

Non-Hermitian operators in QM & *PT*-symmetry

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Based on :

- D.K., H. Bíla, M. Znojil, *J. Phys. A* 39 (2006), 10143; [math-ph/0604055]
- D.K., *submitted*; arXiv:0707.1781 [math-ph]
- D. Borisov, D.K., *submitted*; arXiv:0707.3039 [math-ph]

¿ QM with non-Hermitian operators ?

\mathbb{C}



\mathbb{R}

$$H^* = H$$

\mathbb{I}

$$H^{PT} = H$$

Imaginary Numbers by Yves Tanguy, 1954

(Museo Thyssen-Bornemisza, Madrid)

Insignificant non-Hermiticity

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$$\begin{cases} i\dot{U}(t) = H U(t) \\ U(0) = 1 \end{cases}$$

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Theorem (spectral theorem).

Let $H = H^*$ have discrete spectrum, $H\psi_j = E_j\psi_j$.

Then

$$f(H) = \sum_j f(E_j) \psi_j \langle \psi_j, \cdot \rangle$$

for any complex-valued continuous function f .

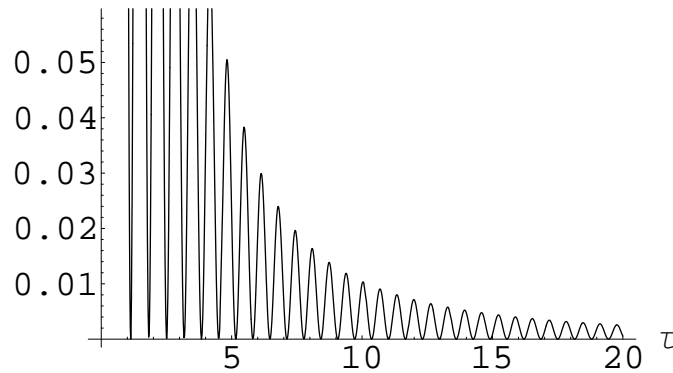
\implies important consequences: **minimax principle**

Technical non-Hermiticity

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Example 1. adiabatic transition probability for $H(t) := \vec{\gamma}(t/\tau) \cdot \vec{\sigma}$, $\tau \rightarrow \infty$

transition probability

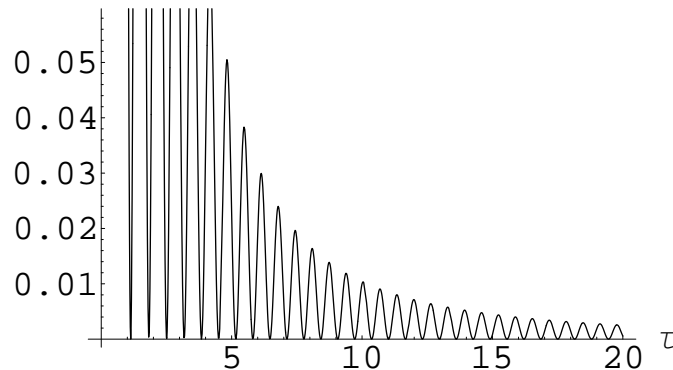


[Berry 1990], [Joye, Kunz, Pfister 1991], [Jakšić, Segert 1993], ...

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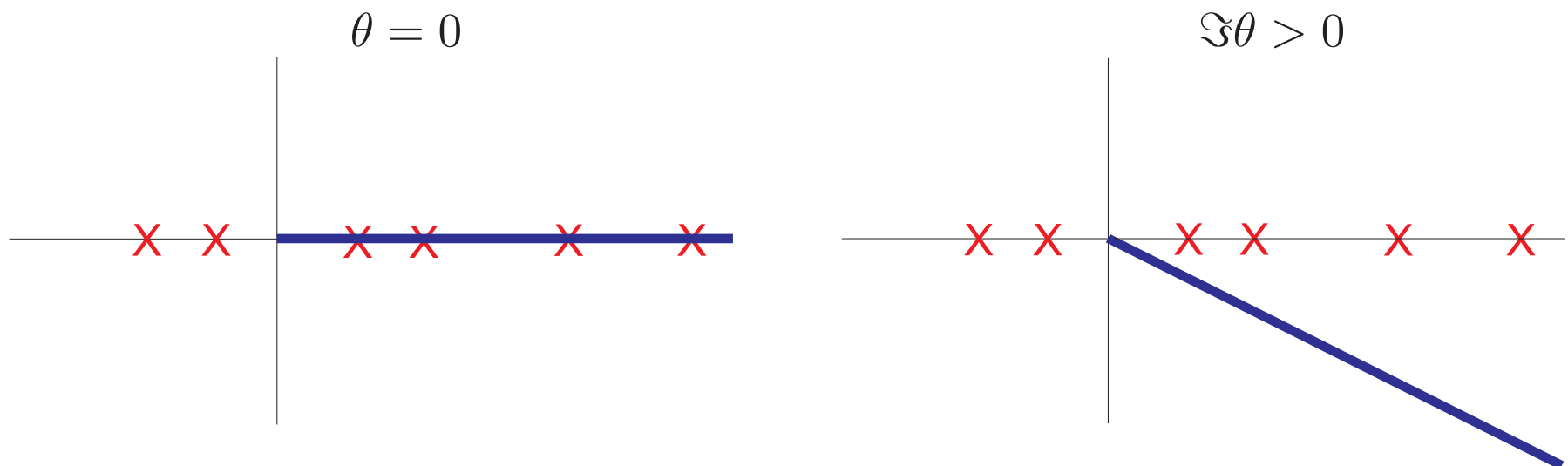
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Example 2. complex scaling $H_\theta := S_\theta(-\Delta + V)S_\theta^{-1}$, $(S_\theta\psi)(x) := e^{\theta/2}\psi(e^\theta x)$

