

# Higgs search at LHC

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# Outline

- **Standard Model Higgs Boson:**

- discovery channels at LHC over the Higgs mass range.

Discussion on the main channels in each Higgs mass region.

Overall view of the statistical significance for the discovery as a function of  $M_H$

- SM Higgs parameters measurement: determining the Higgs sector after the discovery

- **MSSM Higgs:**

- discovery perspectives for the LHC experiments

All the results assume full operative detector (full detector installed, final performances in terms of alignment, calibration, etc.)

# SM Higgs boson - Introduction

- The Higgs boson mass is not predicted by the theory. Both theoretical and experimental limits exist

From direct LEP search:

$$M_H > 114.4 \text{ GeV}$$

From the Electroweak fit of the standard model

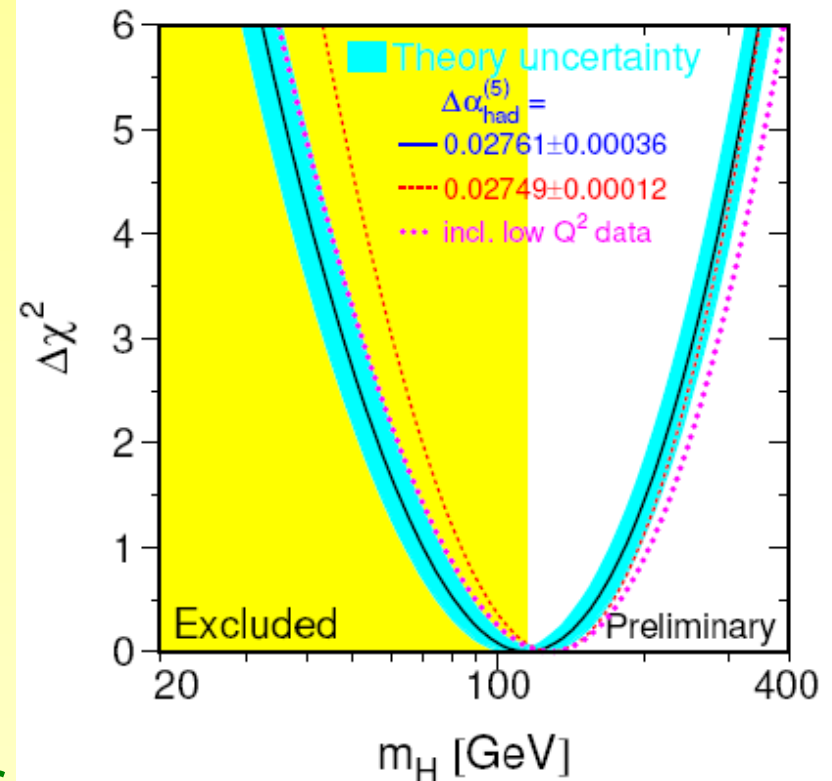
$$M_H < 260 \text{ GeV}$$

$M_H > 1 \text{ TeV}$  is theoretically forbidden

The LHC experiments have to cover

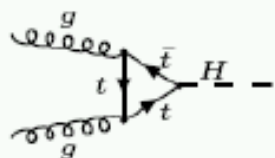
from the LEP limit up to

few 100 GeV ... the TeV scale

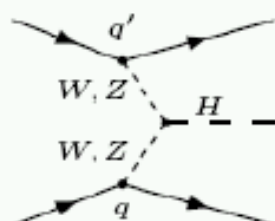


# Production mechanisms & cross-sections

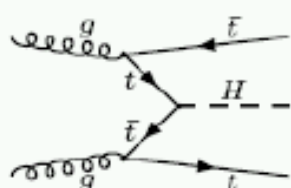
## Production



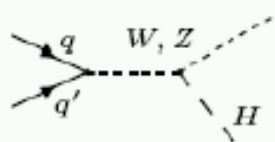
GF:  
Gluon Fusion  
( $gg \rightarrow H$ )



