

# Global controllability for Burgers equation

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# Introduction

## Case of Navier-Stokes Equations

$(\bar{y}, \bar{p})$  : “ideal” solution of Navier-Stokes equations (for example a stationary solution).

$$\begin{cases} \frac{\partial \bar{y}}{\partial t} - \nu \Delta \bar{y} + \bar{y} \cdot \nabla \bar{y} + \nabla \bar{p} = \mathbf{f} & \text{in } \Omega \times (0, T), \\ \operatorname{div} \bar{y} = 0 & \text{in } \Omega \times (0, T), \\ \bar{y} = 0 & \text{on } \Gamma \times (0, T) \\ \bar{y}(0) = \bar{y}_0 & \text{in } \Omega. \end{cases} \quad (1)$$

Consider a solution of the controlled system, starting from a different initial value

$$\begin{cases} \frac{\partial \mathbf{y}}{\partial t} - \nu \Delta \mathbf{y} + \mathbf{y} \cdot \nabla \mathbf{y} + \nabla p = \mathbf{f} + \mathbf{v} \cdot \mathbf{1}_\omega & \text{in } \Omega \times (0, T), \\ \operatorname{div} \mathbf{y} = 0 & \text{in } \Omega \times (0, T), \\ \mathbf{y} = 0 & \text{on } \Gamma \times (0, T) \\ \mathbf{y}(0) = \mathbf{y}_0 & \text{in } \Omega, \end{cases} \quad (2)$$

$\mathbf{1}_\omega$  : characteristic function of a (little) subset  $\omega$  of  $\Omega$ .

## Exact Controllability to Trajectories :

Can we find a control  $\mathbf{v}$  such that

$$\mathbf{y}(T) = \bar{\mathbf{y}}(T) ?$$

i.e can we reach exactly in finite time the “ideal” trajectory  $\bar{\mathbf{y}}$ ?

Local version : same result provided  $\|\mathbf{y}_0 - \bar{\mathbf{y}}_0\|$  is **small enough**.

Last result (Fernandez-Cara, Guerrero, Imanuvilov, Puel, Journal de Math. Pures et Appl., 2004) (dimension 3) : **Local** exact controllability to trajectories.

$$H = \{\mathbf{y} \in L^2(\Omega)^3, \operatorname{div} \mathbf{y} = 0, \mathbf{y} \cdot \boldsymbol{\nu} = 0 \text{ on } \Gamma\}.$$

**Theorem 1** *Let us assume that*

$$\bar{\mathbf{y}}_0 \in H \cap L^4(\Omega)^3, \quad \bar{\mathbf{y}} \in L^\infty(\Omega \times (0, T))^3$$

*and*

$$\frac{\partial \bar{\mathbf{y}}}{\partial t} \in L^2(0, T; L^\sigma(\Omega))^3, \quad \sigma > \frac{6}{5}$$

*then there exists  $\eta > 0$  such that for every  $\mathbf{y}_0 \in H \cap L^4(\Omega)^3$  such that  $\|\mathbf{y}_0 - \bar{\mathbf{y}}_0\|_{L^4(\Omega)^3} \leq \eta$ , there exists a control  $\mathbf{v} \in L^2(0, T; L^2(\omega))^3$  and a solution  $(\mathbf{y}, p)$  of (2) such that*

$$\mathbf{y}(T) = \bar{\mathbf{y}}(T).$$

Among open problems :

Can the result be **global** (at least to achieve 0)?

Open problem except for control on the whole boundary : combining results of Coron for approximate controllability and a local exact controllability result (Fursikov-Imanuvilov or result mentioned above).

Can we use a more “nonlinear” method ?



## Case of Burgers Equations

For 1-d Burgers equation : counter-example due to Guerrero-Imanuvilov.  
Therefore **no global exact controllability**.

