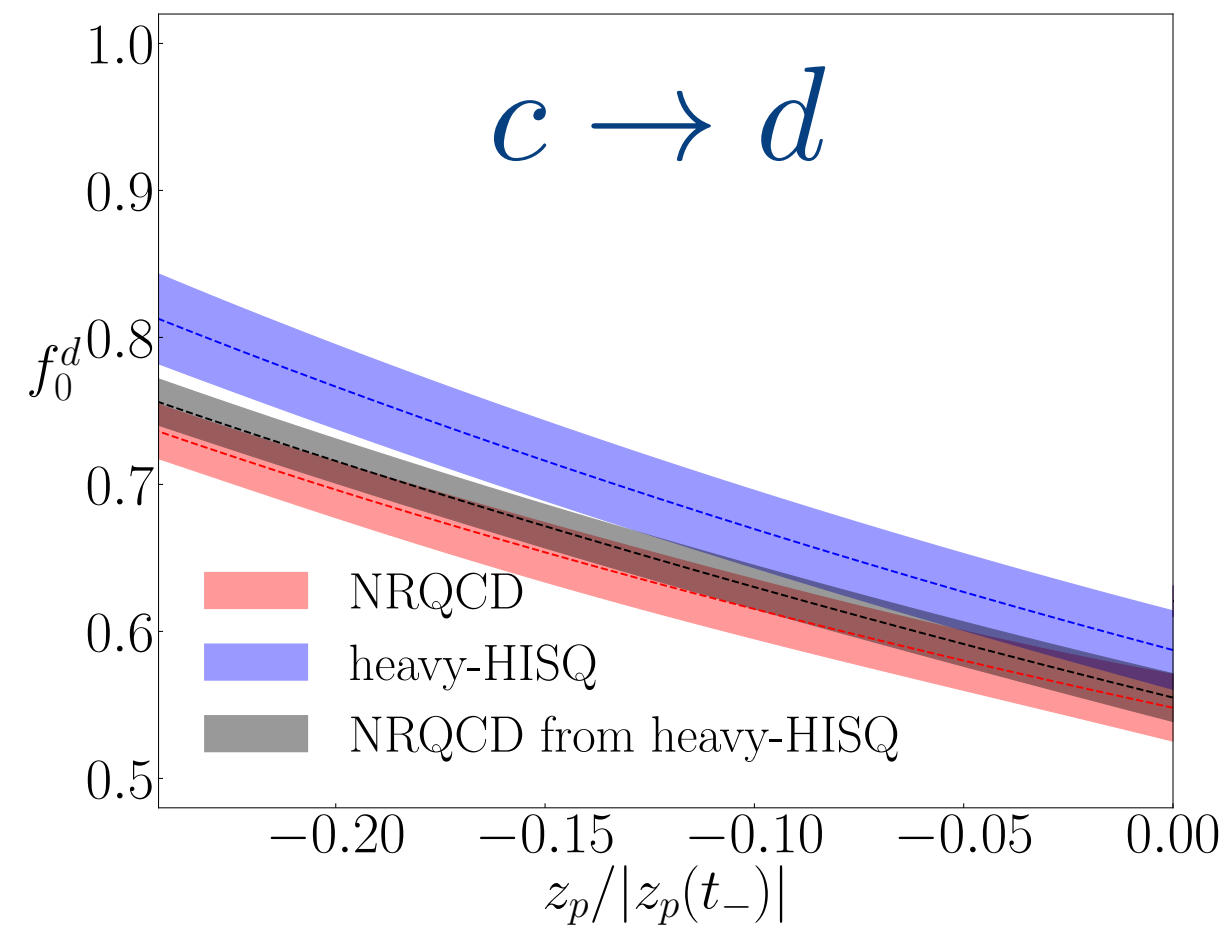
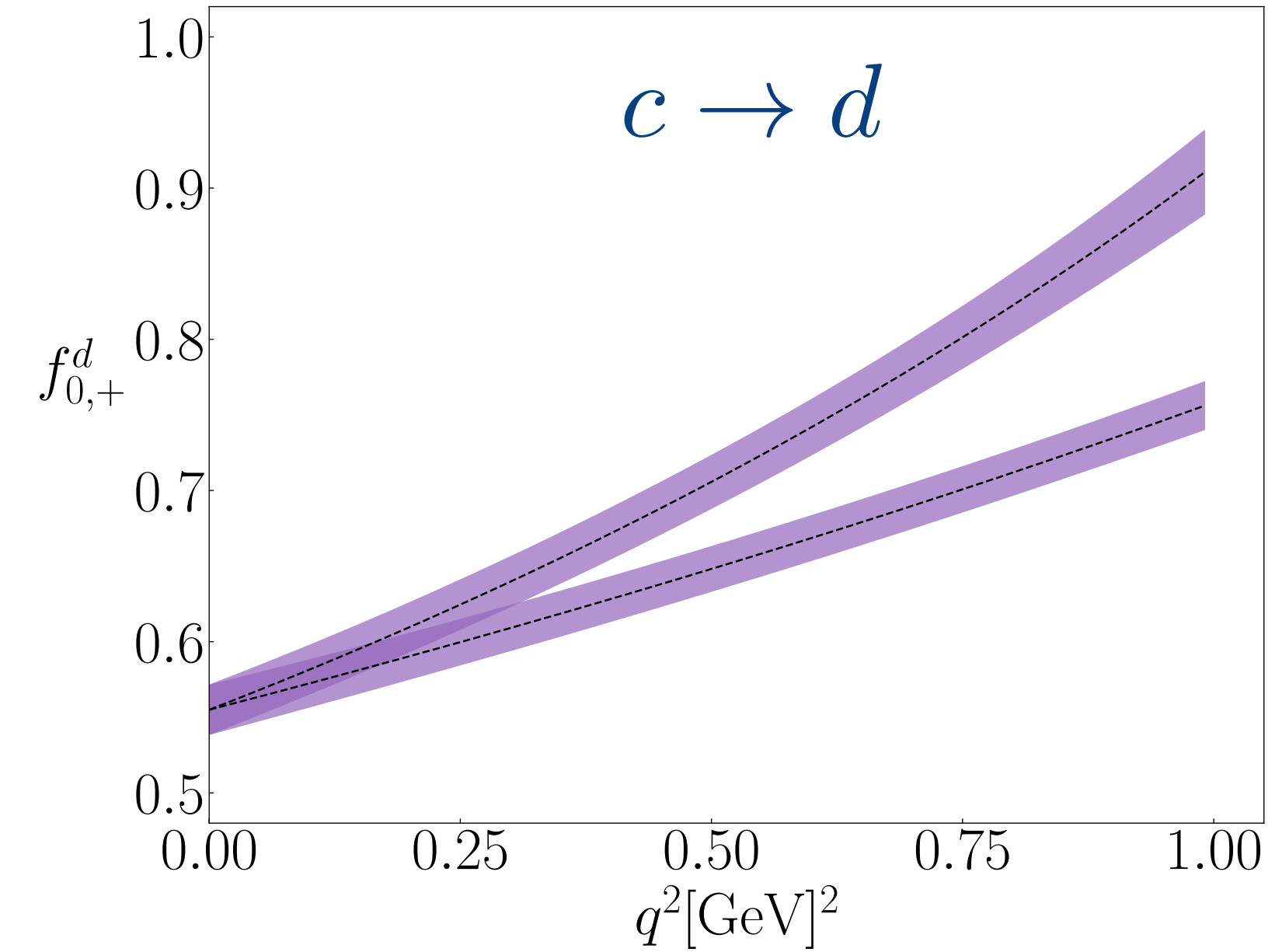
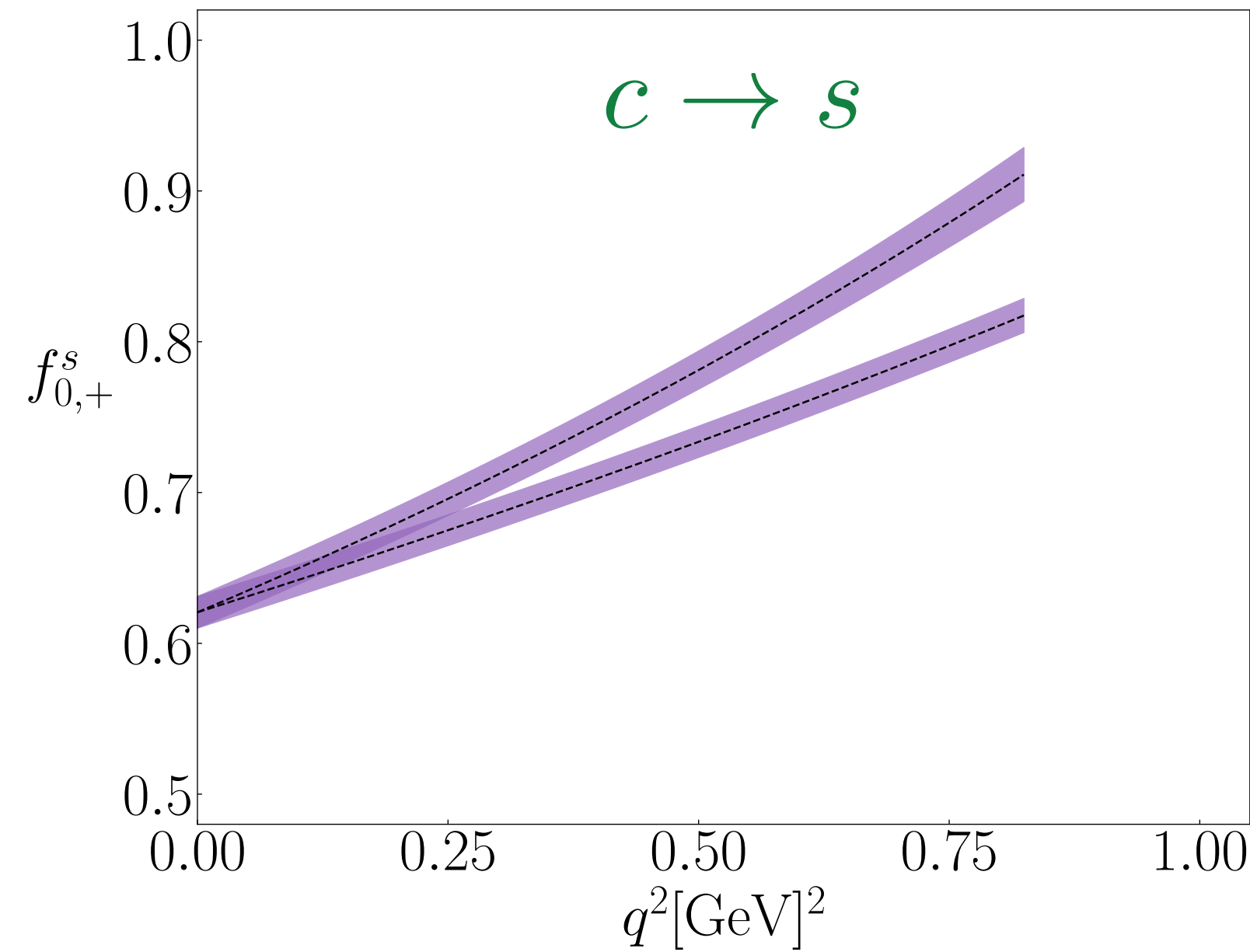
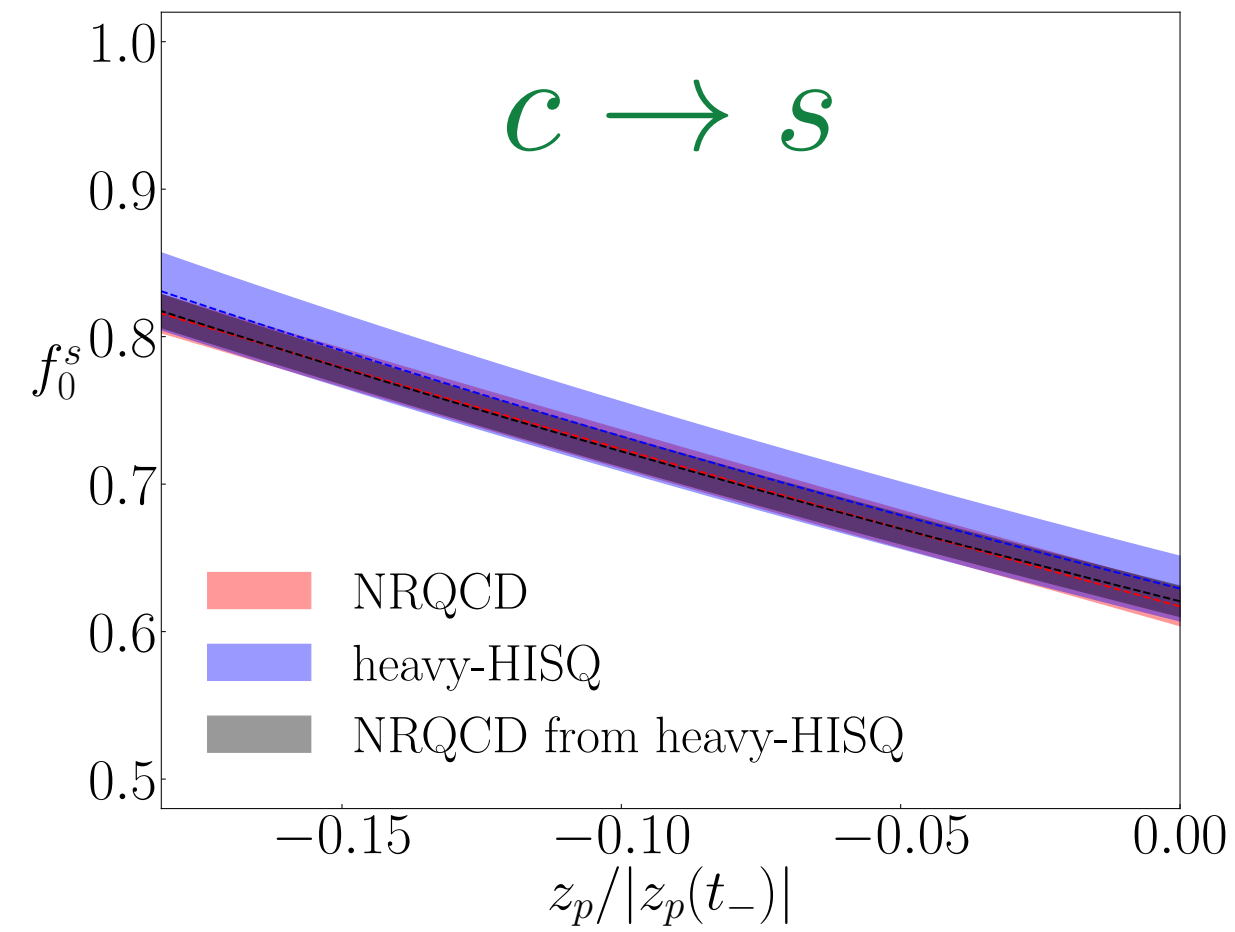
Vector form factor f_+



Charm decays in B_c



Scalar form factor f_0



SM predictions:

$$\Gamma(B_c^+ \rightarrow B_s^0 \bar{\ell} \nu_\ell) = 26.2(1.2) \times 10^9 \text{ s}^{-1}$$

$$\Gamma(B_c^+ \rightarrow \bar{B}^0 \bar{\ell} \nu_\ell) = 1.65(10) \times 10^9 \text{ s}^{-1}$$