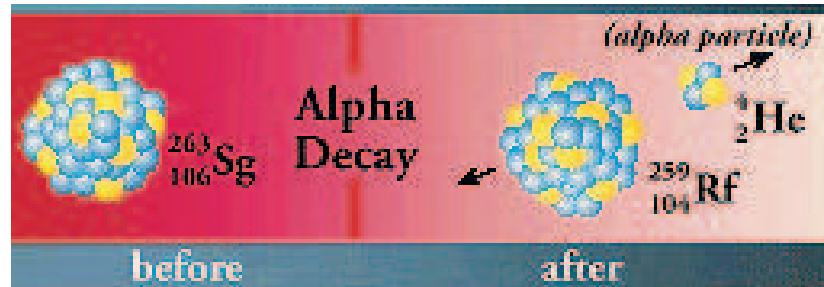


[Aguilar/Balslev, Combes 1971], [Simon 1972], [Van Winter 1974], ...

Approximate non-Hermiticity

open systems

Example 1. radioactive decay



Example 2. dissipative Schrödinger operators in semiconductor physics

Baro, Behrndt, Kaiser, Neidhardt, Rehberg 2002–...

¿ Fundamental non-Hermiticity ?

without violating the “physical axioms” of QM

Non-Hermitian Hamiltonians with real spectra

$$-\Delta + V \quad \text{in} \quad L^2(\mathbb{R})$$

$$V(x) = x^2 + ix^3$$

[Caliceti, Graffi, Maioli 1980]

$$V(x) = ix^3$$

[Bessis, Zinn-Justin]

[Bender, Boettcher 1998]

[Dorey, Dunning, Tateo 2001]

$$V(x) = \begin{cases} i \operatorname{sgn}(x) & \text{if } x \in (-1, 1) \\ \infty & \text{elsewhere} \end{cases}$$

[Znojil 2001]

$$V(x) = \begin{cases} -i \delta(x + \frac{1}{2}) + i \delta(x - \frac{1}{2}) & \text{if } x \in (-1, 1) \\ \infty & \text{elsewhere} \end{cases}$$

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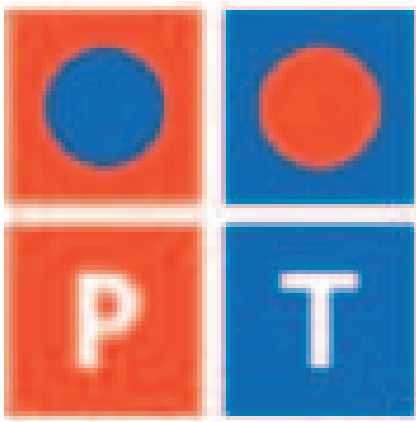
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¿ What is behind the reality of the spectrum ?

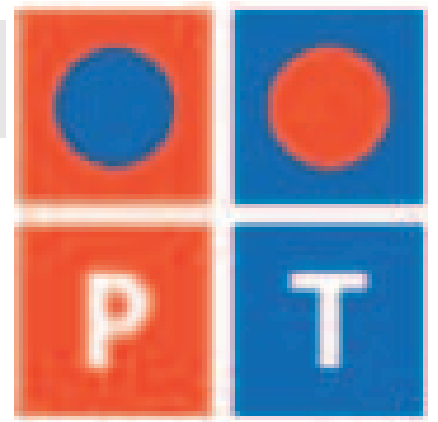
What is \mathcal{PT} -symmetry ?