## Global exact boundary controllability for the 2-d Burgers equation

$$\frac{\partial u}{\partial t} - \Delta u + \frac{\partial u^2}{\partial x_1} + \frac{\partial u^2}{\partial x_2} = f \quad \text{in } Q = (0, T) \times \Omega, \quad (3)$$

$$u|_{\Gamma_0} = 0, \quad u|_{\Gamma_1} = h, \quad (4)$$

$$u(0, \cdot) = u_0, \quad (5)$$

$$u(T, \cdot) = 0. \quad (6)$$

Without loss of generality we may assume that  $\Omega$  is included in the rectangle  $0 \leq x_2 - x_1 \leq A$ ,  $-B \leq x_1 + x_2 \leq B$  with  $A$  and  $B$  two positive constants.

**Theorem 2** *Let us assume that*

$$\Gamma_0 \subset \{x \in \Gamma \mid x_1 - x_2 = 0\} \quad (7)$$

*(or  $\Gamma_0$  is empty which is allowed). Suppose that  $f \in L^2(0, T; L^2(\Omega))$  and that there exists  $T_0 \in (0, T)$  such that  $f(t, x) = 0, \forall t \geq T_0$ . Then for every  $u_0 \in L^2(\Omega)$  there exists a solution  $u \in L^2(0, T; H_{\Gamma_0}^1(\Omega)) \cap C([0, T]; L^2(\Omega))$  such that  $t^2.u \in H^{1,2}(Q) = H^1(0, T; L^2(\Omega)) \cap L^2(0, T; H^2(\Omega) \cap H_{\Gamma_0}^1(\Omega))$  to problem (3)-(5) satisfying (6) (and a corresponding control  $h$ ).*

Proof : related to the return method by Coron but different. Use of a special solution of Burgers equation that we can drive to zero whenever we want.

First of all some existence and regularity results for Burgers equations (good exercises !!)

**Proposition 3** *For every  $f \in L^2(0, T; H^{-1}(\Omega))$  and  $u_0 \in L^2(\Omega)$  there exists a unique solution  $u$  to 2-D Burgers equation with  $u \in L^2(0, T; H_0^1(\Omega)) \cap C([0, T]; L^2(\Omega))$  and we have*

$$\|u\|_{L^2(0, T; H_0^1(\Omega))} + \|u\|_{C([0, T]; L^2(\Omega))} \leq C(\|u_0\|_{L^2(\Omega)} + \|f\|_{L^2(0, T; H^{-1}(\Omega))}).$$

*If  $f \in L^2(0, T; L^2(\Omega))$  and  $u_0 \in H_0^1(\Omega)$  then  $u \in H^{1,2}(Q) = H^1(0, T; L^2(\Omega)) \cap L^2(0, T; H^2(\Omega) \cap H_0^1(\Omega))$  and we have*

$$\|u\|_{H^{1,2}(Q)} \leq C(\|u_0\|_{H_0^1(\Omega)} + \|f\|_{L^2(0, T; L^2(\Omega))} + \|u_0\|_{H_0^1(\Omega)}^5 + \|f\|_{L^2(0, T; L^2(\Omega))}^5).$$

**Proposition 4** *Let us assume that  $f \in L^2(0, T; L^2(\Omega))$  and that  $u_0 \in L^2(\Omega)$ . Then  $t^2.u \in H^{1,2}(Q)$  which implies that for every  $\eta > 0$ ,  $u \in C([\eta, T]; H_0^1(\Omega)) \cap L^2(\eta, T; H^2(\Omega))$  and  $\frac{\partial u}{\partial t} \in L^2(\eta, T; L^2(\Omega))$ . Moreover we have the following estimate*

$$\begin{aligned} \|t^2.u\|_{H^{1,2}(Q)} \leq C(& \|u_0\|_{L^2(\Omega)} + \|f\|_{L^2(0,T;L^2(\Omega))} + \\ & + \|u_0\|_{L^2(\Omega)}^{13} + \|f\|_{L^2(0,T;L^2(\Omega))}^{13}). \end{aligned} \quad (8)$$