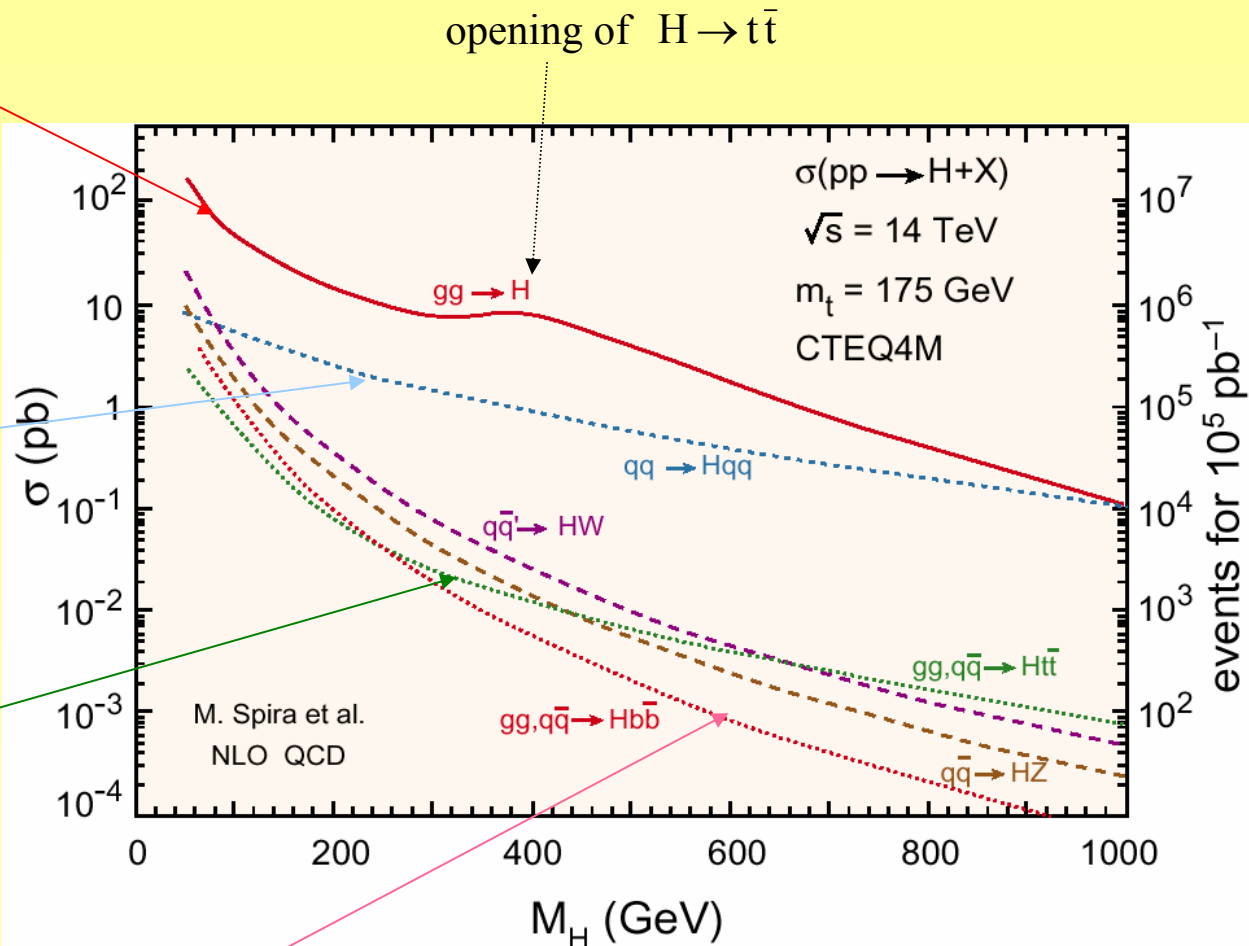
WBF:  
Weak Boson  
Fusion  
( $qq \rightarrow H$ )