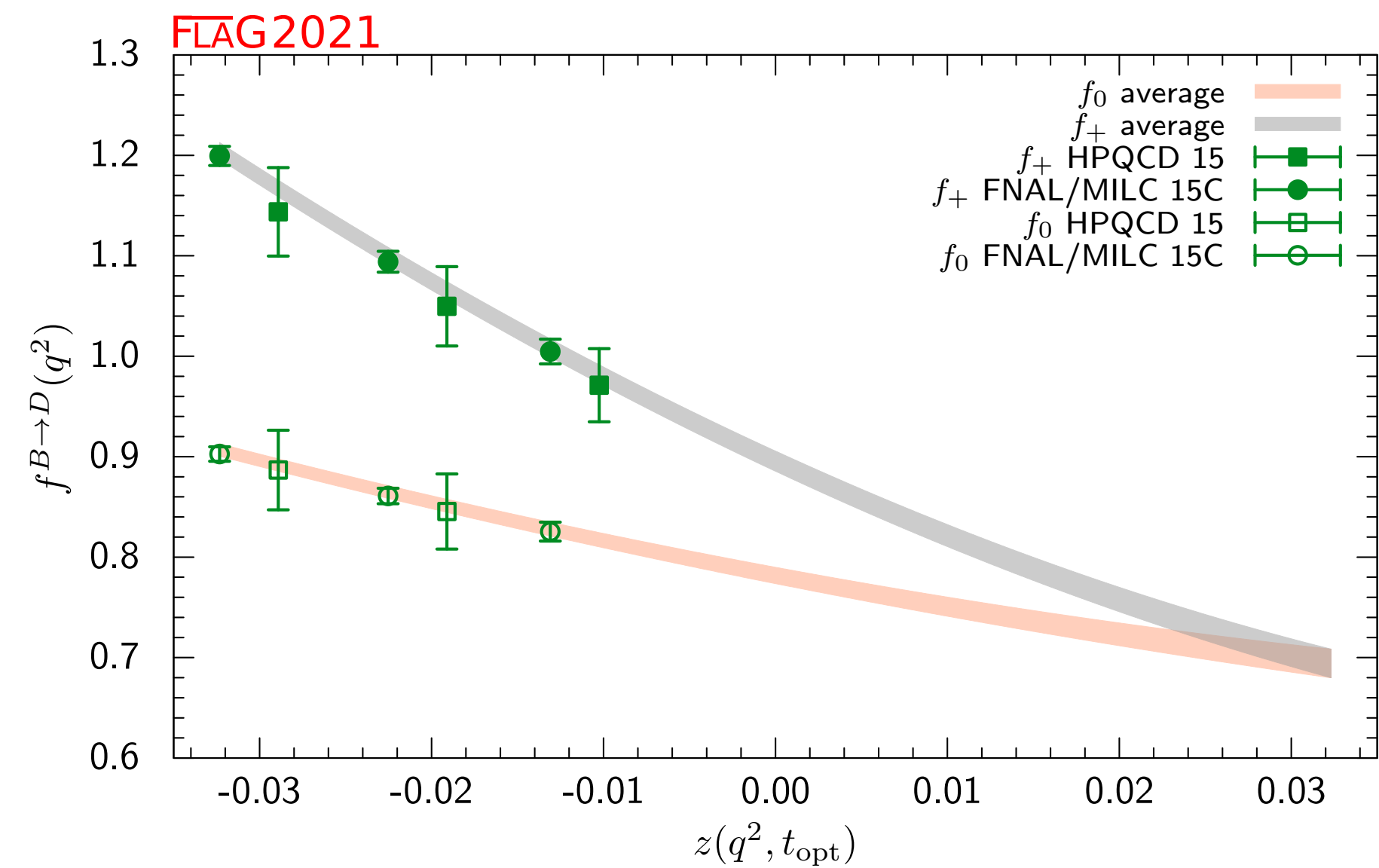
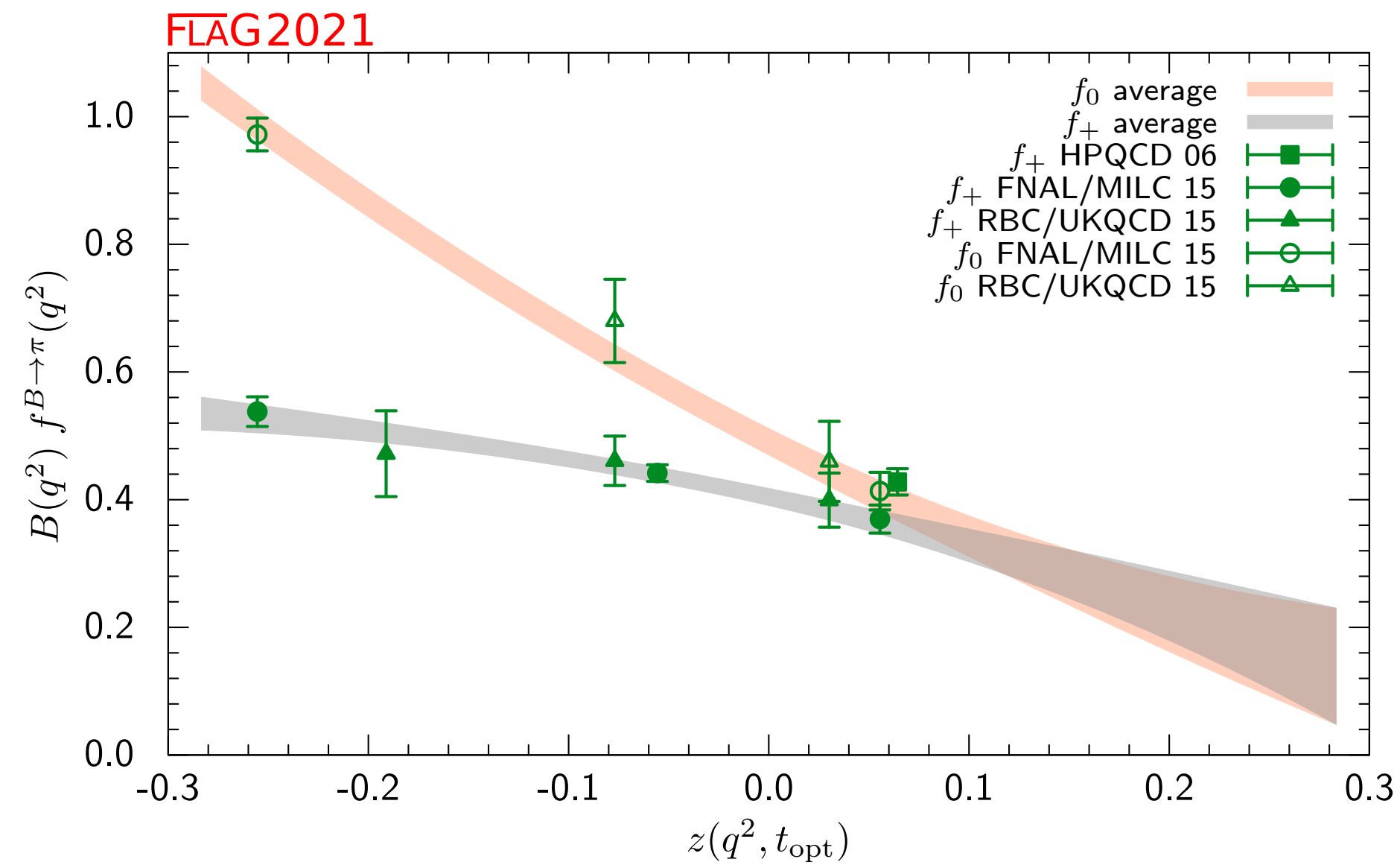
CKM: 3rd column & V_{xb}

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Semileptonic decays

High precision modes

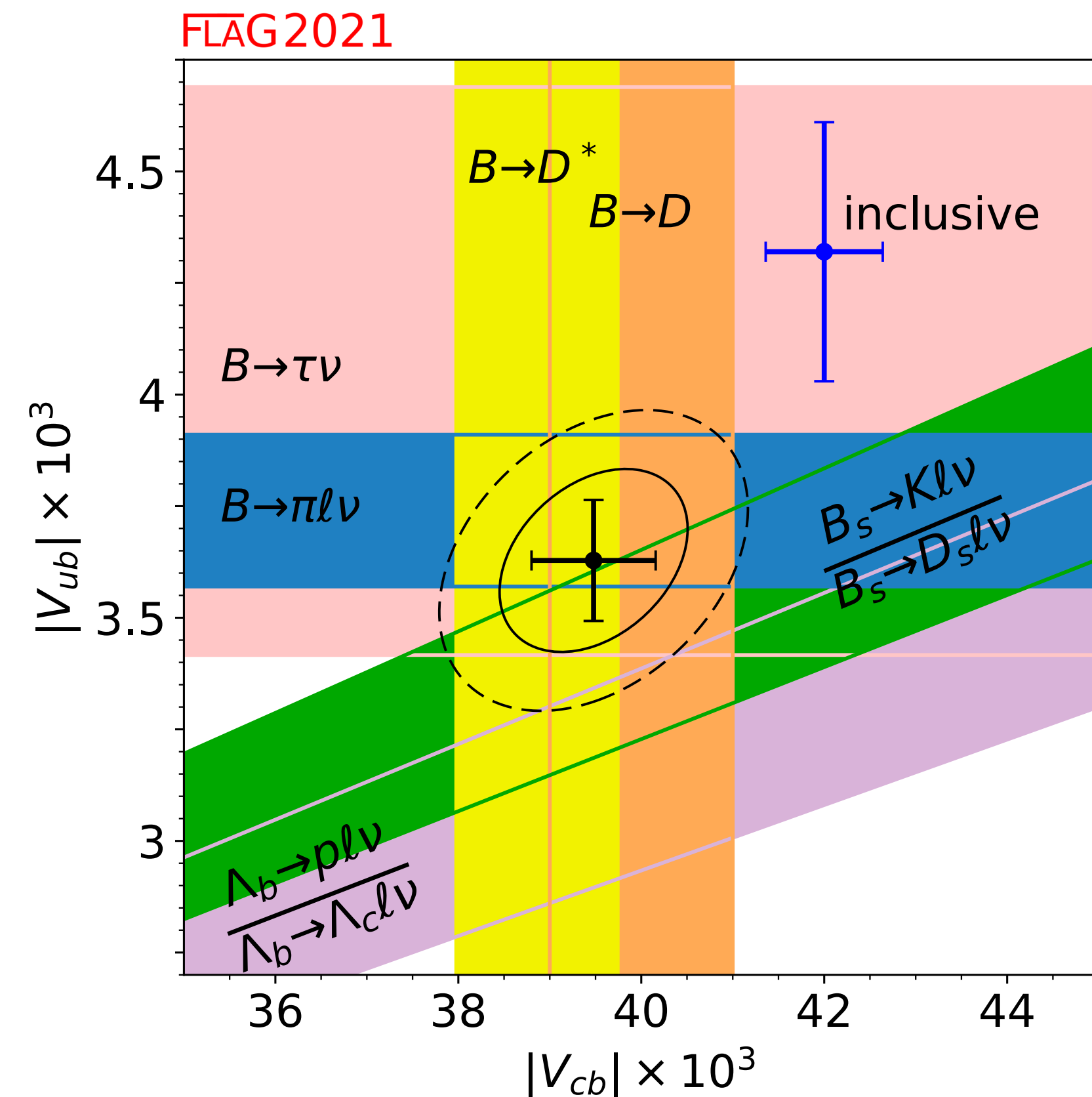
$B \rightarrow \pi$ and $B \rightarrow D$ form factors awaiting updates



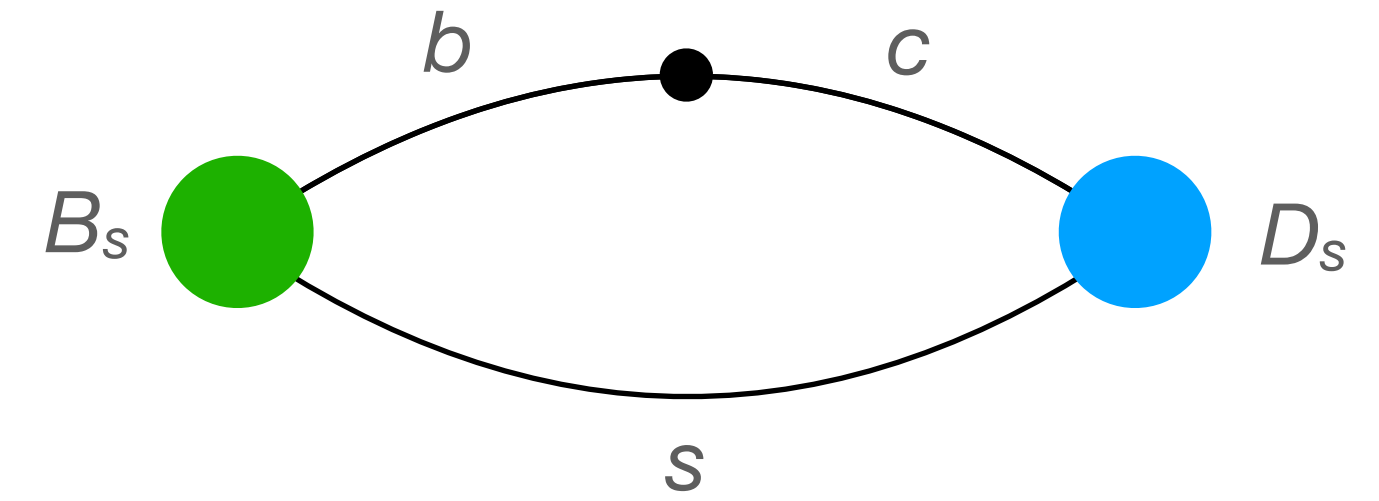
Talk by Alejandro Vaquero next, covering new results for $B \rightarrow D^*$ form factors

Third column

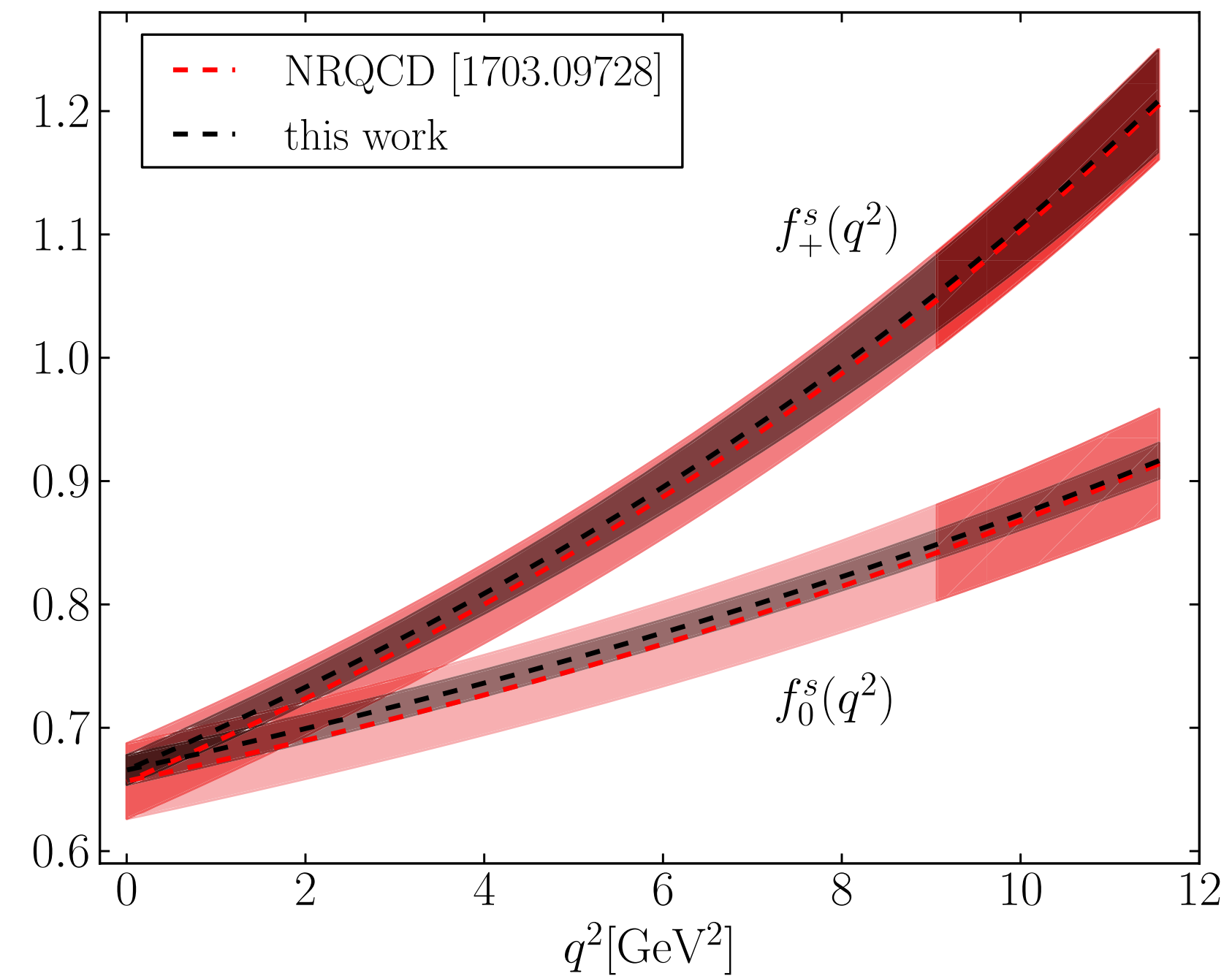
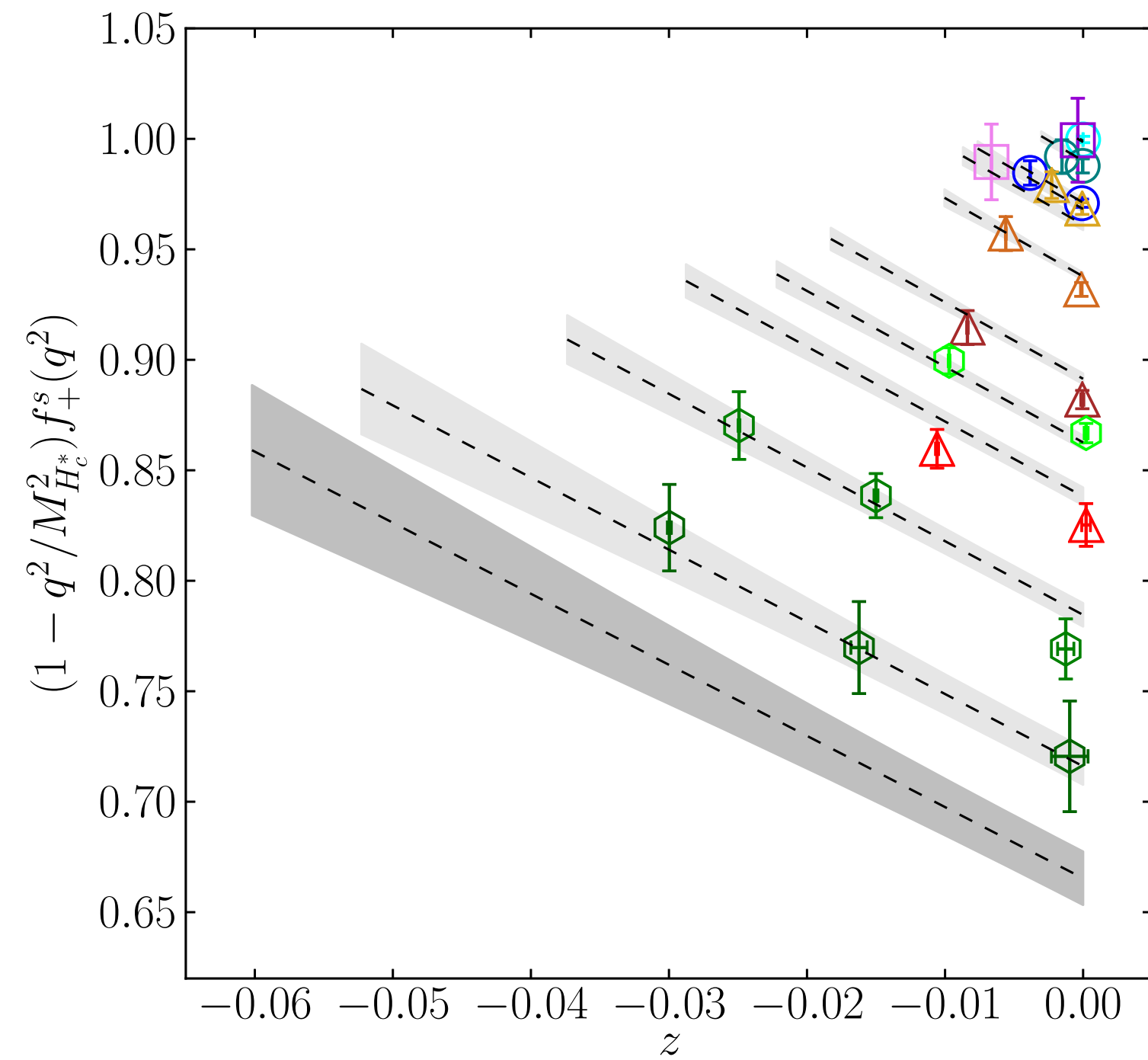
- Persistent tension between inclusive and exclusive determinations
- Much more about $B \rightarrow D^* \ell \nu$ in Alejandro Vaquero's talk on Monday at 11:00
- Details of form factor calculations also in Chris Bouchard's talk on Tuesday at 09:30
- More exclusive modes being measured, requiring corresponding form factors