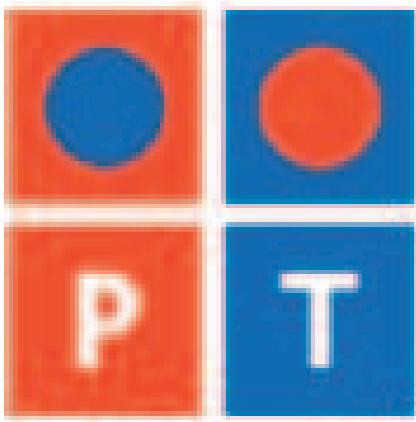
$$[H, \mathcal{PT}] = 0$$

$$(\mathcal{P}\psi)(x) = \psi(-x)$$

$$(\mathcal{T}\psi)(x) = \overline{\psi(x)}$$



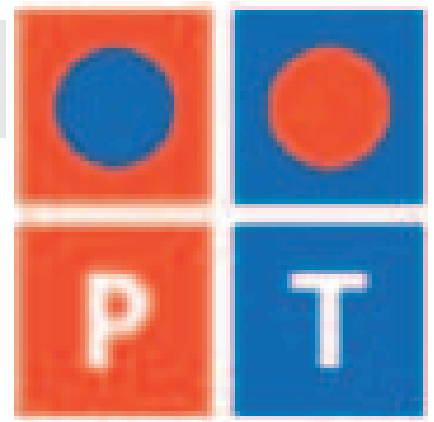
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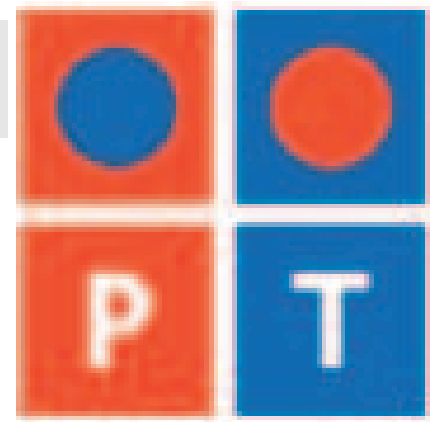
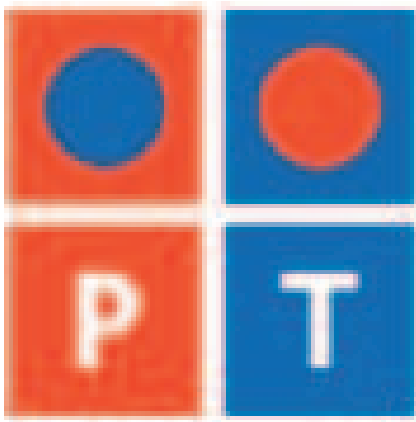
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unbroken \mathcal{PT} -symmetry $:\Leftrightarrow H$ and \mathcal{PT} have the same eigenstates $\Leftrightarrow \sigma(H) \subset \mathbb{R}$

Here we assume that $H = -\Delta + V$ has purely discrete spectrum.

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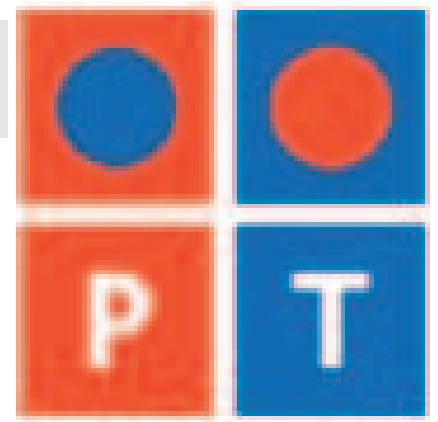
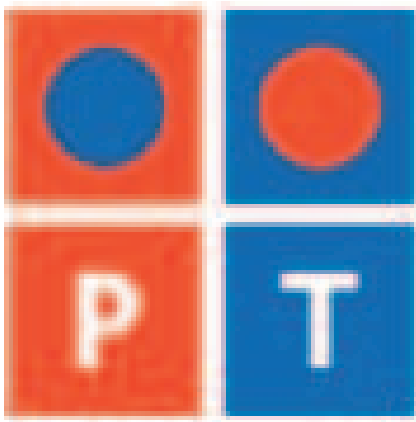
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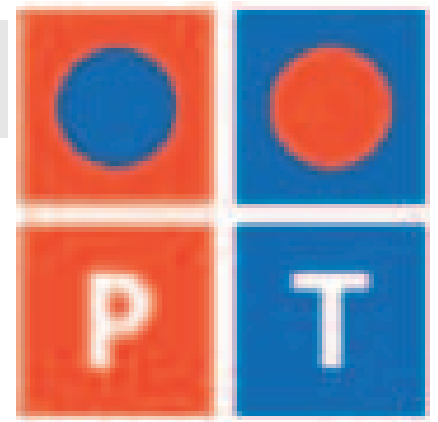
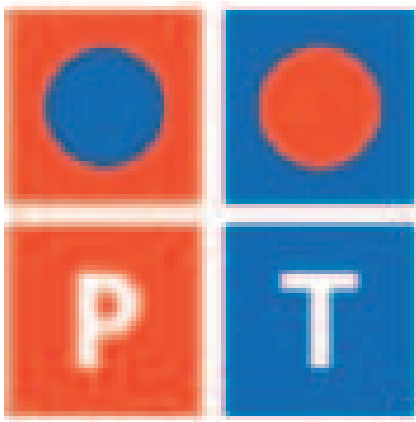
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$\implies H^* := \mathcal{P}H\mathcal{P}$ i.e. H is **pseudo-Hermitian**, i.e. Hermitian w.r.t. $\langle \cdot, \mathcal{P}\cdot \rangle$