$t\bar{t}H$



$WH$   
 $ZH$

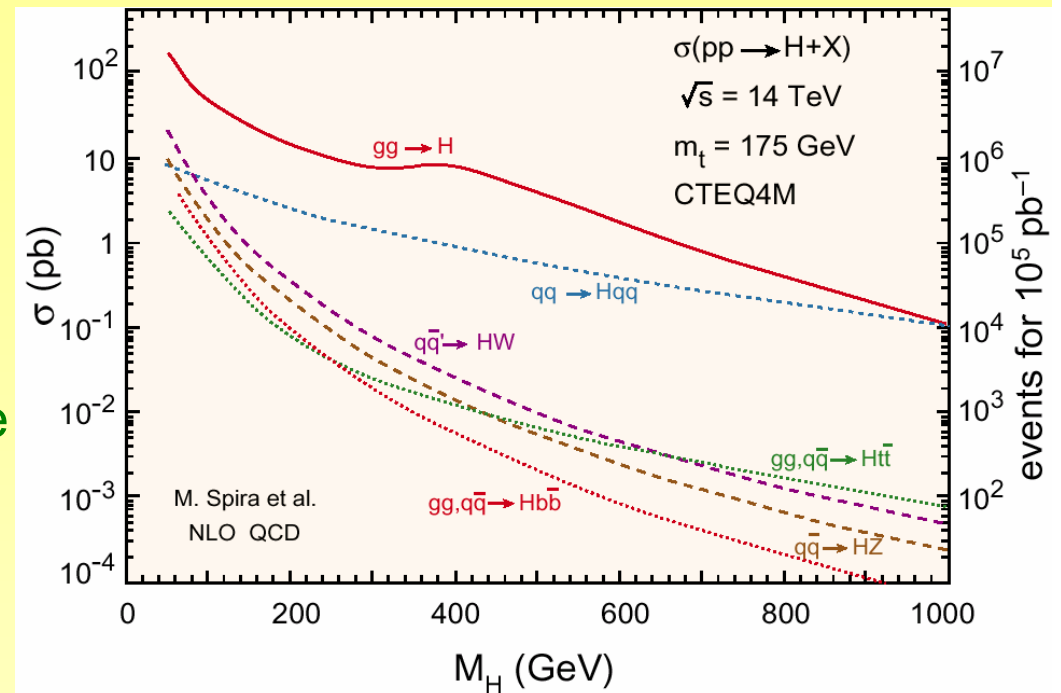


# Cross sections...

Higgs production cross sections available at NLO.

Gluon fusion is the dominant production channel

Vector Boson Fusion (VBF) in the following is the second production channel over the whole mass range



VBF is characterized by energetic jets in the forward region and by lack of hadronic activity in the central region

Though less important in terms of absolute cross section values, associated productions (ttH,WH), provide distinctive signatures. Relevant for largely background dominated decay channels, e.g. bb

# ... and Branching Ratios

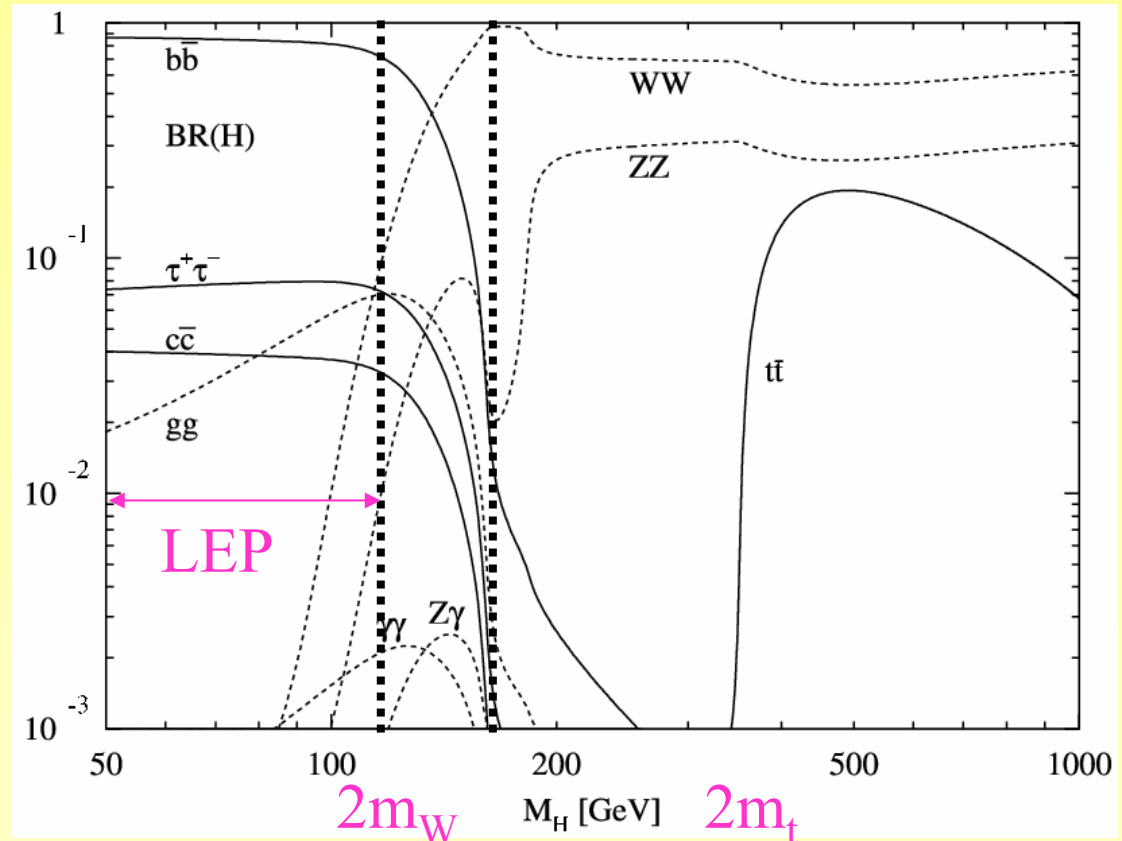
Decays into third generation down-type fermions largely dominating at low mass ( $M_H < 130$  GeV)

