

The Cluster Condensum

Cosmological Framework

- H_0 measurements
- Supernovae
- CMB
- Clusters
- BBN
- Lensing

Concordance
Model!

$$H_0 \approx 70 \text{ km/s/Mpc} \quad \Omega_\lambda \approx 0.7 \quad \Omega_m \approx 0.3$$

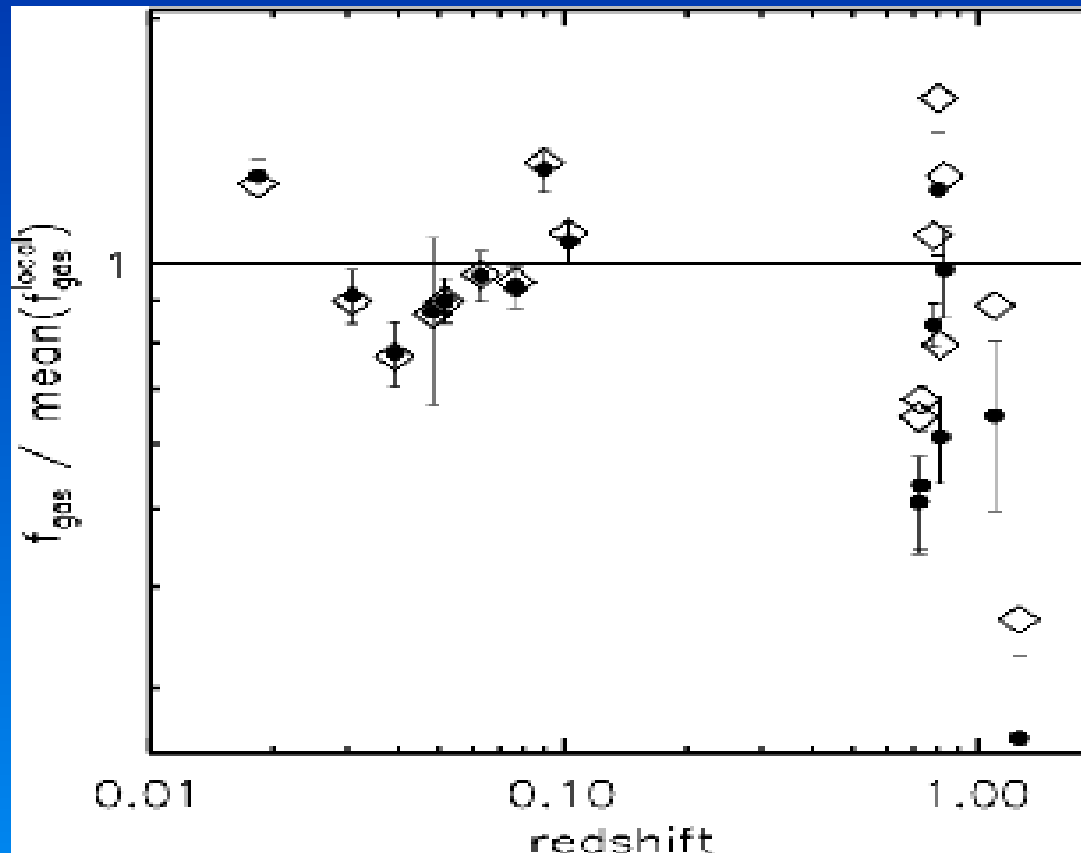
$$\Omega_m$$

From X-ray Clusters

Baryon Fraction evolution

Baryon Fraction

(non)evolution as a test of Ω_m



Chandra

Scaling relations
of
(X-ray) clusters

Basic principle:

Two clusters at two different epochs

z_1 and z_2 with different

masses M_1 and M_2

should be identical once densities

and sizes are scaled properly

Scaling argument for Clusters:

Clusters are geometrically identical

With virial radius-mass relation

$$M = \frac{4\pi}{3} \rho_0 (1+z)^3 (1+\Delta) R_V^3$$

i.e.

$$R_V = \sqrt[3]{\frac{3M}{4\pi\rho_0(1+\Delta)(1+z)}}$$

Mass-temperature relation:

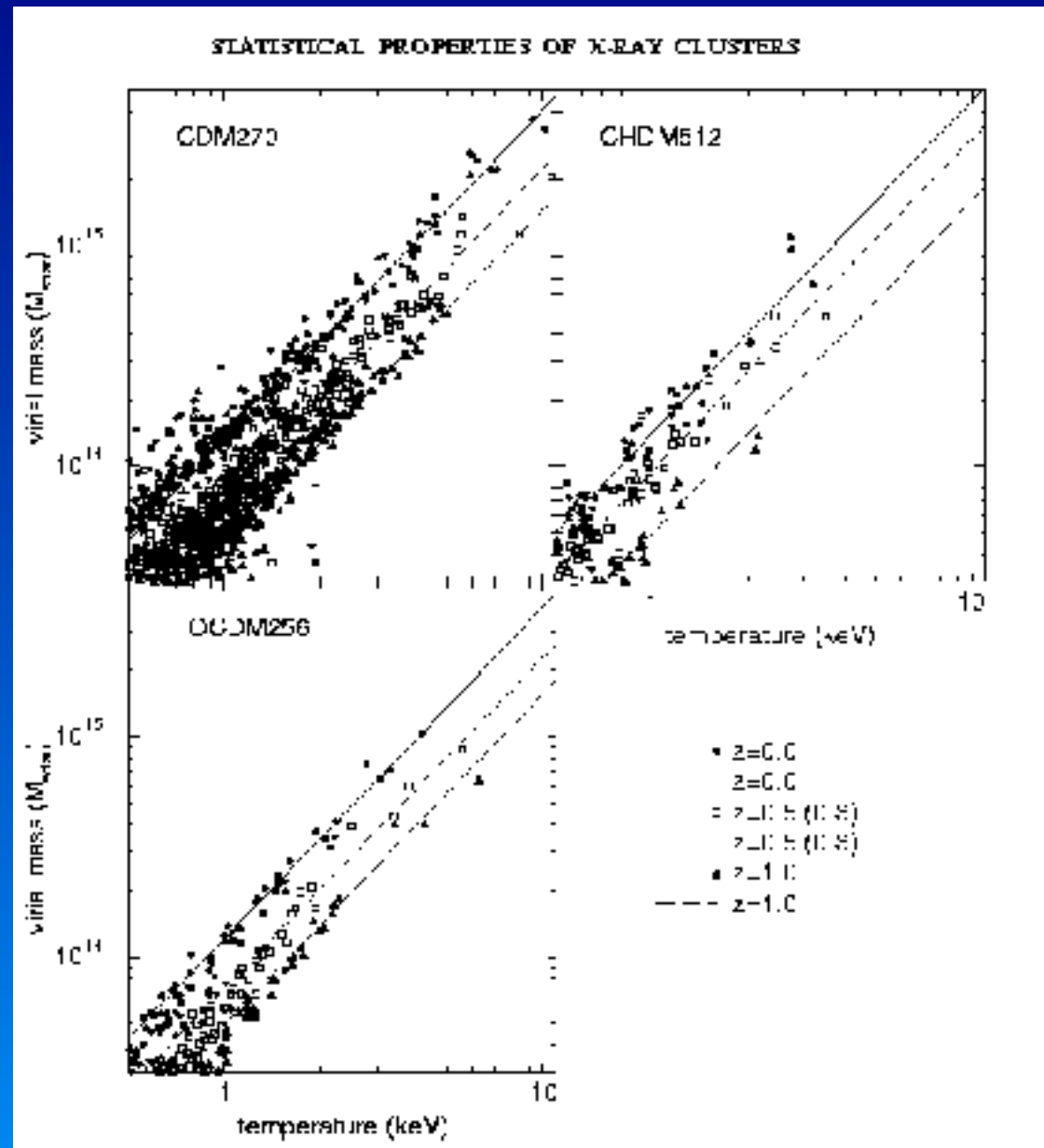
$$T \propto GM/r$$

whatever you do with gravity...

$$T_x \simeq A_{TM} M^{2/3} (1+z) \text{ keV}$$

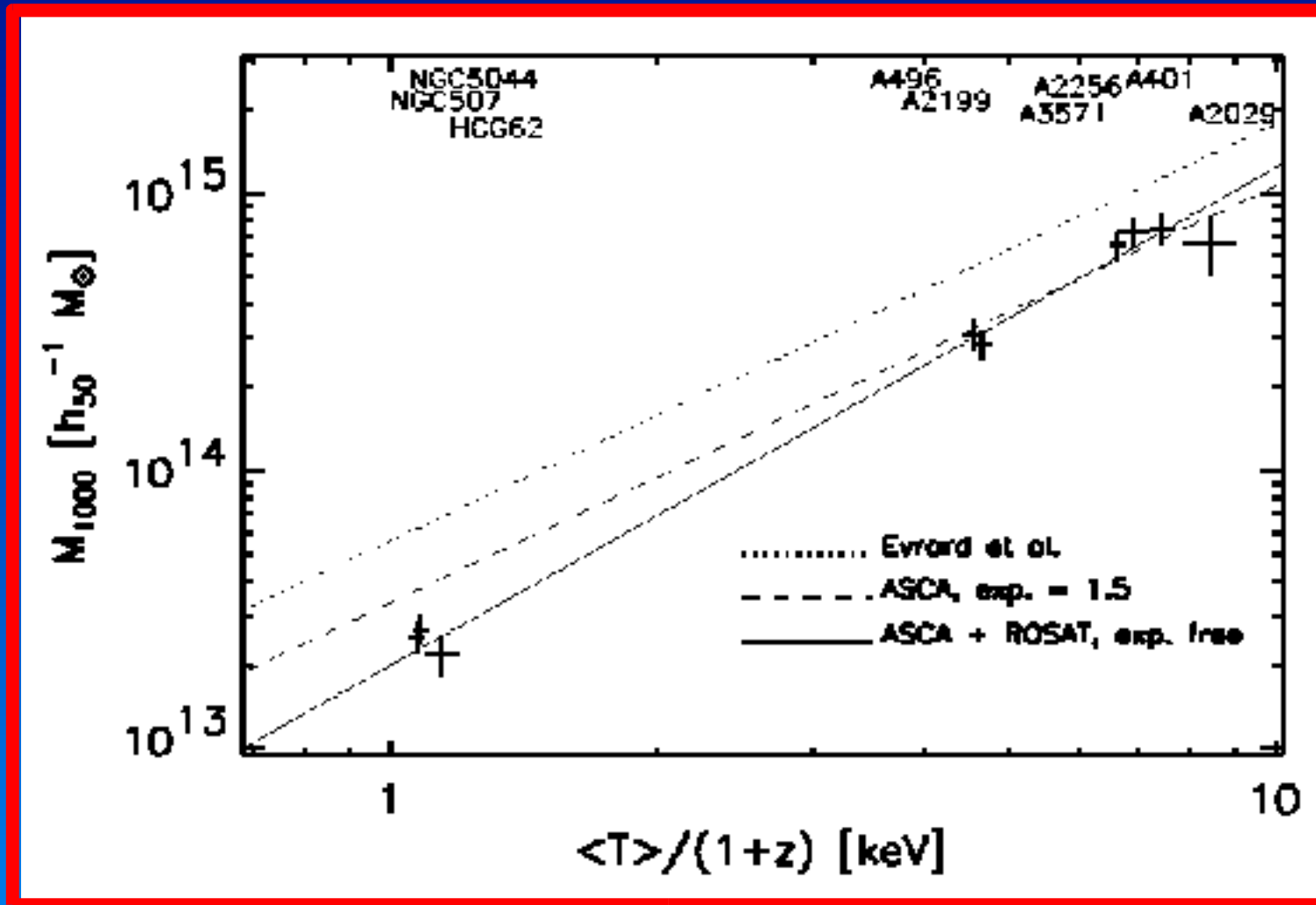
Hydrostatic equilibrium equation

Numerical simulations, Bryan & Norman, 1998



Bryan & Norman (ApJ 495 80 1998)

« Temperature-mass diagram » :



Nevalainen, Markevitch & Forman, 2000

Mass-Luminosity Relation :

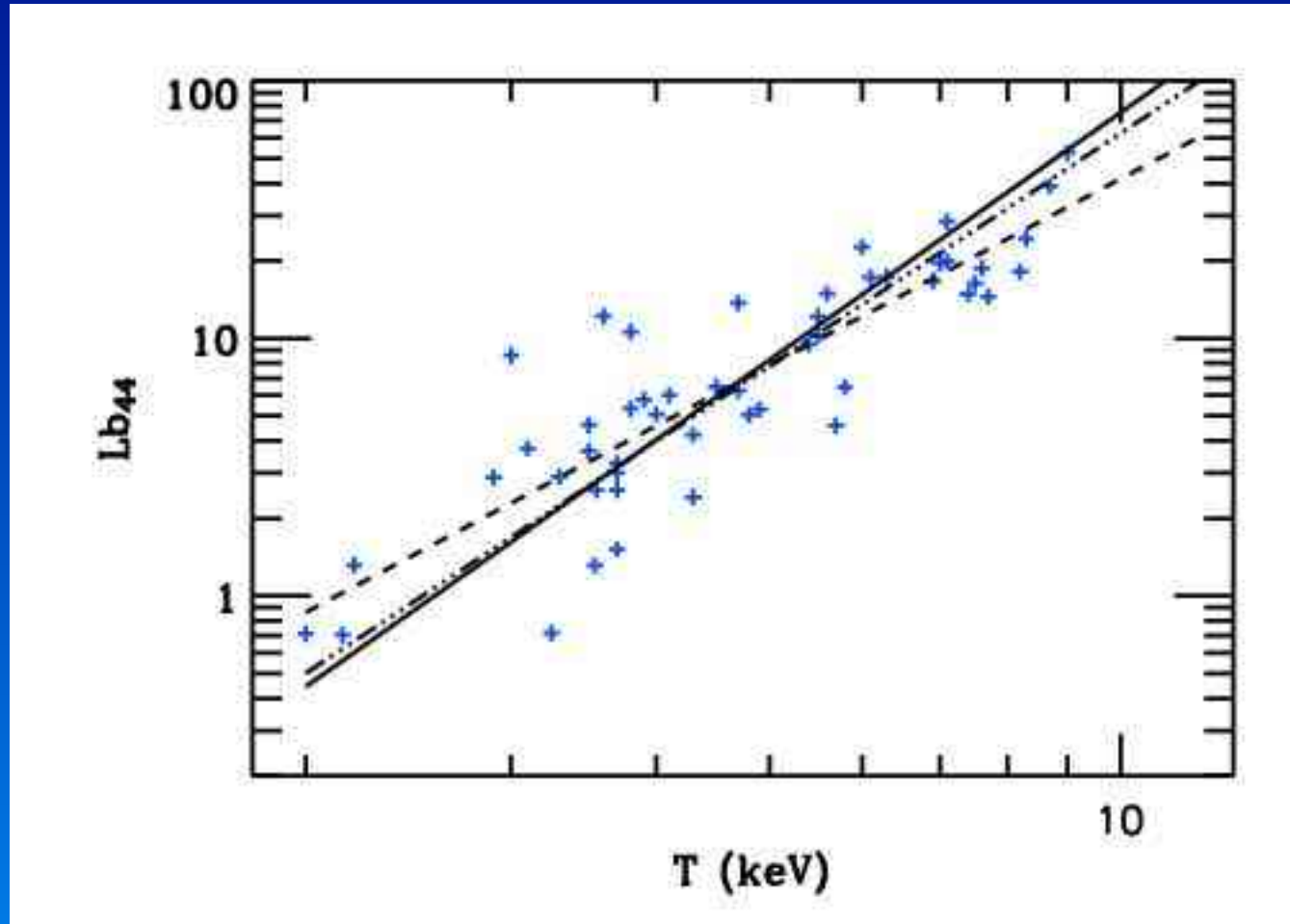
$$L_x \propto n^2 T^{1/2} V$$

...

$$L_x \simeq B M^{4/3} \Omega^{1/6} \Delta^{7/3} (1+z)^{3.5}$$


$$L_x \propto T^2 \Delta^{1/2} (1+z)^{1.5}$$

Observed Temperature -Luminosity Relation



$$L_x \propto T_x^{\sim 3}$$

Conclusion at that point is that

Clusters are not
self similar!

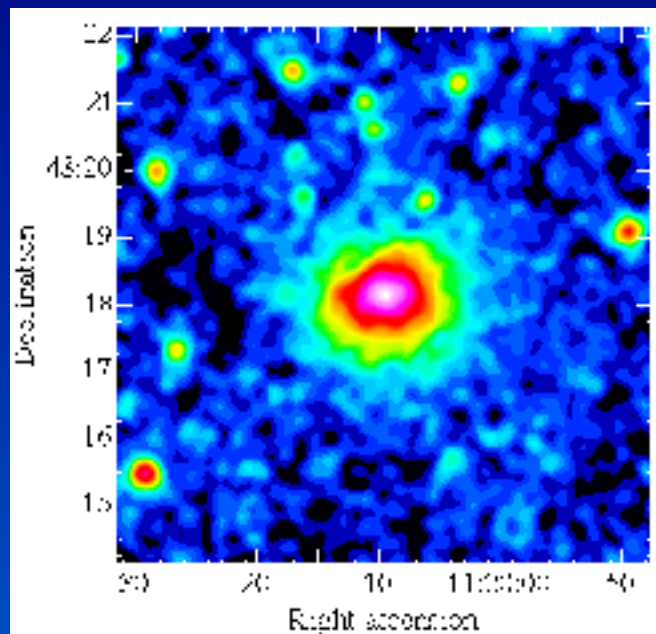
However...

$$\Omega_m$$

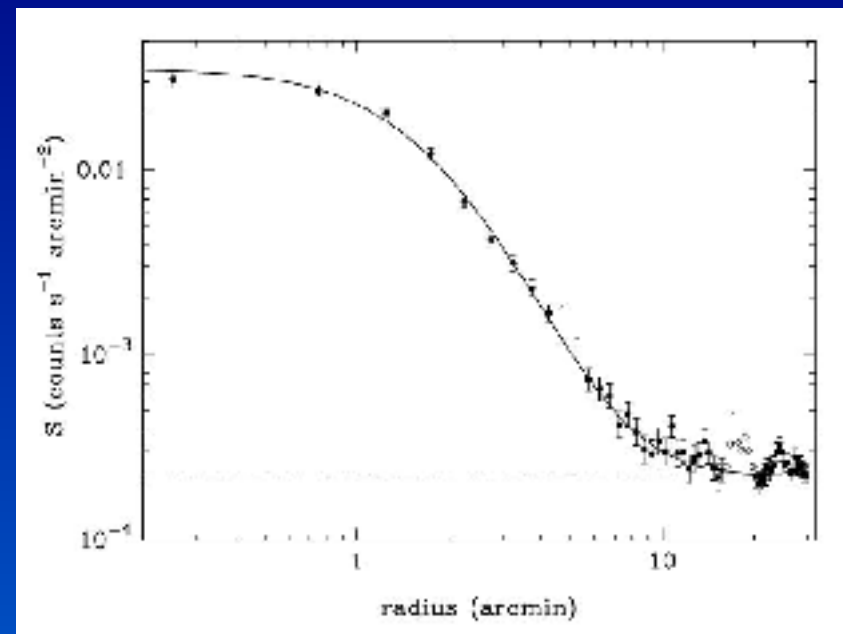
From X-ray Clusters

Baryon Fraction

What do you do with a cluster?

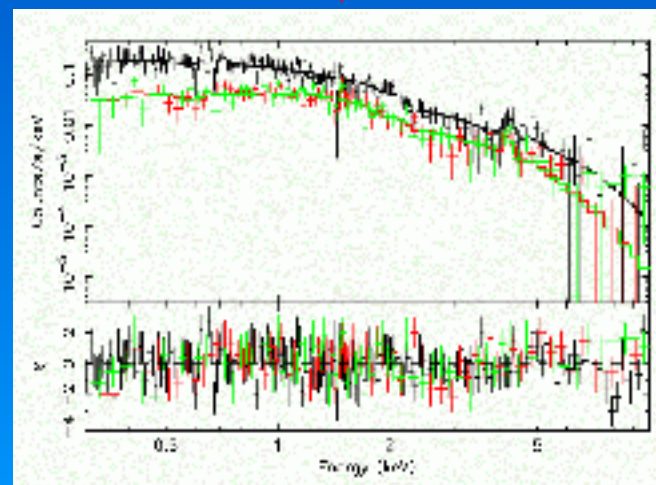


Fit

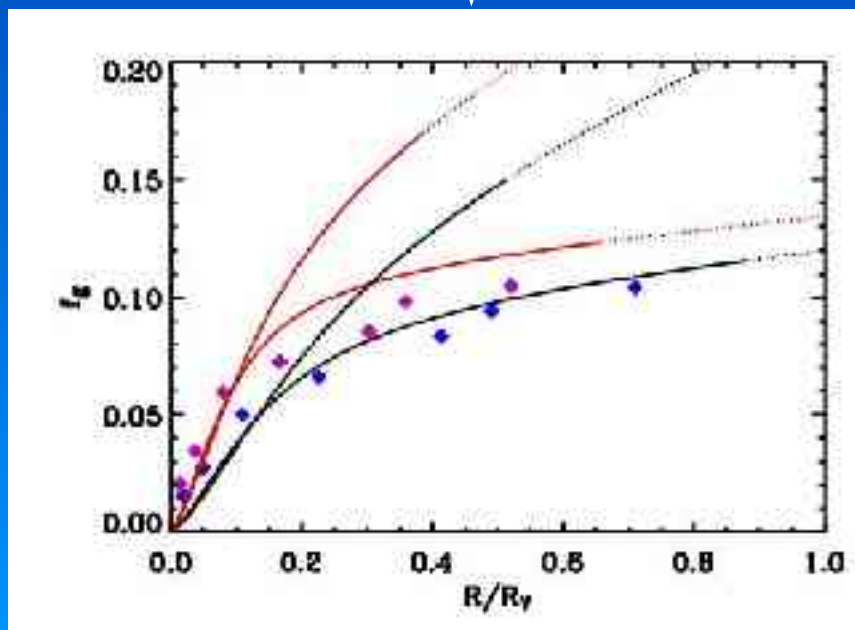


Gas mass $M_g(r)$

X-ray spectrum



$M(r)$



Method :

Ratio :