On the time interval  $(0, T_0)$  set  $h(t, x) = 0$  and leave the system evolve without control. For every  $\eta > 0$ , we have

$$u \in C([\eta, T_0]; H_0^1(\Omega)) \cap L^2(\eta, T_0; H^2(\Omega)), \quad \frac{\partial u}{\partial t} \in L^2(\eta, T_0; L^2(\Omega))$$

and we write

$$u(T_0, \cdot) = u_1 \in H_0^1(\Omega) \subset L^p(\Omega), \quad \forall p, \quad 1 \leq p < +\infty.$$

Now we set

$$\delta_0 = \frac{T - T_0}{4} > 0.$$

We will construct a solution  $u$  in the interval  $(T_0, T_0 + 3\delta_0)$  (and a corresponding control) such that  $u(T_0 + 3\delta_0, \cdot)$  is as small as desired in the norm  $H_0^1(\Omega)$ .



First of all we construct a very specific solution  $U$  of the 2-d Burgers equation.

Let  $w(t, z)$  be a solution to the heat equation

$$\frac{\partial w}{\partial t} - 2\frac{\partial^2 w}{\partial z^2} = 0 \quad z \in (0, A), \quad t > T_0, \quad (9)$$

$$w(t, 0) = 0, \quad w(t, A) = v(t), \quad (10)$$

$$w(T_0, \cdot) = 0, \quad (11)$$

where  $v(\cdot)$  is a boundary control which will be determined later on. This control will be chosen regular so that  $w$  will also be regular.

We now set

$$U(t, x) = w(t, x_2 - x_1). \quad (12)$$

We have

$$\frac{\partial U}{\partial x_1} + \frac{\partial U}{\partial x_2} = 0, \quad \frac{\partial U^2}{\partial x_1} + \frac{\partial U^2}{\partial x_2} = 0$$

so that for every  $N > 0$ ,  $N.U$  is a regular solution of the 2-d Burgers equation

$$\begin{aligned} \frac{\partial(N.U)}{\partial t} - \Delta(N.U) + \frac{\partial(N.U)^2}{\partial x_1} + \frac{\partial(N.U)^2}{\partial x_2} &= 0 \quad \text{in } (T_0, T) \times \Omega, \\ N.U|_{\Gamma_0} &= 0, \\ N.U(T_0, \cdot) &= 0. \end{aligned}$$

Notice that the value of  $N.U$  on  $(T_0, T) \times \Gamma_1$ , which will be a boundary control  $h$  and which depends on  $v$ , does not appear explicitly. If  $\delta$  is any number such that  $0 < \delta \leq \delta_0$ , from the controllability results for the heat equation, we can choose this control  $h$  (and in fact  $v$ ) on  $(T_0 + \delta, T_0 + 2\delta_0)$  such that

$$N.U(T_0 + 2\delta_0, \cdot) = 0.$$

On the interval  $(T_0, T_0 + 2\delta_0)$  we look for  $u$  in the form

$$u = y + N.U, \tag{13}$$

where  $N$  is a large parameter to be determined later on and  $y$  is chosen to vanish on the whole boundary  $\Gamma$ .

Therefore,  $y$  must satisfy the following equation

$$\frac{\partial y}{\partial t} - \Delta y + 2N.U\left(\frac{\partial y}{\partial x_1} + \frac{\partial y}{\partial x_2}\right) + \frac{\partial y^2}{\partial x_1} + \frac{\partial y^2}{\partial x_2} = 0 \quad (14)$$

in  $(T_0, T_0 + 2\delta_0) \times \Omega,$

$$y|_{\Gamma} = 0, \quad (15)$$

$$y(T_0, \cdot) = u_1. \quad (16)$$

**Lemma 5** *There exists a unique solution  $y$  to (14), (15), (16) with  $y \in C([T_0, T_0 + 2\delta_0]; H_0^1(\Omega)) \cap L^2(T_0, T_0 + 2\delta_0; H^2(\Omega)), \frac{\partial y}{\partial t} \in L^2(T_0, T_0 + 2\delta_0; L^2(\Omega))$  and for every  $t_0, t_1$  with  $T_0 \leq t_0 \leq t_1 \leq T_0 + 2\delta_0$  and every  $p \geq 1$  we have*

$$\|y(t_1, \cdot)\|_{L^p(\Omega)} \leq \|y(t_0, \cdot)\|_{L^p(\Omega)}. \quad (17)$$

Proof.