$H \rightarrow b\bar{b}$  decay detection feasible only if the associated production  $t\bar{t}H$  is considered.

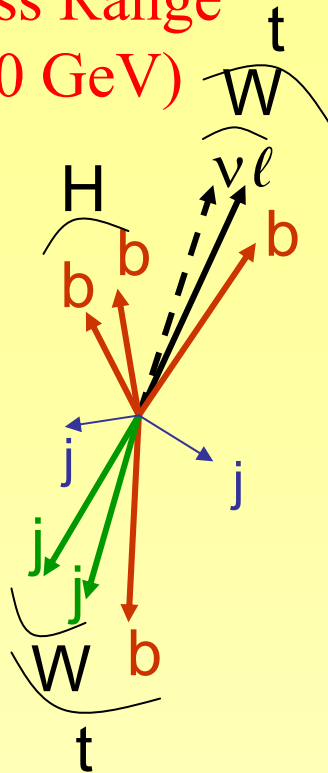
$H \rightarrow \gamma\gamma$  suppressed, but it provides a very clear signature.

Higgs decay into tau pairs important if VBF production is considered

Vector Boson decays important at large  $M_H$  and in the intermediate mass region



## Low Mass Range (until 120 GeV)



## $t\bar{t}H(\rightarrow b\bar{b})$

- $\sigma(pp\rightarrow t\bar{t}H_{120})\times\text{Br}(H\rightarrow b\bar{b})\sim 0.36\text{ pb}$
- $t\bar{t}b\bar{b}$ : 9.5 pb,  $t\bar{t}+jj$  ( $\geq 6$  jets)  $\sim 60$  pb

Tag the top quark to reject background

1)  $t\rightarrow W(\rightarrow \mu\nu)b$  (lepton for the trigger)