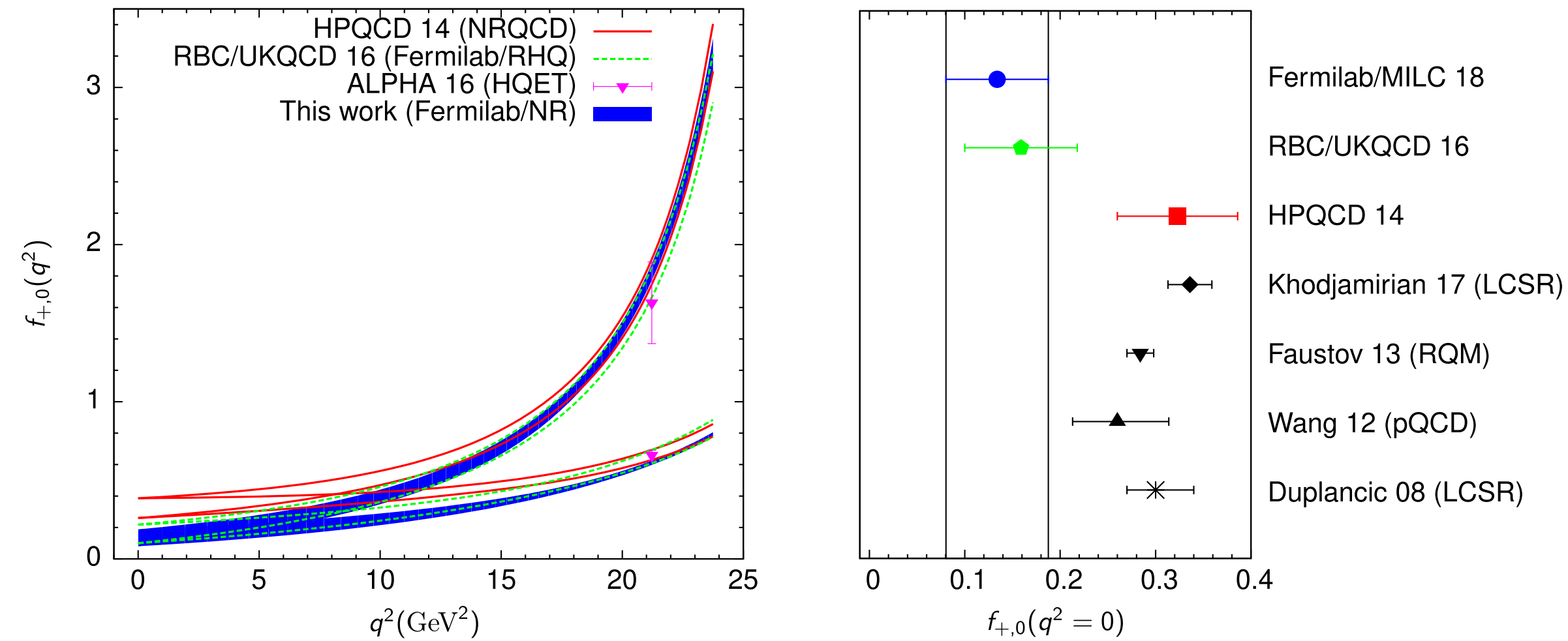
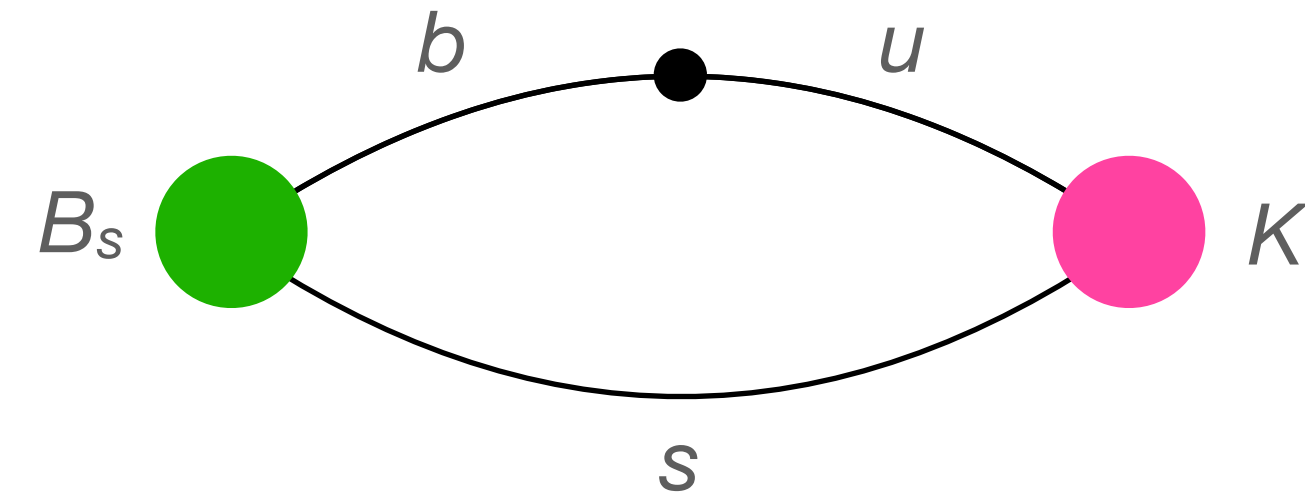
Strange spectator: $b \rightarrow c$



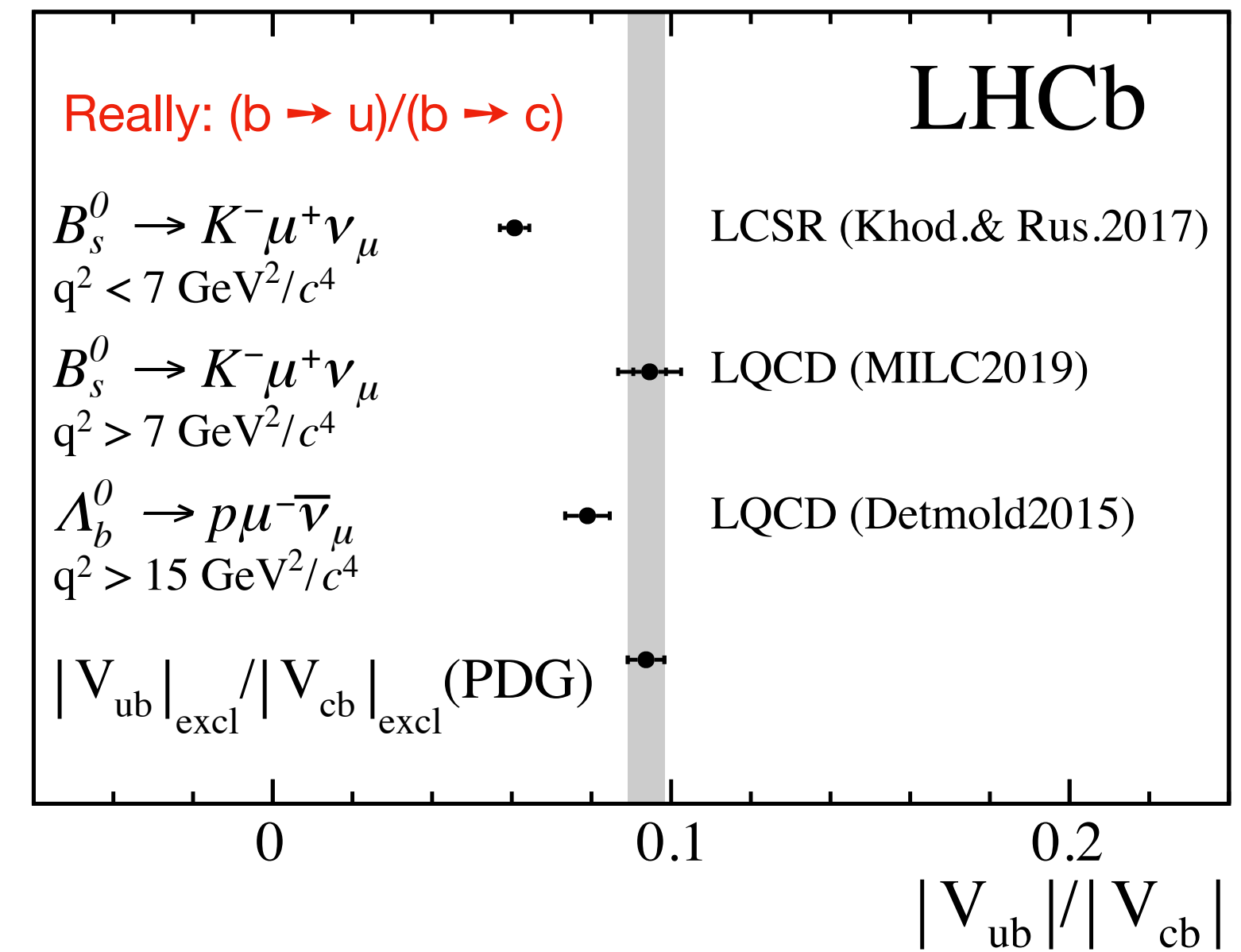
HPQCD: Heavy HISQ on MILC's $n_f = 2+1+1$ HISQ lattices



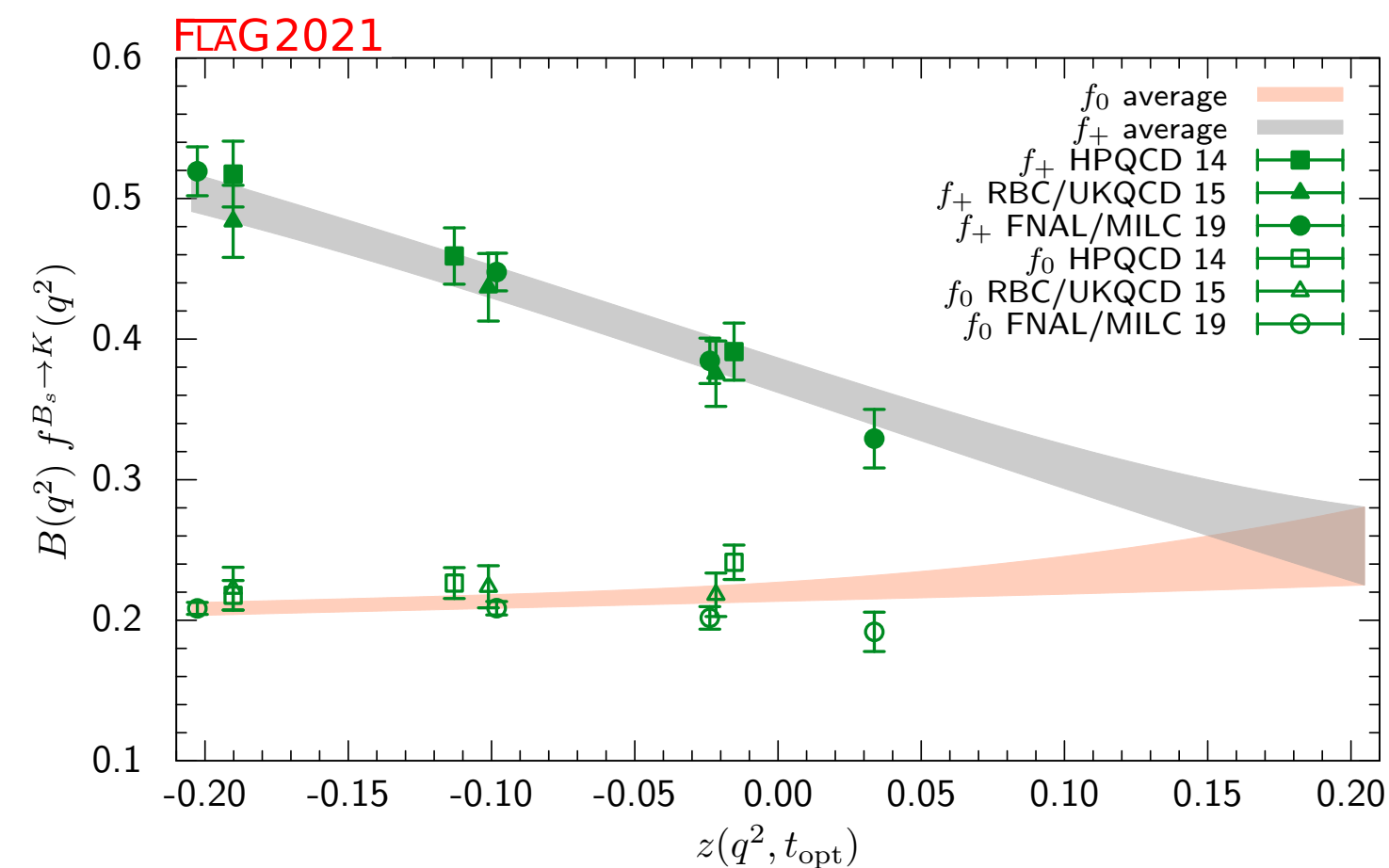
Strange spectator: $b \rightarrow u$



Figures from Bazavov et al., (Fermilab/MILC), *Phys. Rev. D* 100, 034501 (2019)