Existence, uniqueness and regularity of  $y$  is classical as (14) is essentially a Burgers equation. To show that the  $L^p$ -norm of  $y$  is decreasing, multiply equation (14) by  $|y|^{p-2}y$  with  $p \geq 1$ . We obtain

$$\frac{1}{p} \frac{d}{dt} \int_{\Omega} |y|^p dx + (p-1) \int_{\Omega} |y|^{p-2} |\nabla y|^2 dx = 0$$

since

$$\int_{\Omega} U \cdot \left( \frac{\partial y}{\partial x_1} + \frac{\partial y}{\partial x_2} \right) |y|^{p-2} y dx = \frac{1}{p} \int_{\Omega} U \cdot \left( \frac{\partial |y|^p}{\partial x_1} + \frac{\partial |y|^p}{\partial x_2} \right) dx = 0$$

and

$$\int_{\Omega} \left( \frac{\partial y^2}{\partial x_1} + \frac{\partial y^2}{\partial x_2} \right) |y|^{p-2} y dx = \frac{2}{p+1} \int_{\Omega} \left( \frac{\partial |y|^p y}{\partial x_1} + \frac{\partial |y|^p y}{\partial x_2} \right) dx = 0.$$

Let us now define a function  $\beta$  by

$$\beta(x) = C_0 - x_1 - x_2,$$

where  $C_0$  is chosen such that

$$\exists \beta_0 > 0, \forall x \in \Omega, \beta(x) \geq \beta_0.$$

We also write

$$\beta_1 = \max_{x \in \overline{\Omega}} \beta(x).$$

**Lemma 6** *The solution  $y$  of (14), (15), (16) satisfies the following differential inequality*

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega} \beta |y|^2 dx + \int_{\Omega} \beta |\nabla y|^2 dx + \frac{2}{\beta_1} \int_{\Omega} (N.U) \beta |y|^2 dx \leq \frac{4}{3} \int_{\Omega} |u_1|^3 dx. \quad (18)$$

Proof.

Multiply equation (14) by  $\beta y$ . We obtain, as  $\Delta \beta = 0$  and  $\frac{\partial \beta}{\partial x_1} + \frac{\partial \beta}{\partial x_2} = -2$ ,

$$\frac{1}{2} \frac{d}{dt} \int_{\Omega} \beta |y|^2 dx + \int_{\Omega} \beta |\nabla y|^2 dx + 2 \int_{\Omega} (N.U) |y|^2 dx + \frac{4}{3} \int_{\Omega} |y|^2 y dx = 0.$$

Using Lemma 5 with  $p = 3$  we obtain the desired result.

Notice that up to this point **the control  $v$  has not been chosen.**

In the case when  $\Gamma_0$  is empty which means that we can apply a **control on the whole boundary**, we don't have to take the boundary condition  $w(t, 0) = 0$  and we can take  $w$  such that  $\min_{x \in \Omega} U(t, x) \geq \min_{z \in (0, A)} w(t, z) \geq \alpha(t) > 0$  if  $t > T_0$ , which ensures that  $U$  has a strictly positive minimum when  $t > T_0$ .

When  $\Gamma_0$  is not empty, due to the boundary condition  $w(t, 0) = 0$  we cannot have a strictly positive minimum for  $U$  over  $\Omega$ .



Let us now make a choice for  $w$  and  $v$ . On the interval  $(T_0, T_0 + \delta)$ , where  $0 < \delta \leq \delta_0$ , we set

$$w(t, z) = \frac{1}{\sqrt{(t - T_0)}} \left( e^{-\frac{(z-5A)^2}{8(t-T_0)}} - e^{-\frac{(z+5A)^2}{8(t-T_0)}} \right). \quad (19)$$

We can see that  $w$  satisfies (9), (10) with a suitable control  $v$  and (11).