2)  $t\rightarrow W(\rightarrow jj)b$

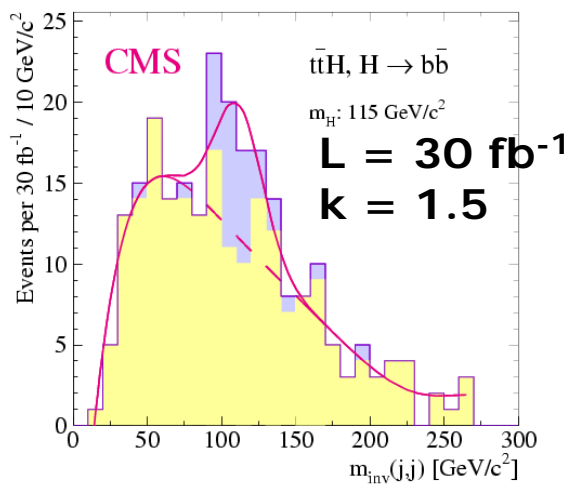
-6 jets

-need **good b-tagging** to reduce combinatorics

-pairing likelihood w/ 6 variables  $\rightarrow$  best

combination

$$\Delta = (m_{\ell\nu b} - m_{\text{top}})^2 + (m_{jjb} - m_{\text{top}})^2$$



Reducible background:  $t\bar{t}+jets, W+jets$  production

Irreducible background:  $t\bar{t}b\bar{b}$  production

CMS applies  $K_{ttH}=1.5 \rightarrow S/\sqrt{B} = 5.3$

ATLAS (no K,  $t\bar{t}b\bar{b}$  at ME, CTEQ5L)  $\rightarrow S/\sqrt{B} = 2.8$

# Higgs decay into $\gamma\gamma$

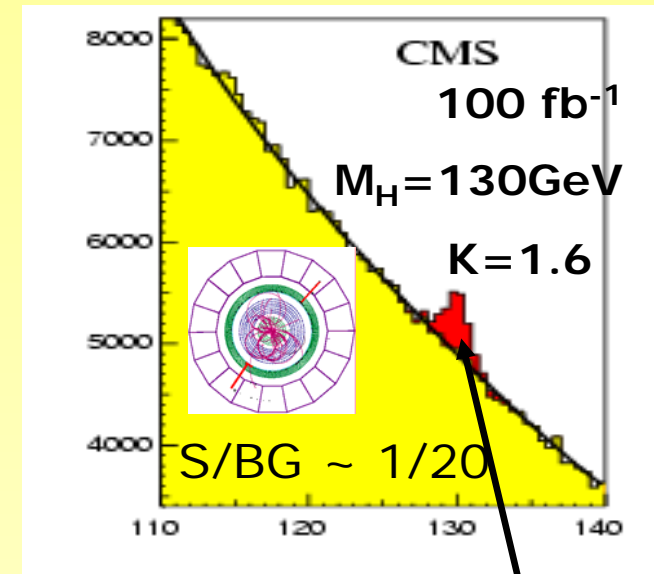
The Higgs decay into  $\gamma\gamma$  is a challenge for the EM calorimeters.

Key points: calorimeter resolution (and linearity), id of photons vs rejection against  $\pi_0$

Calorimeter segmentation is critical to reject jets.  
For  $\sim 80\%$  efficiency, the jet rejection factor is 1000 to 4000, depending on the  $E_T$

Dedicated algorithm for the recovery of photon conversion ( $\sim 1/3$ ) in the material in front of the EM calorimeter

Main backgrounds: irreducible  $\gamma\gamma$  dominant.  $\gamma j$  and  $jj$  together are half the irreducible background



$\sigma_M$ :  $\sim 1\text{ GeV}$



# Higgs decay into ZZ

Final states investigated:  $2e2\mu$ ,  $4\mu$ ,  $4e$

Irreducible background coming from direct ZZ production

Reducible background coming mainly from  $t\bar{t}, Zbb$

The reducible background are strongly reduced by isolation criteria on the leptons and by b-veto

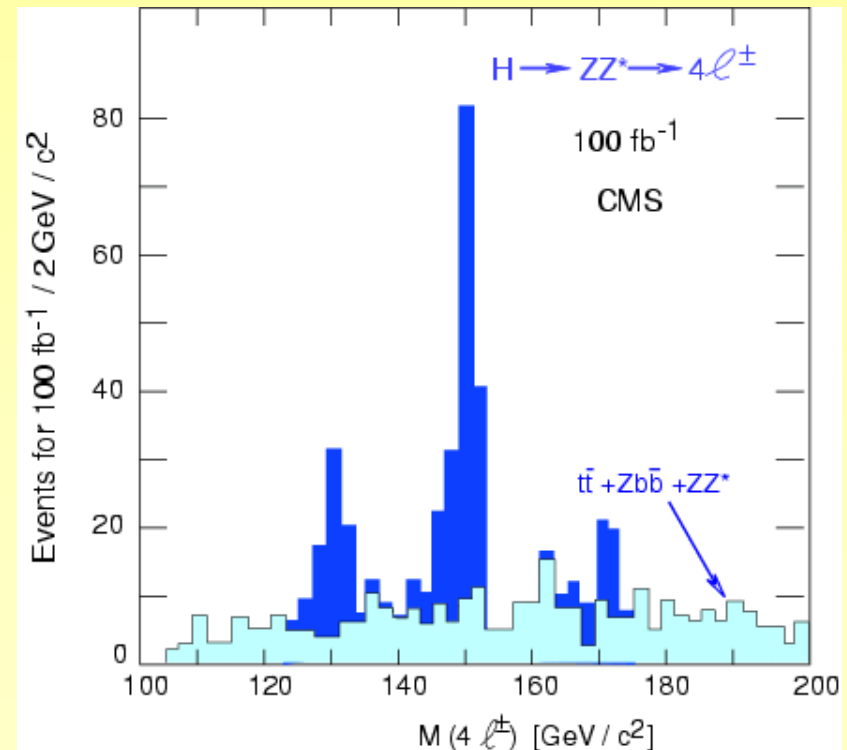
The channel is useful for the Higgs discovery in the ranges

$$130 \text{ GeV} < M_H < 150 \text{ GeV}$$

and

$$180 \text{ GeV} < M_H < 600 \text{ GeV}$$

( $H \rightarrow WW$  is opened at  $\sim 160 \text{ GeV}$ )