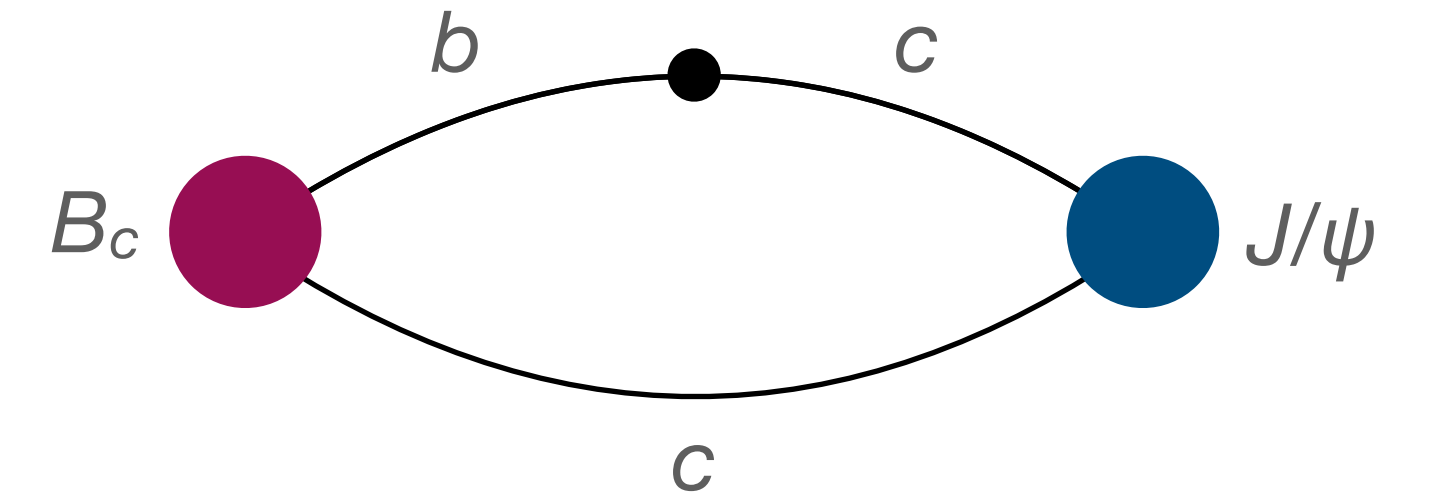


LHCb, *Phys. Rev. Lett.* 126, 081804 (2021)

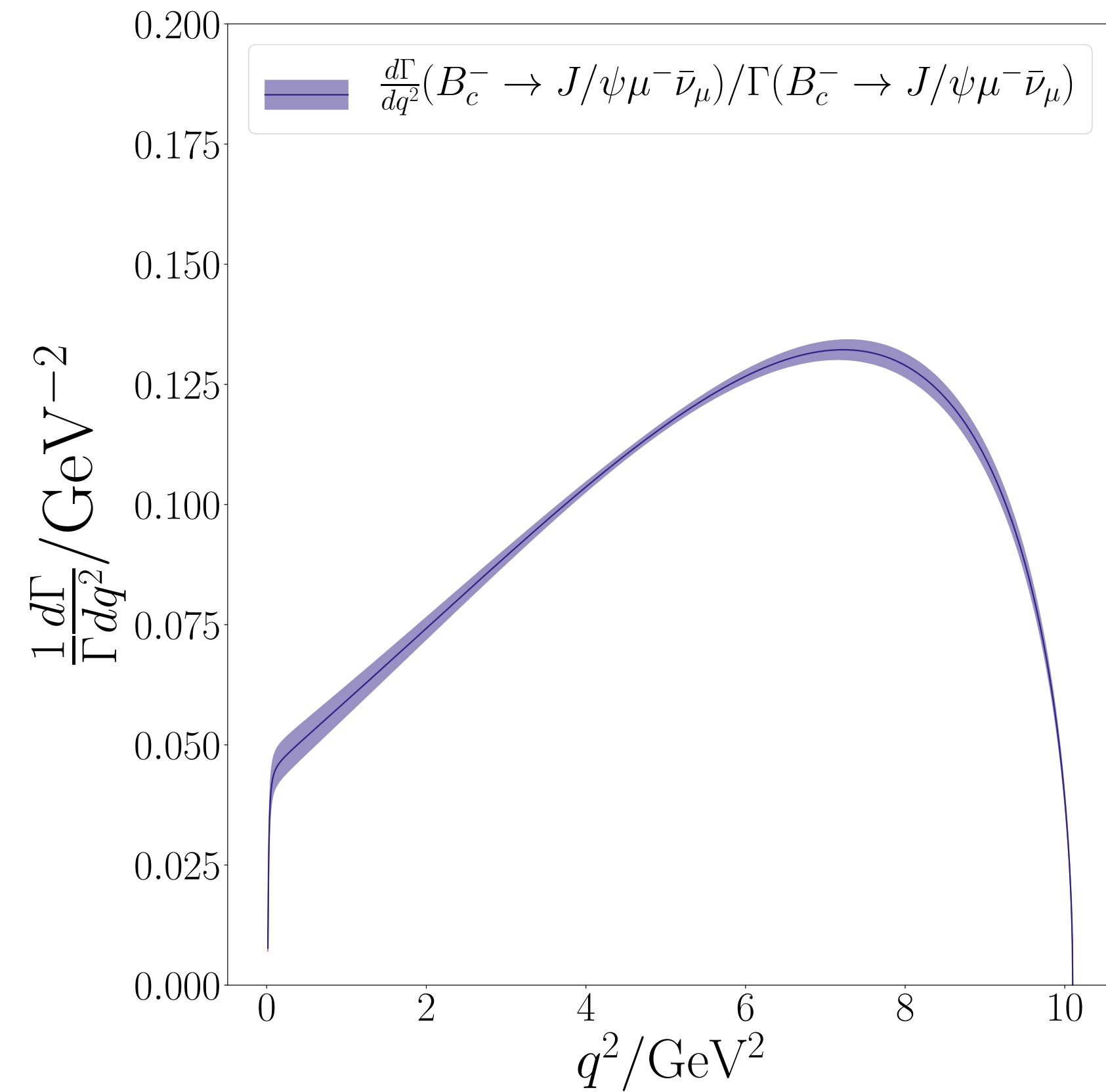
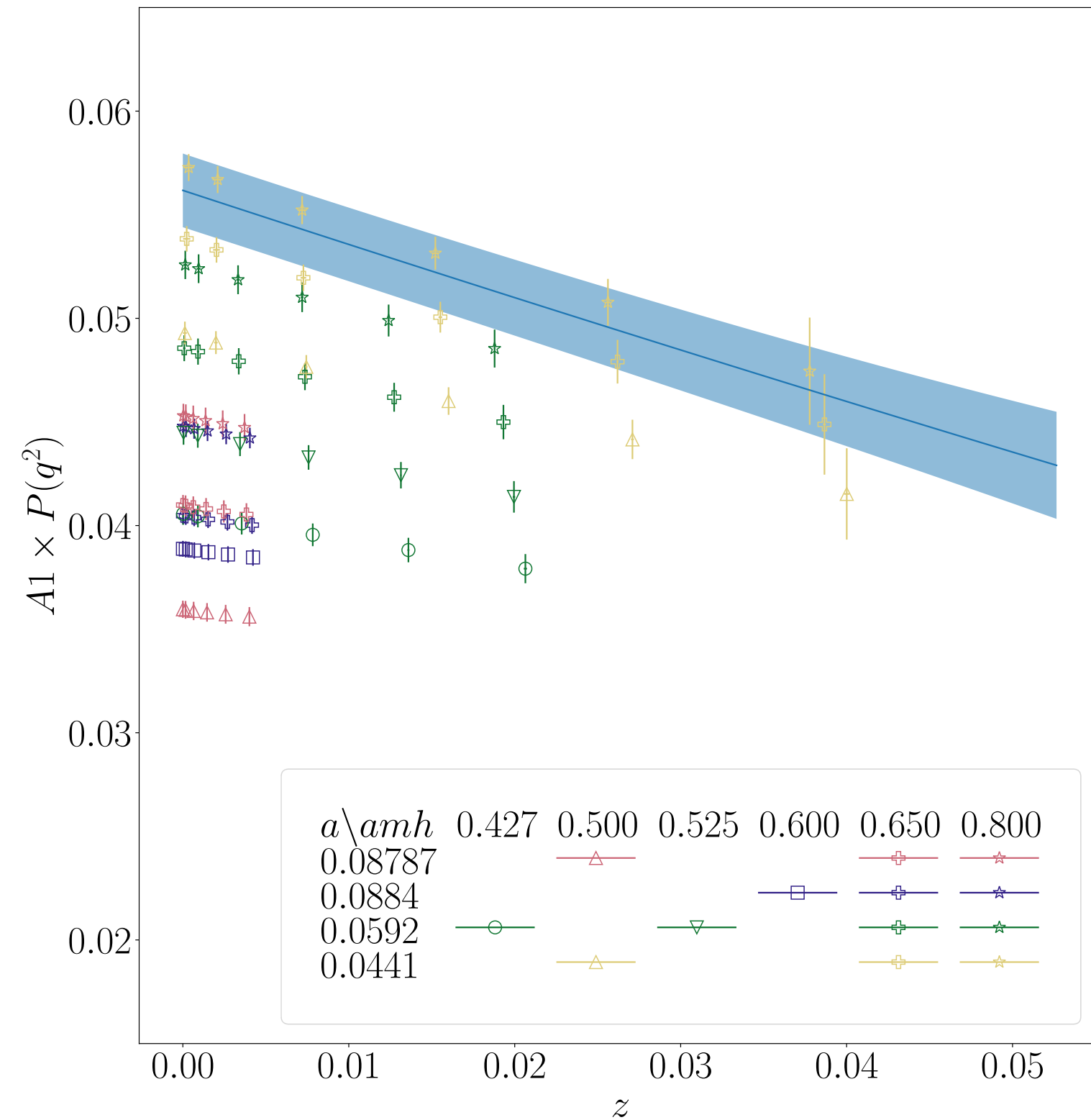


FLAG2021: Fit to $B_s \rightarrow K$ form factor data has $\chi^2/\text{d.o.f.} = 1.54$