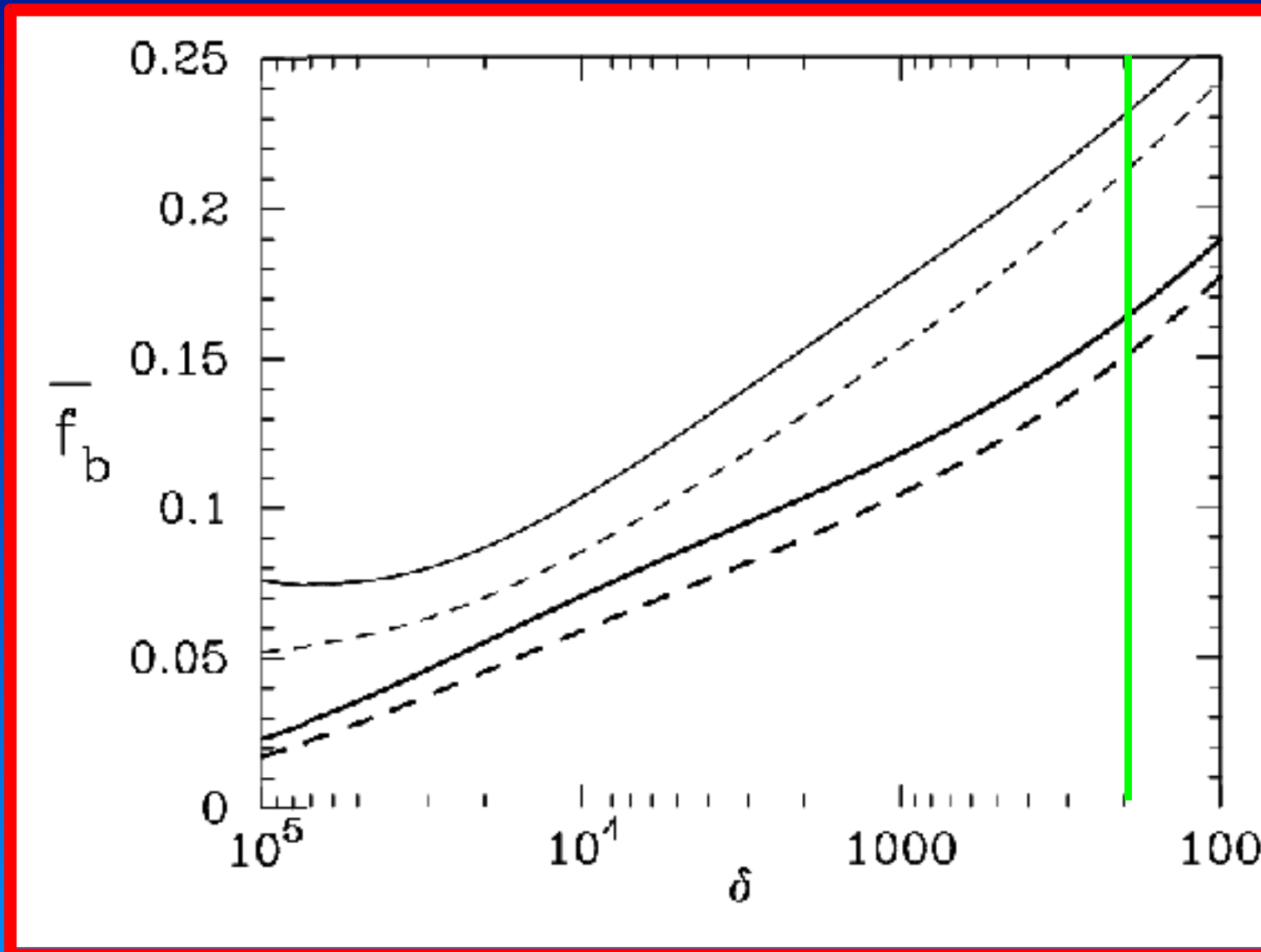
$$f_b = M_b / M_t$$

Observations $\Rightarrow M_g, T_x \Rightarrow M_t$

i.e. : $f_b \approx 15.-20. h_{50}^{-3/2} \%$

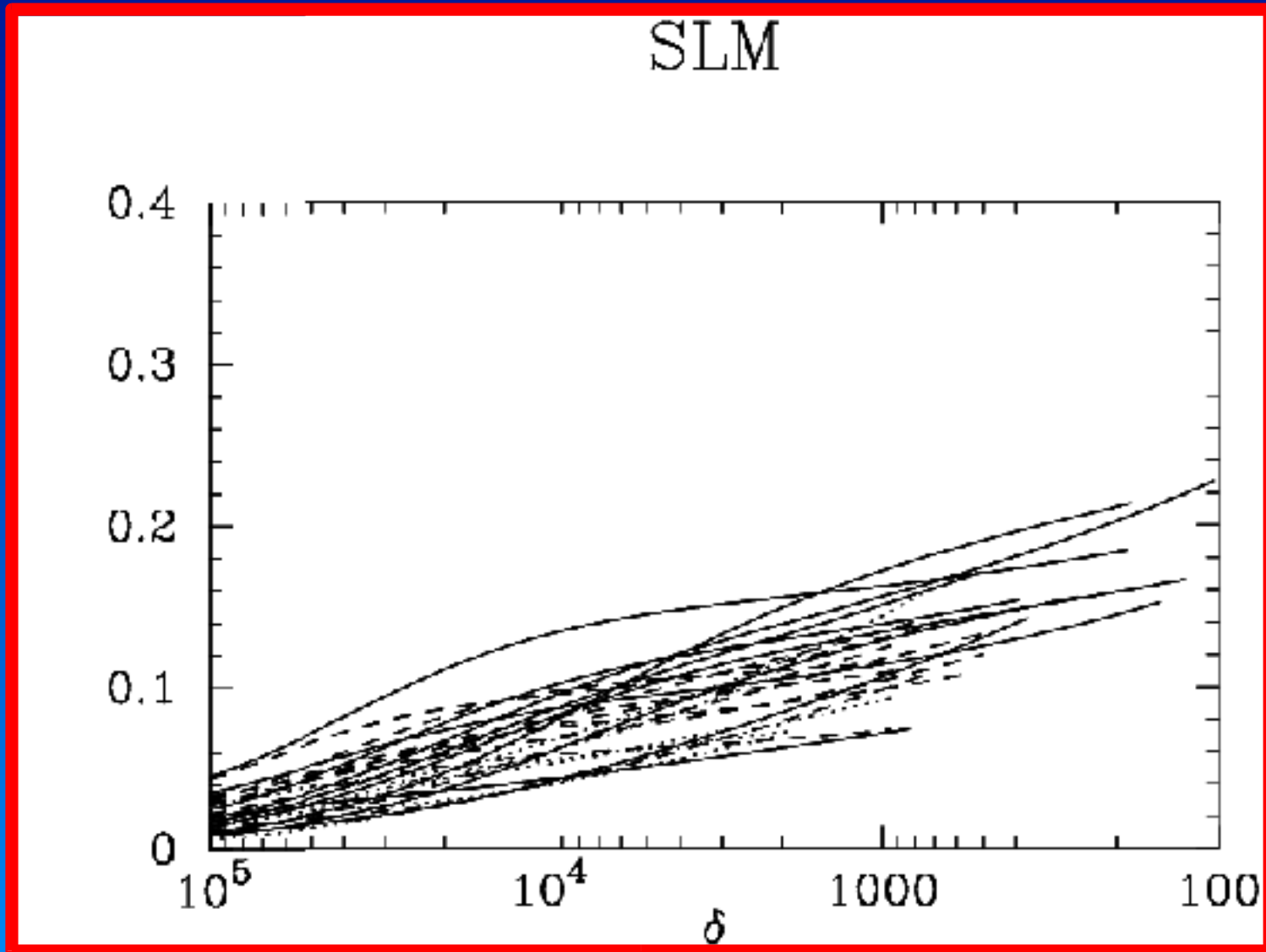
Cosmology : $f_b = \Omega_{bbn} / \Omega_0$

Average baryon fraction :

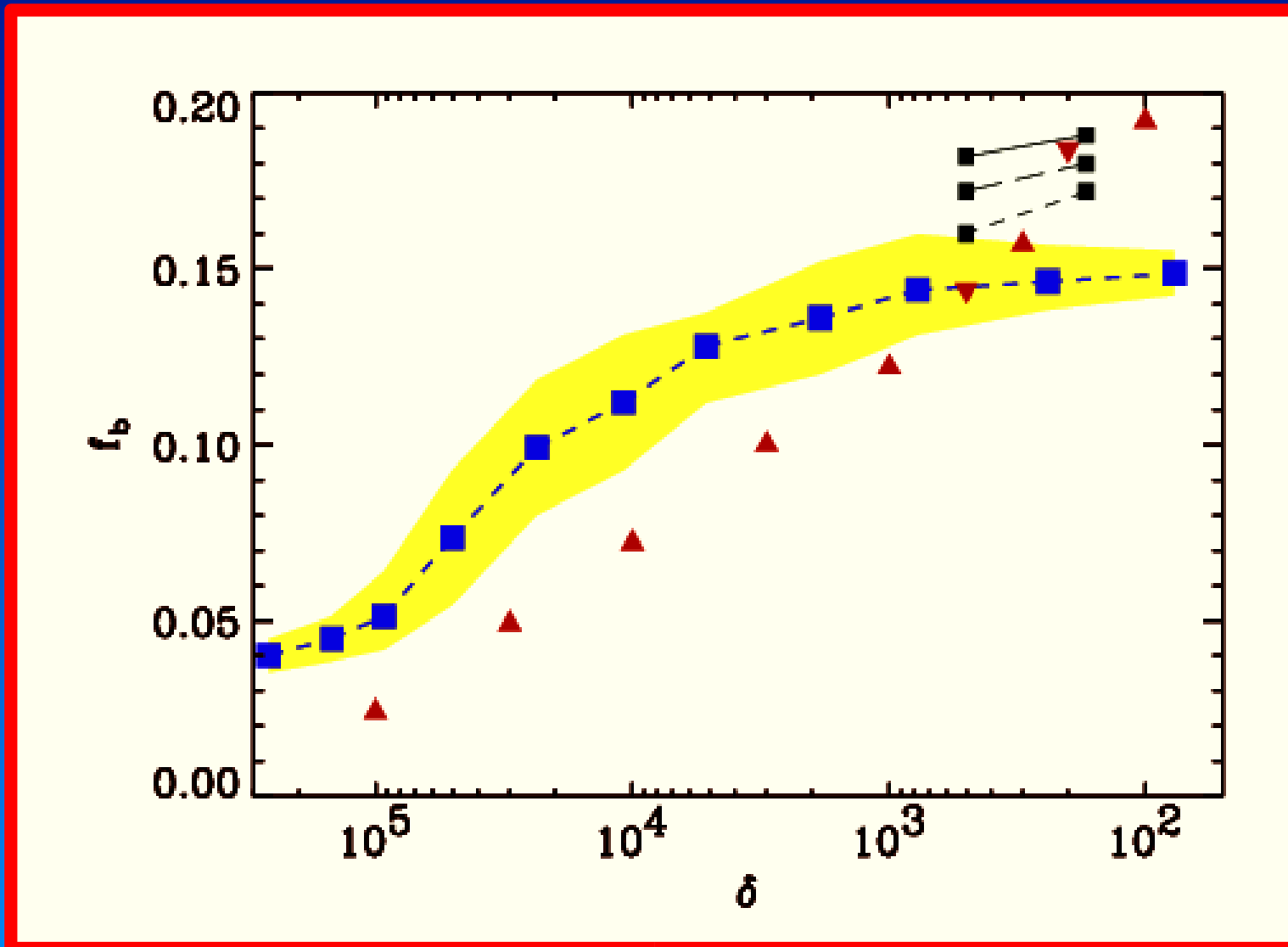


Roussel et al, 2000

Scaling of baryon fractions :



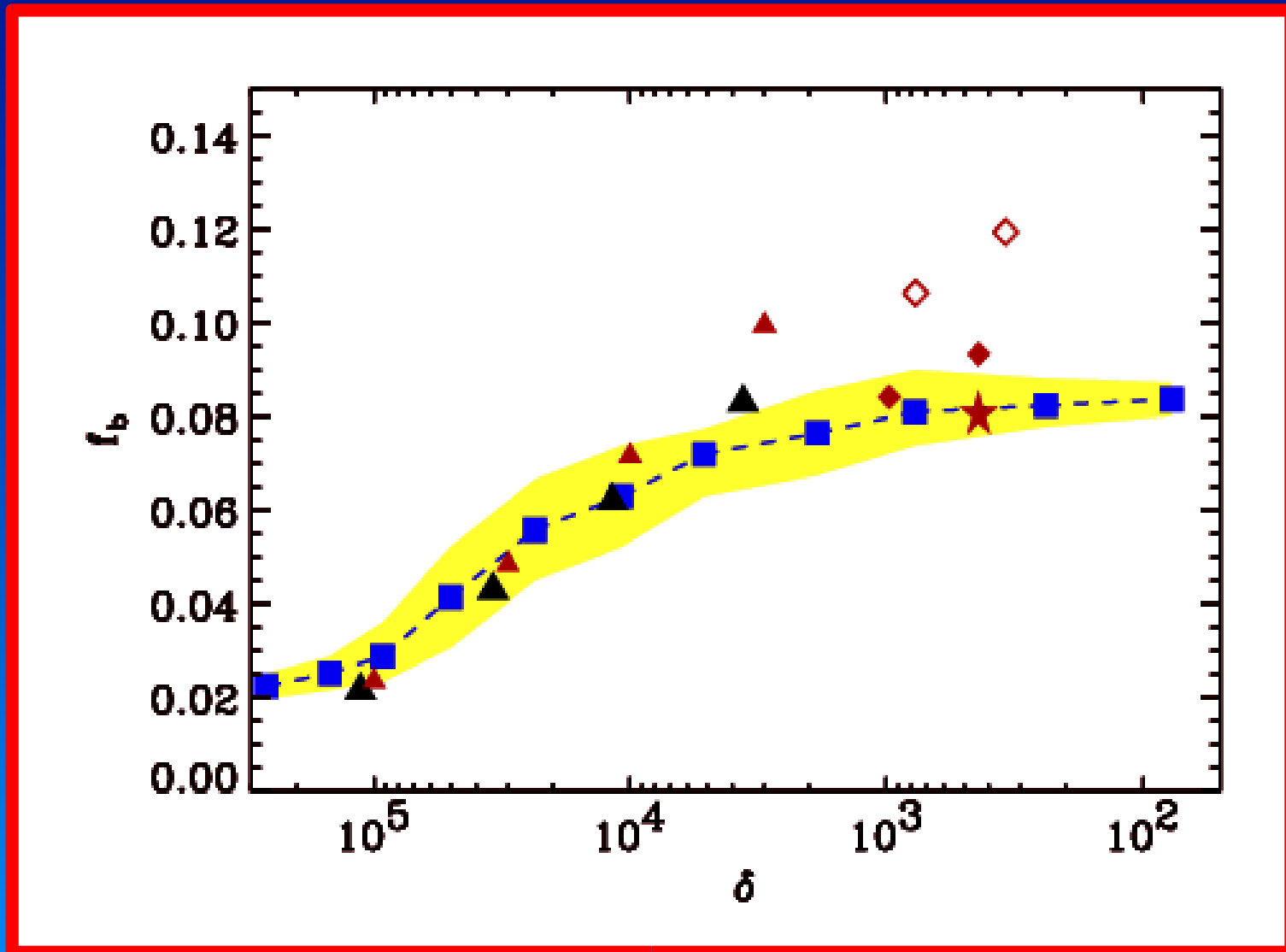
Comparison with simulations:



Three correcting factors :

- ◆ **Mass estimator**
- ◆ **Clumping of the gas**
- ◆ **Outer emission**

Actual comparison :



Conclusion :

→ baryon fraction favors low Ω_m :

→ however outer region of clusters
not so well known...

→ a baryon fraction of the order of $10. h_{50}^{-3/2} \%$
or less could be consistent with data...

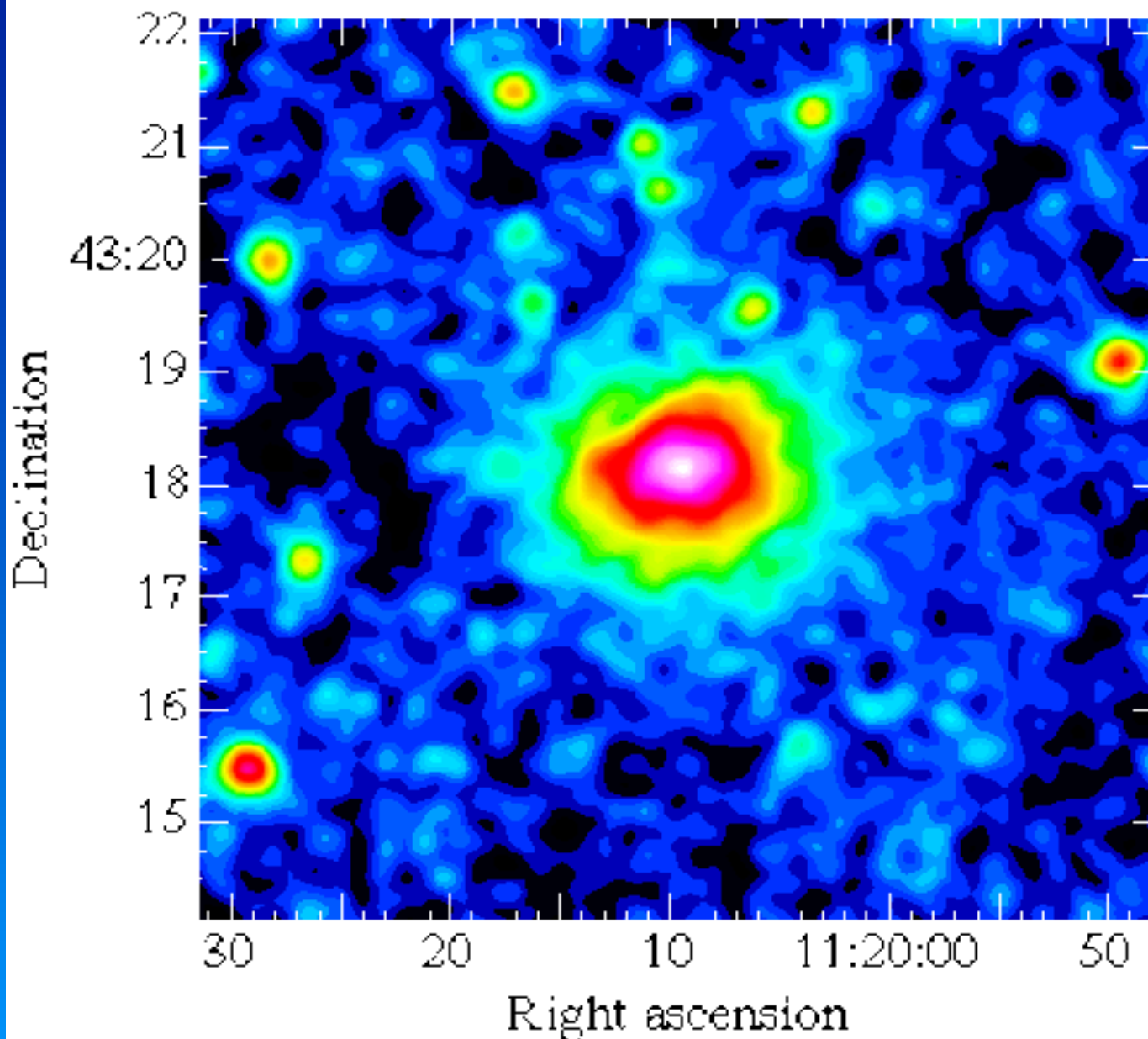
XMM

Ω -Project

**X-ray properties of distant
SHARC clusters
for Cosmology**

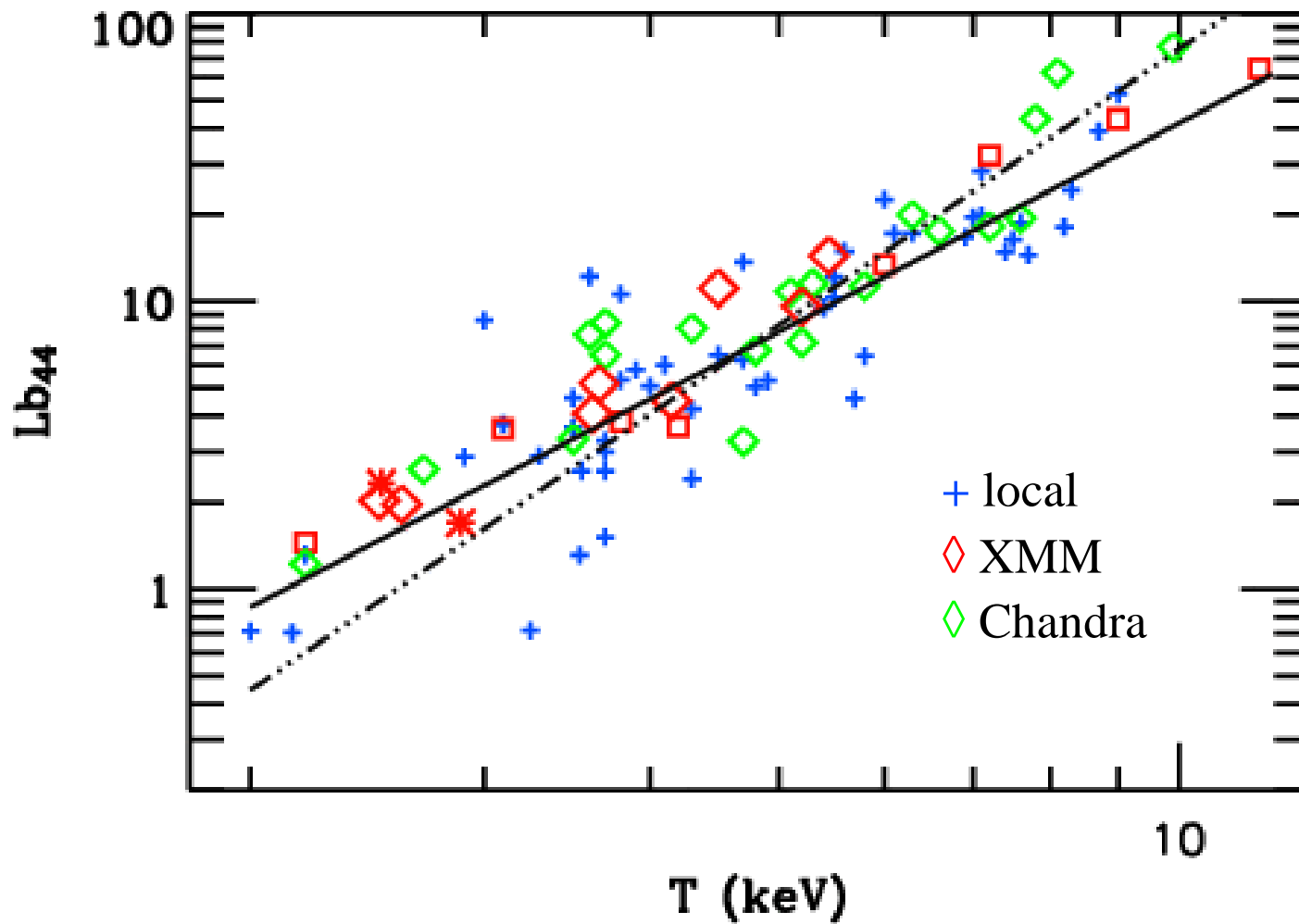
Z ≈ 0.6

X
M
M



R
X
J
1
1
2
0

XMM Lx-Tx evolution



Conclusion on evolution:

❖ remarkable convergence

$$\left(\frac{L_x}{T_x}\right)_z = \left(\frac{L_x}{T_x}\right)_{z=0} (1+z)^\beta$$

with $\beta = 0.65 \pm 0.28$

D.Lumb et al., 2003

in full agreement with ASCA (Sadat et al., 1998; Novicki et al., 2003;...), Chandra (Vikhlinin et al, 2002), new XMM analysis (Kotov & Vikhlinin, 2005)