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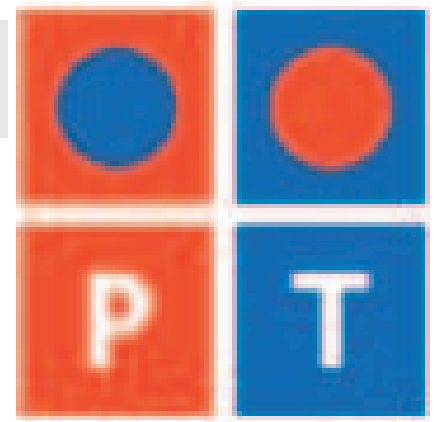
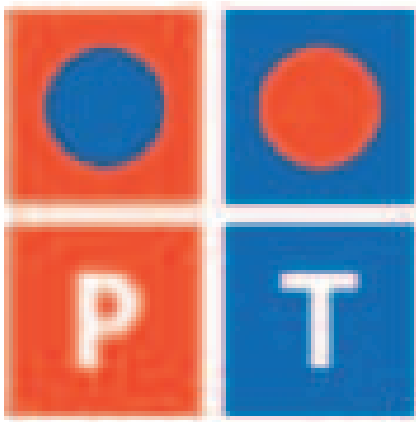
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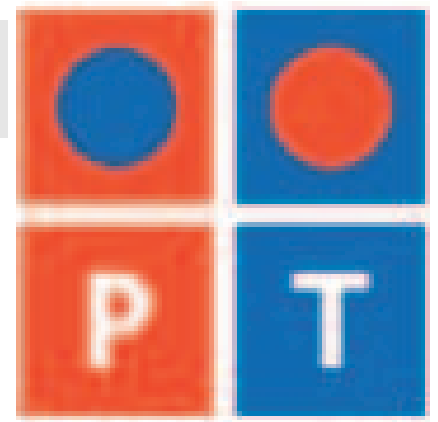
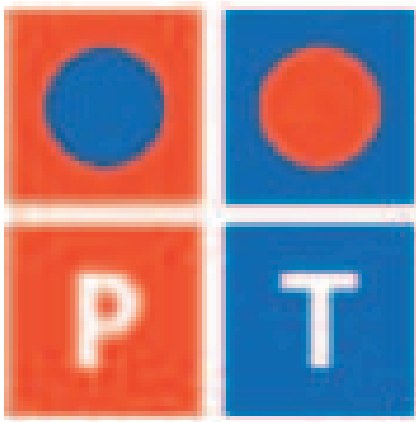
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Albeverio-Fei-Kurasov, Bender-Boettcher, Caliceti-Graffi-Sjöstrand, Jones-Mateo, Langer-Tretter, Mostafazadeh, Scholtz-Geyer-Hahne, Znojil, ...

Attempts to calculate the metric

1. Perturbative

For instance, [Bender 2004] for $\frac{1}{2}p^2 + \frac{1}{2}x^2 + \varepsilon ix^3$: $\Theta^{-1} = \exp(\varepsilon Q_1 + \varepsilon^3 Q_3 + \dots)$

$$Q_1 = -\frac{4}{3}\mu^{-4}p^3 - 2\mu^{-2}S_{1,2},$$

$$Q_3 = \frac{128}{15}\mu^{-10}p^5 + \frac{40}{3}\mu^{-8}S_{3,2} + 8\mu^{-6}S_{1,4} - 12\mu^{-8}p,$$

$$Q_5 = -\frac{320}{3}\mu^{-16}p^7 - \frac{544}{3}\mu^{-14}S_{5,2} - \frac{512}{3}\mu^{-12}S_{3,4} \\ - 64\mu^{-10}S_{1,6} + \frac{24736}{45}\mu^{-14}p^3 + \frac{6368}{15}\mu^{-12}S_{1,2},$$

$$Q_7 = \frac{553984}{315}\mu^{-22}p^9 + \frac{97792}{35}\mu^{-20}S_{7,2} + \frac{377344}{105}\mu^{-18}S_{5,4} \\ + \frac{721024}{315}\mu^{-16}S_{3,6} + \frac{1792}{3}\mu^{-14}S_{1,8} - \frac{2209024}{105}\mu^{-20}p^5 \\ - \frac{2875648}{105}\mu^{-18}S_{3,2} - \frac{390336}{35}\mu^{-16}S_{1,4} + \frac{46976}{5}\mu^{-18}p.$$

$S_{0,0} = 1$, $S_{0,3} = x^3$, $S_{1,1} = \frac{1}{2}(xp + px)$, $S_{1,2} = \frac{1}{3}(x^2p + xpx + px^2)$, and so on.

Other models: [Mostafazadeh, Batal 2004], [Scholtz, Geyer 2006], ...

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2. Formal

– Is Θ well defined ?

– Is Θ bounded ?

– Does $\Theta D(H) \subseteq D(H^*)$ hold ?

(NB $H^*\Theta = \Theta H$)

What is \mathcal{PT} -symmetry ?

a mathematical approach

Special case of J -self-adjointness

[Edmunds, Evans 1987]

$$H^* = JHJ$$

where J is a conjugation operator:
$$\begin{cases} (J\phi, J\psi) = (\psi, \phi) \\ J^2\psi = \psi \end{cases} \quad \forall \phi, \psi \in \mathcal{H}$$

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N.B.
$$\begin{cases} \sigma_p(H) = \{\lambda \mid H - \lambda \text{ is not injective}\} \\ \sigma_c(H) = \{\lambda \mid H - \lambda \text{ is not surjective} \ \& \ \mathfrak{R}(H - \lambda) \text{ is dense}\} \\ \sigma_r(H) = \{\lambda \mid H - \lambda \text{ is injective} \ \& \ \mathfrak{R}(H - \lambda) \text{ is not dense}\} \end{cases}$$

Proof. $\lambda \in \sigma_r(H) \Leftrightarrow \bar{\lambda} \in \sigma_p(H^*) \Leftrightarrow \lambda \in \sigma_p(H)$

q.e.d.