For  $0 < a \leq z \leq A$  we have

$$\begin{aligned} w(t, z) \geq w(t, a) &= \frac{2}{\sqrt{(t - T_0)}} e^{-\frac{(a^2+25A^2)}{8(t-T_0)}} \sinh\left(\frac{5Aa}{4(t - T_0)}\right) \\ &\geq \frac{5Aa}{2(t - T_0)^{\frac{3}{2}}} e^{-\frac{(a^2+25A^2)}{8(t-T_0)}}. \end{aligned}$$

At the same time we also have

$$\exists C_0 > 0, \forall a \in (0, A), \forall t \in (T_0, T_0 + \delta), \forall z, 0 \leq z \leq a, w(t, z) \leq w(t, a) \leq C_0 a.$$

We will write

$$\Omega_a = \{x \in \Omega, 0 \leq x_2 - x_1 \leq a\}$$

and we have

$$|\Omega_a| \leq Ca,$$

and

$$\min_{x \in \Omega \setminus \Omega_a} U(t, x) \geq w(t, a) \geq \frac{5Aa}{2(t - T_0)^{\frac{3}{2}}} e^{-\frac{(a^2 + 25A^2)}{8(t - T_0)}}.$$

Therefore, from (18), we obtain

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\Omega} \beta |y|^2 dx + \int_{\Omega} \beta |\nabla y|^2 dx + \frac{5NAa}{\beta_1(t-T_0)^{\frac{3}{2}}} e^{-\frac{(a^2+25A^2)}{8(t-T_0)}} \int_{\Omega} \beta |y|^2 dx \\
& \leq \frac{4}{3} \int_{\Omega} |u_1|^3 dx + 2N \int_{\Omega_a} w(t, a) |y|^2 dx \\
& \leq \frac{4}{3} \int_{\Omega} |u_1|^3 dx + 2Nw(t, a) |\Omega_a|^{\frac{1}{3}} \left( \int_{\Omega} |y|^3 dx \right)^{\frac{2}{3}} \\
& \leq \frac{4}{3} \int_{\Omega} |u_1|^3 dx + CNa^{\frac{4}{3}} \left( \int_{\Omega} |u_1|^3 dx \right)^{\frac{2}{3}}.
\end{aligned}$$

We now take

$$a = \frac{1}{N^{\frac{3}{4}}}$$

which implies the following differential inequality

$$\frac{d}{dt} \int_{\Omega} \beta |y|^2 dx \leq -\frac{10N^{\frac{1}{4}}A}{\beta_1(t-T_0)^{\frac{3}{2}}} e^{-\frac{26A^2}{8(t-T_0)}} \int_{\Omega} \beta |y|^2 dx + C(\|u_1\|_{L^3(\Omega)}).$$

Using Gronwall Lemma, integrating this inequality on  $(T_0, T_0 + \delta)$ , we obtain

$$\int_{\Omega} \beta |y(T_0 + \delta, x)|^2 dx \leq \left( \int_{\Omega} \beta |u_1|^2 dx \right) e^{-N^{\frac{1}{4}}g(\delta)} + \delta C(\|u_1\|_{L^3(\Omega)})$$

where for  $\delta$  small enough

$$g(\delta) = \int_{T_0}^{T_0+\delta} \frac{10A}{\beta_1(t-T_0)^{\frac{3}{2}}} e^{-\frac{26A^2}{8(t-T_0)}} dt \geq C e^{-\frac{A^2}{\delta}} > 0$$

This implies

$$\int_{\Omega} |y(T_0 + \delta, x)|^2 dx \leq \frac{\beta_1}{\beta_0} \|u_1\|_{L^2(\Omega)}^2 e^{-N^{\frac{1}{4}} g(\delta)} + \frac{\delta}{\beta_0} C(\|u_1\|_{L^3(\Omega)})$$

and, choosing first  $\delta$  sufficiently small then  $N$  sufficiently large we have proved the following