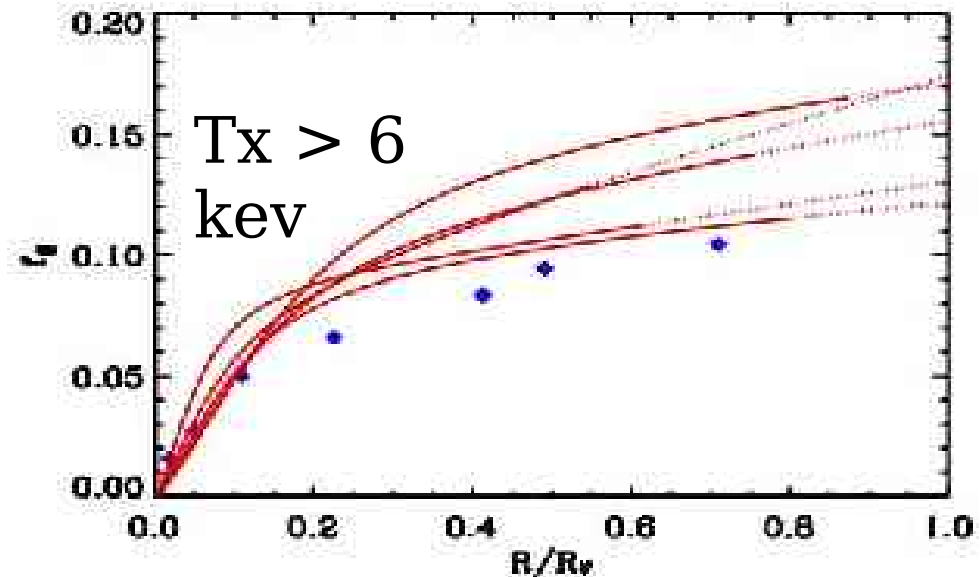
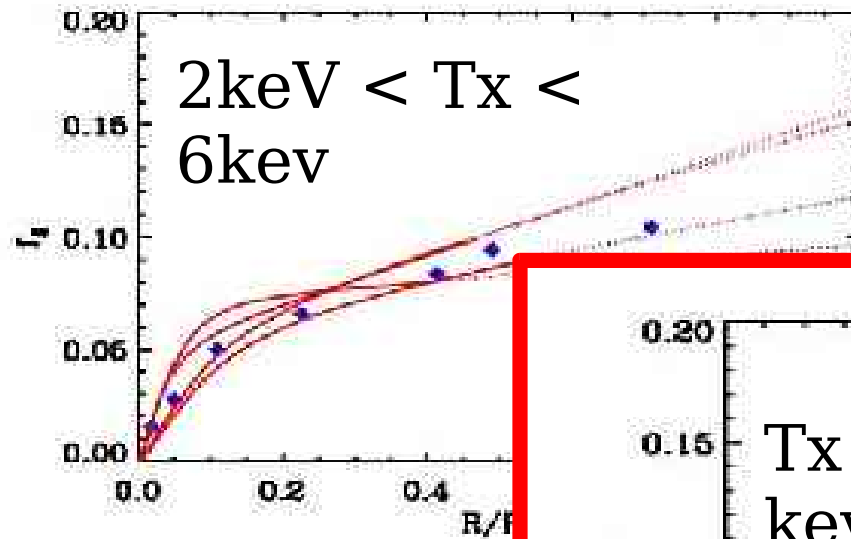
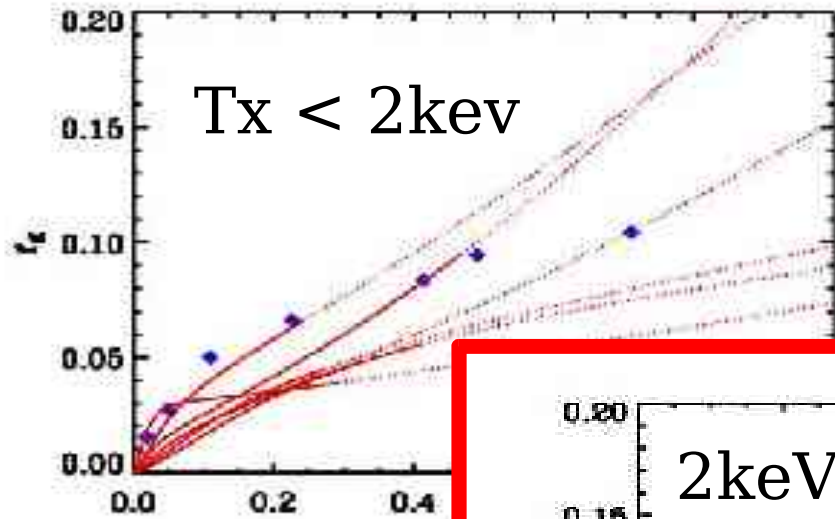
$$\Omega_m$$

From X-ray Clusters

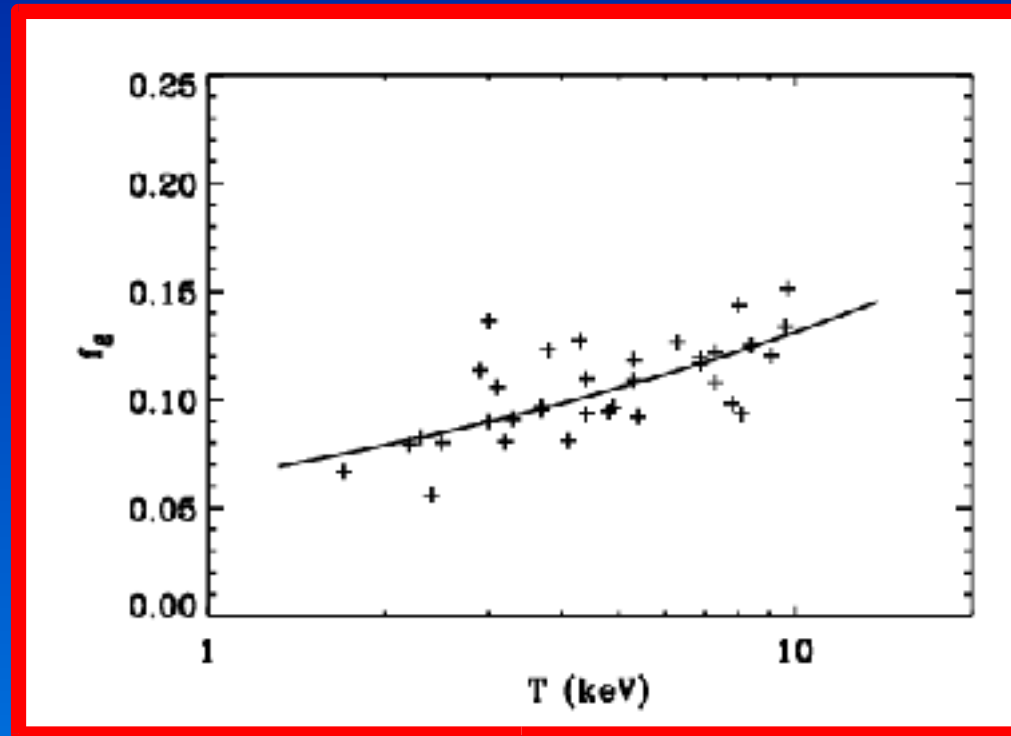
Baryon Fraction evolution
in the XMM Ω -project

(Sadat et al., 2005, A&A)

Baryon Fraction @ $z = 0$

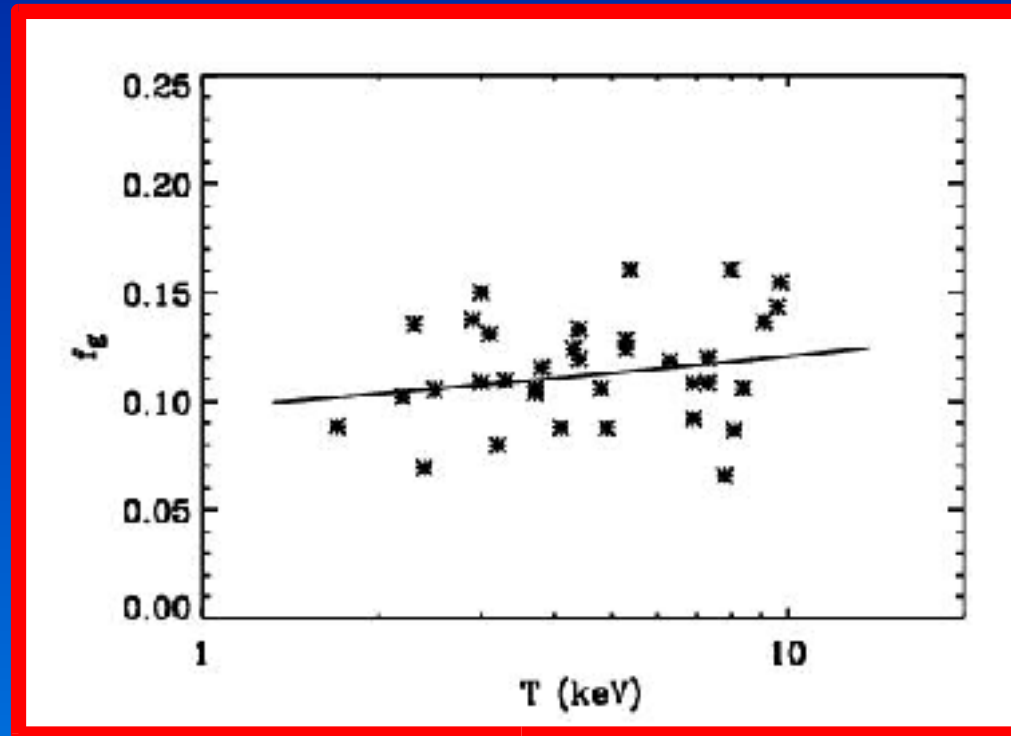


Baryon Fraction @ $z = 0$



R_{2000} in Vikhlinin, Forman, Jones 1999 ($\sim 35-45\%$
Rv)

Baryon Fraction @ $z = 0$

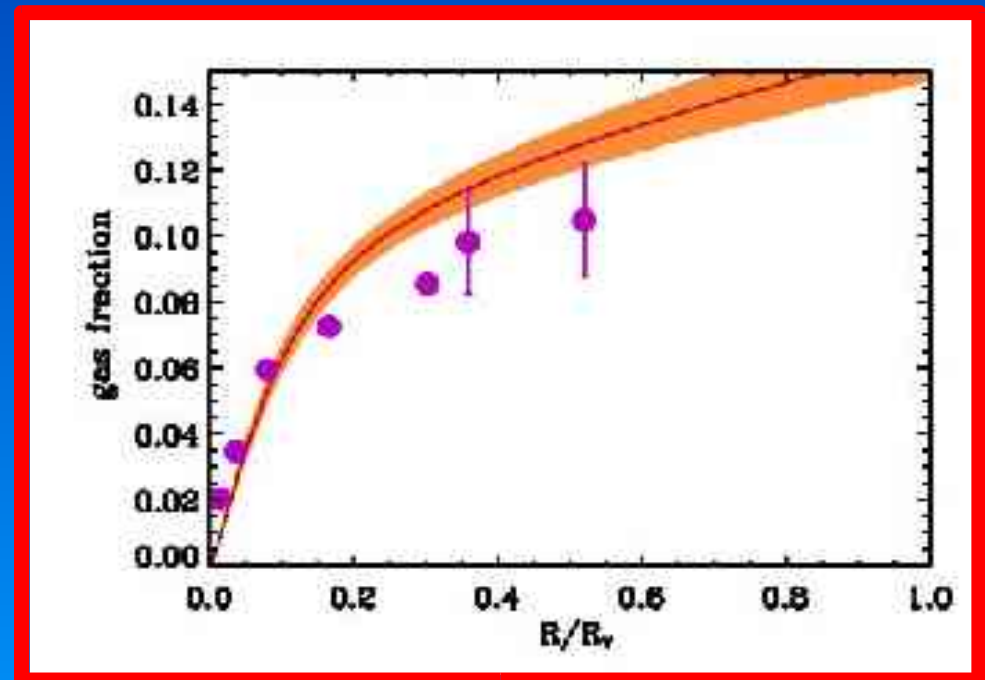
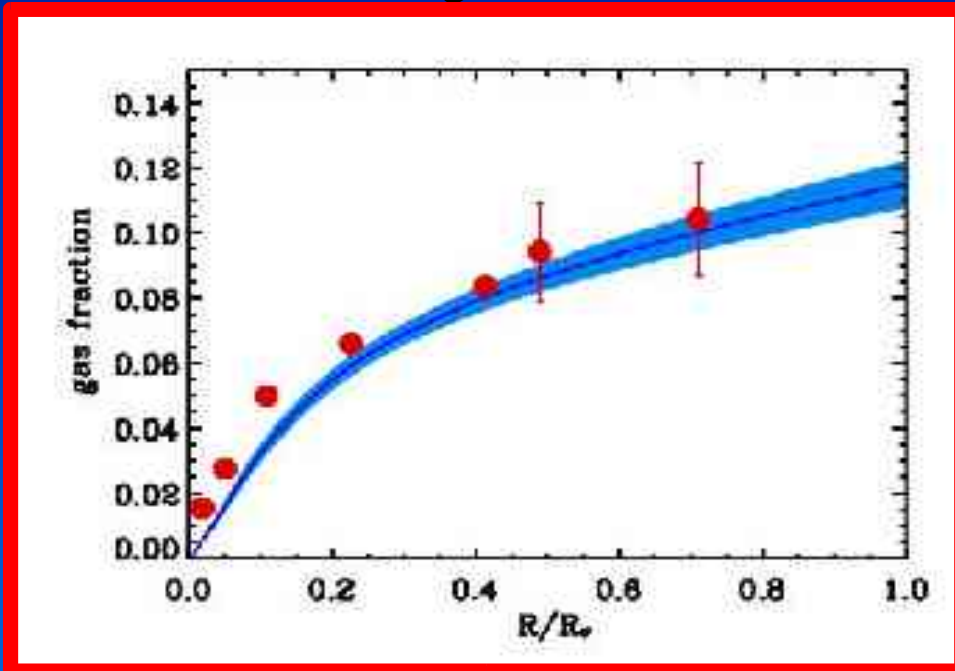


R_v in Vikhlinin, Forman, Jones 1999

Baryon Fraction @ $z = 0.6$

Internal structure

is complex...

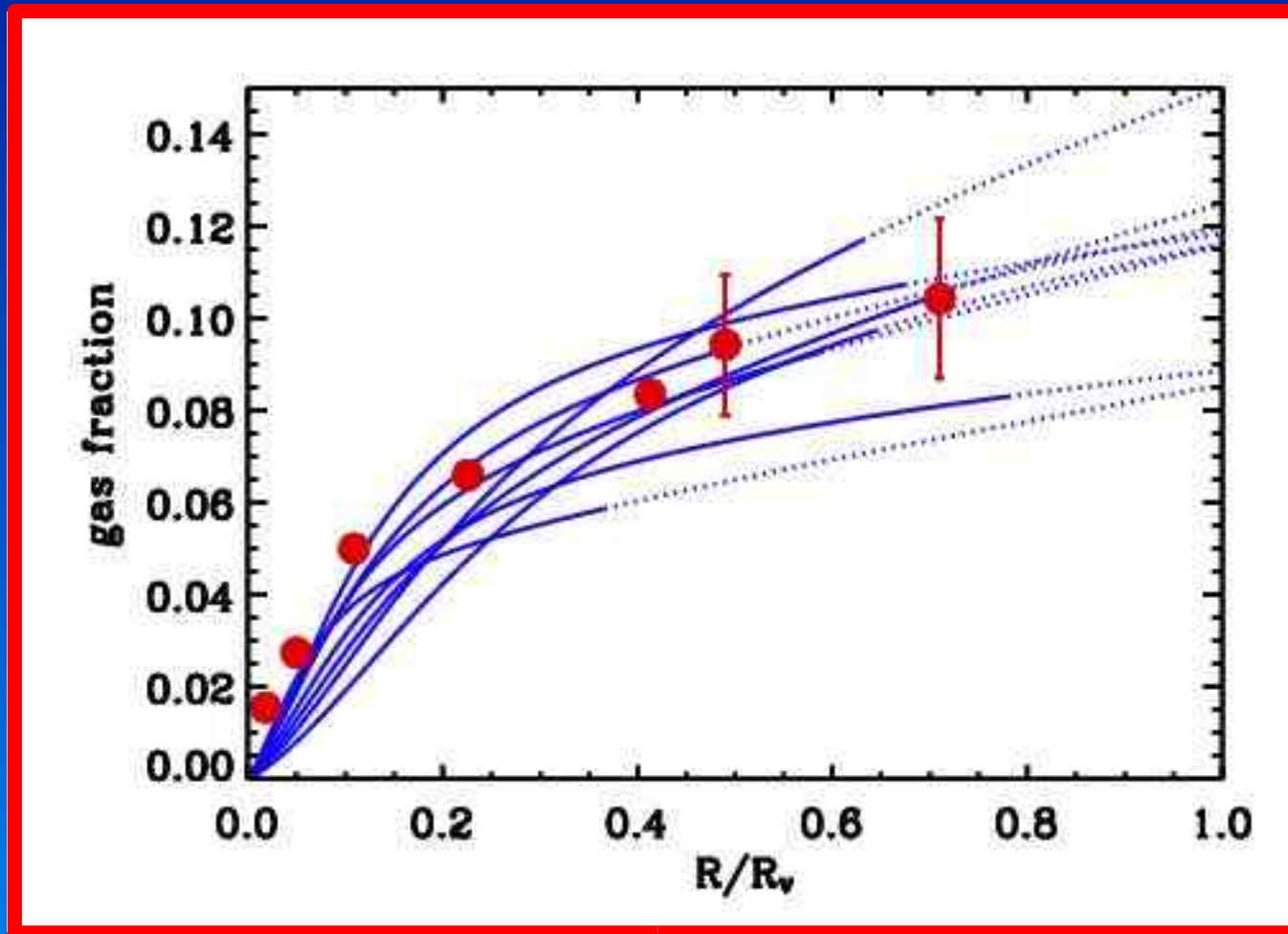


Conclusion at that point is that

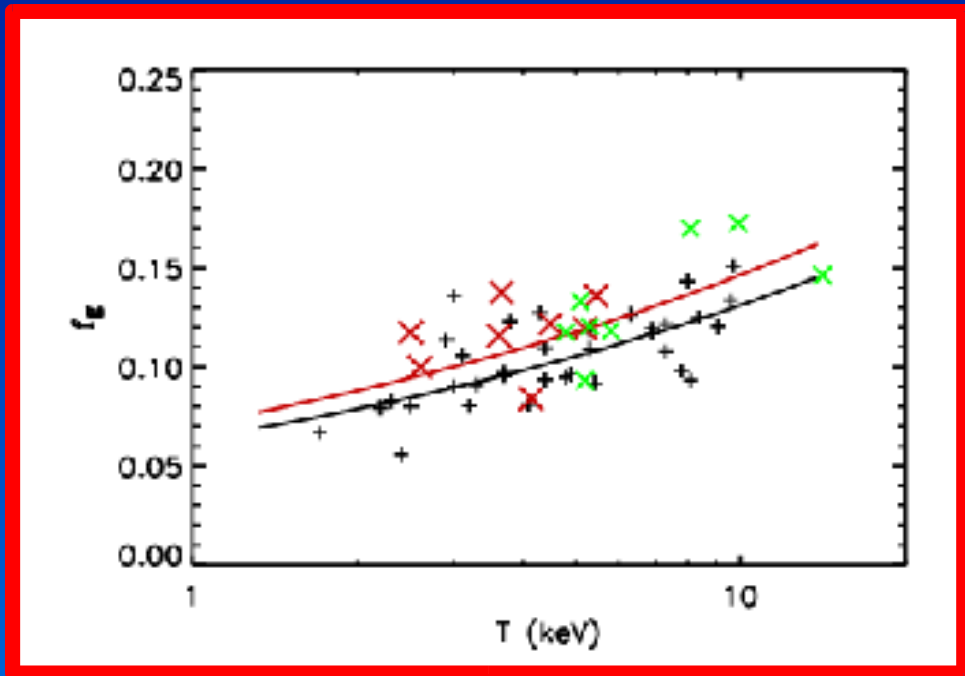
Clusters are not
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However...