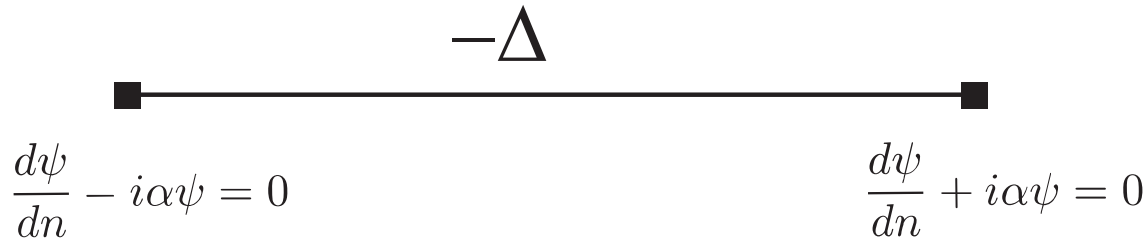
The simplest \mathcal{PT} -symmetric model

[D.K., Břila, Znojil 2006]

$$\mathcal{H} := L^2(0, \pi), \quad H_\alpha \psi := -\psi'', \quad D(H_\alpha) := \left\{ \psi \in W^{2,2}(0, \pi) \left| \begin{array}{l} \psi'(0) + i\alpha\psi(0) = 0 \\ \psi'(\pi) + i\alpha\psi(\pi) = 0 \end{array} \right. \right\}$$

$\alpha \in \mathbb{R}$

$-\Delta$


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Theorem 2. $\sigma(H_\alpha) = \{\alpha^2\} \cup \{n^2\}_{n=1}^\infty$

Corollary. The spectrum of H_α is $\begin{cases} \text{always real,} \\ \text{simple if } \alpha \notin \mathbb{Z} \setminus \{0\}. \end{cases}$

Biorthonormal eigenbasis

$$\alpha \notin \mathbb{Z} \setminus \{0\}$$

$$H_\alpha \psi_n^\alpha = E_n^\alpha \psi_n^\alpha$$

$$H_\alpha^* \phi_n^\alpha = E_n^\alpha \phi_n^\alpha$$

$$\psi_n^\alpha(x) = \begin{cases} A_0^\alpha \exp(-i\alpha x) \\ A_n^\alpha \left(\cos(nx) - i \frac{\alpha}{n} \sin(nx) \right) \end{cases}$$

$$\phi_n^\alpha(x) = \begin{cases} B_0^\alpha \exp(i\alpha x) \\ B_n^\alpha \left(\cos(nx) + i \frac{\alpha}{n} \sin(nx) \right) \end{cases}$$

$$A_n^\alpha = \begin{cases} \sqrt{\frac{1}{\pi}} \frac{i2\pi\alpha}{1 - \exp(-i2\pi\alpha)} \\ \sqrt{\frac{2}{\pi}} \frac{n^2}{n^2 - \alpha^2} \end{cases}$$

Special normalisation:

$$B_n^\alpha = \begin{cases} \sqrt{\frac{1}{\pi}} \\ \sqrt{\frac{2}{\pi}} \end{cases}$$

Theorem 3.

$$\langle \phi_n^\alpha, \psi_m^\alpha \rangle = \delta_{nm}$$

and

$$\psi = \sum_{n=0}^{\infty} \psi_n^\alpha \langle \phi_n^\alpha, \psi \rangle = \sum_{n=0}^{\infty} \phi_n^\alpha \langle \psi_n^\alpha, \psi \rangle$$

Corollary. $f(H) = \sum_j f(E_j) \psi_j \langle \phi_j, \cdot \rangle$

for $f(E) = E^N$, etc.

Calculation of the metric operator

[D.K. 2007]

$$\Theta \equiv \sum_{n=0}^{\infty} \phi_n^\alpha \langle \phi_n^\alpha, \cdot \rangle = \phi_0^\alpha \langle \phi_0^\alpha, \cdot \rangle + \sum_{n=1}^{\infty} \left(\chi_n^N + i \frac{\alpha}{n} \chi_n^D \right) \left\langle \left(\chi_n^N + i \frac{\alpha}{n} \chi_n^D \right), \cdot \right\rangle =: \Theta_1 + \Theta_2$$

$$\Theta_2 = \sum_{n=1}^{\infty} \left\{ \chi_n^N \langle \chi_n^N, \cdot \rangle + \frac{\alpha^2}{n^2} \chi_n^D \langle \chi_n^D, \cdot \rangle + \frac{\alpha}{n^2} p \chi_n^D \langle \chi_n^D, \cdot \rangle + \frac{\alpha}{n^2} p^* \chi_n^N \langle \chi_n^N, \cdot \rangle \right\}$$