**Proposition 7** *Given  $u_1$  in  $H_0^1(\Omega)$  (in fact  $u_1 \in L^3(\Omega)$  would be enough), for every  $\delta_0 > 0$  and for every  $\epsilon_0 > 0$ , there exists  $\delta$  with  $0 < \delta \leq \delta_0$  and there exists  $N$  sufficiently large such that*

$$\|y(T_0 + \delta, \cdot)\|_{L^2(\Omega)} \leq \epsilon_0.$$

Now we choose the control  $v$  on the time interval  $(T_0 + \delta, T_0 + 2\delta_0)$  in (10) such that  $w$  satisfies

$$w(T_0 + 2\delta_0, \cdot) = 0.$$

This is possible using classical results on null controllability for the heat equation. Then we also have

$$U(T_0 + 2\delta_0, \cdot) = 0.$$

Therefore,

$$\|u(T_0 + 2\delta_0, \cdot)\|_{L^2(\Omega)} = \|y(T_0 + 2\delta_0, \cdot)\|_{L^2(\Omega)} \leq \|y(T_0 + \delta, \cdot)\|_{L^2(\Omega)} \leq \epsilon_0.$$

Notice that  $\epsilon_0$  can be chosen as small as we wish. At this point we only know that the  $L^2(\Omega)$ -norm of  $u(T_0 + 2\delta_0, \cdot)$  is as small as we wish.

On the interval  $(T_0 + 2\delta_0, T_0 + 3\delta_0)$  we let the system evolve freely and we take the boundary control equal zero. Then using the regularizing effect of Burgers equation we see that at time  $T_0 + 3\delta_0$  we have

$$\|u(T_0 + 3\delta_0, \cdot)\|_{H_0^1(\Omega)} \leq \epsilon_1,$$

where  $\epsilon_1$  can be taken as small as we wish provided  $\epsilon_0$  is small enough.

Therefore, on the time interval  $(T_0 + 3\delta_0, T)$  we can use a result of local exact controllability to trajectories for 2-d Burgers equations (not completely trivial !) to find a boundary control  $h$  such that

$$u(T, \cdot) = 0.$$

**A situation without global controllability**



Theorem 2 was proved under the restrictive assumption (7) on the boundary  $\Gamma_0$ . The next result shows that without this assumption the global controllability property may fail.

Let us suppose that the geometrical situation is such that there exists a function  $\rho(x) \in C^2(\bar{\Omega})$  such that

$$\rho|_{\Gamma_1} = 0, \quad \rho(x) > 0 \text{ in } \Omega, \quad \frac{\partial \rho}{\partial x_1} + \frac{\partial \rho}{\partial x_2} < 0 \quad \forall x \in \bar{\Omega}. \quad (20)$$

Of course this cannot occur in the situation considered in the previous section, but there are many cases where such a function  $\rho$  exists, for example when  $\Omega = \{(x_1, x_2), 0 < x_2 - x_1 < 1, -1 < x_1 + x_2 < 1\}$  and  $\Gamma_1 = \{(x_1, x_2), 0 < x_2 - x_1 < 1, x_1 + x_2 = 1\}$ .

For a function  $v$  defined on  $\Omega$  or  $(0, T) \times \Omega$  we set

$$v^+ = \max(v, 0), \quad v^- = (-v)^+.$$

**Theorem 8** *Suppose that condition (20) holds true. Let  $f \in L^2(0, T; L^2(\Omega))$  and  $u_0 \in H_0^1(\Omega)$  such that  $u_0^- \neq 0$ . Then there exists a time  $T_0(u_0^-, f) > 0$  such that for each  $T \leq T_0(u_0^-, f)$  there is no solution to problem (3)-(5) in the space  $u \in H^{1,2}(Q)$  satisfying (6).*