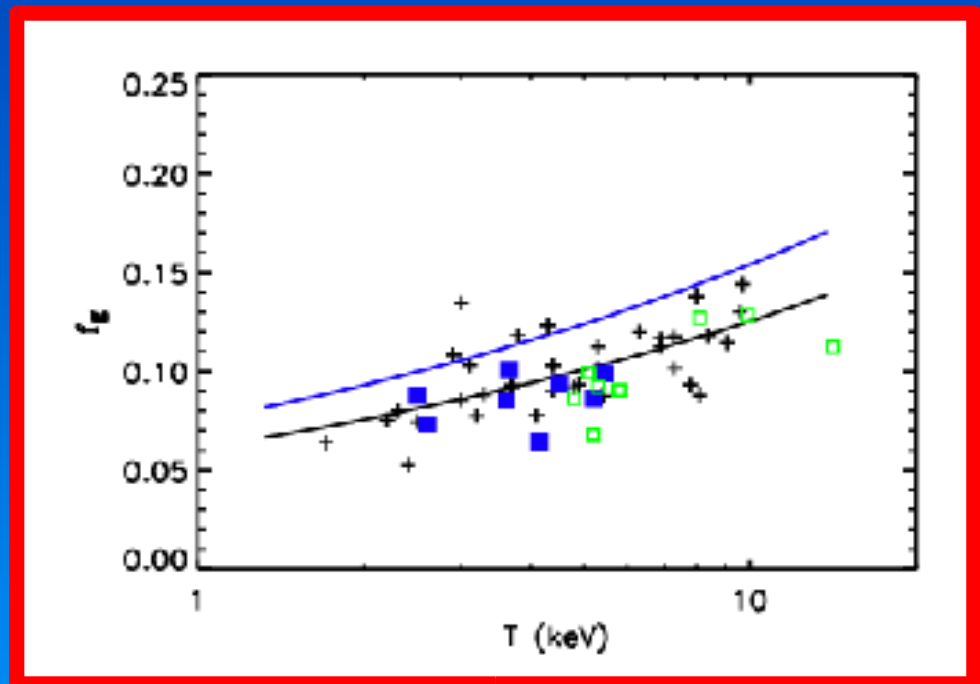
Baryon Fraction @ $z = 0.6$



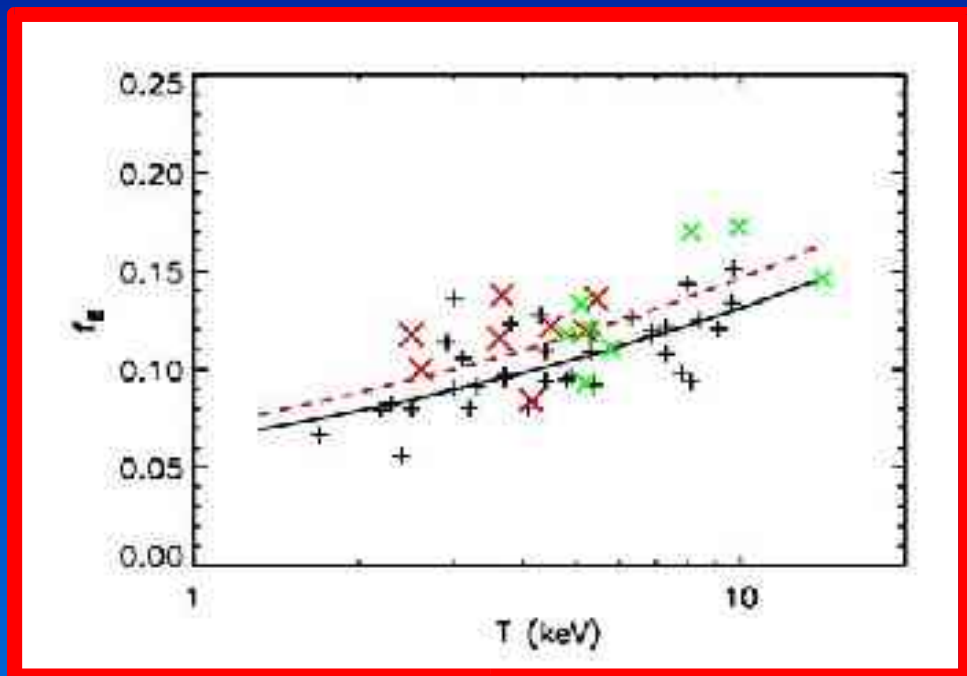
Baryon Fraction @ $z = 0.6$



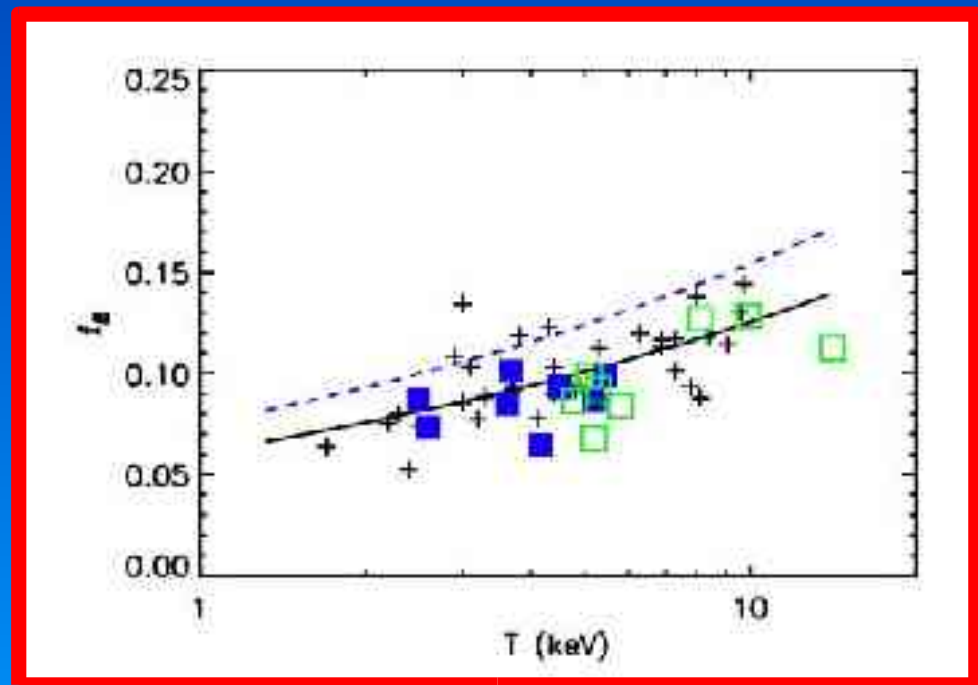
$\Delta = 2000$



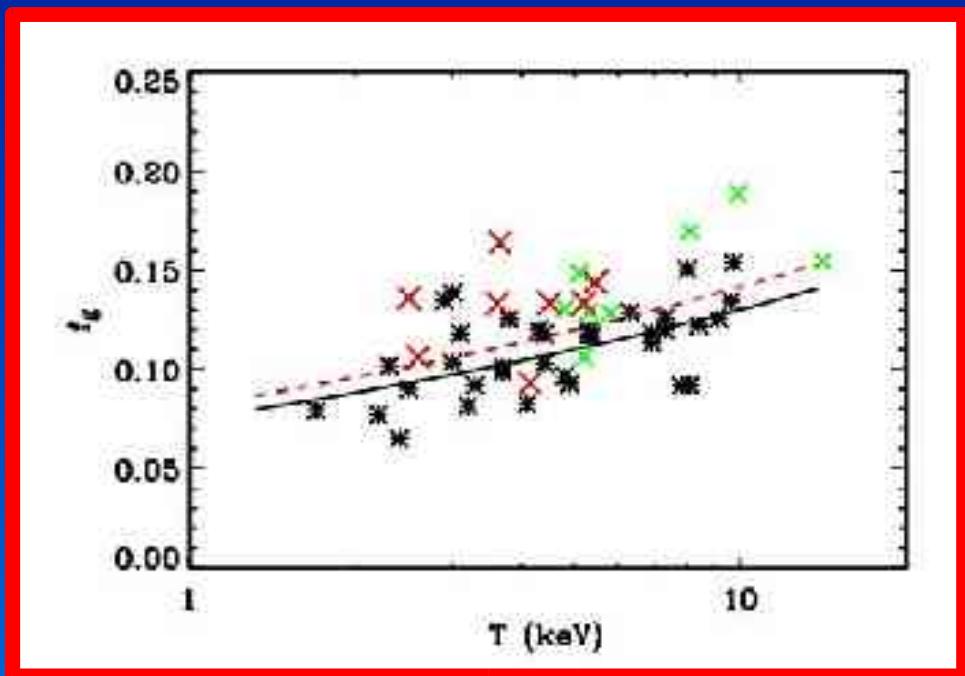
Baryon Fraction @ $z = 0.6$



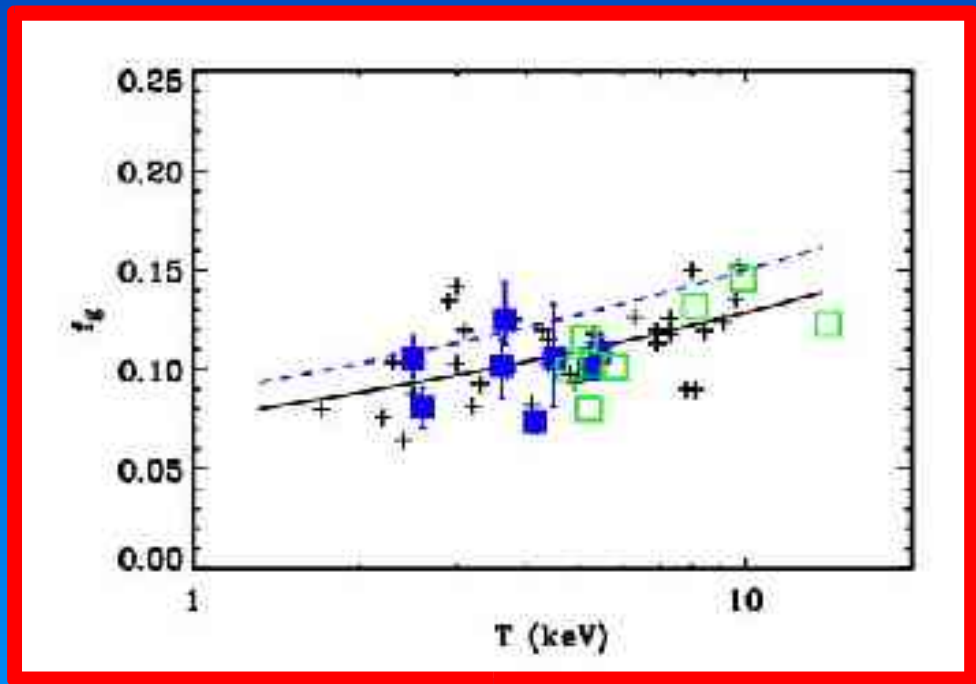
$$\Delta = 2000$$



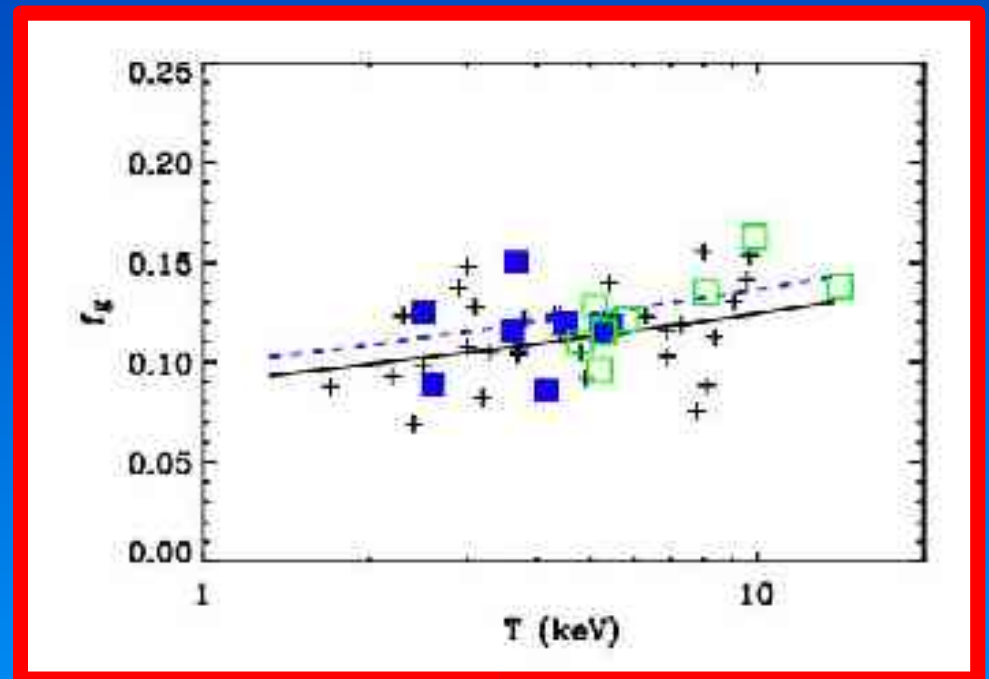
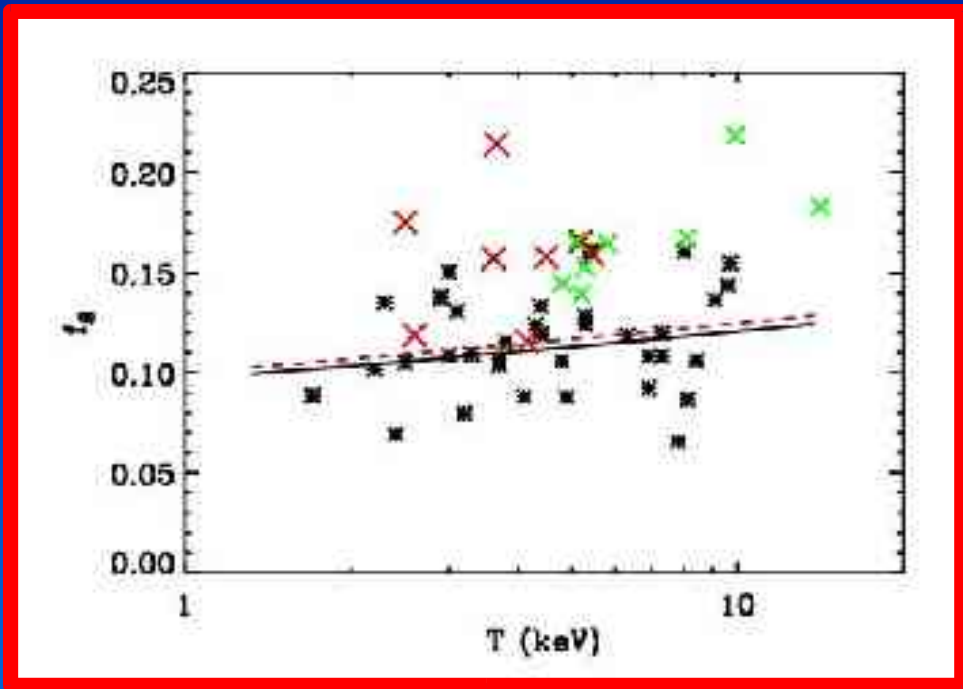
Baryon Fraction @ $z = 0.6$



$\Delta = 1000$



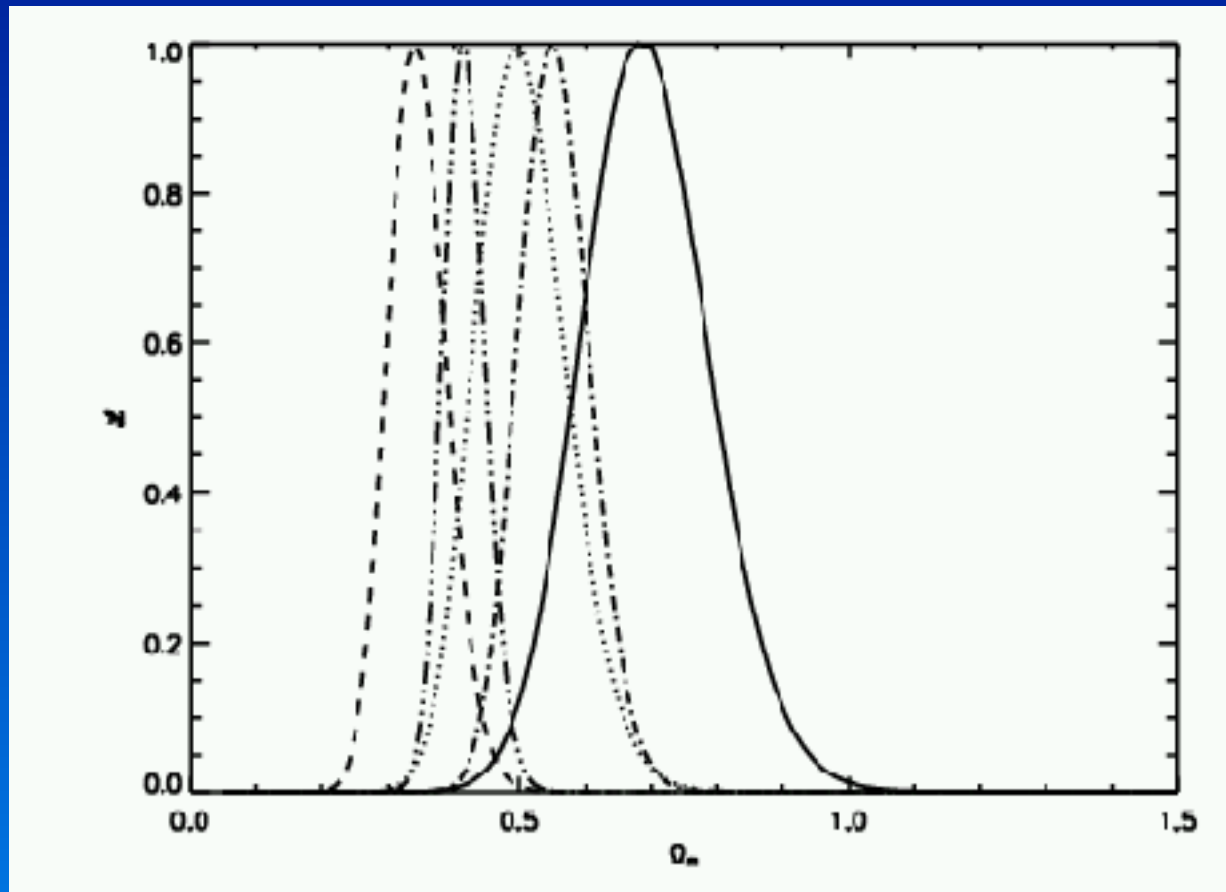
Baryon Fraction @ $z = 0.6$



$$\Delta_v$$

... favors high Ω_m !

Likelihood analysis:



(Ferramacho et al., 2006, submitted)

Theory of the mass function

Basics

Matter = random field:

$$\rho(x) = \bar{\rho}(1 + \delta(x))$$

δ mathematically ill-behaved...

1. The field is smoothed = ρ is convolved:

$$\tilde{\delta} = \delta * W_R$$

with a window function:

$$\int W_R(u) du = 1$$

the smoothed field:

$$\bar{\delta}(x) = \int \delta(x + u) W_R(u) du$$

Variance of the field :

$$\overline{\tilde{\delta}^2(x)} = \sigma^2(R)$$

Ex: top-hat window:

$$\begin{aligned} W_R(u) &= 1/V & \text{for } |u| < R \\ W_R(u) &= 0. & \text{for } |u| \geq R \end{aligned}$$

Mass associated:

$$M(R) = \frac{4\pi}{3} R^3 \bar{\rho}$$

2. Nonlinear model

- linear overdensity

$$\tilde{\delta}(x) \sim 1$$

Ex: spherical model $\delta_{\text{NL}}(\mathbf{z}, \Omega_m, \Omega_\lambda \dots)$

3. Mass function:

- trivially (!): dV will be in an object with mass $> M$ if included in a NL fluctuation of $\tilde{\delta}_R$ with radius $> R$

$$\int_M^{+\infty} m n(m) dm = \bar{\rho} \int \mathcal{F}_\delta(\delta) s(\delta) d\delta \sim \bar{\rho} \int_{\delta_{NL}}^{+\infty} \mathcal{F}_\delta(\delta) d\delta$$

or (sharp threshold):

$$\int_M^{+\infty} mn(m)dm = \bar{\rho} \int_{\delta_{NL}}^{+\infty} \mathcal{F}_\delta(\delta)d\delta = \bar{\rho} \int_{\nu_{NL}}^{+\infty} \mathcal{F}(\nu)d\nu$$

with:

$$\delta = \nu\sigma(R) = \nu\sigma(M) \quad \text{and}$$

$$\nu_{NL} = \frac{\delta_{NL}}{\sigma(M)}$$

and just take the derivative...

the mass function:

$$N(M) = -\frac{\rho}{M^2 \sigma(M)} \delta_{NL} \frac{\ln \sigma}{\ln M} \mathcal{F}(\nu_{NL})$$

normalization condition:

$$\frac{1}{\bar{\rho}} \int_0^{+\infty} m n(m) dm = \int_0^{+\infty} \mathcal{F}(\nu) d\nu = 1$$

Press and Schechter (1974) used:

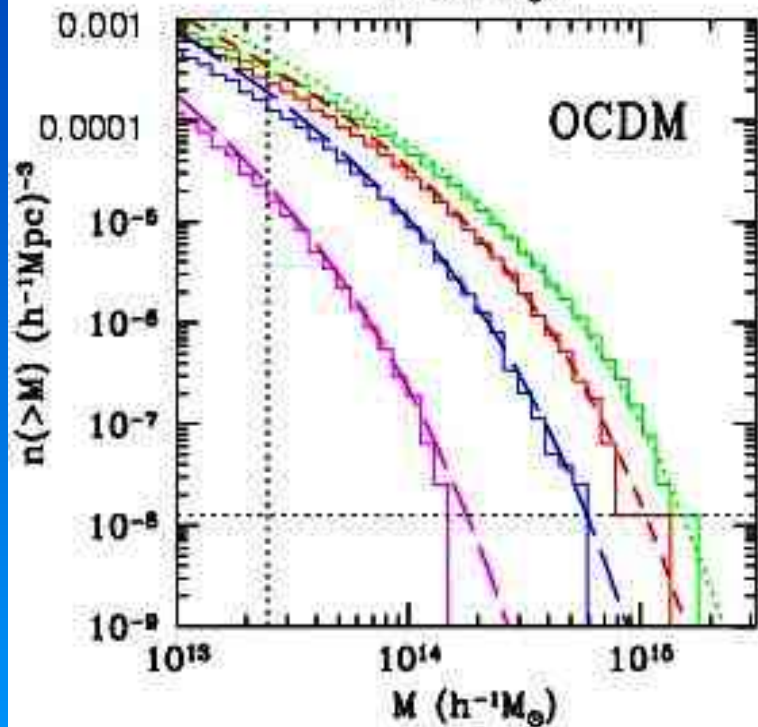
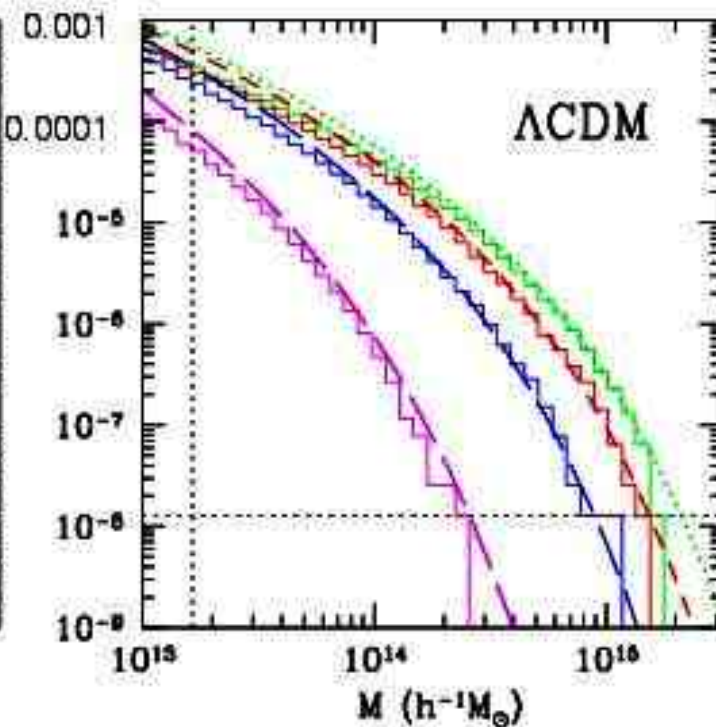
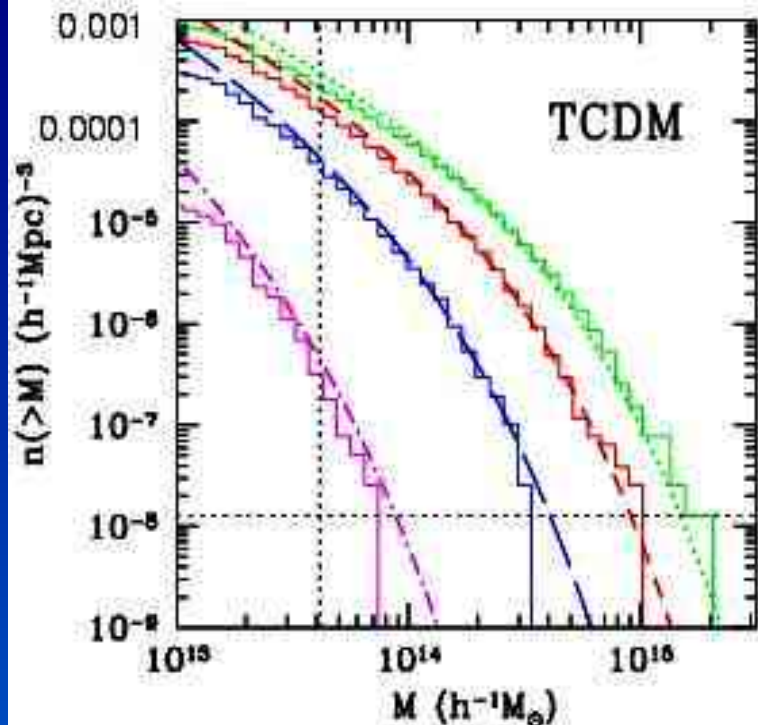
$$\mathcal{F}(\nu) = \sqrt{\frac{2}{\pi}} \exp\left(-\frac{\nu^2}{2}\right)$$

major recent improvements:

$$\mathcal{F}(\nu) = \sqrt{\frac{2A}{\pi}} C \exp(-0.5A\nu^2) (1. + (1./ (A\nu)^2)^Q)$$