B_c semileptonic decays: $b \rightarrow c$

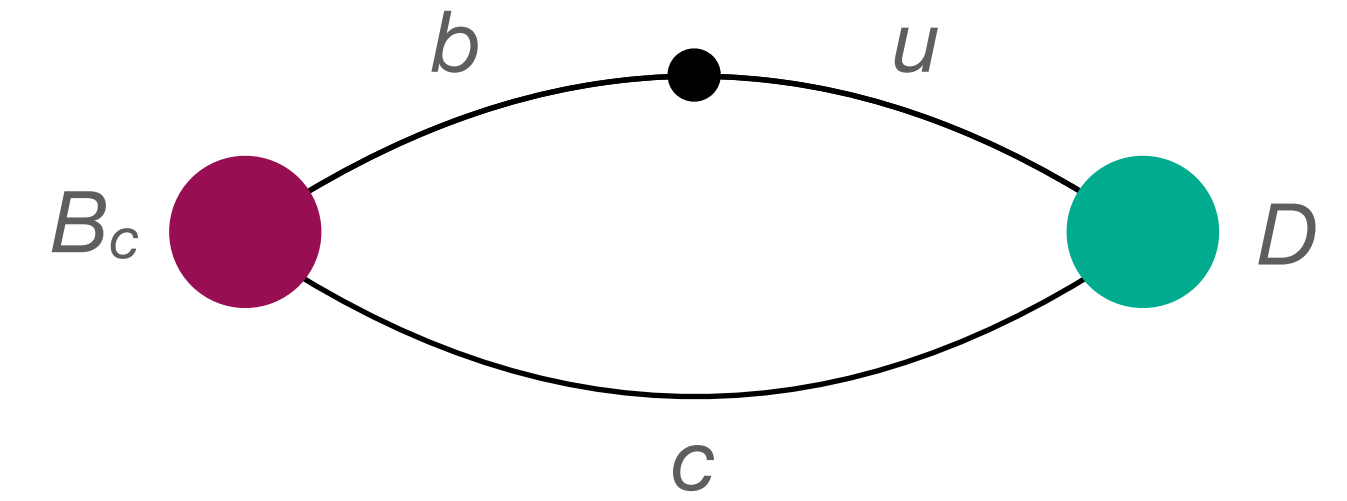


HPQCD: Heavy HISQ on MILC's $n_f = 2+1+1$ HISQ lattices

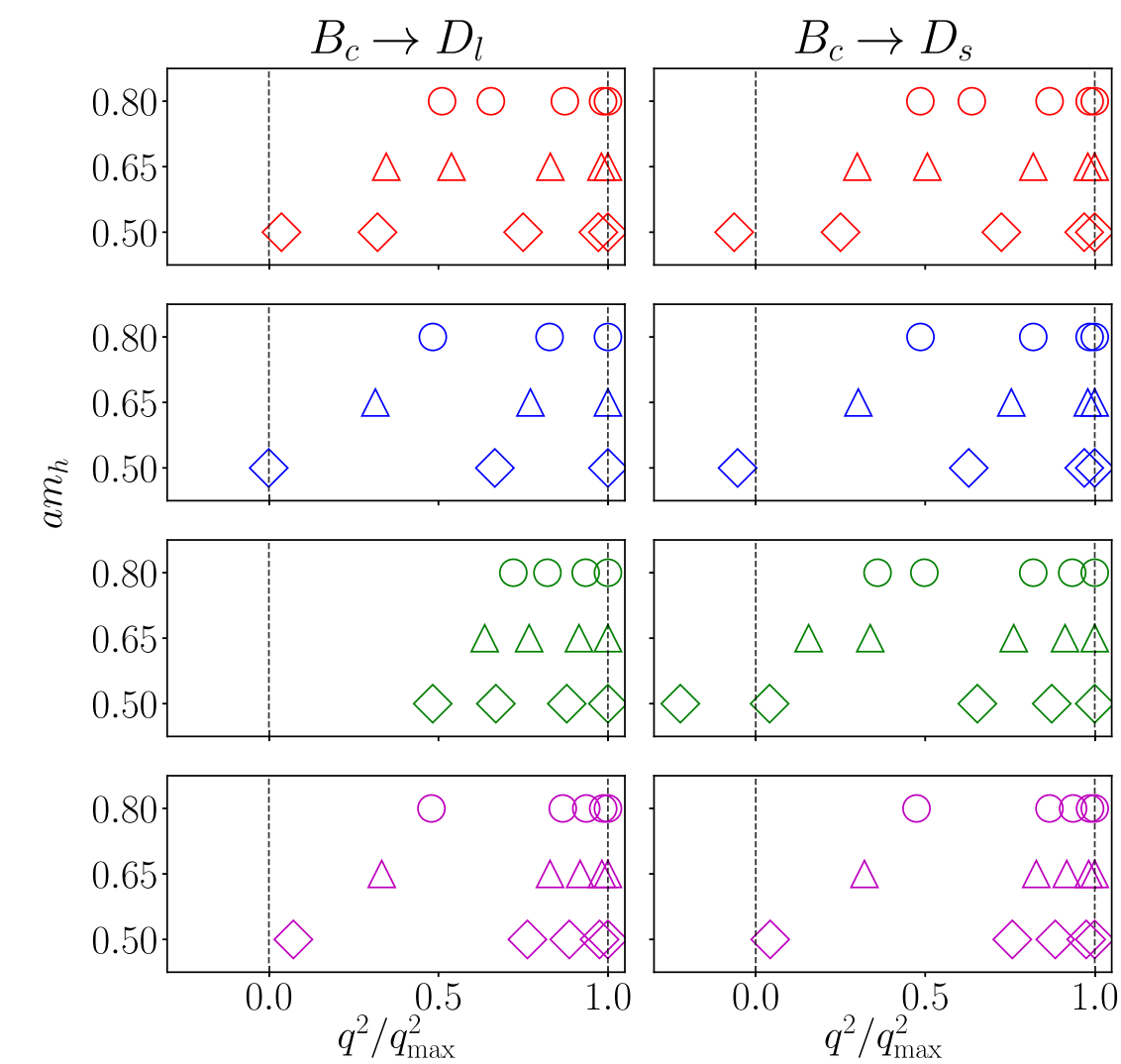
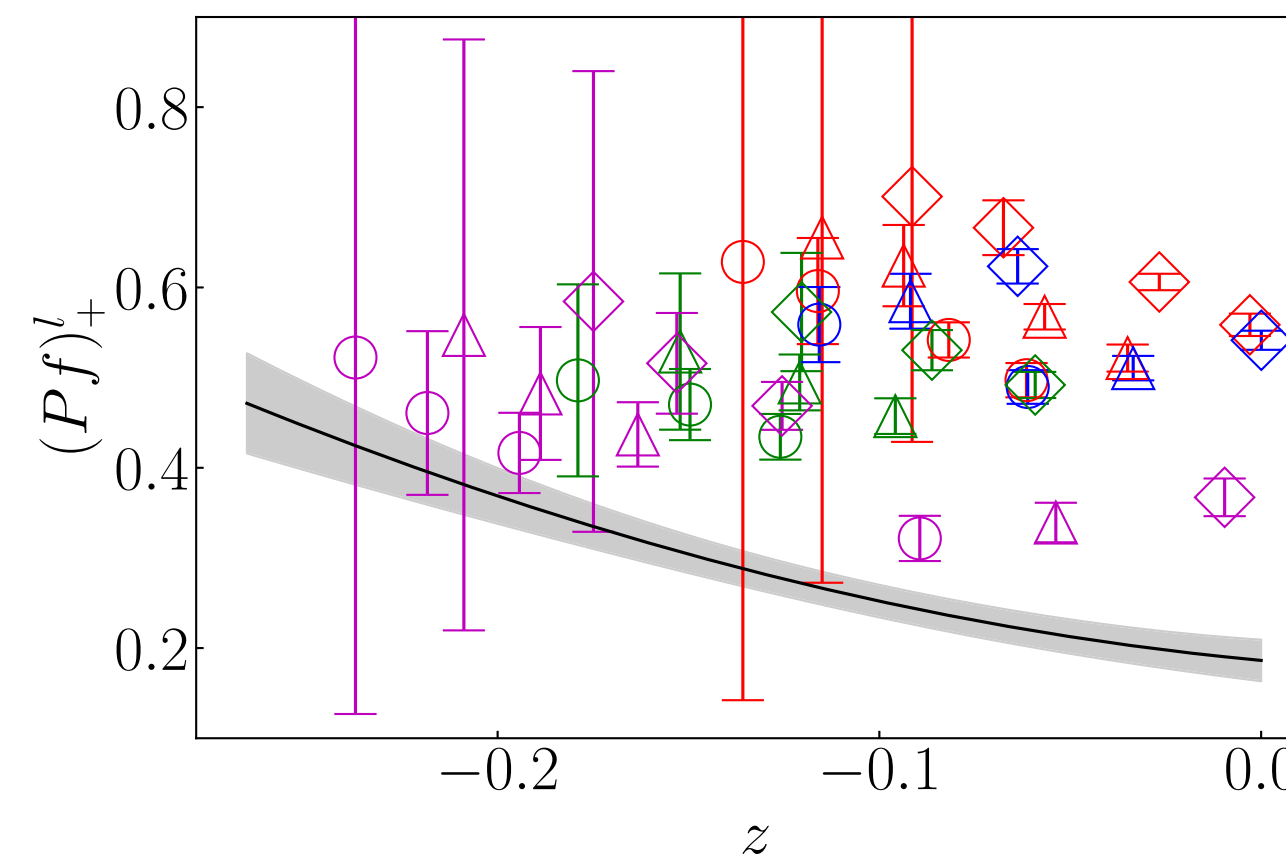
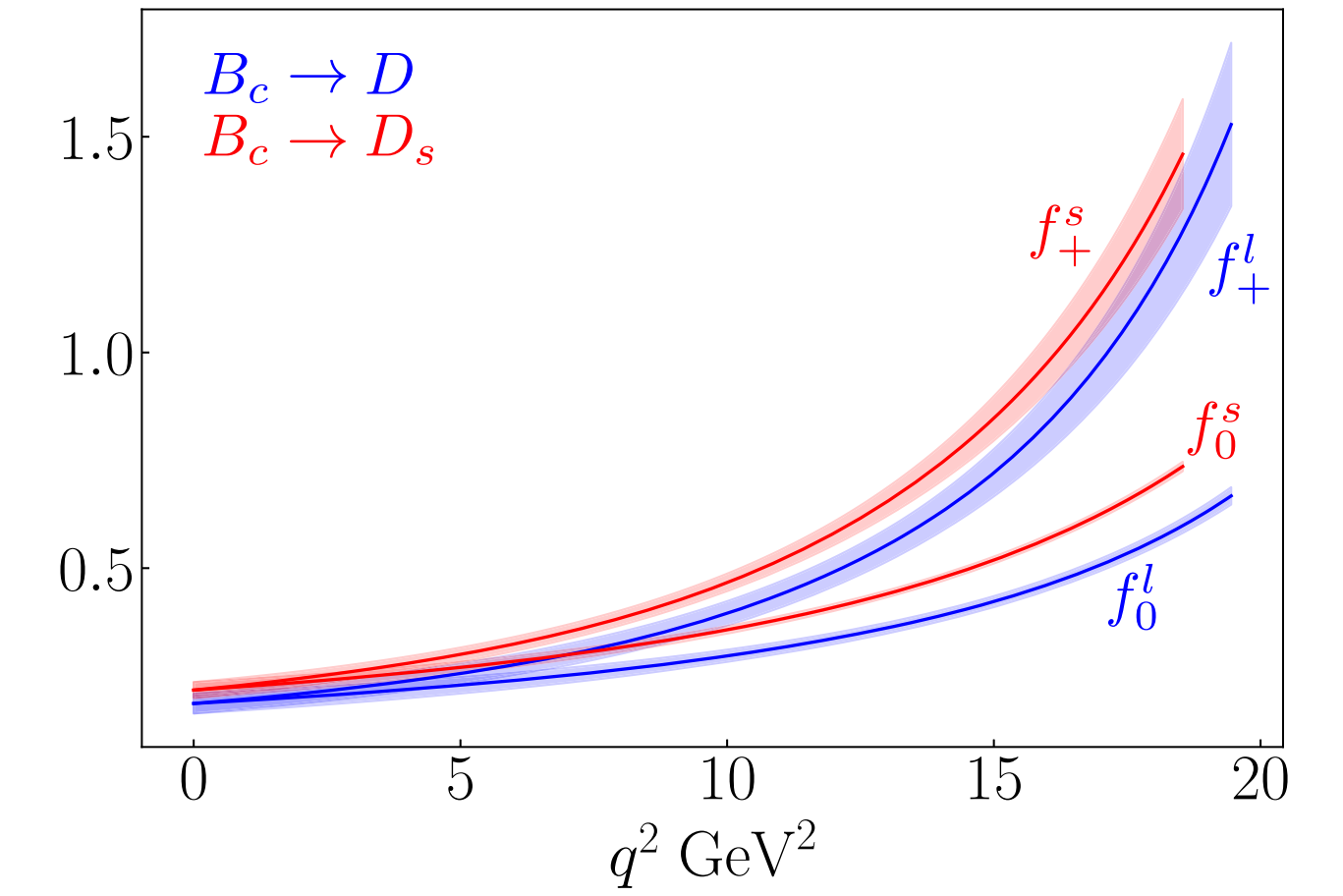
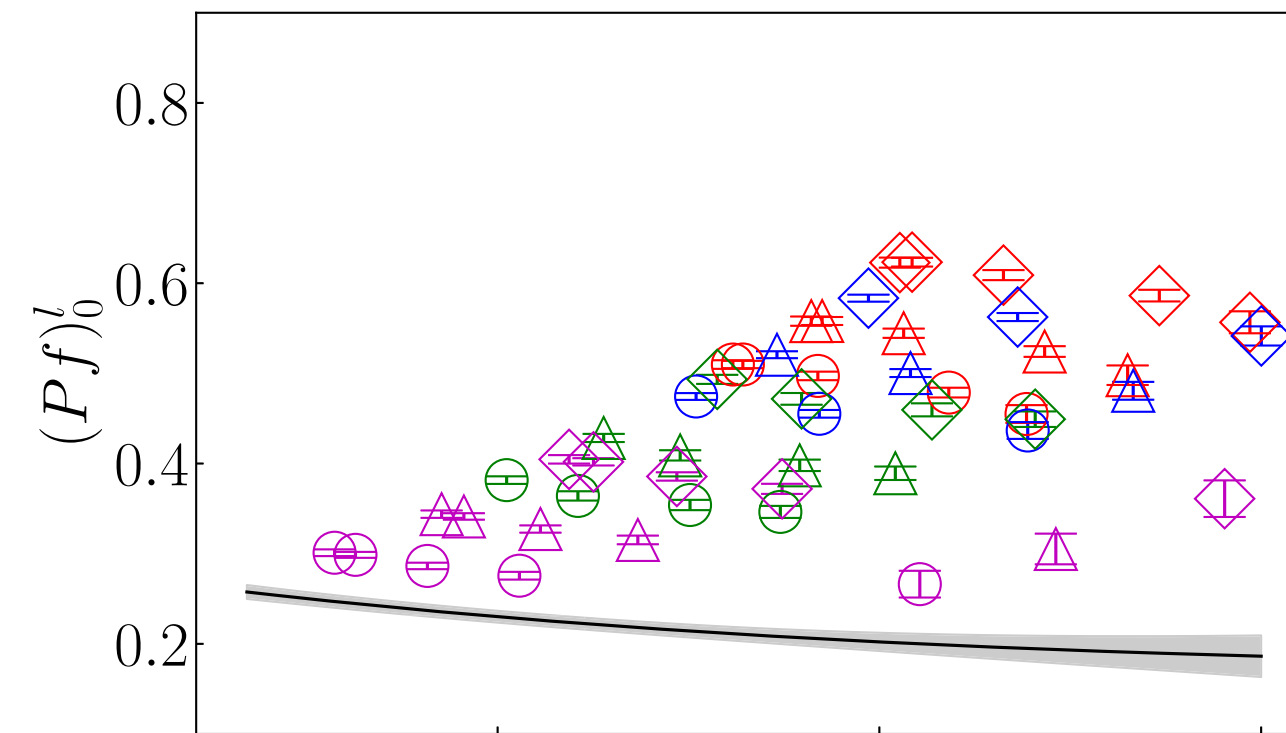
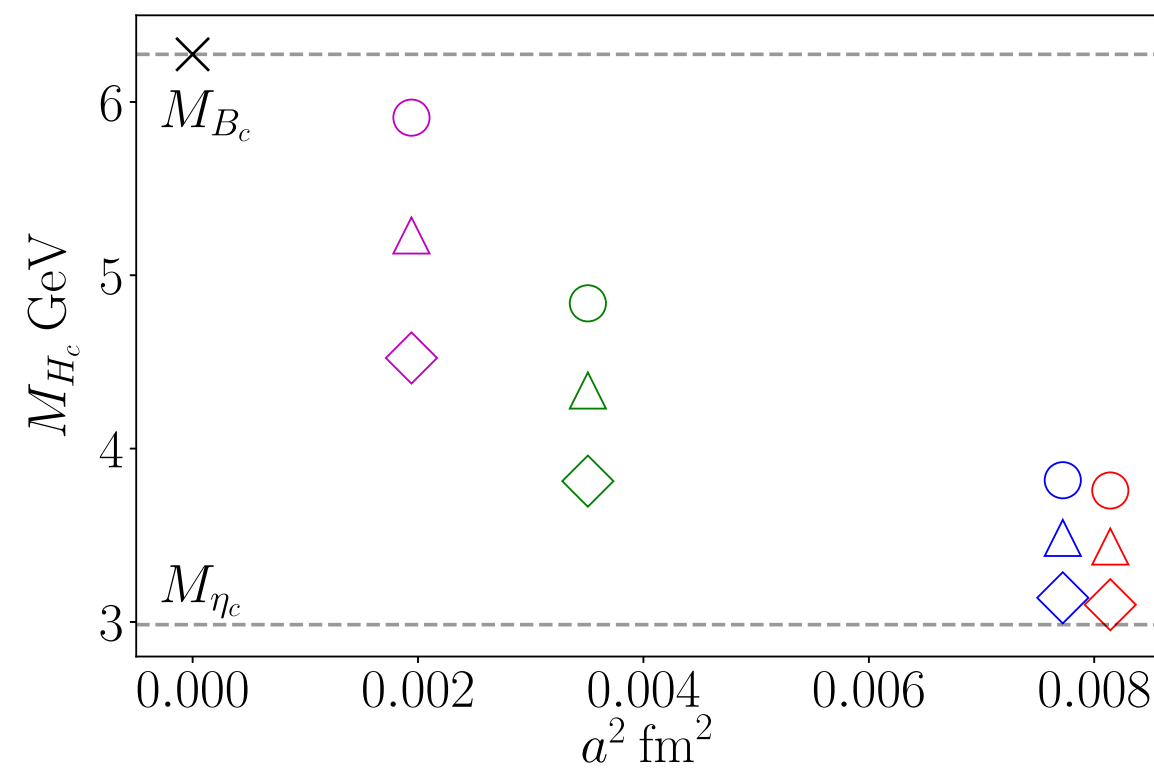


$$R(J/\psi)_{\text{SM}} = 0.2601(36)$$

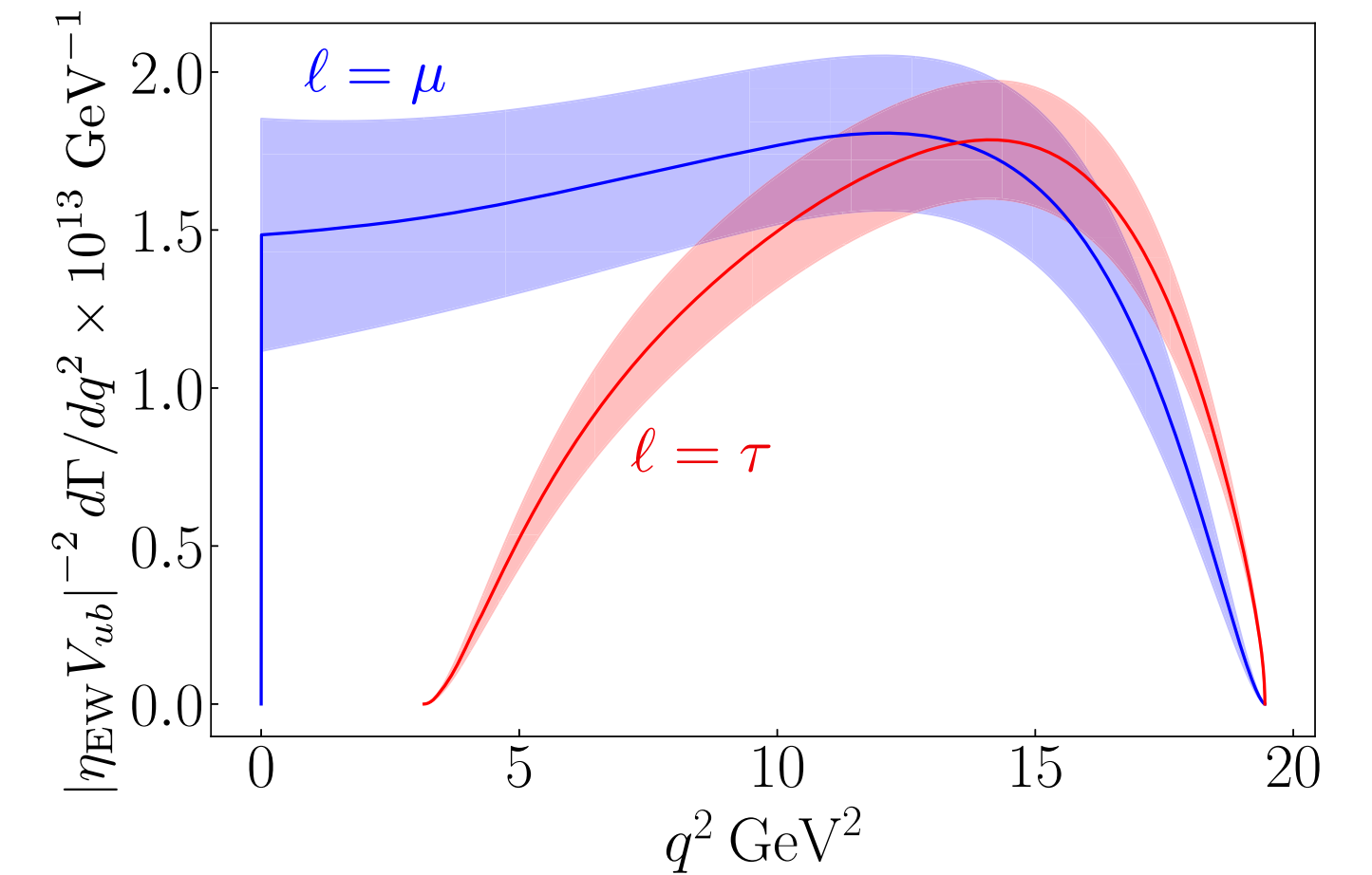
B_c semileptonic decays: $b \rightarrow u$



HPQCD: Heavy HISQ on MILC's $n_f = 2+1+1$ HISQ lattices



LHCb is working on an analysis of $\mathcal{B}(B_c \rightarrow D^{(*)} \ell \nu) / \mathcal{B}(B_c \rightarrow J/\psi \ell \nu)$



Inclusive decays

- Gambino and Hashimoto ([PRL 125, 032001 \(2020\)](#)) proposal to determine inclusive decay rates from lattice QCD

- Want

$$W^{\mu\nu}(p, q) = \sum_{X_c} (2\pi)^3 \delta^{(4)}(p - q - r) \frac{1}{2E_{B_s}} \langle \bar{B}_s(\mathbf{p}) | J^{\mu\dagger} | X_c(\mathbf{r}) \rangle \langle X_c(\mathbf{r}) | J^\nu | \bar{B}_s(\mathbf{p}) \rangle$$

- Can compute

$$C_{\mu\nu}^{JJ}(t; \mathbf{q}) = \sum_{\mathbf{x}} \frac{e^{i\mathbf{q}\cdot\mathbf{x}}}{2m_{B_s}} \langle \bar{B}_s(\mathbf{0}) | J_\mu^\dagger(\mathbf{x}, t) J_\nu(\mathbf{0}, 0) | \bar{B}_s(\mathbf{0}) \rangle$$