**Proof.** We argue by contradiction. Let  $u_0 \in H_0^1(\Omega)$  and  $f \in L^2(0, T; L^2(\Omega))$  be given functions. Suppose that there exists a solution  $u$  to (3)-(6). Then we consider the function  $y(t, x) = u(t, x) - u_0(x)$  which satisfies the following system of equations

$$\begin{aligned} \frac{\partial y}{\partial t} - \Delta y + \frac{\partial y^2}{\partial x_1} + \frac{\partial y^2}{\partial x_2} + 2\frac{\partial(yu_0)}{\partial x_1} + 2\frac{\partial(yu_0)}{\partial x_2} &= q \quad \text{in } (0, T) \times \Omega, \\ y|_{\Gamma_0} &= 0, \quad y|_{\Gamma_1} = h \quad y(0, \cdot) = 0, \\ y(T, \cdot) &= -u_0, \end{aligned}$$

where

$$q = \Delta u_0 - \frac{\partial u_0^2}{\partial x_1} - \frac{\partial u_0^2}{\partial x_2} + f.$$

We set

$$\rho_1(x) = \rho(x)^4.$$

Multiplying the equation by  $\rho_1 y^+$  and integrating by parts we obtain

$$\begin{aligned}
& \frac{1}{2} \frac{d}{dt} \int_{\Omega} \rho_1 |y^+|^2 dx + \int_{\Omega} \left( \rho_1 |\nabla y^+|^2 - \frac{\Delta \rho_1}{2} |y^+|^2 - \frac{2}{3} \left( \frac{\partial \rho_1}{\partial x_1} + \frac{\partial \rho_1}{\partial x_2} \right) (y^+)^3 \right) dx \\
& + \int_{\Gamma_0} \frac{1}{2} \frac{\partial \rho_1}{\partial n} |y^+|^2 d\sigma - 2 \int_{\Omega} \left( \left( \frac{\partial y^+}{\partial x_1} + \frac{\partial y^+}{\partial x_2} \right) \rho_1 u_0 y^+ - u_0 \left( \frac{\partial \rho_1}{\partial x_1} + \frac{\partial \rho_1}{\partial x_2} \right) |y^+|^2 \right) dx \\
& = \int_{\Omega} f \rho_1 y^+ dx - \int_{\Omega} \nabla u_0 \cdot \nabla y^+ \rho_1 dx - \int_{\Omega} \nabla u_0 \cdot \nabla \rho_1 y^+ dx \\
& \quad + \int_{\Omega} u_0^2 y^+ \left( \frac{\partial \rho_1}{\partial x_1} + \frac{\partial \rho_1}{\partial x_2} \right) dx + \int_{\Omega} u_0^2 \rho_1 \left( \frac{\partial y^+}{\partial x_1} + \frac{\partial y^+}{\partial x_2} \right) dx \\
& \leq \int_{\Omega} f \rho_1 y^+ dx - \int_{\Omega} \nabla u_0 \cdot \nabla y^+ \rho_1 dx - \int_{\Omega} \nabla u_0 \cdot \nabla \rho_1 y^+ dx \\
& \quad + \int_{\Omega} u_0^2 \rho_1 \left( \frac{\partial y^+}{\partial x_1} + \frac{\partial y^+}{\partial x_2} \right) dx.
\end{aligned}$$

By (20) we have  $\int_{\Gamma_0} \frac{1}{2} \frac{\partial \rho_1}{\partial \bar{n}} |y^+|^2 d\sigma = 0$ . Again using (20) we may assume that for some positive constant  $M$  we have  $-\frac{2}{3} \left( \frac{\partial \rho_1}{\partial x_1} + \frac{\partial \rho_1}{\partial x_2} \right) > M \rho_1^{\frac{3}{4}}$  for all  $x \in \bar{\Omega}$ . Then denoting by  $C_i$  various constants independent of  $y$  and  $u_0$  we have

$$\begin{aligned} \int_{\Omega} \left( -\frac{\Delta \rho_1}{2} |y^+|^2 - \frac{2}{3} \left( \frac{\partial \rho_1}{\partial x_1} + \frac{\partial \rho_1}{\partial x_2} \right) (y^+)^3 \right) dx &\geq \int_{\Omega} \left( -C_0 \rho_1^{\frac{1}{2}} |y^+|^2 + M \rho_1^{\frac{3}{4}} (y^+)^3 \right) dx \\ &\geq -C_1 \left( \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx \right)^{\frac{2}{3}} + M \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx \geq \frac{3M}{4} \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx - C_2. \end{aligned}$$