with

$$A = 0.707 \quad C = 0.3222 \quad Q = 0.3$$

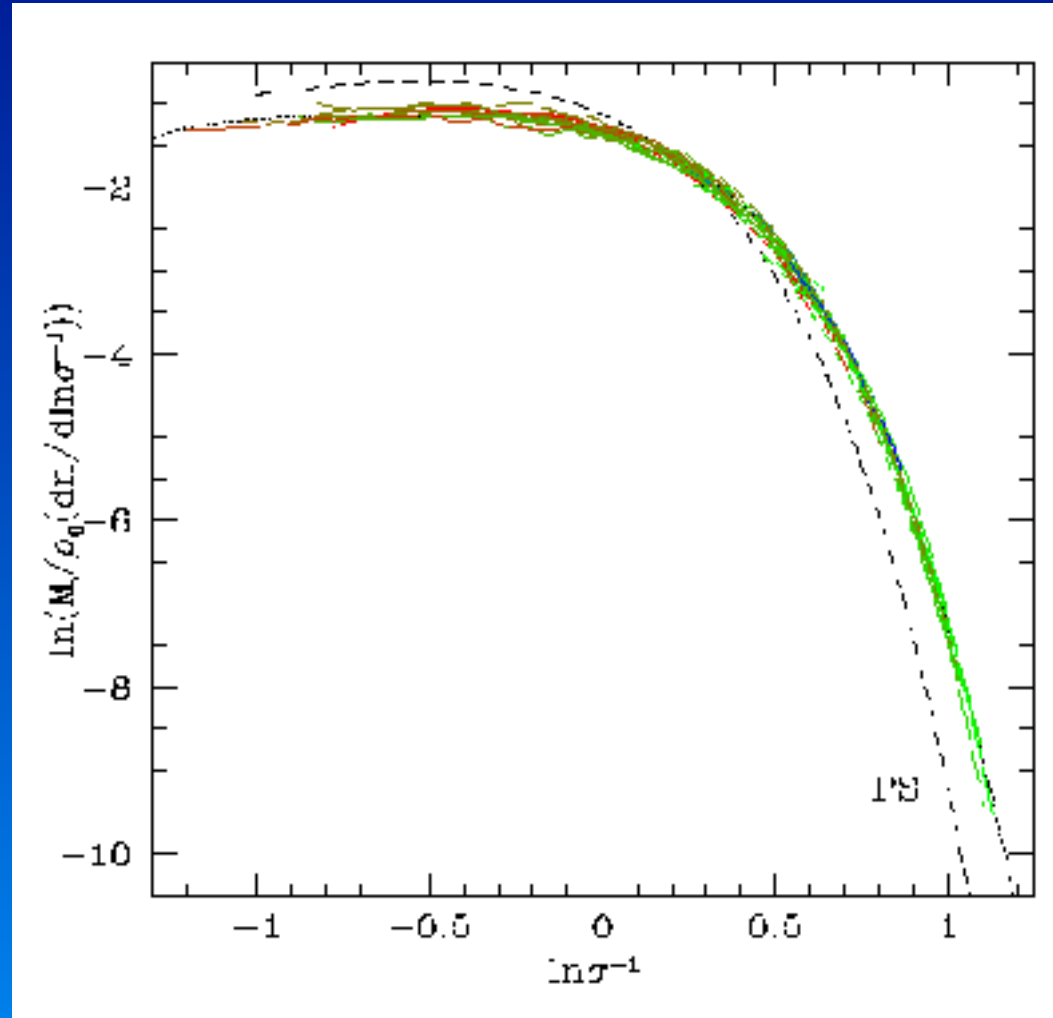


$L_{\text{box}} = 250 \text{ h}^{-1}\text{Mpc}$

$N_{\text{part}} = N_{\text{gr}} = 128^3$

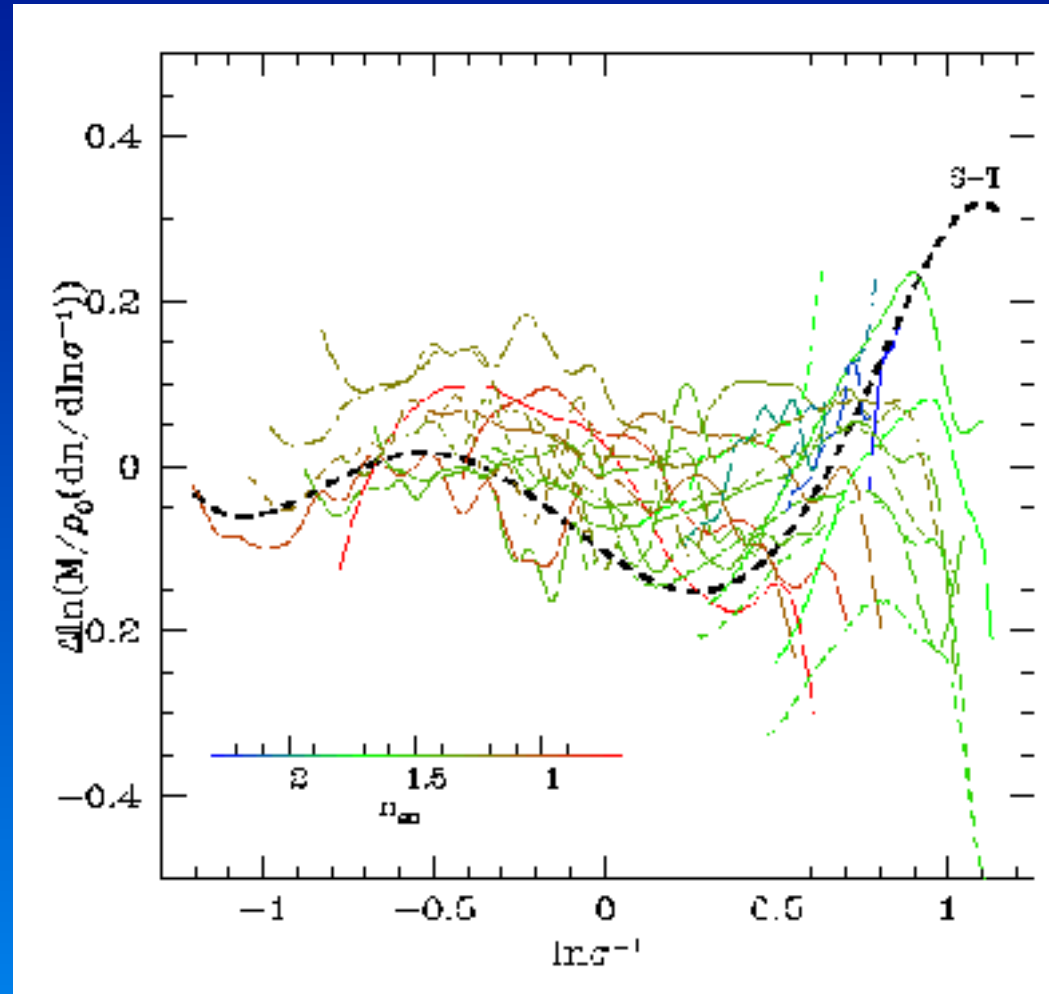
- $z=0$
- $z=0.21$
- $z=0.55$
- $z=1.40$

Checking N(M)



Jenkins et al., 2001 MNRAS, 321, 372

Checking N(M) (2)



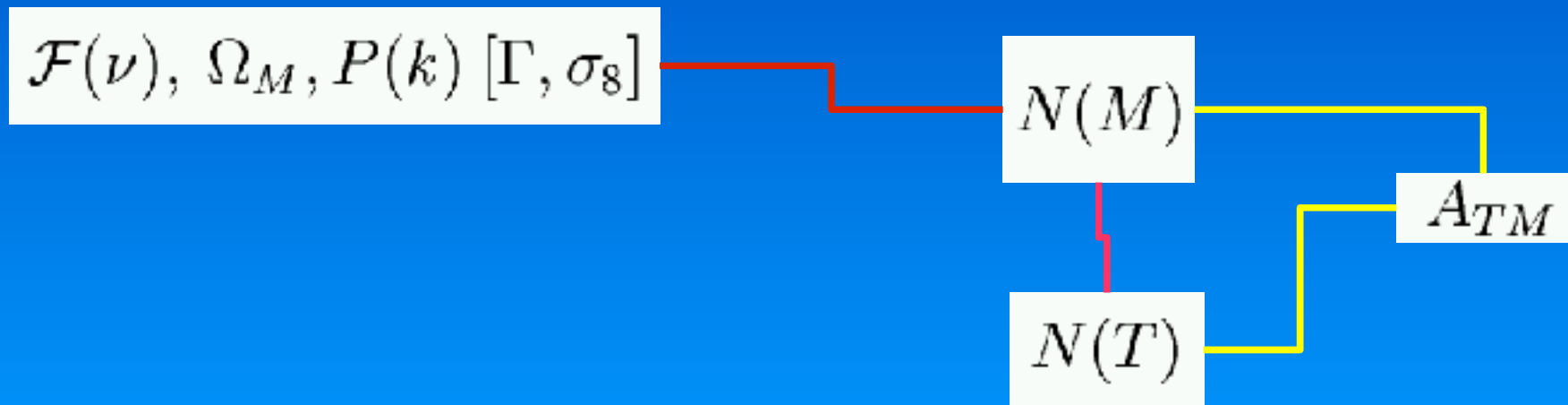
Jenkins et al., 2001 MNRAS, 321, 372

Conclusion

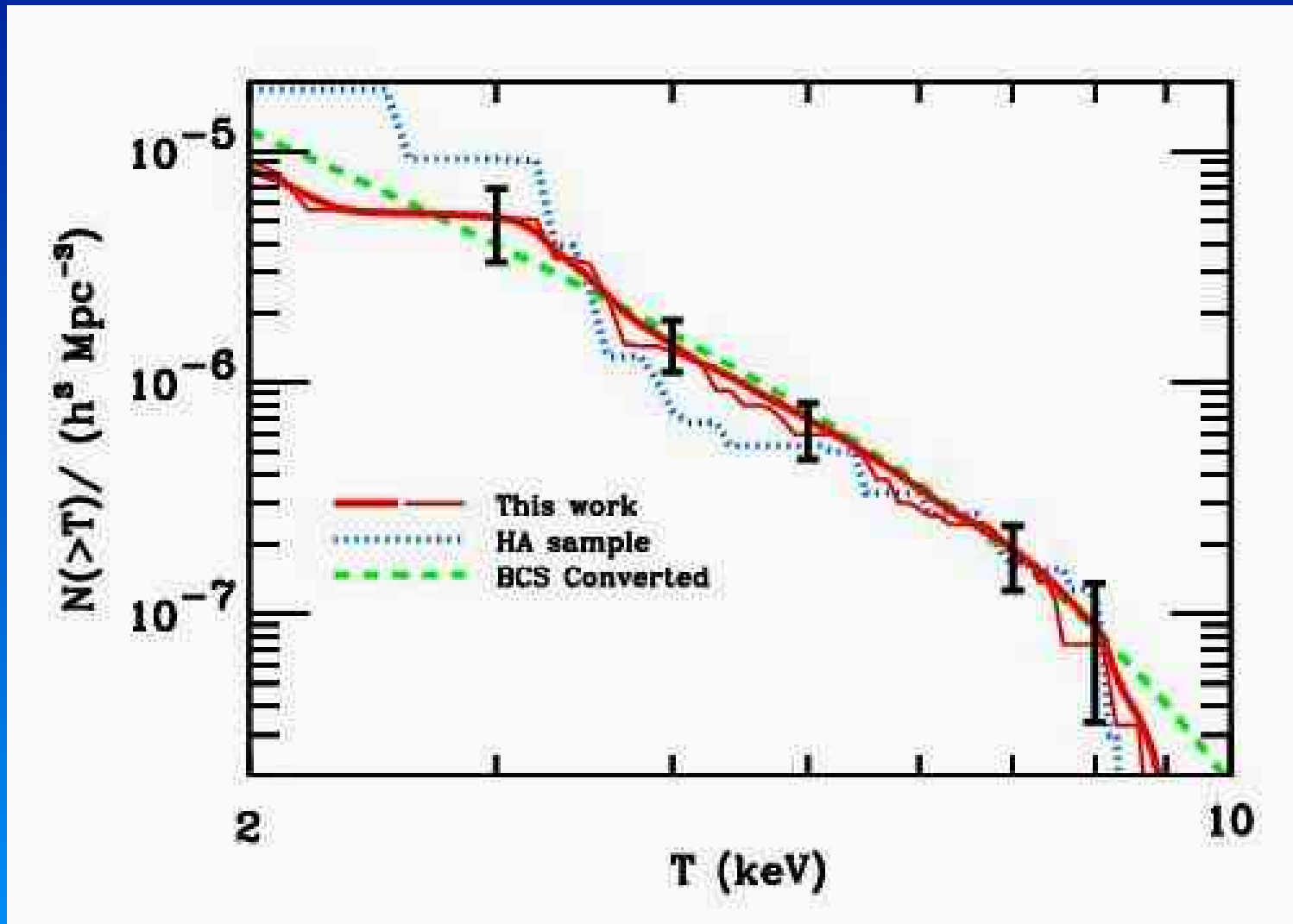
Reliable description
of the mass function in
numerical simulations...

σ_8

from **X**-ray clusters:

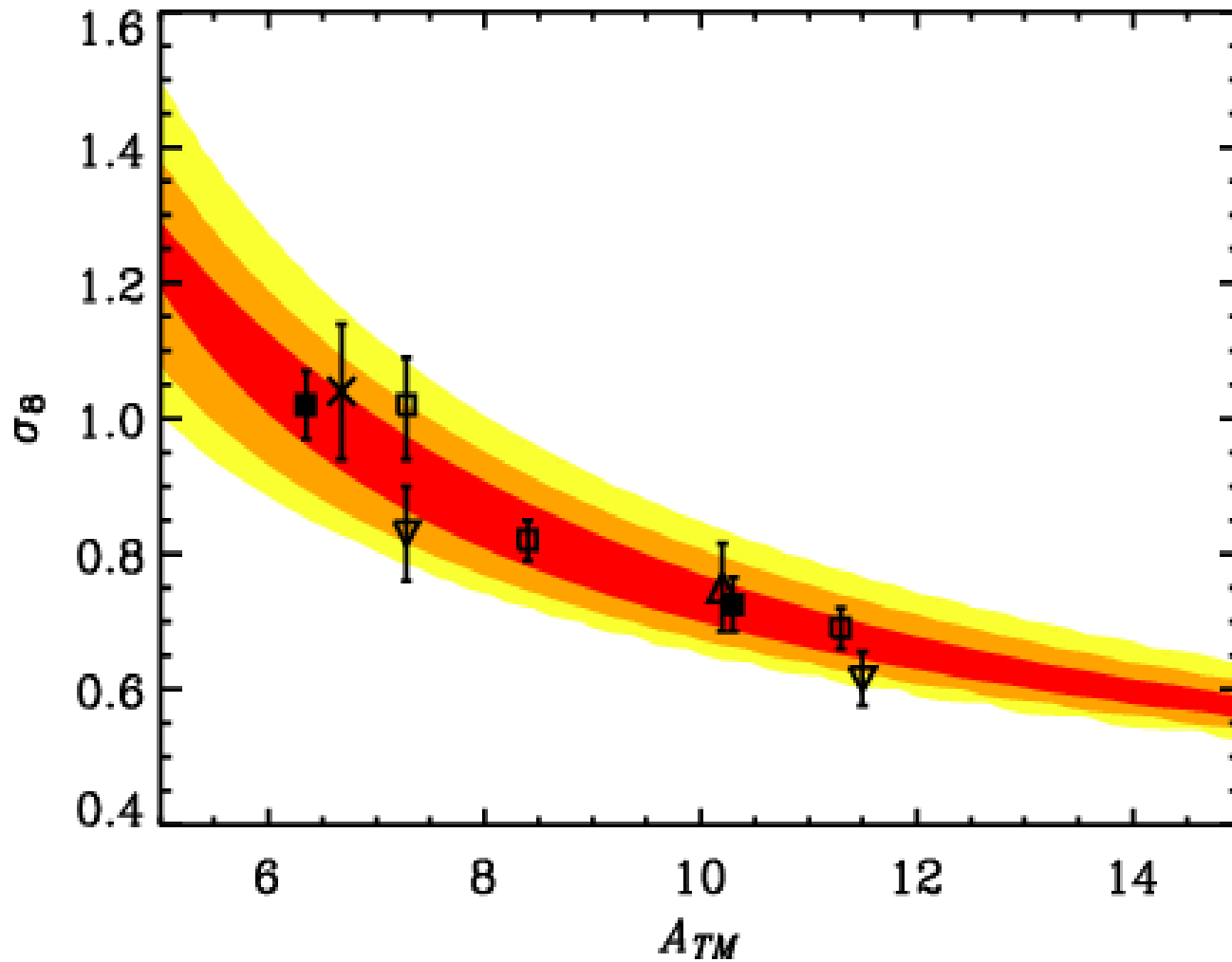


A revised estimation of the local $N(T)$

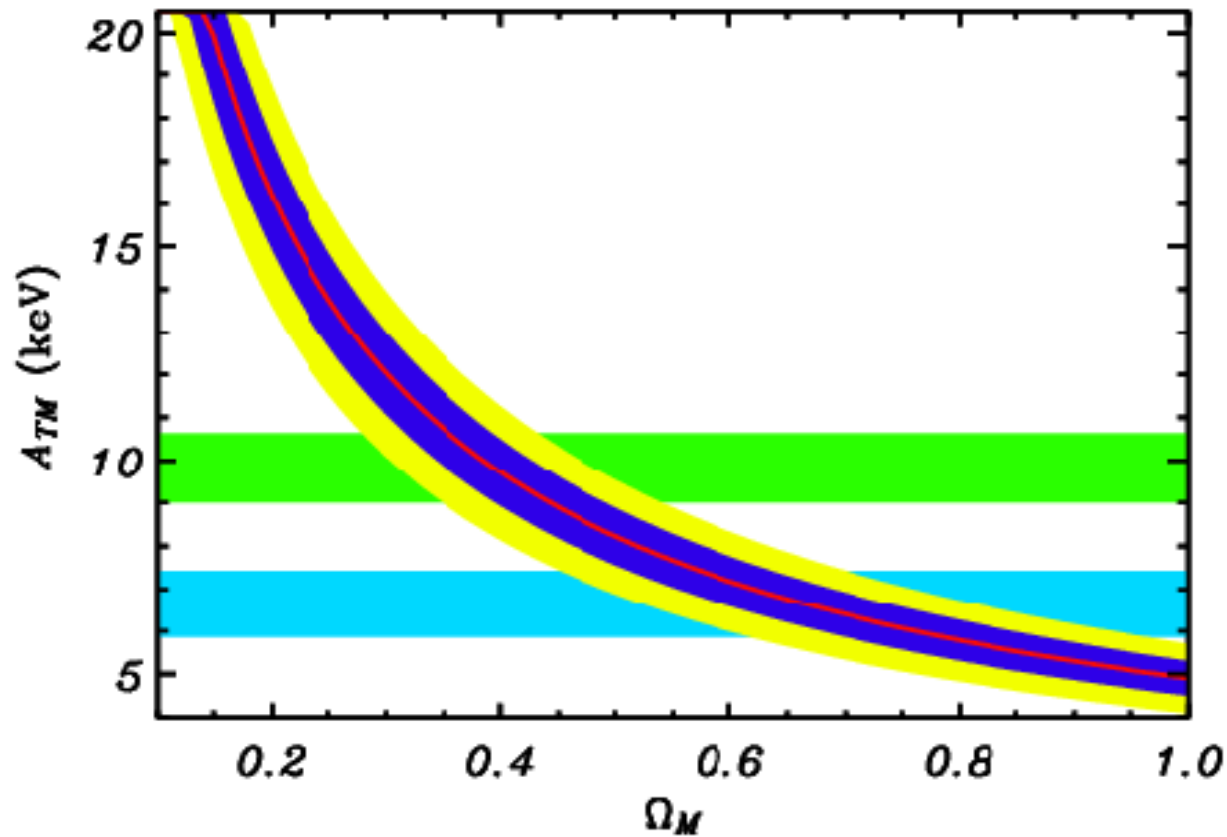


σ_8

from X-ray clusters:

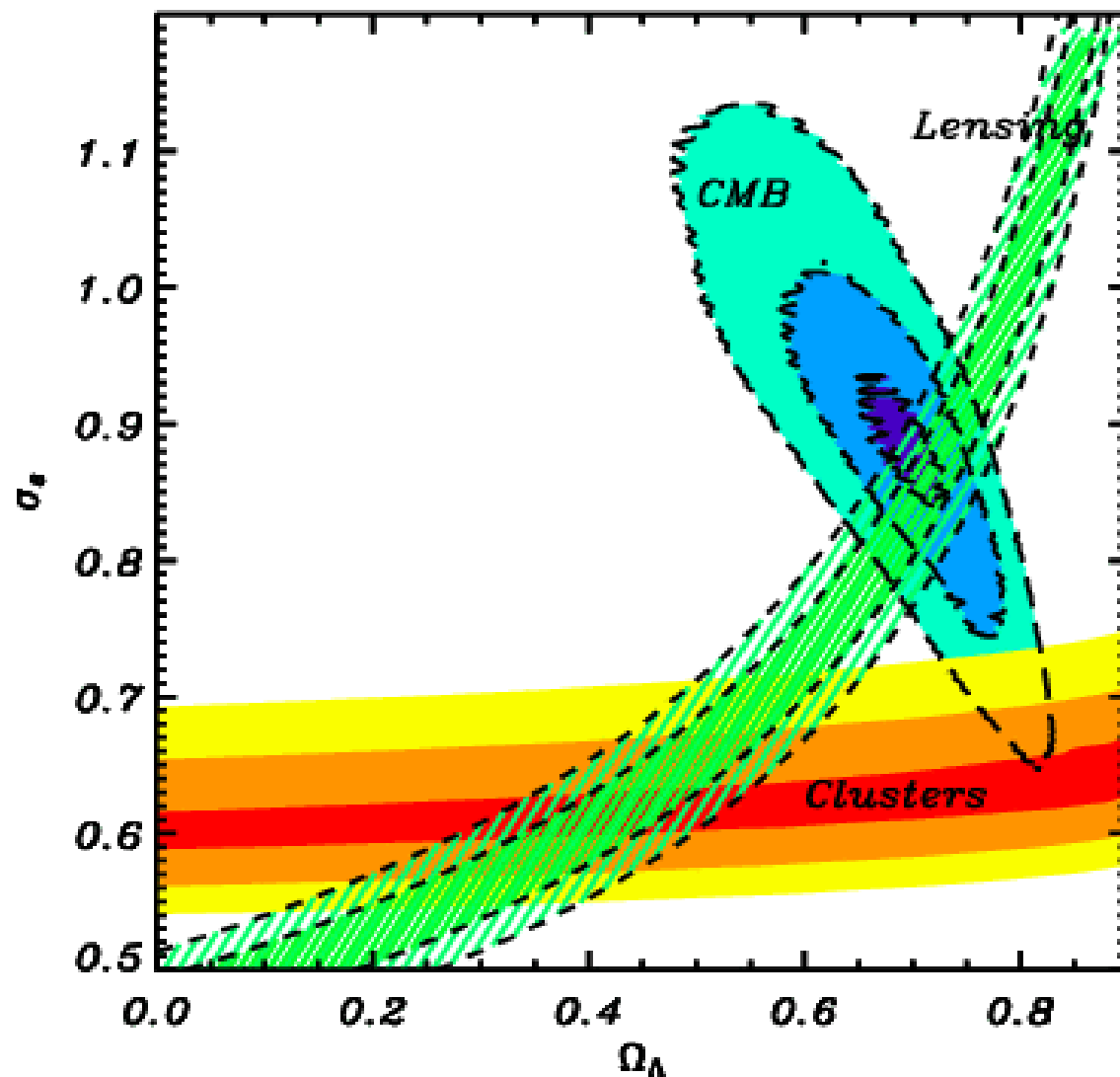


Self-consistent baryon fraction



σ_8 all together:

Blanchard & Douspis 2004

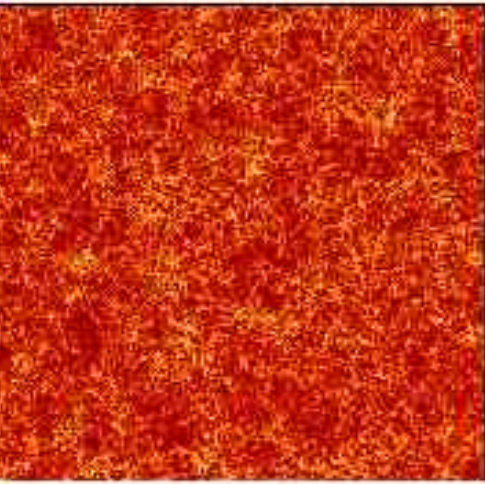


$$\Omega_m$$

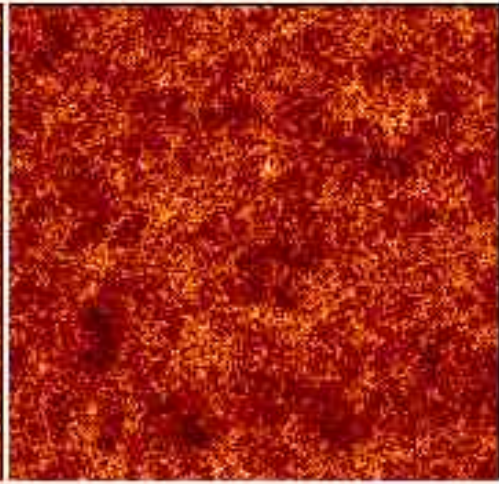
From X-ray Clusters

Number evolution

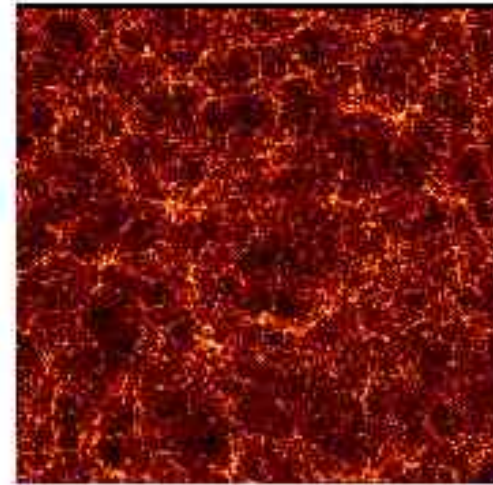
Λ CDM



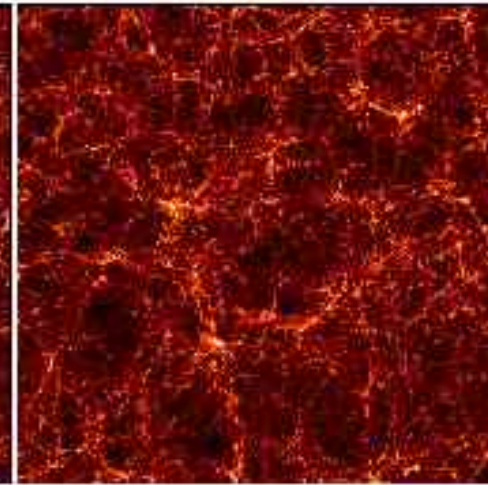
r CDM



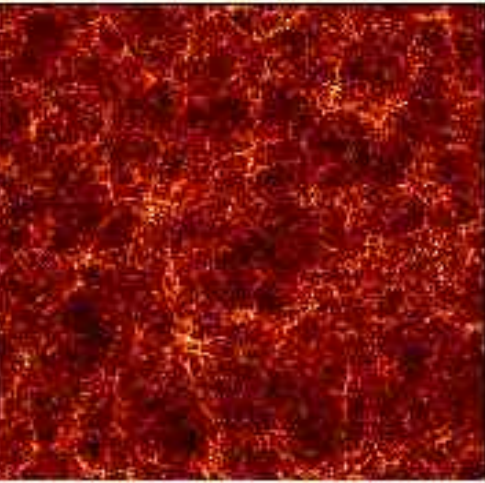
Λ CDM



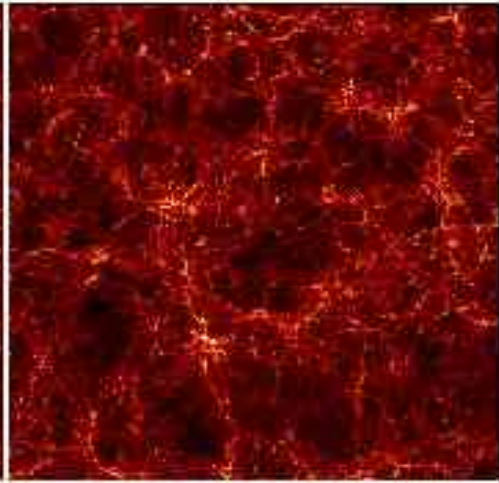
r CDM



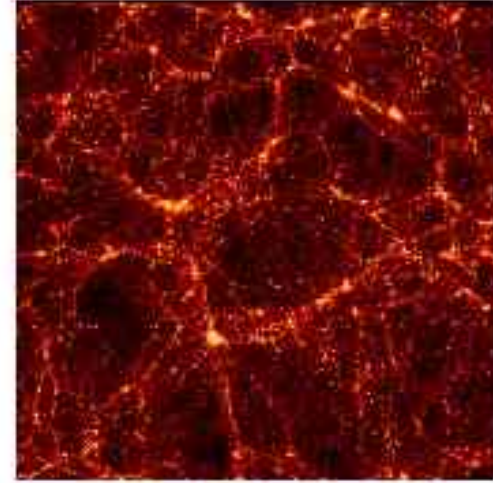
Λ CDM



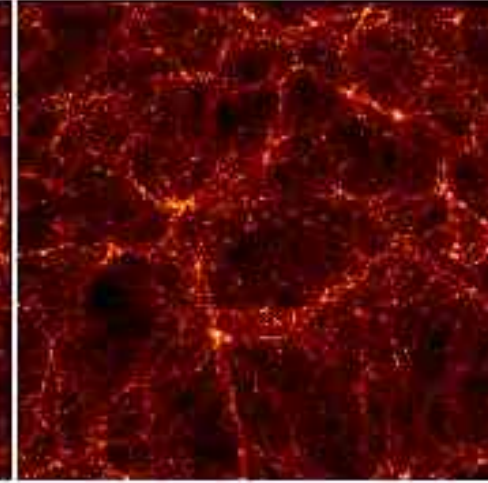
r CDM



Λ CDM

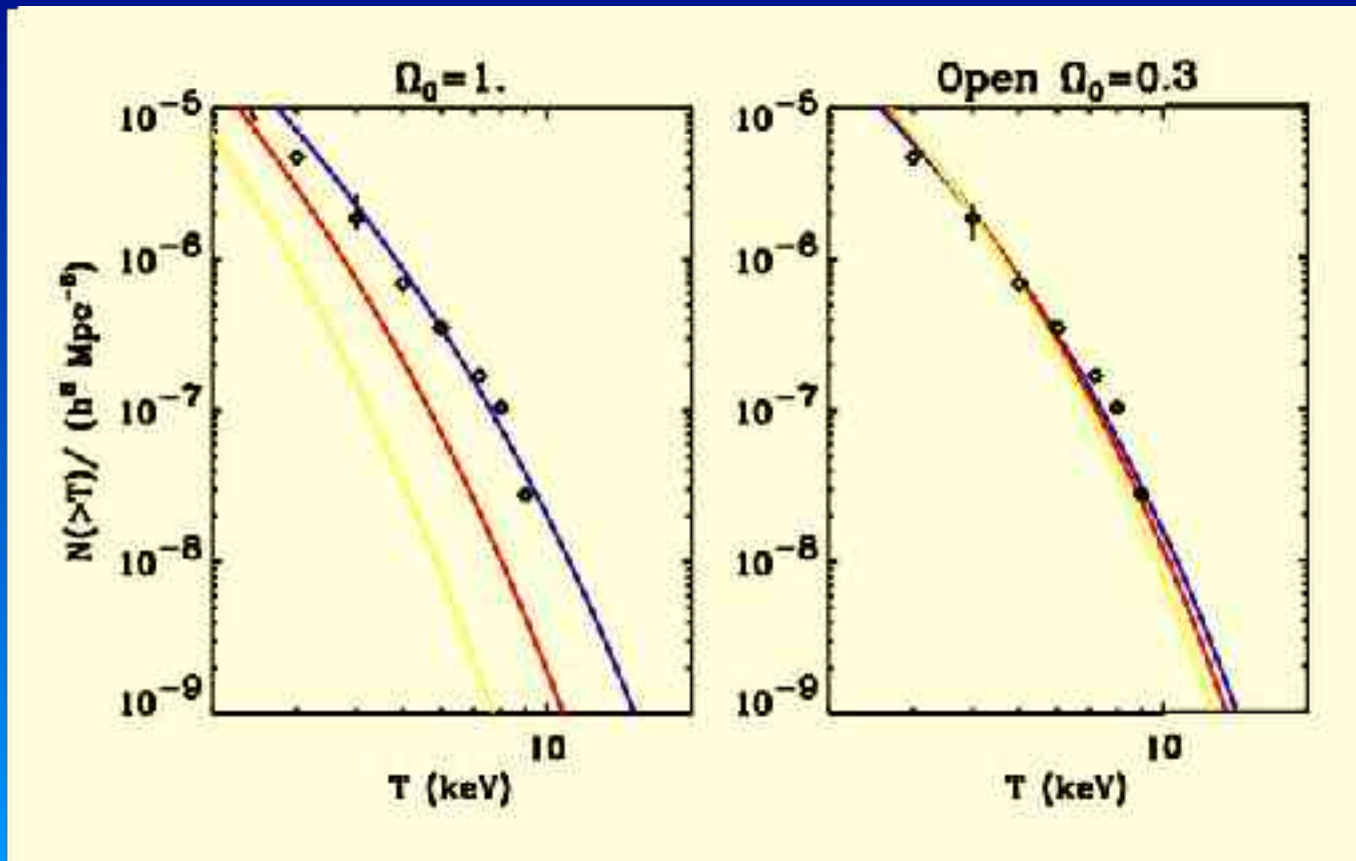


r CDM

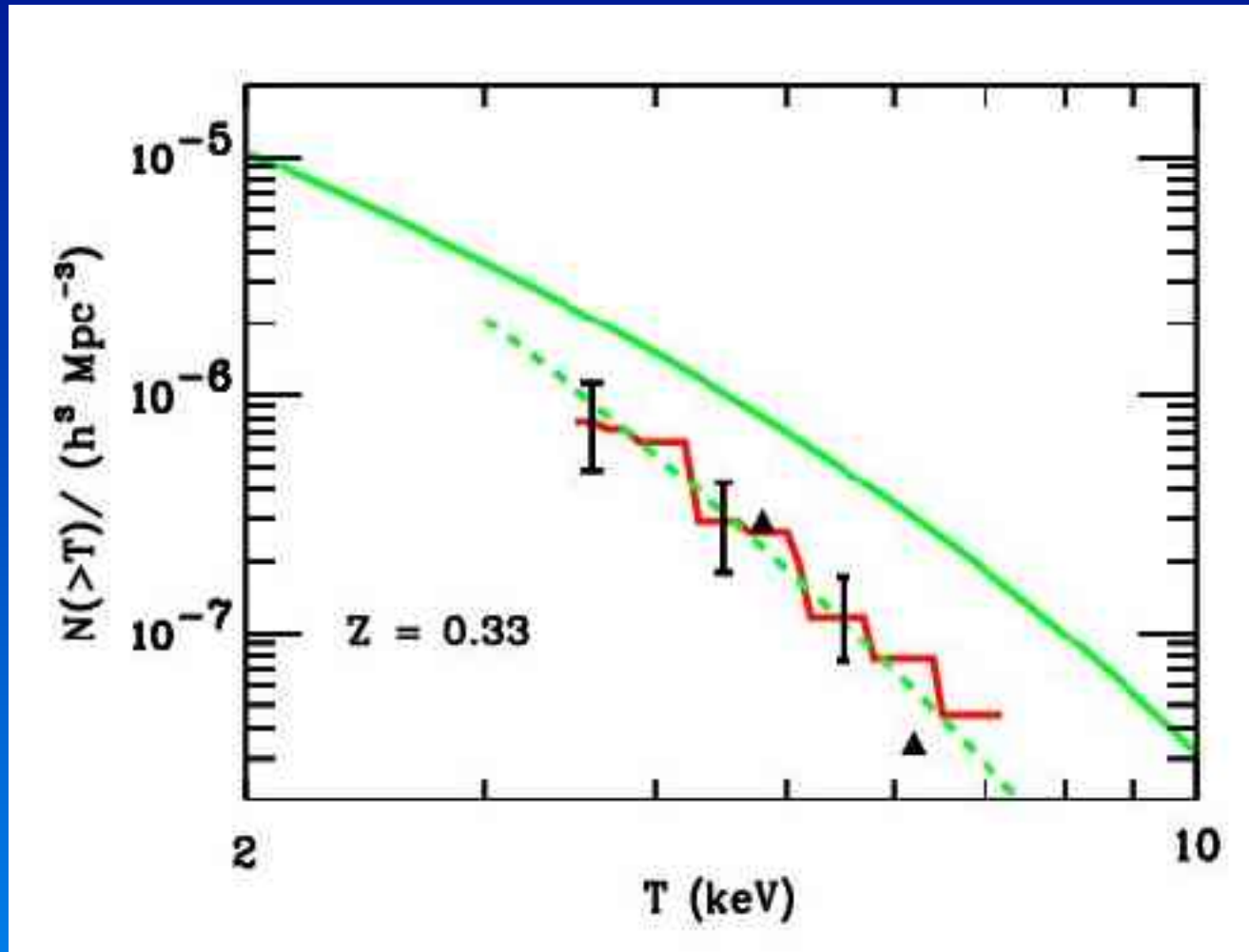


VIRGO simulations

Principle

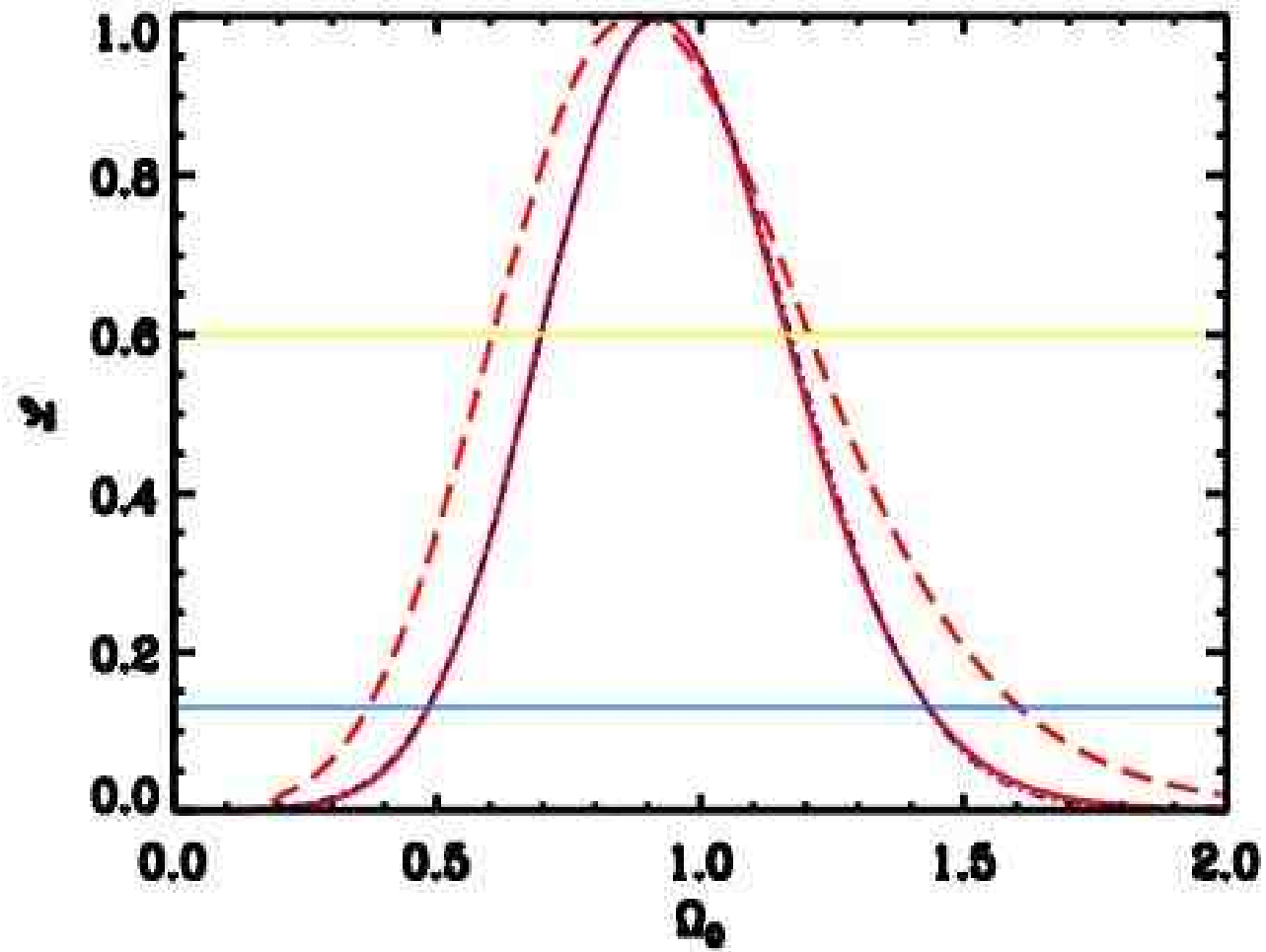


Estimated $N(T)$ at $z \approx 0.33$



Using Henri's sample (1997)

Likelihood on Ω_m



Blanchard et al (2000)

$$\Omega_m$$

From X-ray Clusters

Vauclair et al, 2003

A&A 412, L37

Number counts:

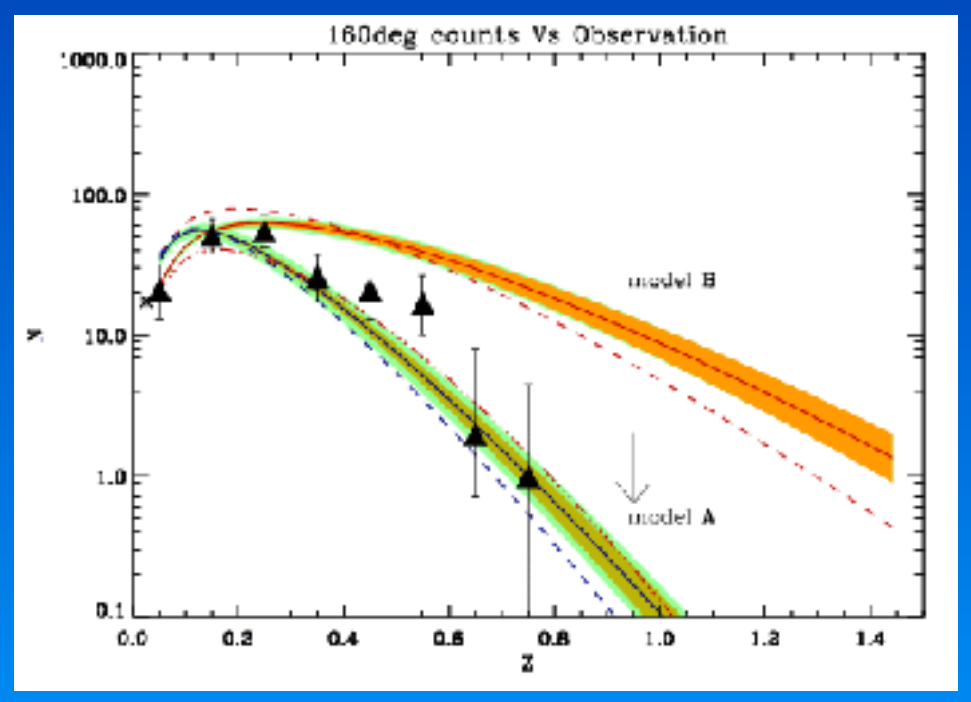
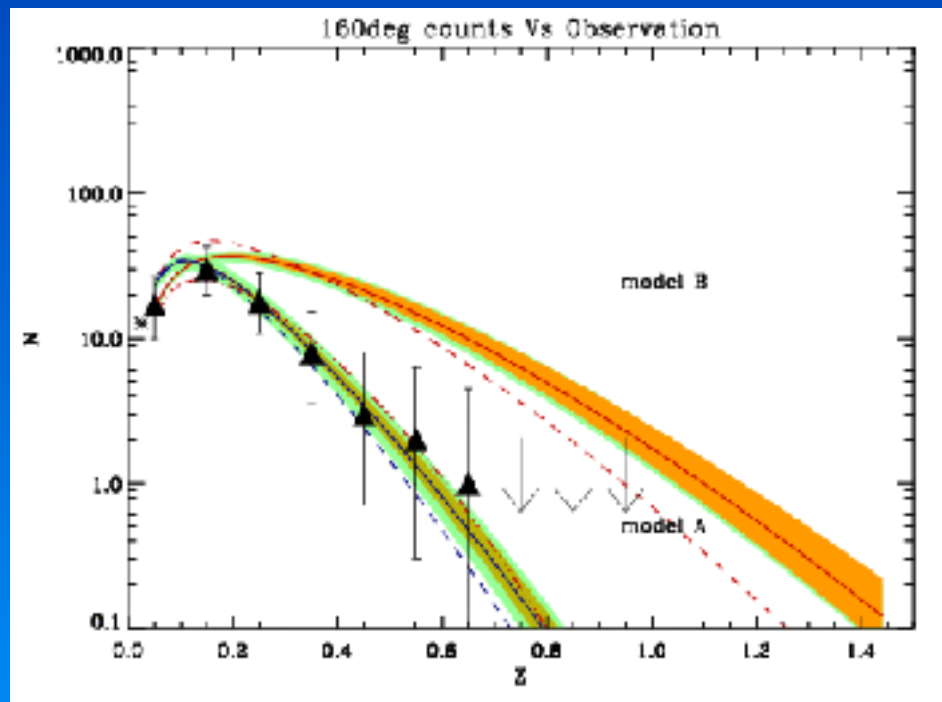
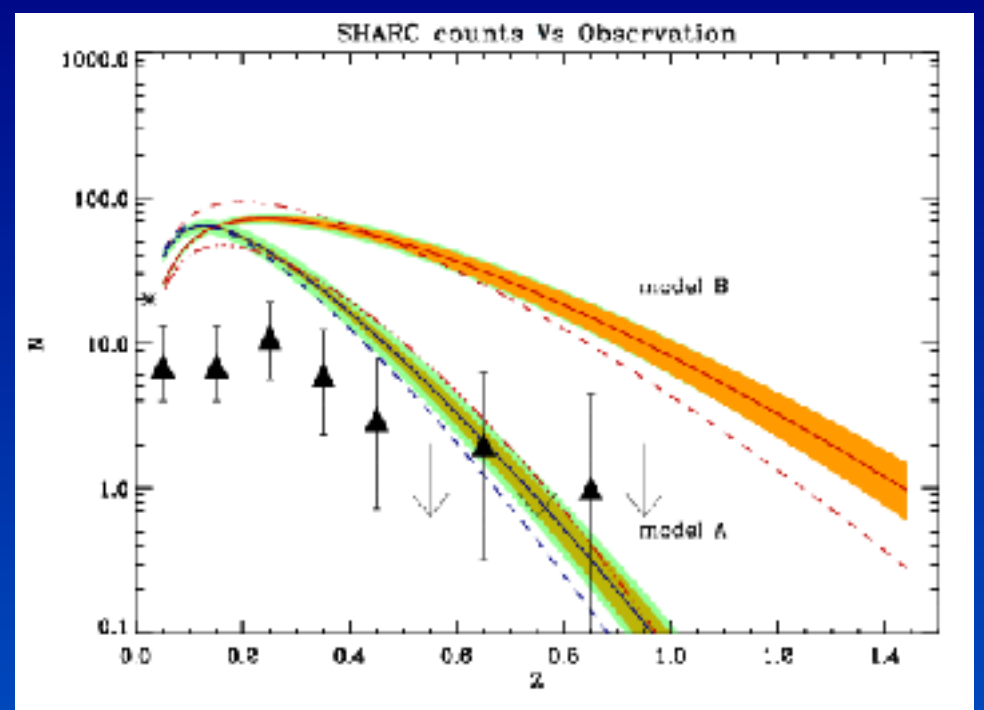
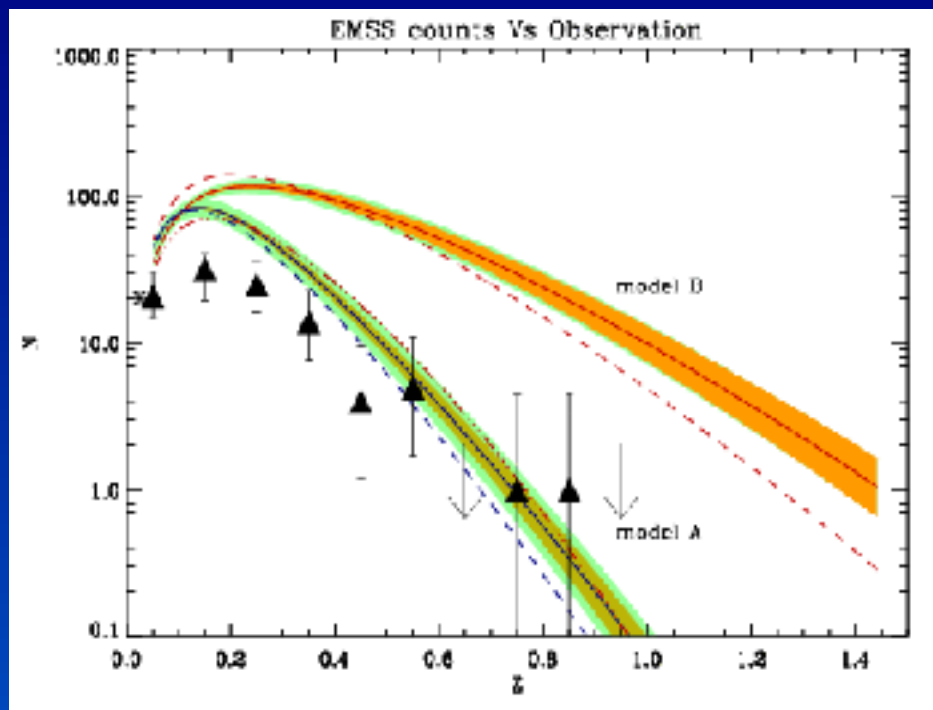
~300 clusters with