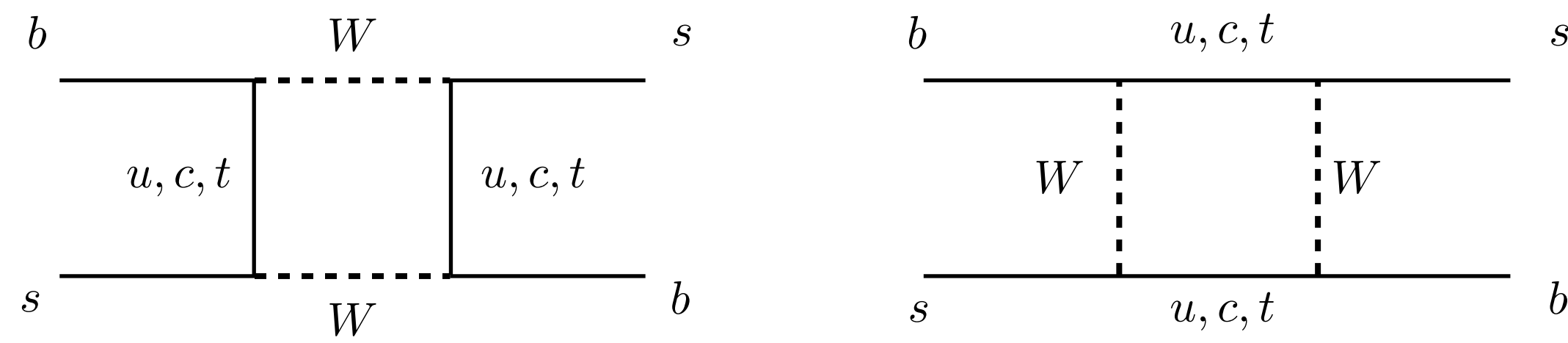
$$\begin{aligned} M_{\mu\nu}(t; \mathbf{q}) &= \int d^3x \frac{e^{i\mathbf{q}\cdot\mathbf{x}}}{2m_B} \langle \bar{B}(\mathbf{0}) | J_\mu^\dagger(\mathbf{0}, 0) e^{-t\hat{H} + i\hat{\mathbf{P}}\cdot\mathbf{x}} J_\nu(\mathbf{0}, 0) | \bar{B}(\mathbf{0}) \rangle \\ &= \frac{(2\pi)^3}{2m_B} \langle \bar{B}(\mathbf{0}) | J_\mu^\dagger(\mathbf{0}, 0) e^{-t\hat{H}} \delta^3(\hat{\mathbf{P}} + \mathbf{q}) J_\nu(\mathbf{0}, 0) | \bar{B}(\mathbf{0}) \rangle \\ &= \int_0^\infty d\omega W_{\mu\nu}(\omega, \mathbf{q}) e^{-\omega t} . \end{aligned}$$

- Can be related to integral over W
- Progress reported at Lattice 2022 by A. Smecca, A. Barone, R. Kellermann

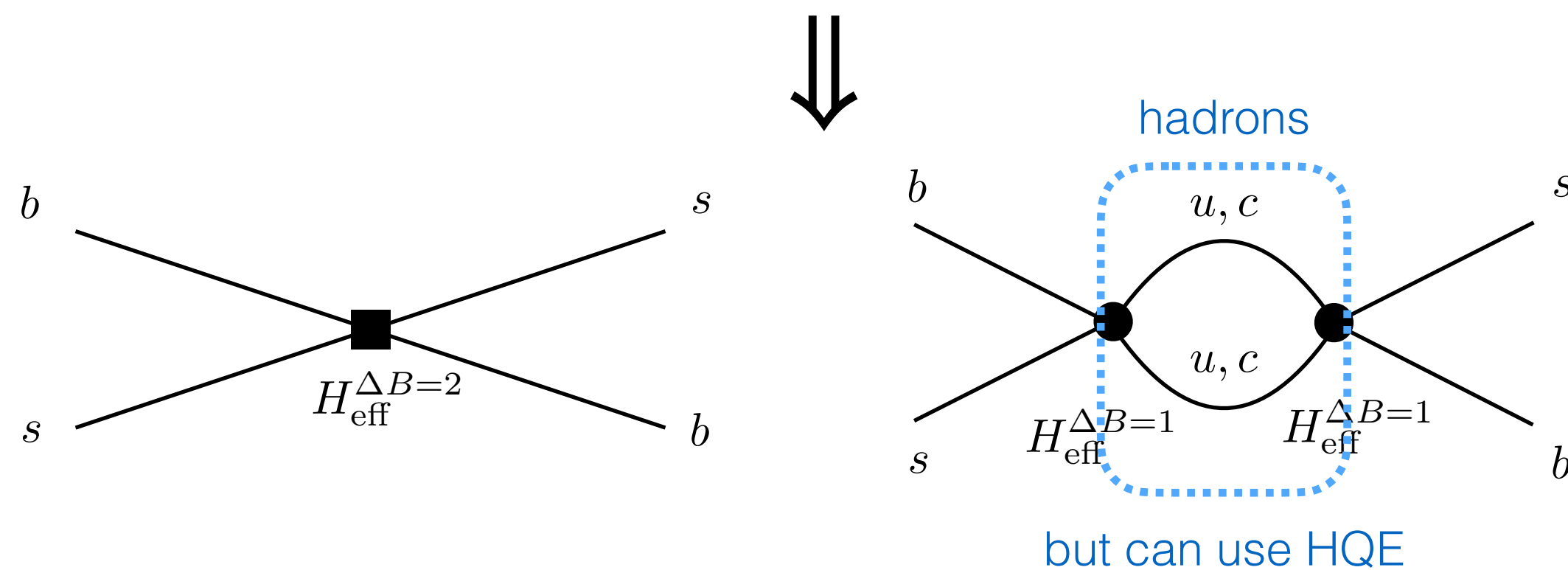
CKM: 3rd row & B mixing

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

B-mixing in the Standard Model



Integrate out W and t



$$\Delta M_s = \frac{1}{2m_{B_s}} \langle \bar{B}_s | H_{\text{eff}}^{\Delta B=2} | B_s \rangle$$

$$H_{\text{eff}}^{\Delta B=2} = \frac{G_F^2 m_W^2}{4\pi^2} (V_{ts}^* V_{tb})^2 \sum_{i=1}^5 C_i Q_i$$

$$Q_1 = (\bar{b}^\alpha \gamma^\mu (1 - \gamma^5) s^\alpha) (\bar{b}^\beta \gamma_\mu (1 - \gamma^5) s^\beta)$$

$$Q_2 = (\bar{b}^\alpha (1 - \gamma^5) s^\alpha) (\bar{b}^\beta (1 - \gamma^5) s^\beta)$$

$$Q_3 = (\bar{b}^\alpha (1 - \gamma^5) s^\beta) (\bar{b}^\beta (1 - \gamma^5) s^\alpha)$$

$$Q_4 = (\bar{b}^\alpha (1 - \gamma^5) s^\alpha) (\bar{b}^\beta (1 + \gamma^5) s^\beta)$$

$$Q_5 = (\bar{b}^\alpha (1 - \gamma^5) s^\beta) (\bar{b}^\beta (1 + \gamma^5) s^\alpha)$$

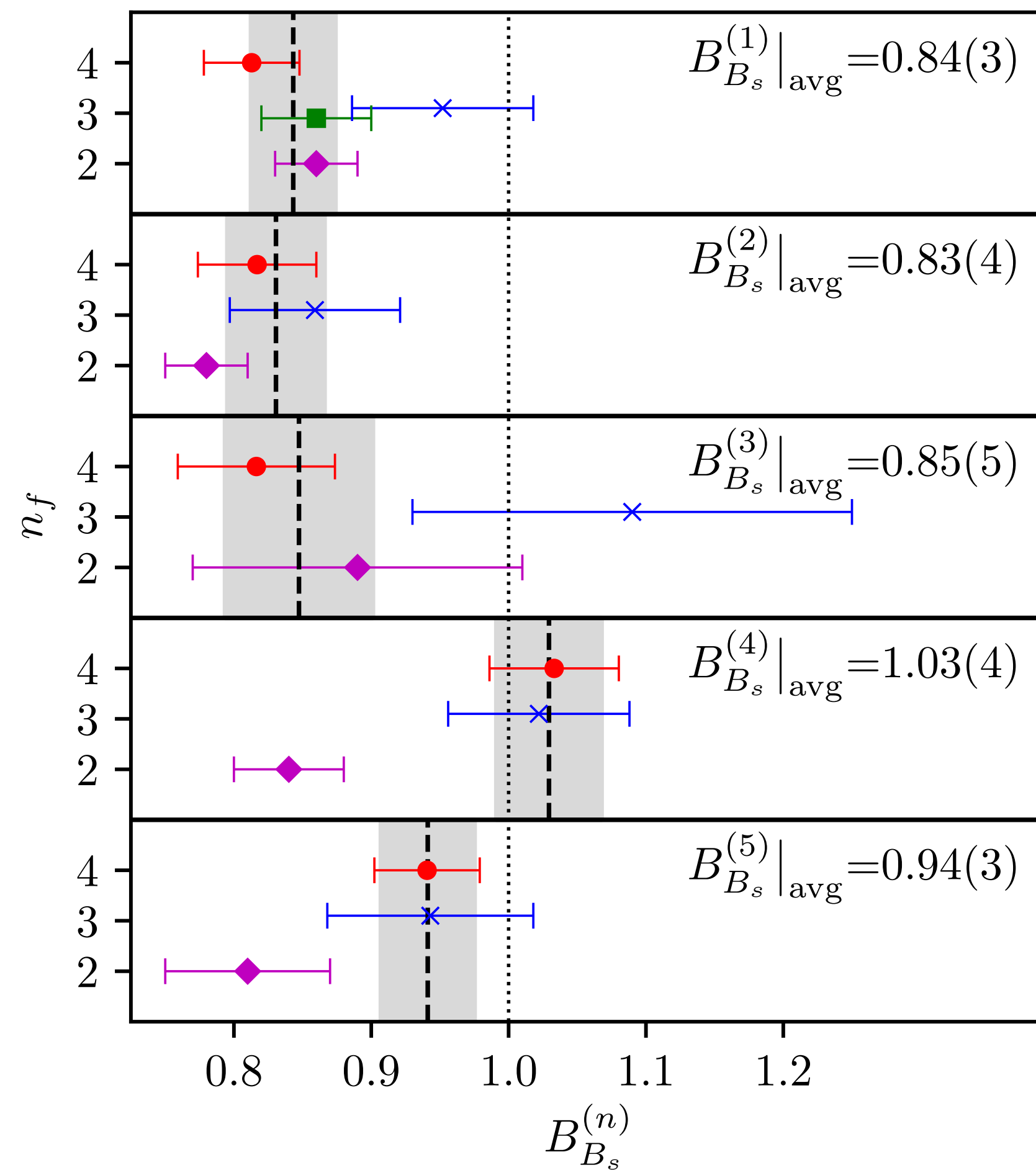
In the SM, only Q_1 enters ΔM_s

B mixing: recent LQCD results

- ETMC (Carrasco et al., [JHEP 03 \(2014\) 016](#)) — ratio method (incl. static limit), twisted mass valence and sea ($n_f = 2$), RI-MOM
- Fermilab/MILC (Bazavov et al., [PRD 93, 113016 \(2016\)](#)) — Fermilab b , AsqTad valence and sea ($n_f = 2+1$), mostly non-perturbative
- RBC/UKQCD (Boyle et al., [arXiv:1812.08791](#)) — domain wall valence (extrapolating to the b) and sea ($n_f = 2+1$), ratios only
- HPQCD (Dowdall et al., [PRD 100, 094508 \(2019\)](#)) — NRQCD b , HISQ valence and sea ($n_f = 2+1+1$, MILC), perturbative
- Results presented for “bag factors”

$$B_{B_q}^{(i)}(\mu) \equiv \frac{\langle B_q | O_i^q | \bar{B}_q \rangle_{\overline{\text{MS}}}^{(\mu)}}{\eta_i^q(\mu) f_{B_q}^2 M_{B_q}^2},$$