$p\psi := -i\psi', \quad D(p) := W_0^{1,2}(0, \pi)$

$$= I - \chi_0^N \langle \chi_0^N, \cdot \rangle + \alpha^2 (-\Delta_D)^{-1} + \alpha p (-\Delta_D)^{-1} + \alpha p^* (-\Delta_N^\perp)^{-1}$$

$$-\Delta_N^\perp := (I - P_0^N)(-\Delta_N)(I - P_0^N)$$

$$P_0^N := \chi_0^N \langle \chi_0^N, \cdot \rangle$$

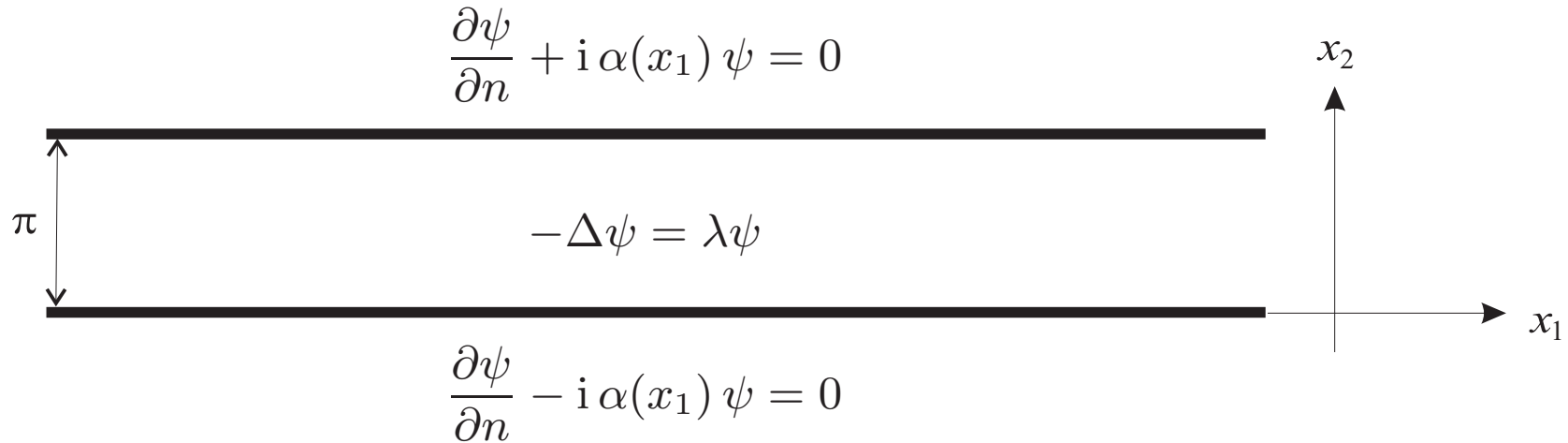
Theorem 4. Θ is bounded, symmetric, non-negative and satisfies

$$\forall \psi \in D(H_\alpha), \quad H_\alpha^* \Theta \psi = \Theta H_\alpha \psi$$

Moreover, Θ is positive if $\alpha \notin \mathbb{Z} \setminus \{0\}$.

\mathcal{PT} -symmetric waveguide

[Borisov, D.K. 2007]



$$\mathcal{H} := L^2(\Omega), \quad \Omega := \mathbb{R} \times (0, \pi)$$

$$H_\alpha \psi := -\Delta \psi, \quad \mathfrak{D}(H_\alpha) := \left\{ \psi \in W^{2,2}(\Omega) \mid \partial_2 \psi + i \alpha \psi = 0 \text{ on } \partial \Omega \right\}, \quad \alpha : \mathbb{R} \rightarrow \mathbb{R}$$

\mathcal{PT} -symmetric waveguide

[Borisov, D.K. 2007]

$$\frac{\partial \psi}{\partial n} + i \alpha(x_1) \psi = 0$$
$$-\Delta \psi = \lambda \psi$$
$$\frac{\partial \psi}{\partial n} - i \alpha(x_1) \psi = 0$$

$$\mathcal{H} := L^2(\Omega), \quad \Omega := \mathbb{R} \times (0, \pi)$$

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Theorem. Let $\alpha \in W^{1,\infty}(\mathbb{R})$. Then H_α is an m -sectorial operator satisfying

$$H_\alpha^* = H_{-\alpha} = \mathcal{T} H_\alpha \mathcal{T} \quad (\mathcal{T}\text{-self-adjointness})$$

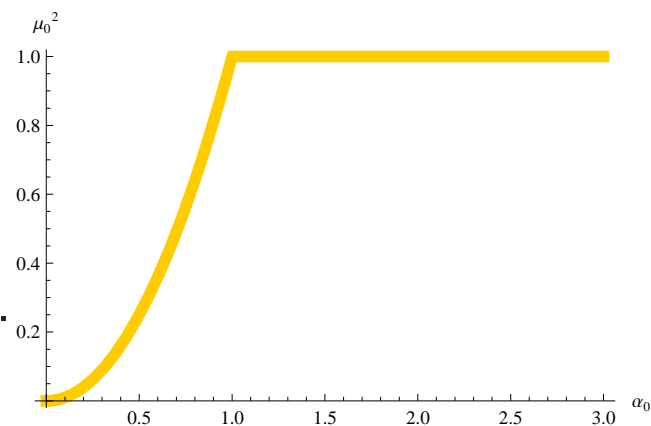
Corollary. $\sigma_r(H_\alpha) = \emptyset$

Spectral analysis

Stability of the continuous spectrum

Theorem. Let $\alpha - \alpha_0 \in C_0(\mathbb{R}) \cap W^{1,\infty}(\mathbb{R})$.