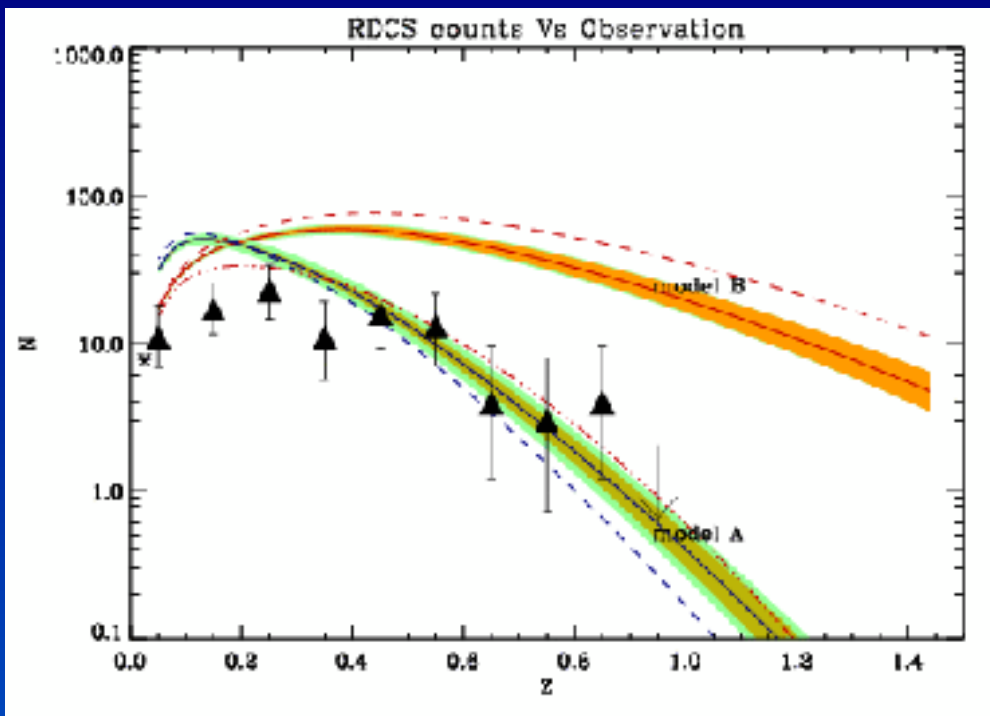
$z > 0.3$

Method:

$$f_x \rightarrow L_x \rightarrow s, T_x$$

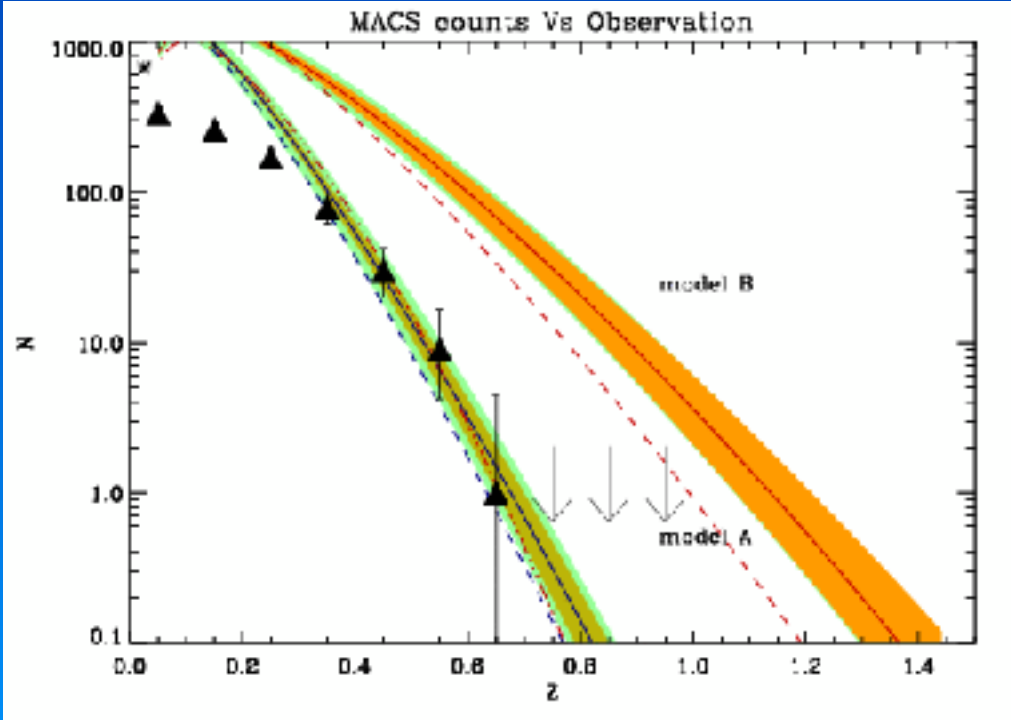
$$\begin{aligned} N(> f_x) &= \int_0^{+\infty} \int_0^{+\infty} s(T, z) N(T, z) dT dV(z) \\ &\sim \int_0^{+\infty} N(> T(z)) dV(z) \end{aligned}$$



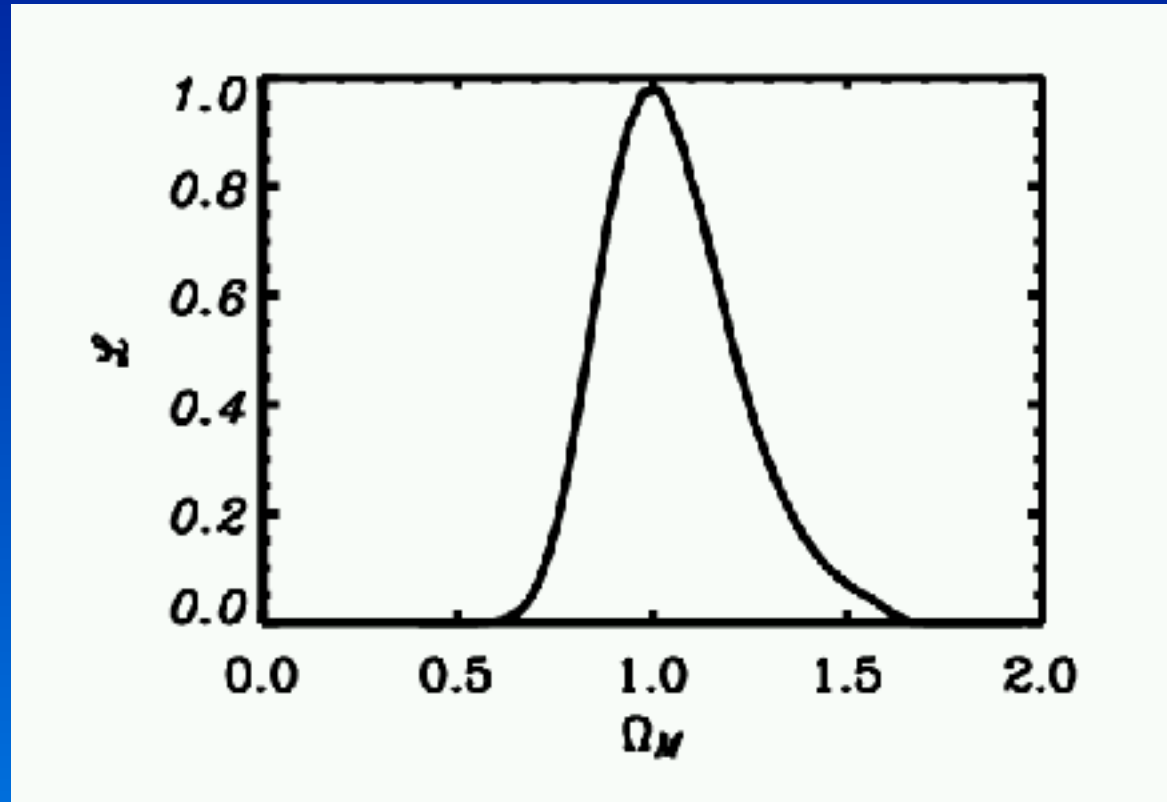


RDCS: 50 deg²
 $f_x \approx 3 \cdot 10^{-14}$ erg/s/cm²

MACS: 22 000 deg²
 $f_x \approx 10^{-12}$ erg/s/cm²



Likelihood analysis:



(Vauclair S., 2004)

$$\Omega_m = 0.99 \pm 0.15 \pm 0.15$$

Conclusions ?

- **Strong Evolution** in the abundance of x-ray clusters appears from all existing surveys in a **very consistent** way.
- This provides a strong argument in favor of a high matter density... 😊
- Consistent with f_b evolution ... 😊

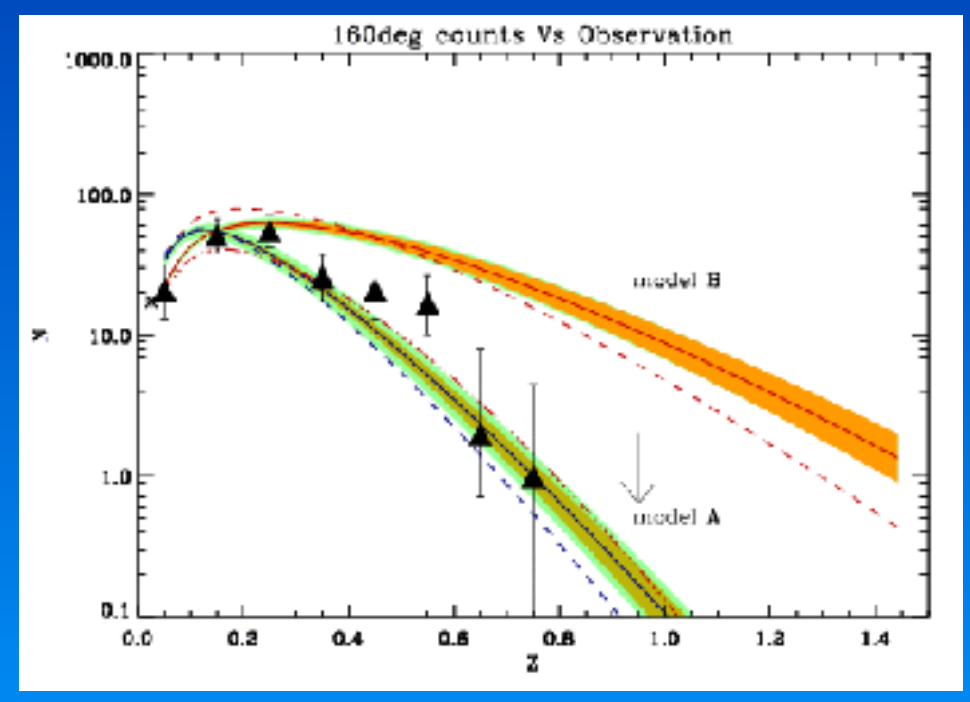
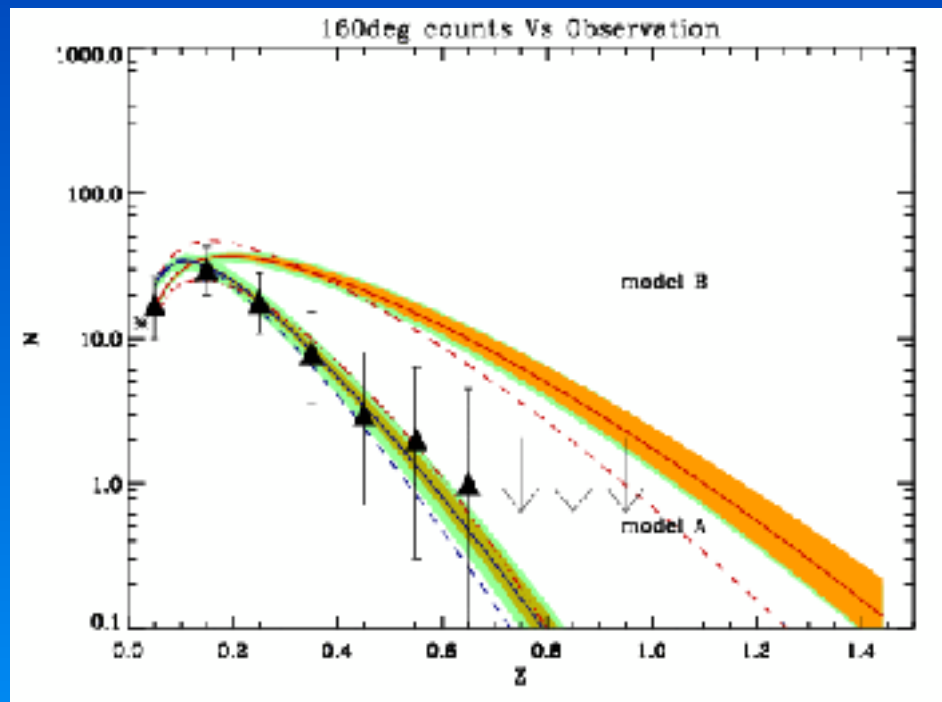
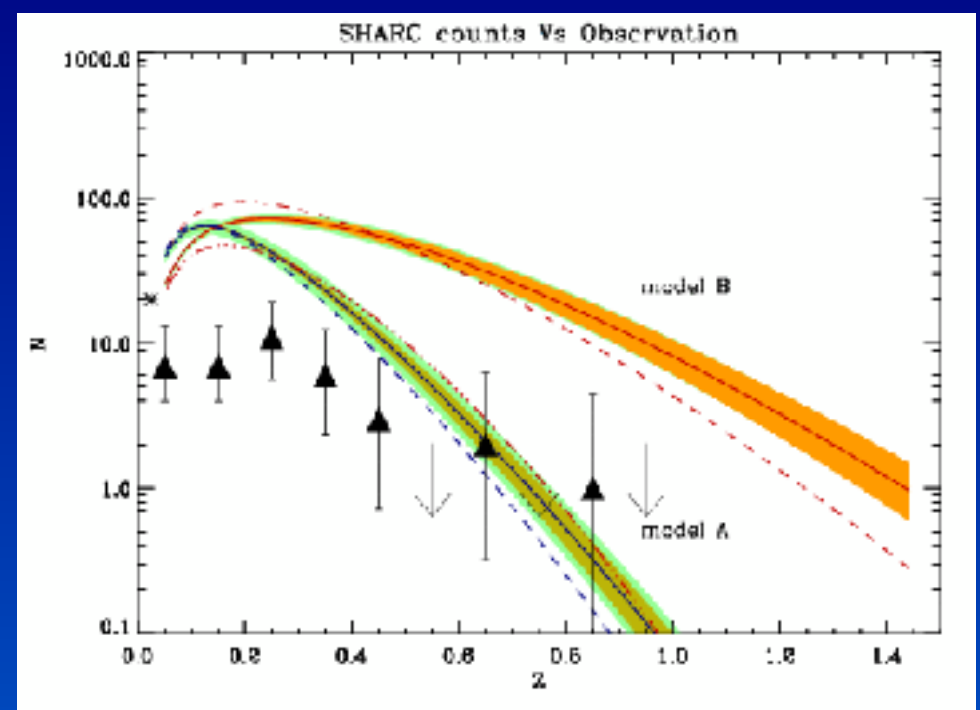
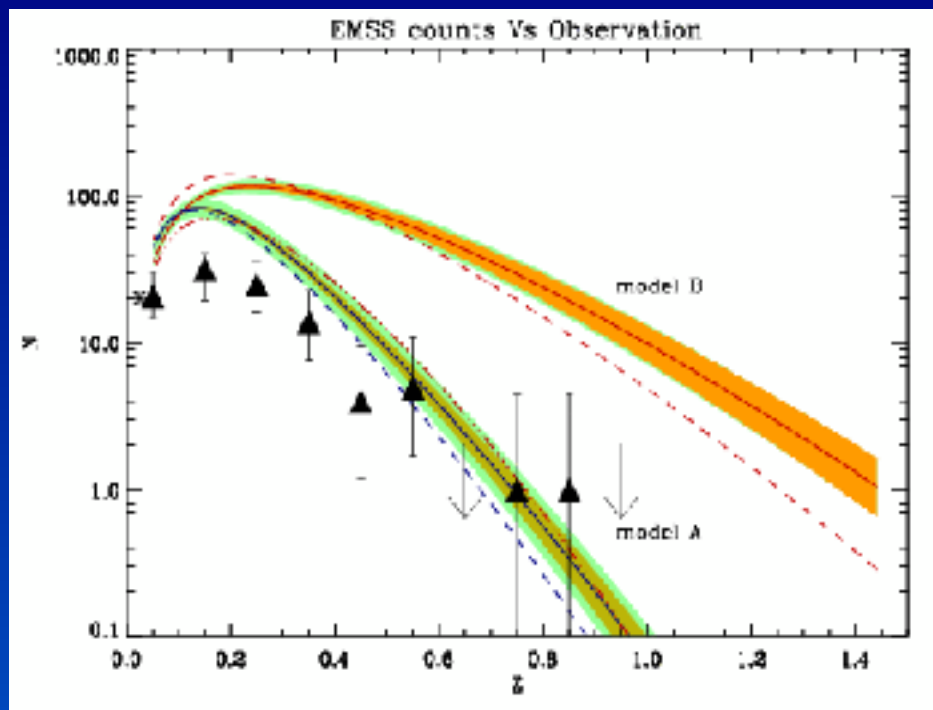
Yet an other
degeneracy?

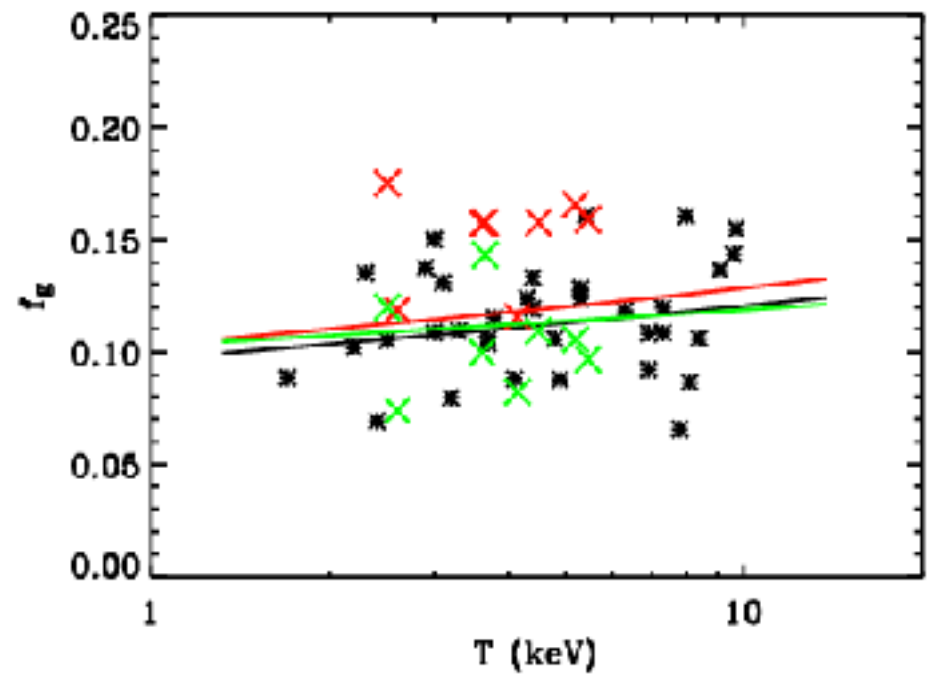
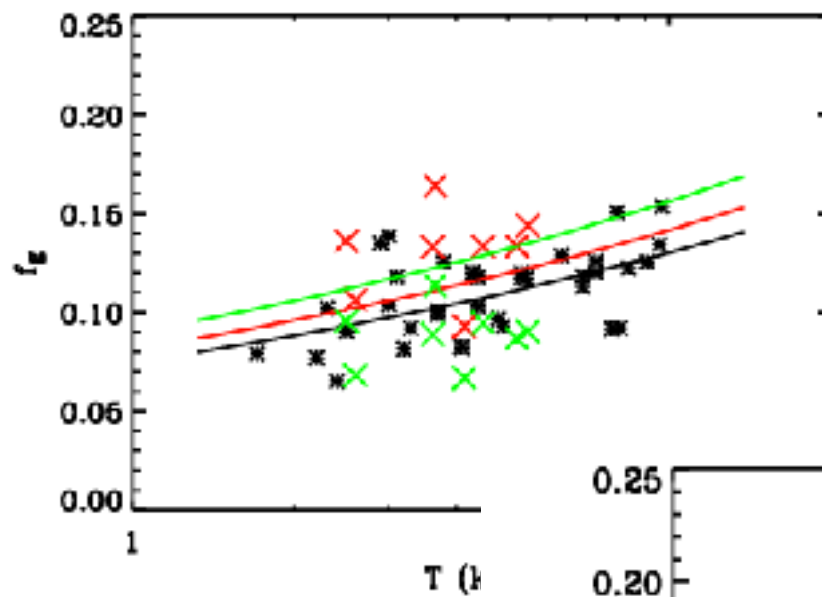
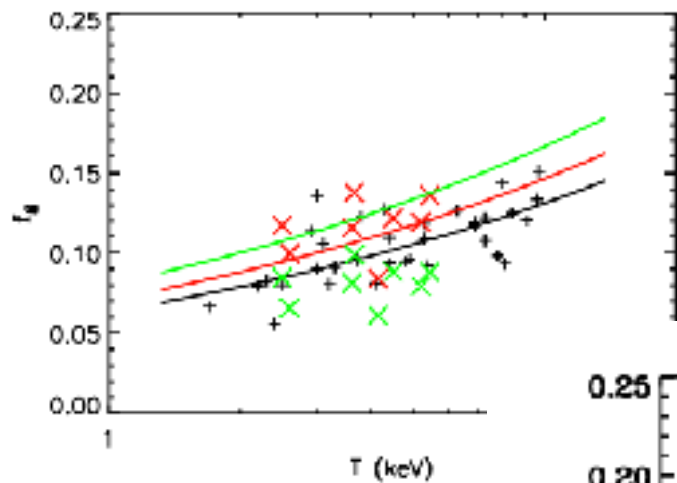
Kill the Mass-Temperature Relation :

$$T \propto GM/r + \dots \propto GM/r / (1+z)$$

i.e. ~ forget gravity...