Comparisons — bag factors

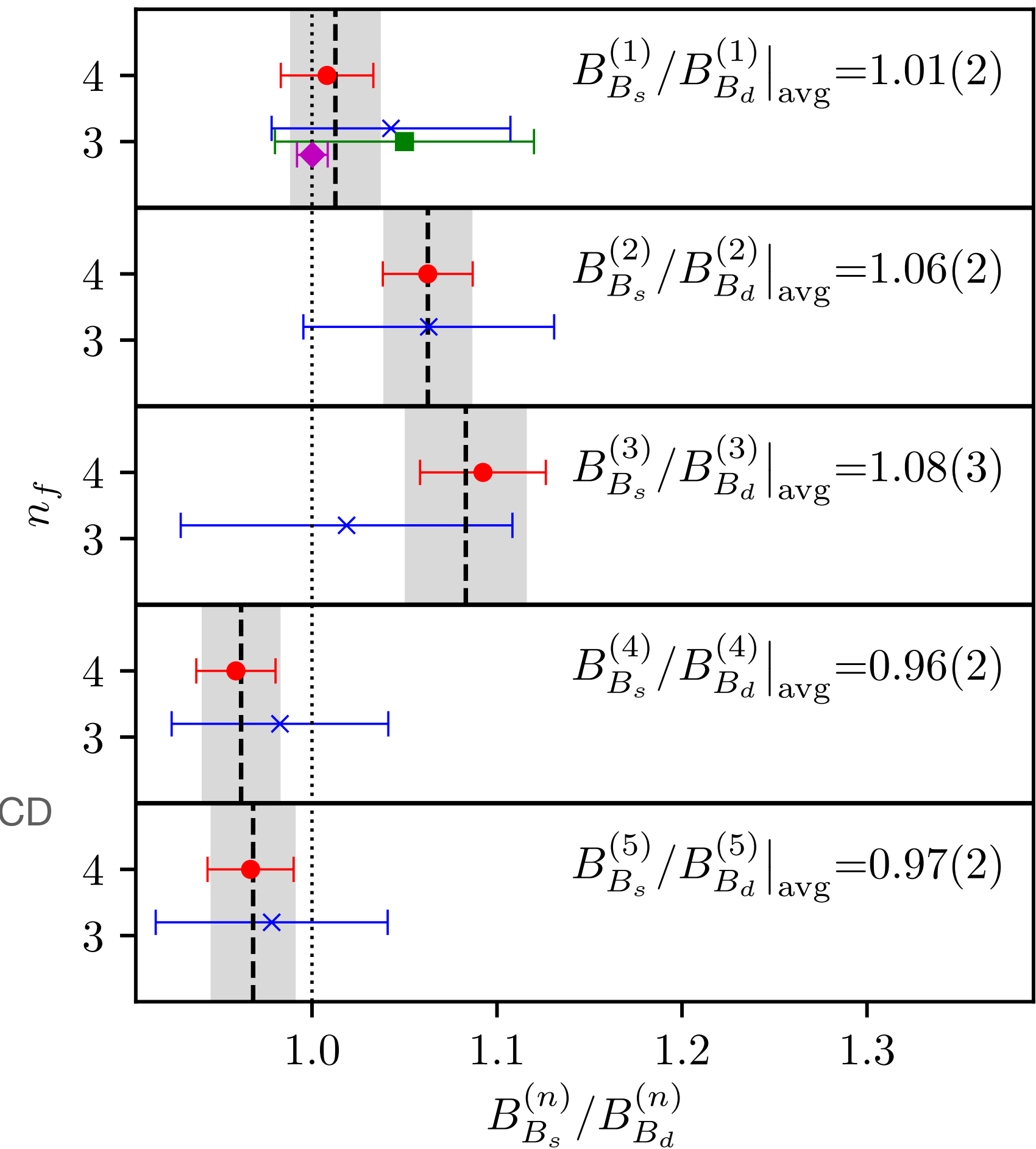


Legend

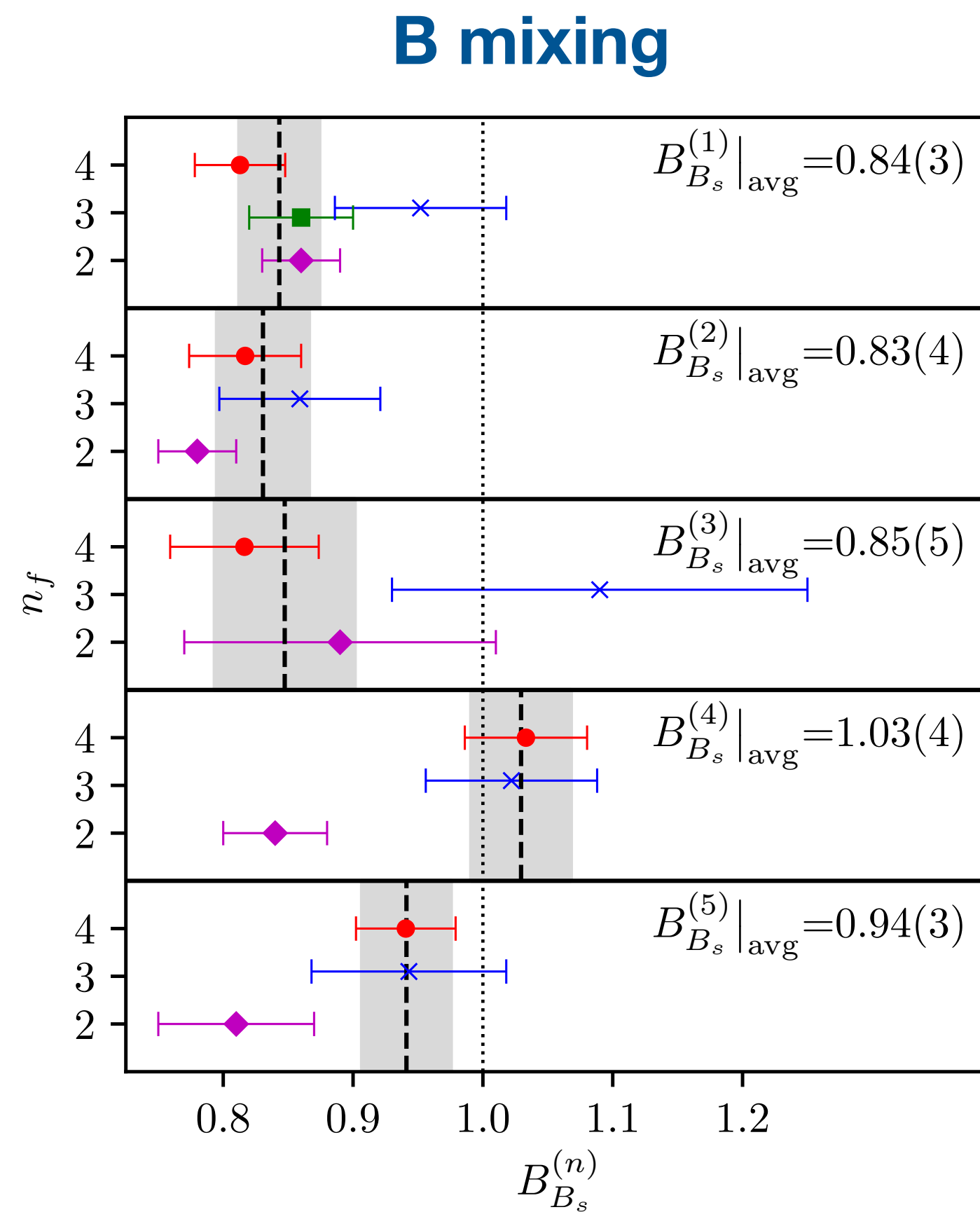
- Red circle: HPQCD
- Blue cross: Fermilab/MILC
- Green square: HPQCD
- Purple diamond: ETMC

Legend

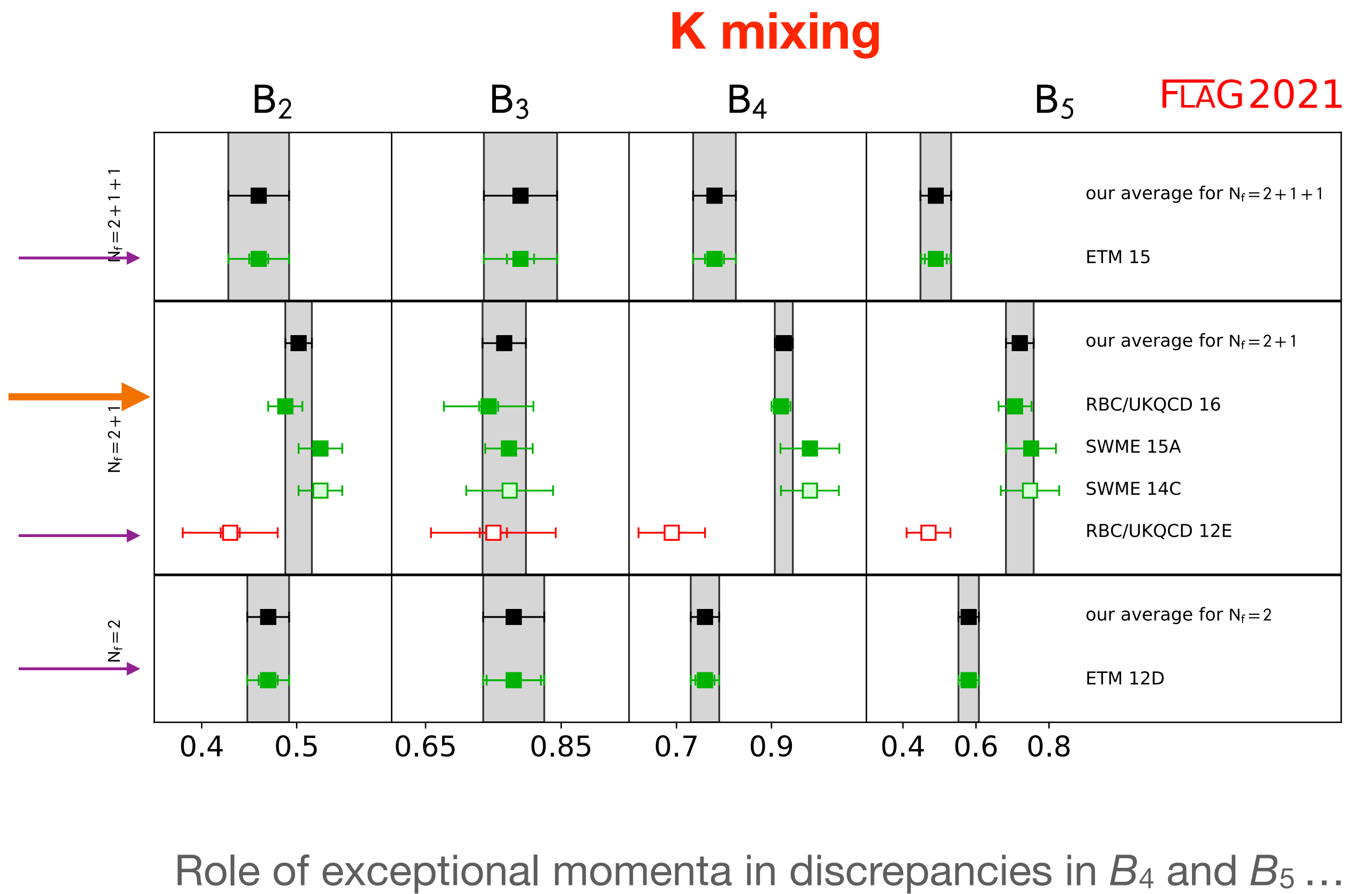
- Red circle: HPQCD
- Blue cross: Fermilab/MILC
- Green square: HPQCD
- Purple diamond: RBC-UKQCD



Similarities between K and B mixing



← Purple: RI-MOM
 RI-SMOM →
 Rest: perturbative



Operator matching

RI-MOM schemes

Calculate (gauge-fixed) quark matrix elements nonperturbatively
 Define renormalization constants by imposing conditions at large momentum. The same conditions can be applied in continuum perturbation theory.

$$O(\mu) = Z_O(\mu a, g(a)) O(a).$$

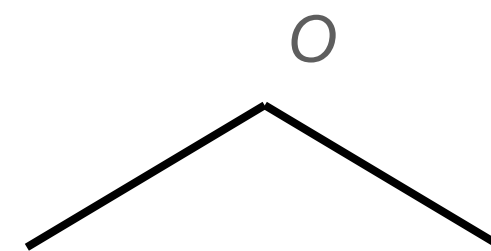
Z_O is found, by imposing the renormalization condition

$$Z_O(\mu a, g(a)) Z_\psi^{-1}(\mu a, g(a)) \Gamma_O(pa) |_{p^2=\mu^2} = 1,$$

$$\Gamma_O(pa) = \frac{1}{12} \text{tr} (\Lambda_O(pa) \hat{P}_O),$$

where

$$\Lambda_O(pa) = S(pa)^{-1} G_O(pa) S(pa)^{-1}.$$



Must take care not to allow large contributions from low energy

(i) symmetric or nonexceptional momentum configuration:

RI-SMOM $p_1^2 = p_2^2 = q^2 = -\mu^2, \quad \mu^2 > 0,$
 $q = p_1 - p_2;$

(ii) asymmetric or exceptional momentum configuration:

RI-MOM $p_1^2 = p_2^2 = -\mu^2, \quad \mu^2 > 0,$
 $p_1 = p_2, \quad q = 0,$

STURM, AOKI, CHRIST, IZUBUCHI, SACHRAJDA, AND SONI

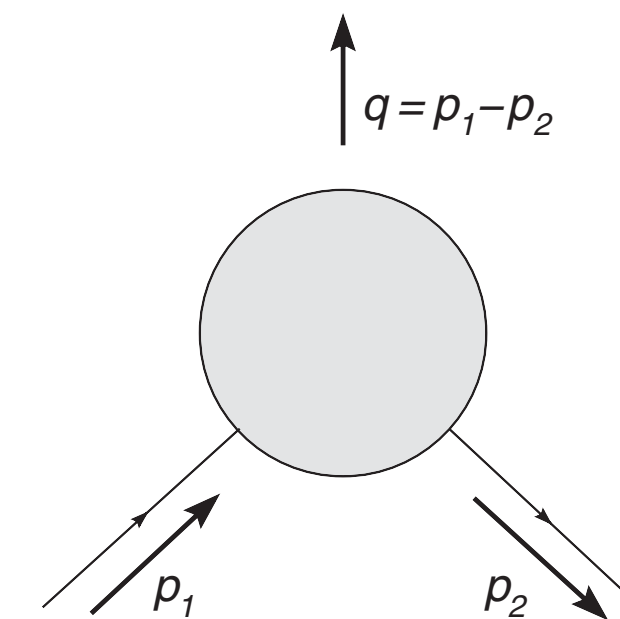
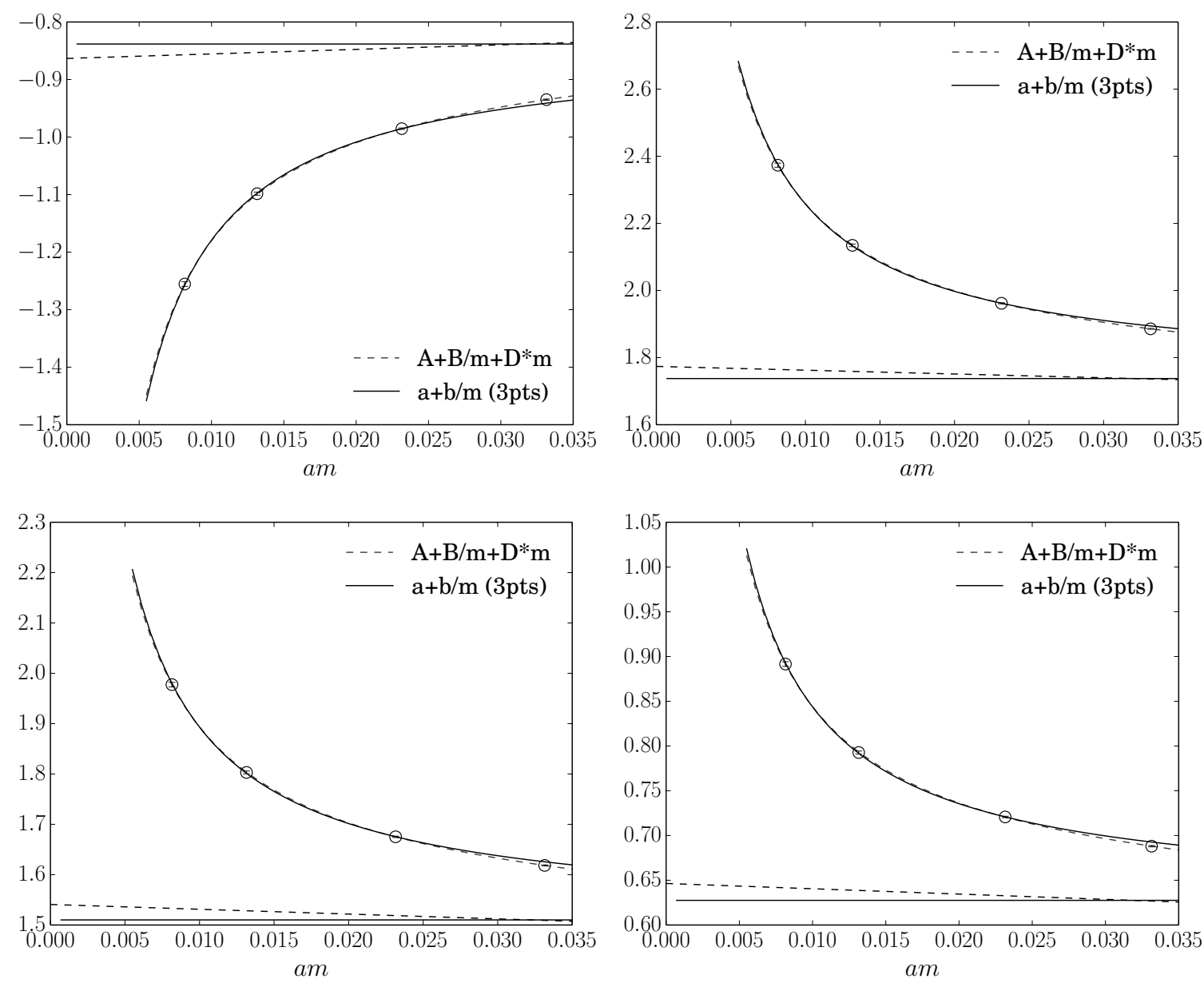


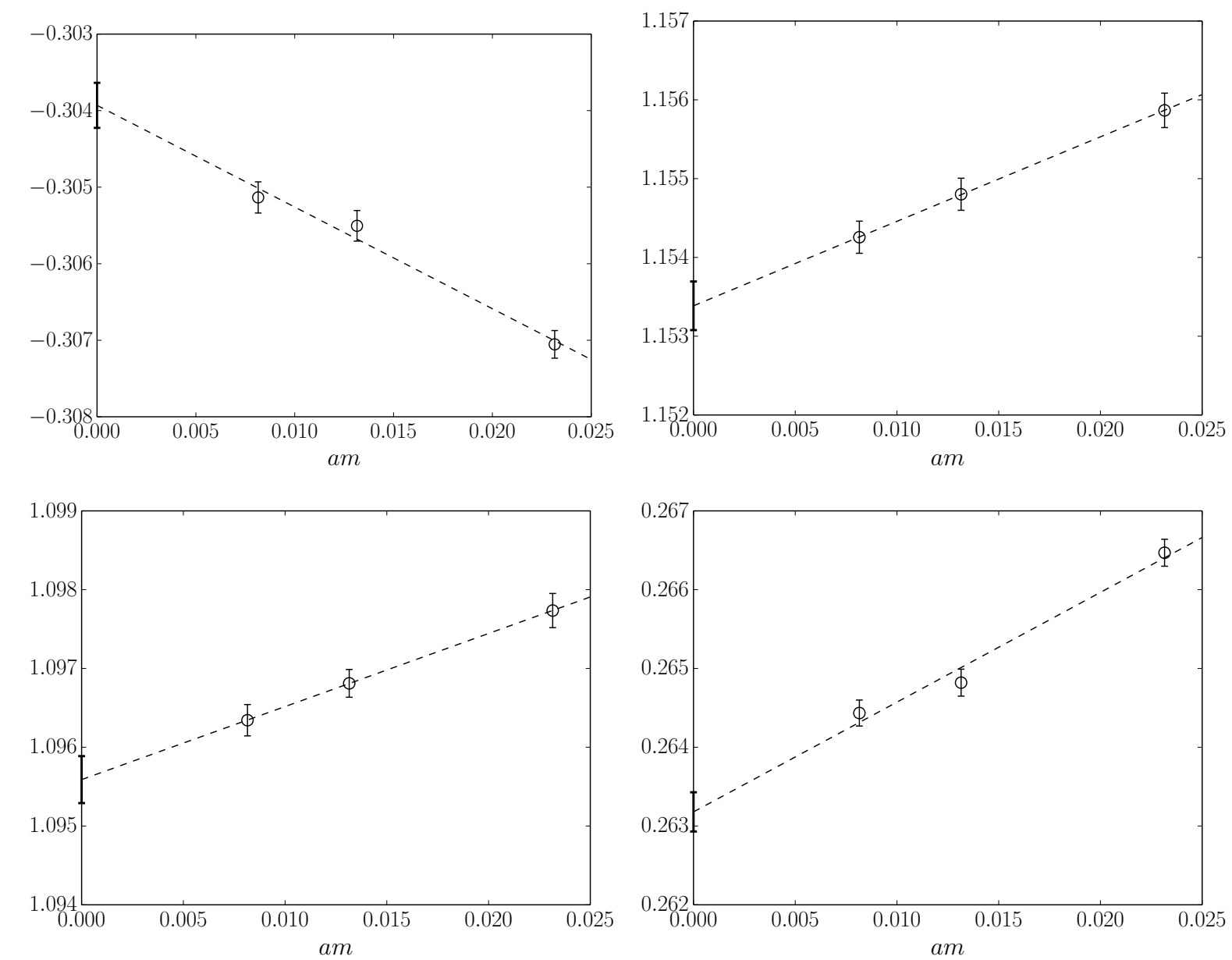
FIG. 1. Momentum flow of a generic diagram required for the renormalization procedure with **nonexceptional momenta**. The gray bubble stands for an operator insertion and higher order corrections.