Then $\sigma_c(H_\alpha) = [\mu_0^2, \infty)$ where $\mu_0 := \min\{|\alpha_0|, 1\}$.

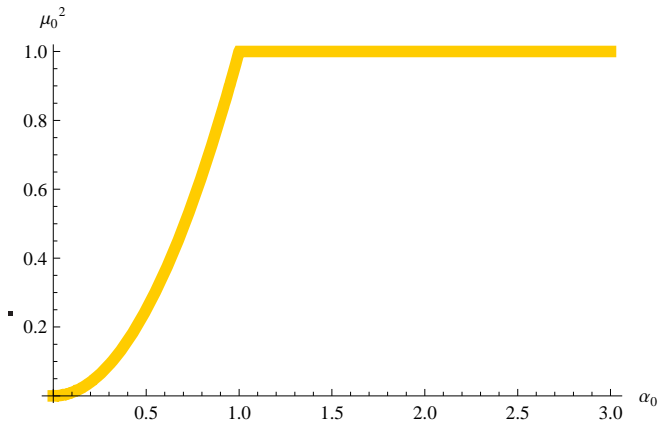


Spectral analysis

Stability of the continuous spectrum

Theorem. Let $\alpha - \alpha_0 \in C_0(\mathbb{R}) \cap W^{1,\infty}(\mathbb{R})$.

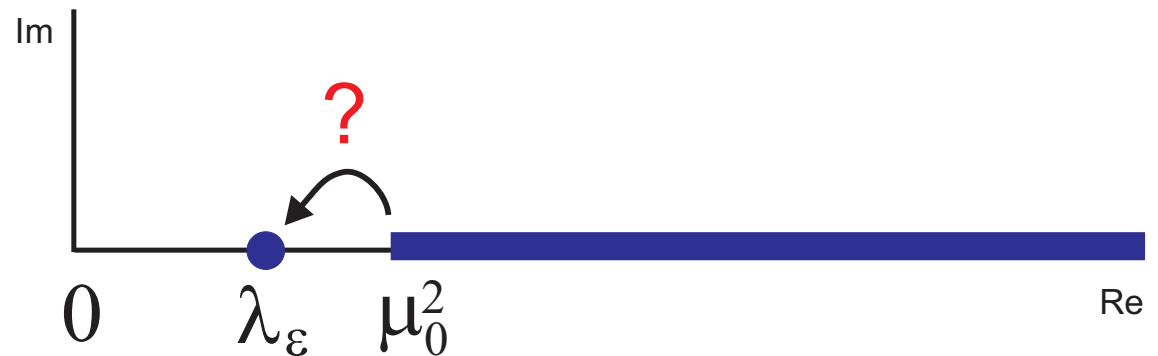
Then $\sigma_c(H_\alpha) = [\mu_0^2, \infty)$ where $\mu_0 := \min\{|\alpha_0|, 1\}$.



Weakly-coupled bound states

$$\alpha(x_1) = \alpha_0 + \varepsilon \beta(x_1)$$

$$\varepsilon \rightarrow 0+ \quad \beta \in C_0^2(\mathbb{R})$$



Theorem. Let $|\alpha_0| < 1$.

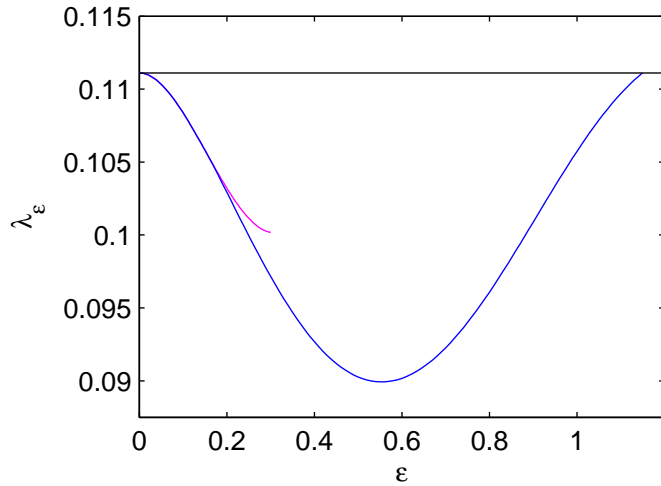
1. $\alpha_0 \langle \beta \rangle < 0 \implies \exists! \lambda_\varepsilon = \mu_0^2 - \varepsilon^2 \alpha_0^2 \langle \beta \rangle^2 + 2\varepsilon^3 \alpha_0 \langle \beta \rangle \tau + \mathcal{O}(\varepsilon^4) \in \mathbb{R}$
2. $\alpha_0 \langle \beta \rangle > 0 \implies \text{no}$
3. $\alpha_0 = 0 \implies \text{no}$