$$T_x \propto M^{2/3}$$





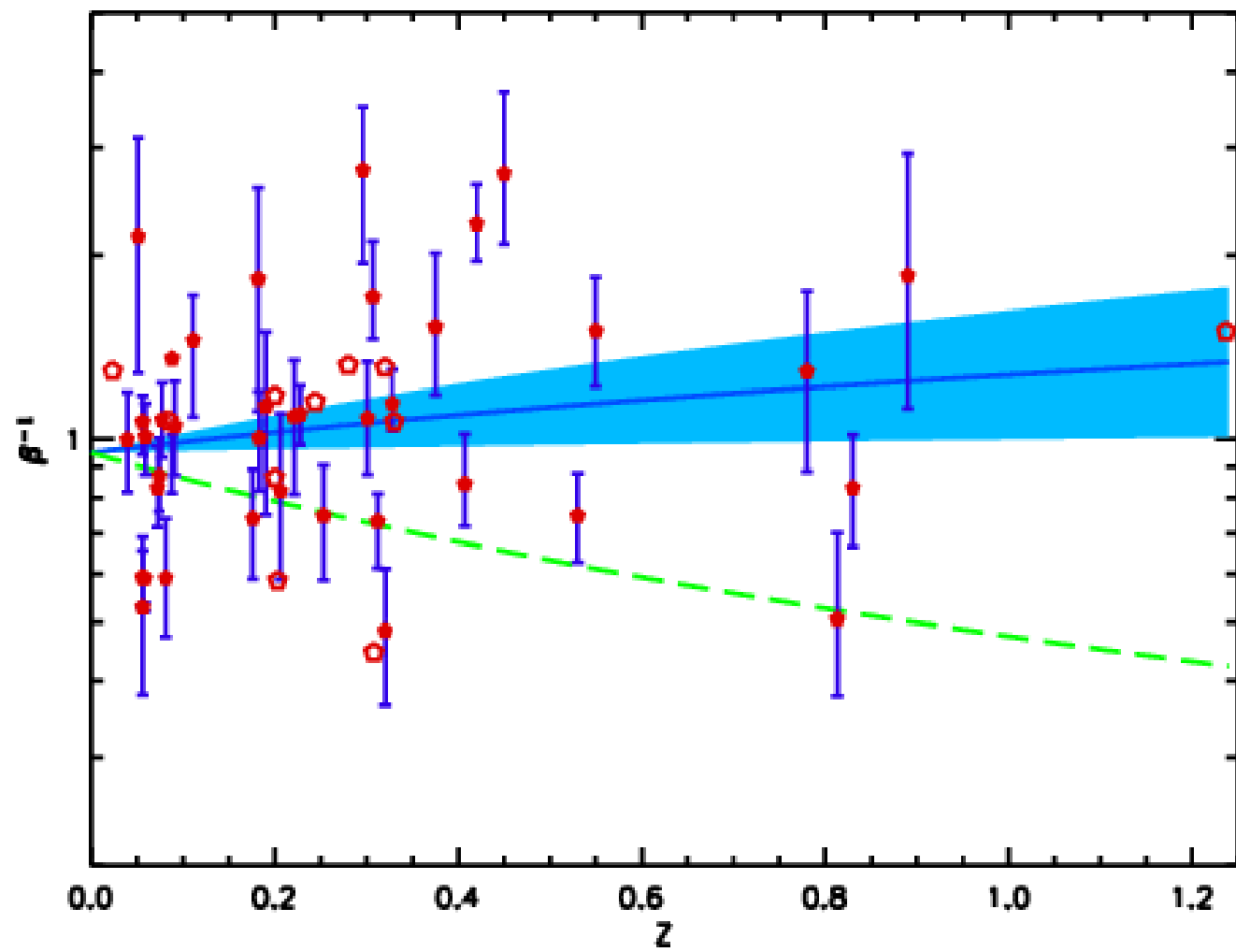
Concordance
works !

Breaking the degeneracy...

$$T \propto GM/r / (1+z) \propto \sigma^2 / (1+z)$$

➤ **Testable... i.e.**

$$\beta^{-1} \propto T/\sigma^2 \propto 1/(1+z)$$



Conclusion:

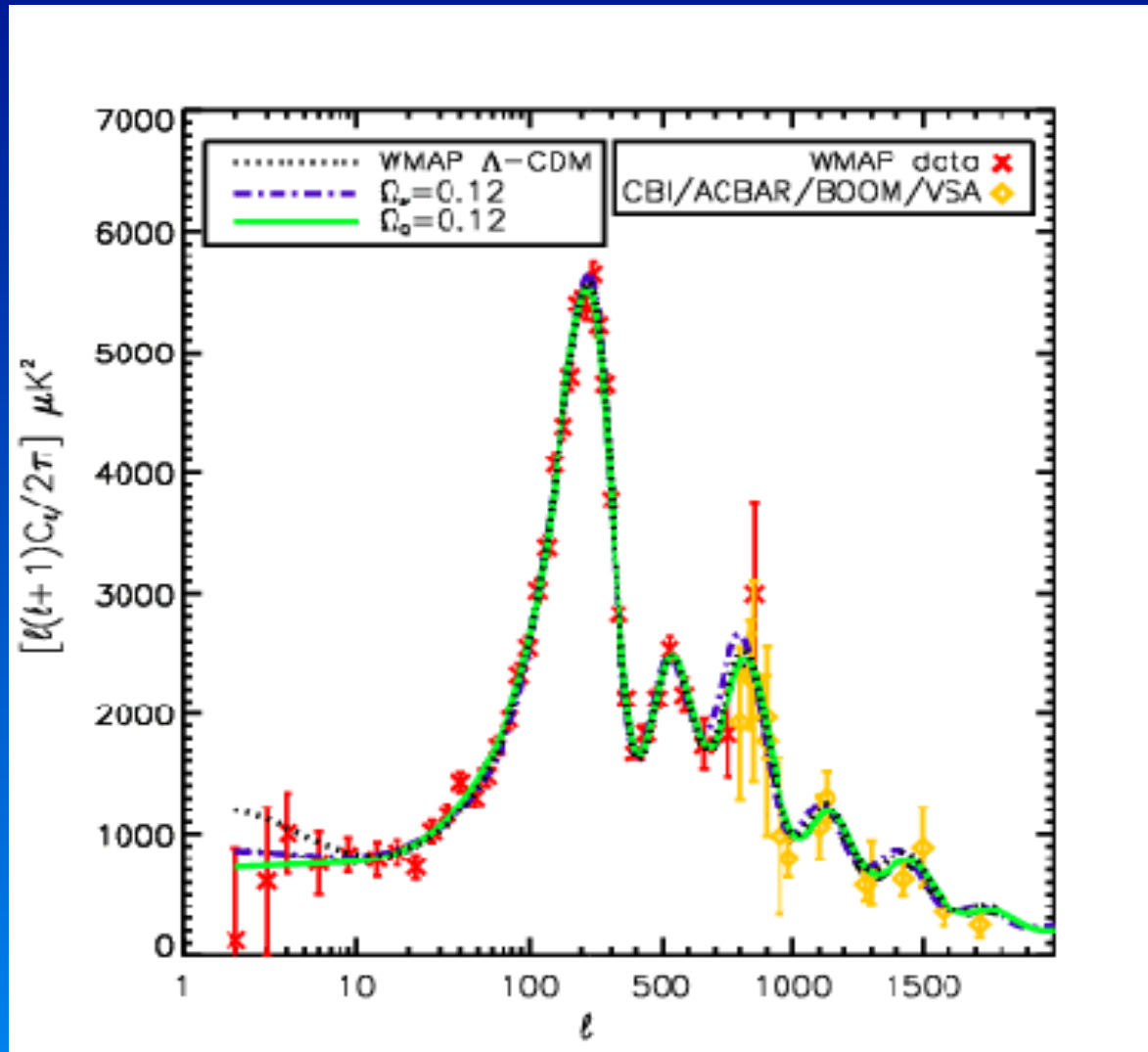
- **Strong Evolution** in the abundance of x-ray clusters appears from all existing surveys in a very **consistent** way.
- **This provides a strong argument in favor of a high matter density... ☺**
- **Consistent with f_b evolution ... ☺**
- **Consistent with f_b amplitude ... ☺**

An Einstein-de Sitter universe, why not?

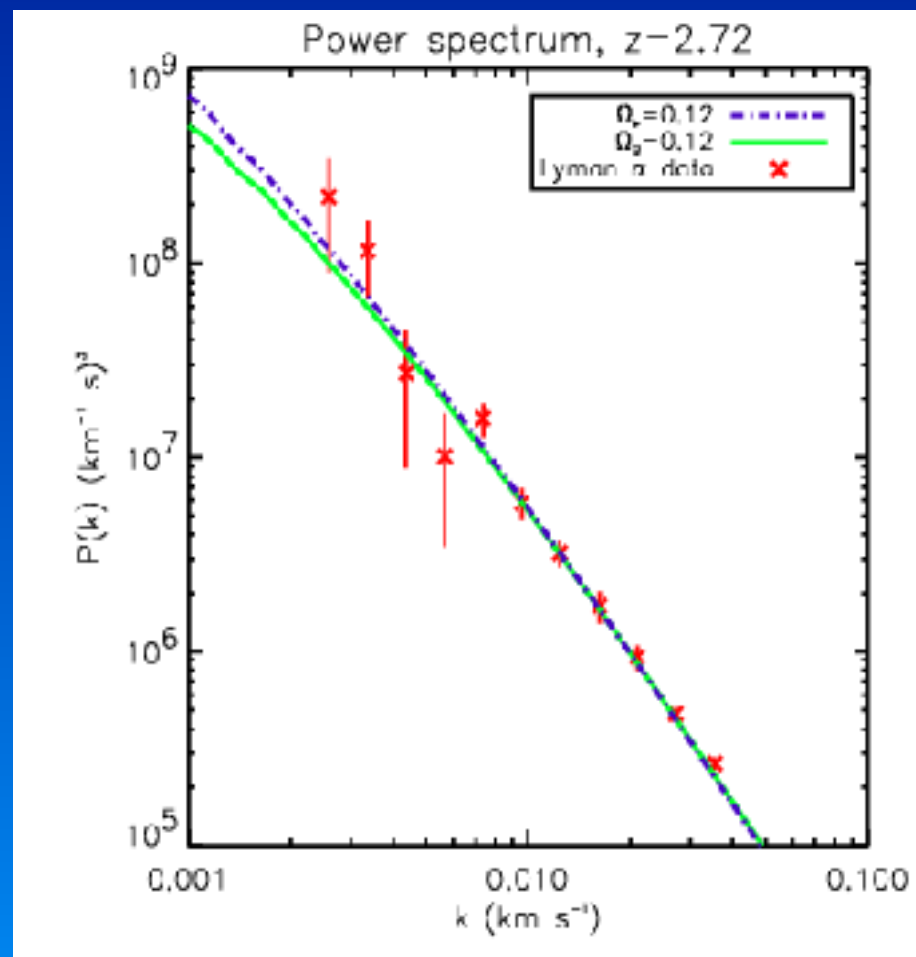
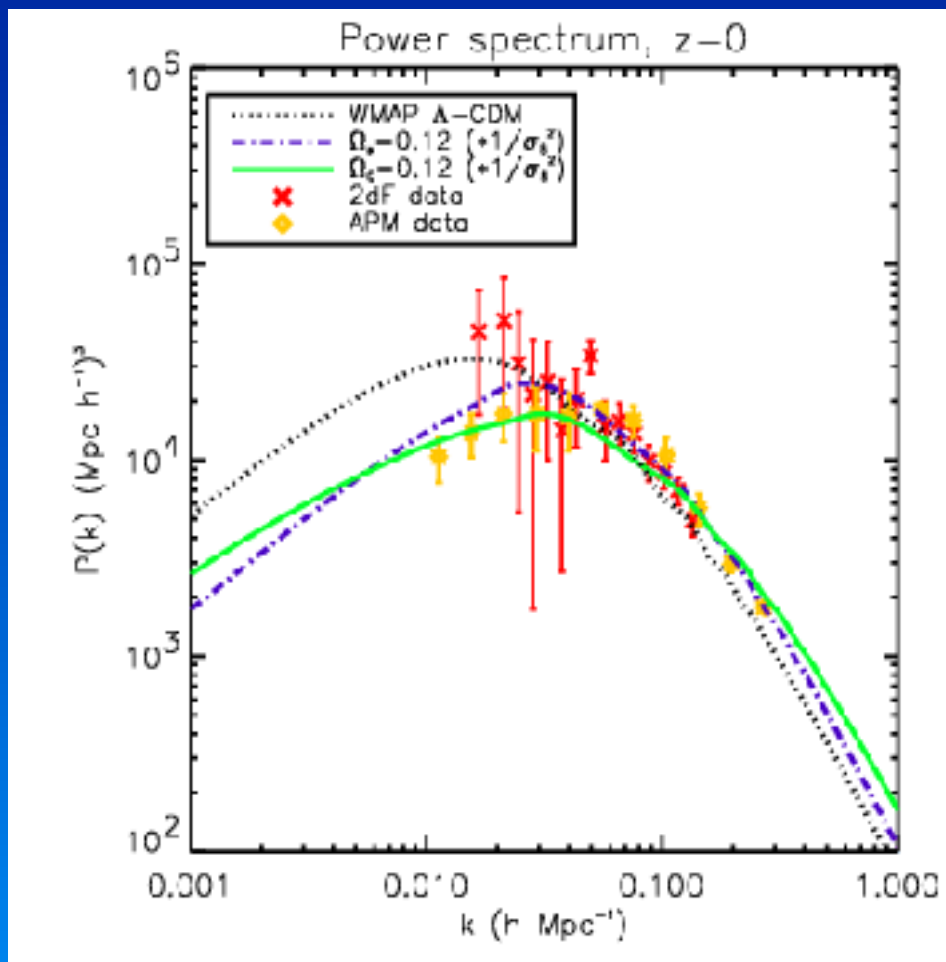
An Einstein-de Sitter universe,

why not ???

C.M.B.

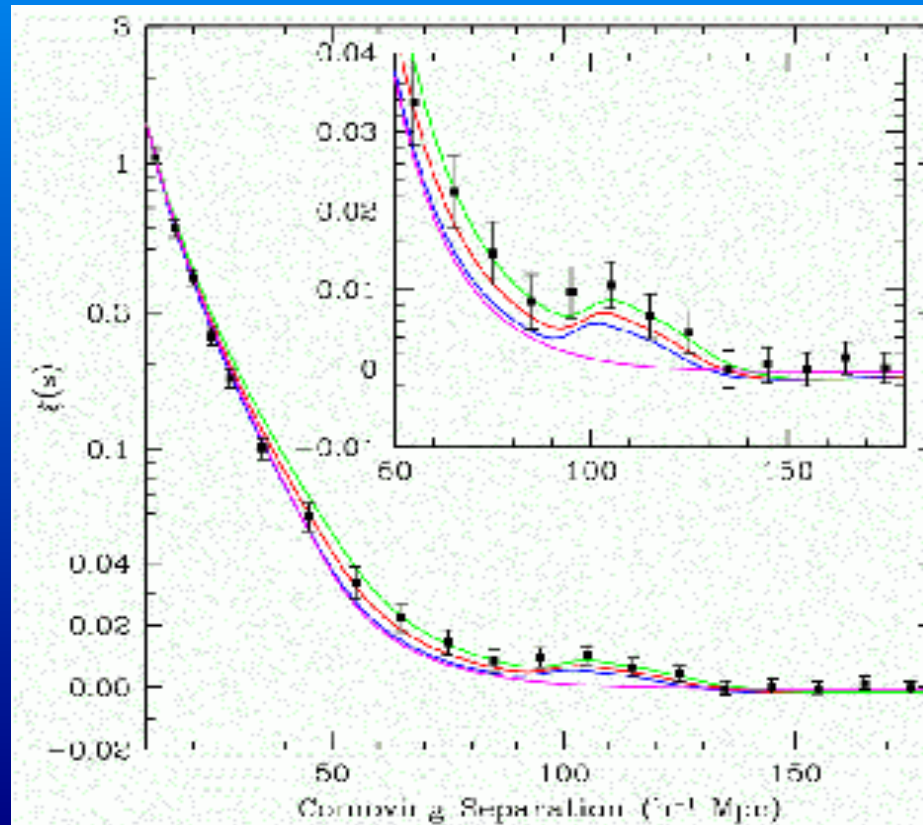


Blanchard, Douspis, Rowan-Robinson, Sarkar 2003

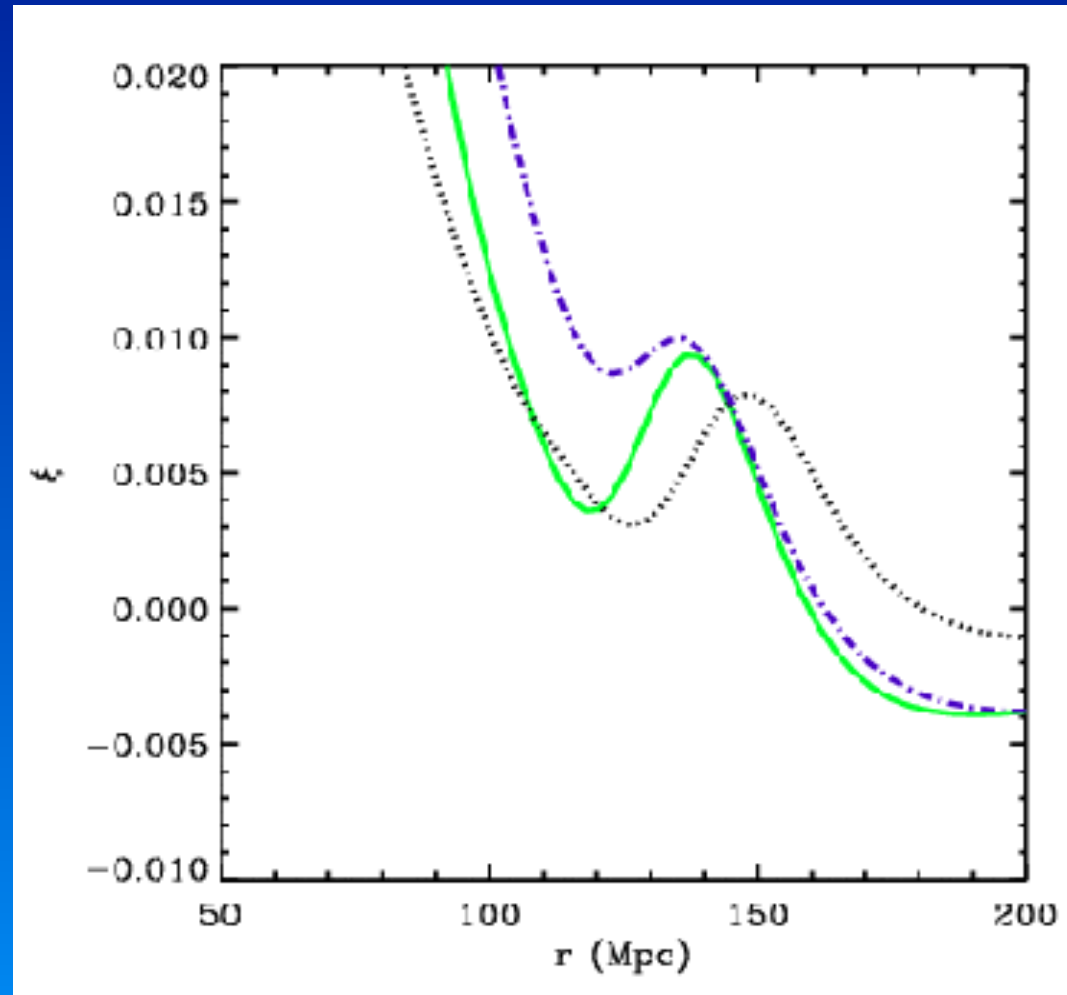


LSS correlation

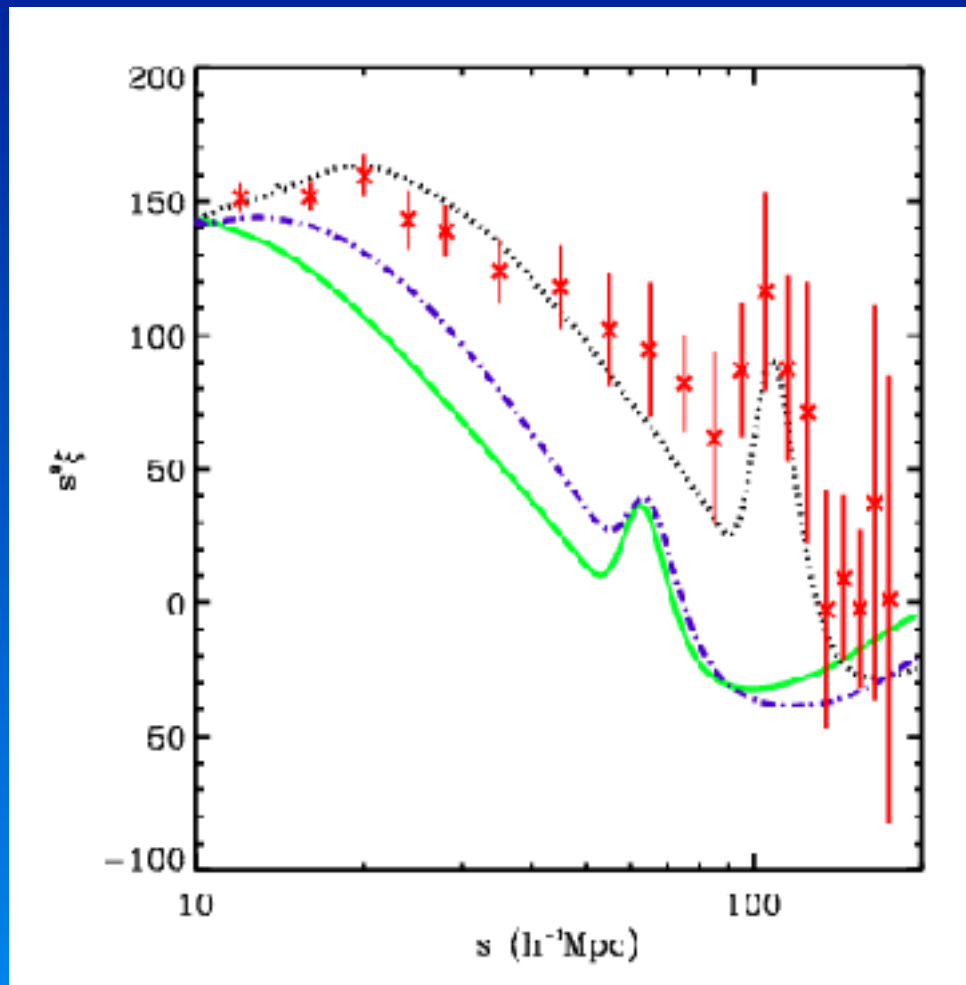
SLOAN red galaxies



Acoustic peak ?



Blanchard, Douspis, Rowan-Robinson, Sarkar 2005



Blanchard, Douspis, Rowan-Robinson, Sarkar 2005

Conclusion

- **Clusters do not fit well in the concordance model**
i.e. favoring an **Einstein de Sitter** ...
 - **SLOAN LSS excludes EdS However**
 - **Troubles from clusters ?**