Avoiding exceptional momenta

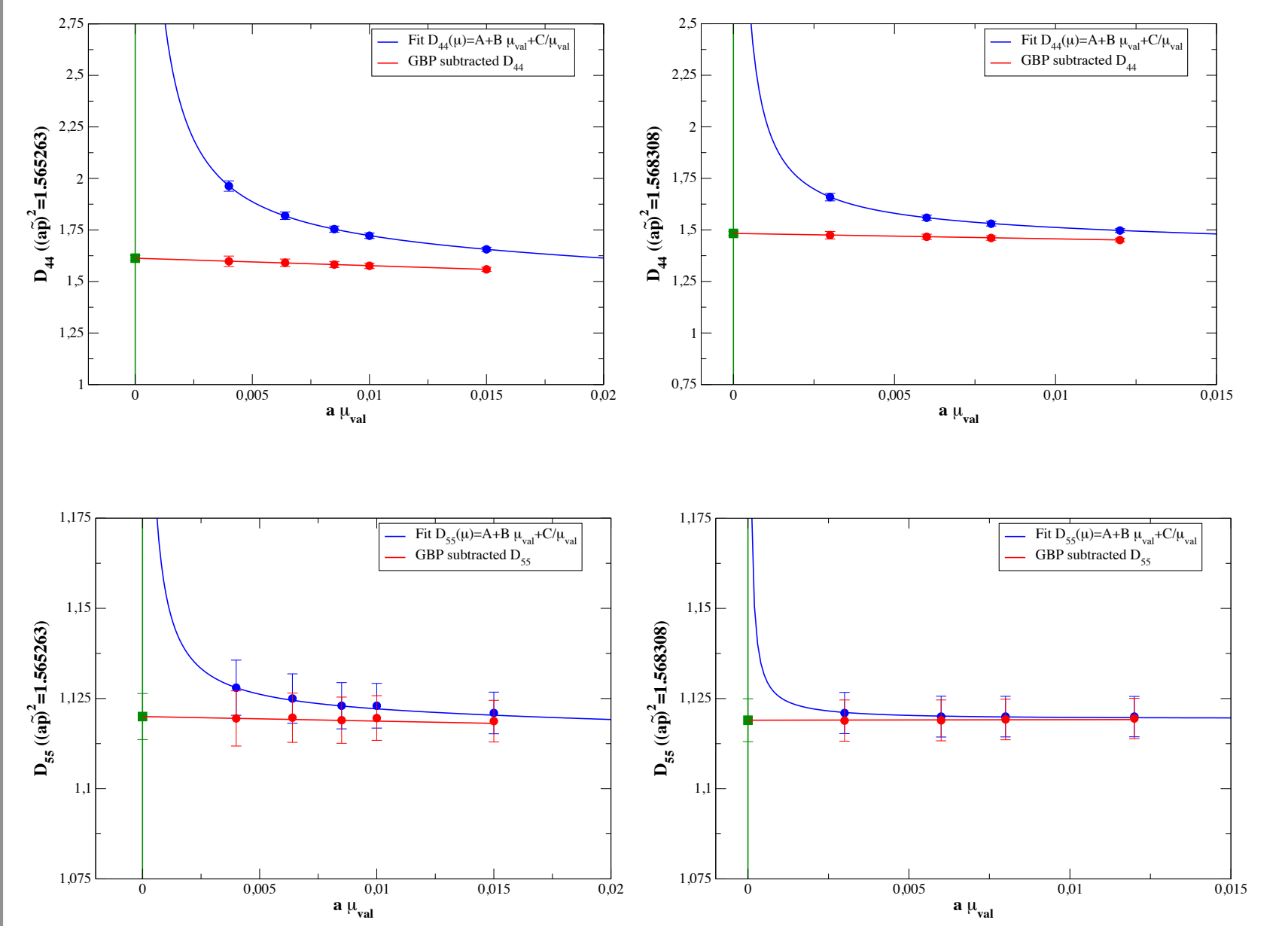
RI-MOM



RI-SMOM



RI-MOM

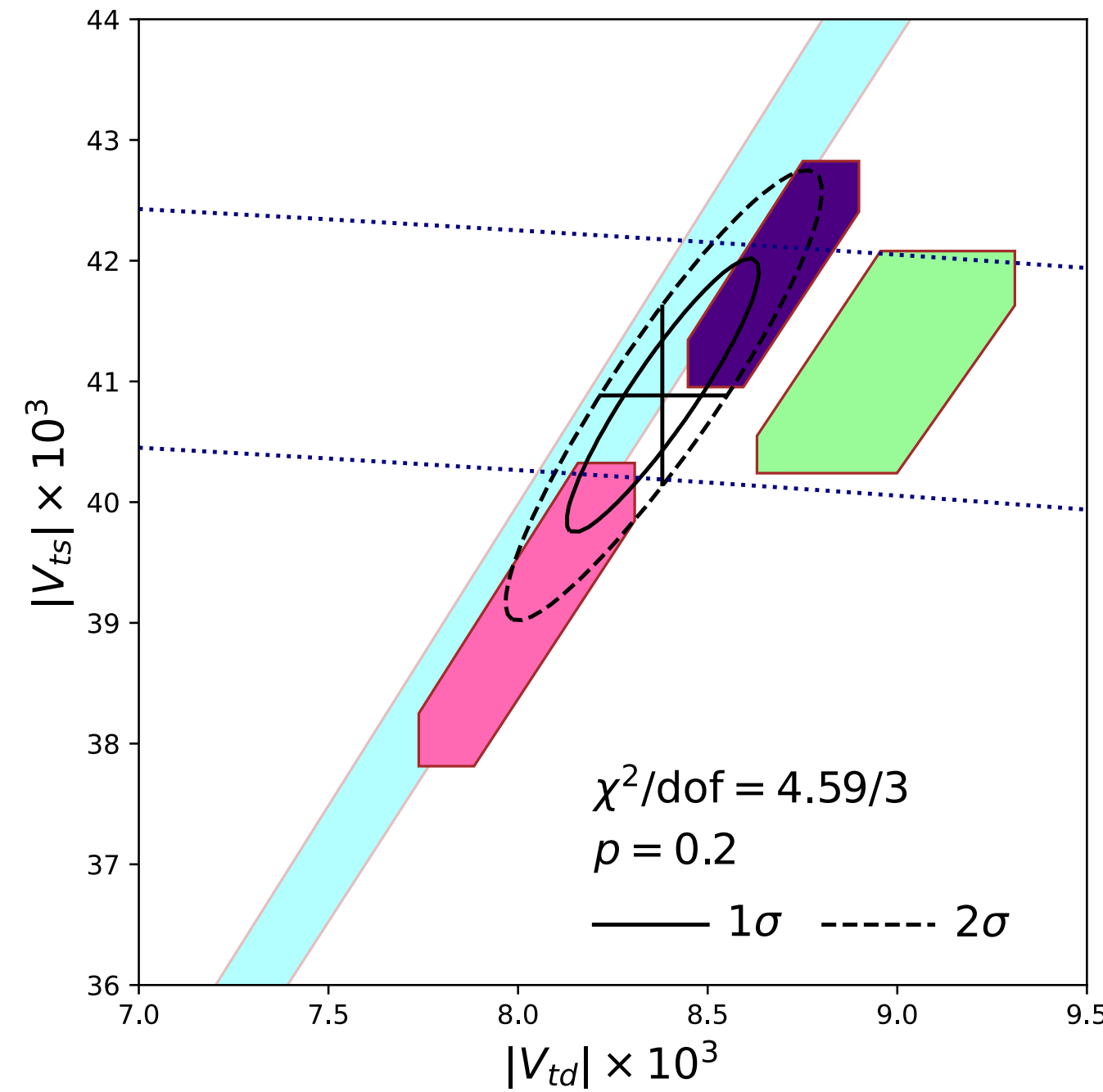
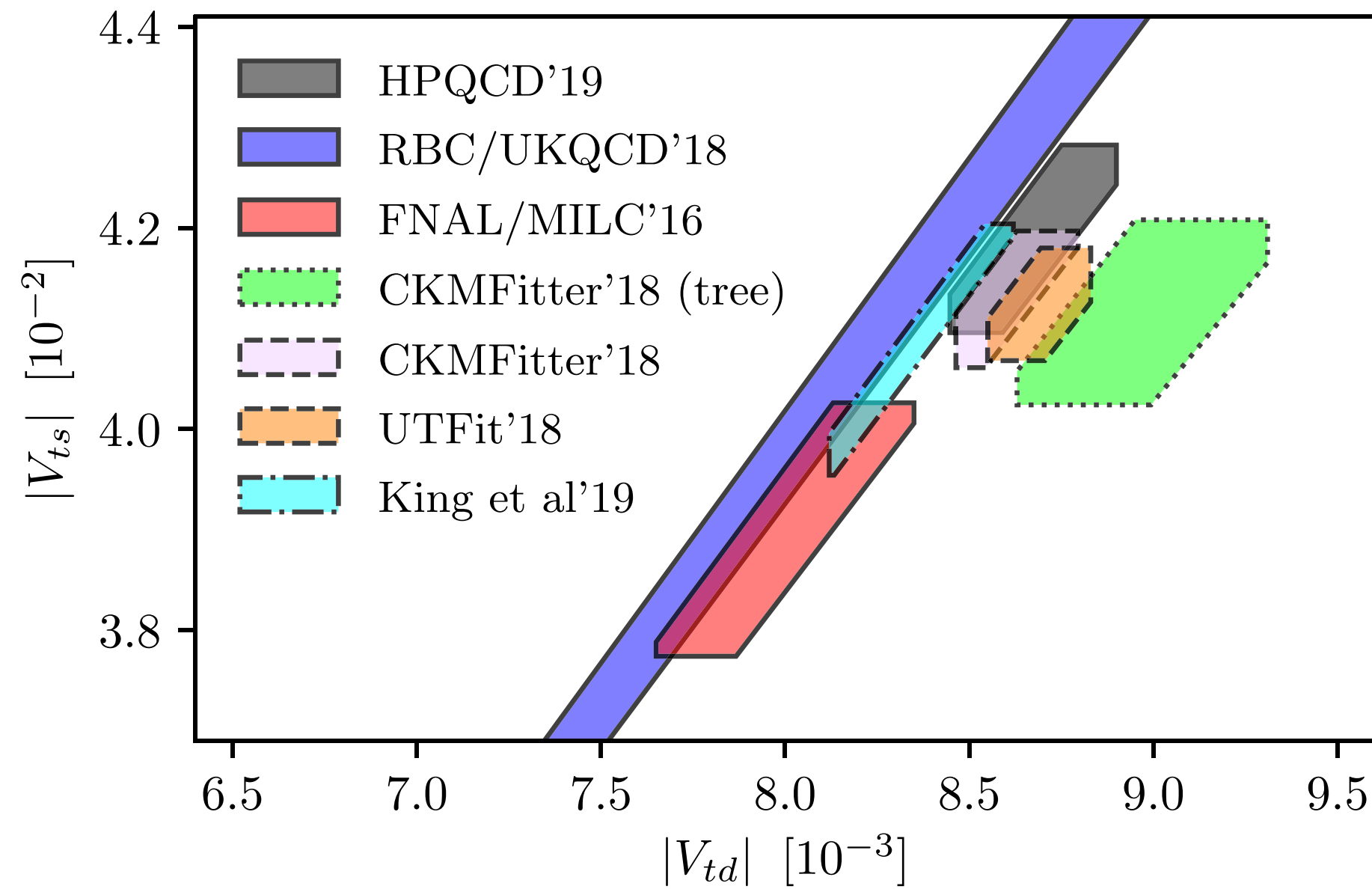


Figs. 3 & 4 from Boyle, Garron, Hudspith, Lehner, Lytle, JHEP 10 (2017) 054

Bertone et al. (ETM), JHEP 03 (2013) 089

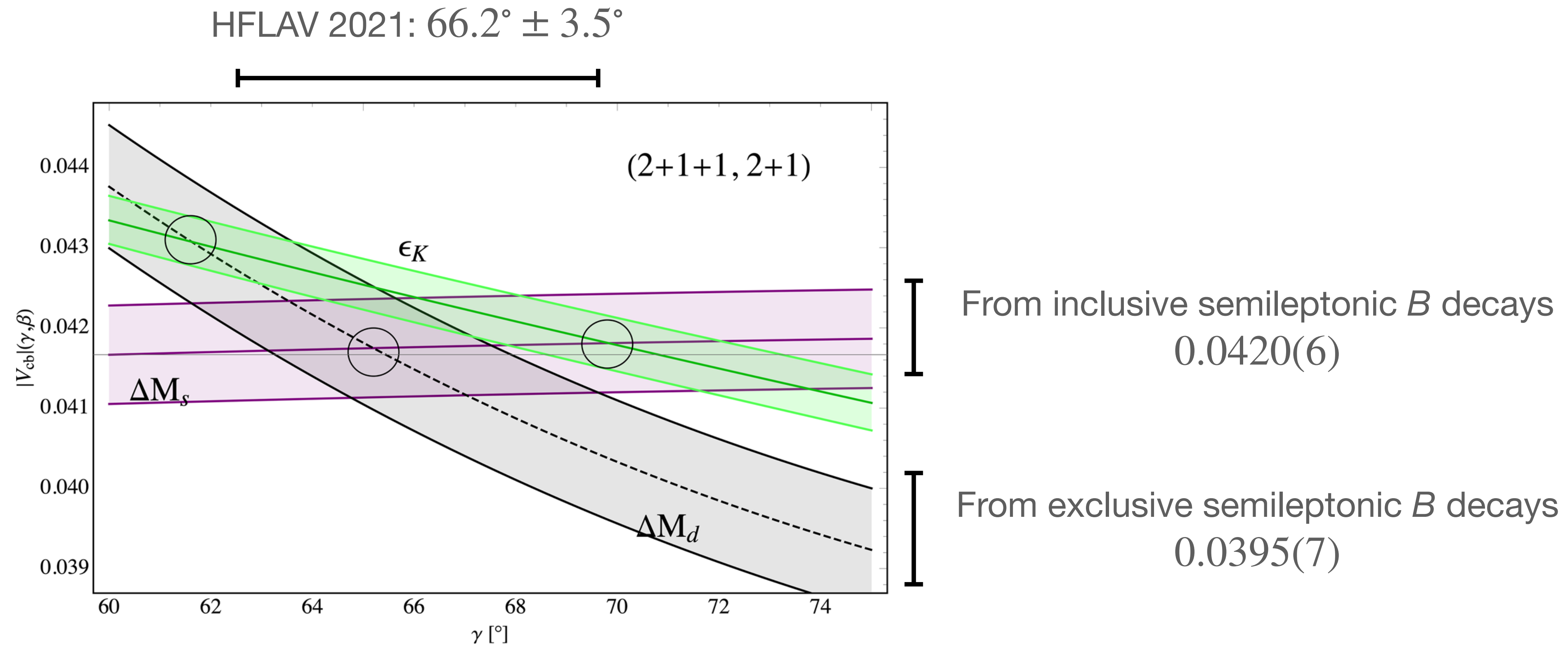
Third row constraints from B mixing

1 σ hexagons



- General agreement of LQCD results, along with sum rules
- Tension with exclusive $|V_{ub}/V_{cb}|$, angle γ , and unitarity

B Mixing, unitarity, and V_{cb}



Buras and Venturini, *Eur. Phys. J. C* 82, 615 (2022)

Conclusions

- LQCD: lots of progress in precision, general agreement between methods
- CKM: persistent inclusive/exclusive tension. Growing tension between exclusive $b \rightarrow c$ and B -mixing (+unitarity)
- Future: lots of “next generation calculations” in progress, errors being reduced with a range of formulations
New ideas for including QED effects, contributing to inclusive decays, improving dispersive constraints on ff shapes
- Advertisement of other talks
 - 11:00 Monday: Alejandro Vaquero, $B \rightarrow D^* l \nu$ and related anomalies
 - 16:30 Monday: Enrico Lunghi, B and K physics, role of lattice QCD
 - 09:30 Tuesday: Chris Bouchard, $B \rightarrow K l l$ and $D \rightarrow K l \nu$ form factors and phenomenology
 - 11:00 Tuesday: Simon Kuberski, Heavy flavor physics with $O(a)$ improved Wilson quarks