$$\langle \beta \rangle := \int_{\mathbb{R}} \beta(x_1) dx_1$$

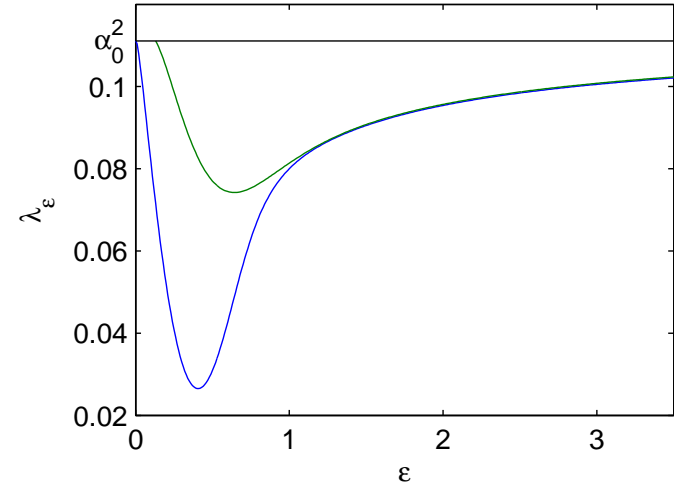
Numerical results

[Tater 2007 (last week)]

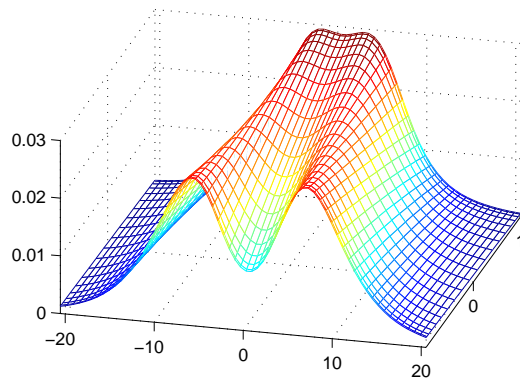
$$\alpha(x) = \alpha_0 - \varepsilon \exp(-x^2)$$



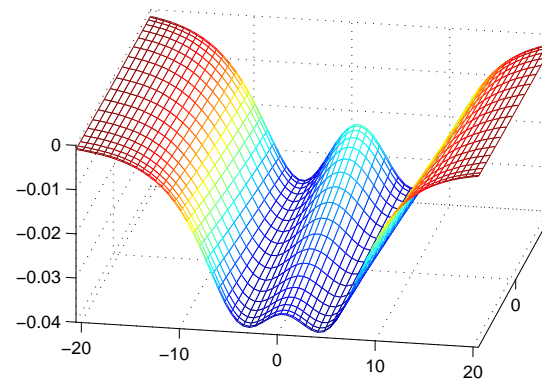
$$\alpha(x) = \alpha_0 - \varepsilon \exp(-0.025x^2)$$



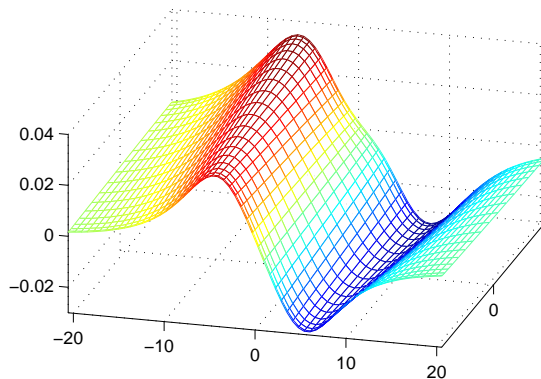
$$\Re(\psi_1) \quad \varepsilon = 0.65$$



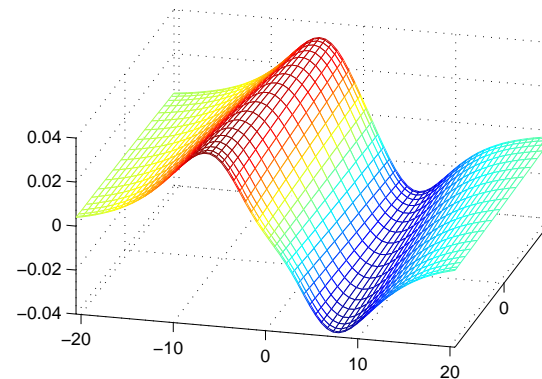
$$\Im(\psi_1) \quad \varepsilon = 0.65$$



$$\Re(\psi_2) \quad \varepsilon = 0.65$$



$$\Im(\psi_2) \quad \varepsilon = 0.65$$



Conclusions

Ad \mathcal{PT} -symmetry:

- no extension of QM
- rather an alternative (pseudo-Hermitian) representation
- overlooked for over 70 years
- i rigorous formulation is still missing !
- i phenomenological relevance ?

Conclusions

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- rather an alternative (pseudo-Hermitian) representation
- overlooked for over 70 years
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Ad \mathcal{PT} -symmetric waveguide:

- rigorous treatment
- ! reality of the total spectrum ?
- ! non-perturbative proof of the existence of the point spectrum ?
- ! calculation of the metric operator ?
- ! physical motivation ?