Then we have

$$\begin{aligned}
2 \int_{\Omega} \left( \frac{\partial y^+}{\partial x_1} + \frac{\partial y^+}{\partial x_2} \right) \rho_1 u_0 y^+ dx &\leq \frac{1}{4} \int_{\Omega} \rho_1 |\nabla y^+|^2 dx + C_3 \int_{\Omega} u_0^2 \rho_1 |y^+|^2 dx \\
&\leq \frac{1}{4} \int_{\Omega} \rho_1 |\nabla y^+|^2 dx + C_4 \|u_0\|_{H_0^1(\Omega)}^2 \cdot \left( \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx \right)^{\frac{2}{3}} \\
&\leq \frac{1}{4} \int_{\Omega} \rho_1 |\nabla y^+|^2 dx + \frac{M}{4} \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx + C_5 \|u_0\|_{H_0^1(\Omega)}^6.
\end{aligned}$$

Also

$$\begin{aligned}
2 \int_{\Omega} u_0 \left( \frac{\partial \rho_1}{\partial x_1} + \frac{\partial \rho_1}{\partial x_2} \right) |y^+|^2 dx &\leq C_6 \|u_0\|_{H_0^1(\Omega)} \cdot \left( \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx \right)^{\frac{2}{3}} \\
&\leq \frac{M}{4} \int_{\Omega} \rho_1^{\frac{3}{4}} (y^+)^3 dx + C_7 \|u_0\|_{H_0^1(\Omega)}^3.
\end{aligned}$$

We also obtain

$$\begin{aligned}
& \int_{\Omega} f \rho_1 y^+ dx - \int_{\Omega} \nabla u_0 \cdot \nabla y^+ \rho_1 dx \\
& - \int_{\Omega} \nabla u_0 \cdot \nabla \rho_1 y^+ dx + \int_{\Omega} u_0^2 \rho_1 \left( \frac{\partial y^+}{\partial x_1} + \frac{\partial y^+}{\partial x_2} \right) dx \\
& \leq C_8 (\|f\|_{L^2(Q)}^2 + \|u_0\|_{H_0^1(\Omega)}^2 + \|u_0\|_{H_0^1(\Omega)}^4) \\
& \quad + \frac{1}{2} \int_{\Omega} \rho_1 |y^+|^2 dx + \frac{1}{4} \int_{\Omega} \rho_1 |\nabla y^+|^2 dx.
\end{aligned}$$

Using all these inequalities we obtain

$$\begin{aligned}
& \frac{d}{dt} \int_{\Omega} \rho_1 |y^+|^2 dx + \int_{\Omega} \rho_1 |\nabla y^+|^2 dx + \int_{\Omega} \frac{M}{2} \rho_1^{\frac{3}{4}} (y^+)^3 dx \\
& \leq C_9 (1 + \|f\|_{L^2(Q)}^2 + \|u_0\|_{H_0^1(\Omega)}^2 + \|u_0\|_{H_0^1(\Omega)}^6) + \int_{\Omega} \rho_1 |y^+|^2 dx.
\end{aligned}$$

Applying Gronwall's inequality we obtain, as  $y^+(0, \cdot) = 0$ ,

$$\sup_{t \in (0, T)} \int_{\Omega} \rho_1 |y^+|^2 dx \leq C_{10} (1 + \|f\|_{L^2(Q)}^2 + \|u_0\|_{H_0^1(\Omega)}^2 + \|u_0\|_{H_0^1(\Omega)}^6) T e^T.$$

Since the right hand side goes to zero as  $T$  goes to zero and  $y^+(T) = u_0^-$ , we immediately arrive to a contradiction and the proof of Theorem 8 is complete.