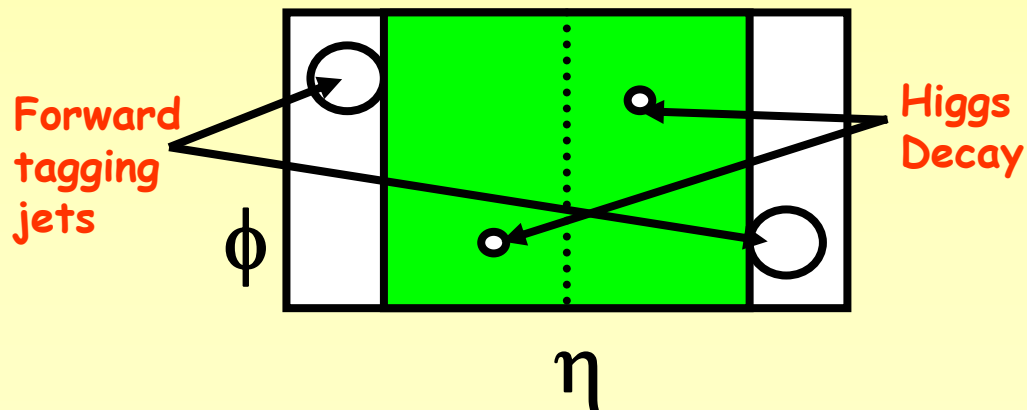
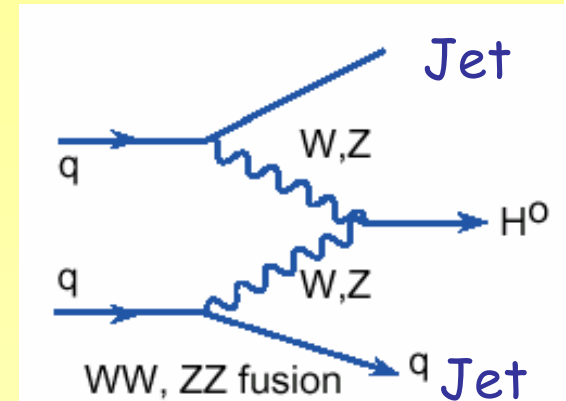


# Vector Boson Fusion

During the last years a lot of work has been put on the VBF production channels.

One expects two hard jets in the forward and backward regions of the detector  $\rightarrow$  forward jet tagging

Lack of color exchange between partons in the initial state  $\rightarrow$  reduced hadronic activity in the central region  
 $\rightarrow$  central jet veto (pile-up effect critical)



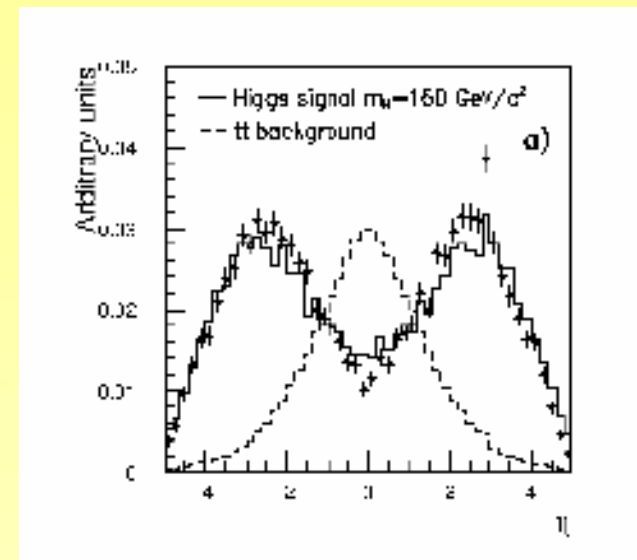
Extremely useful for Higgs discovery in the low and intermediate mass region

It can be used for several Higgs decay channel  $\rightarrow$  useful to measure Higgs coupling

# Vector Boson Fusion (2)

Implementation of fwd jet tagging and jet veto algorithms:

- Associated jets are “forward”. They are well separated in pseudorapidity
- Veto any jet in the central region (except decay products of the Higgs)



Forward jets are the “signature” of VBF

Central jet veto effective for QCD background rejections, in particular against the inclusive tt production, which is a common background for all the VBF channels

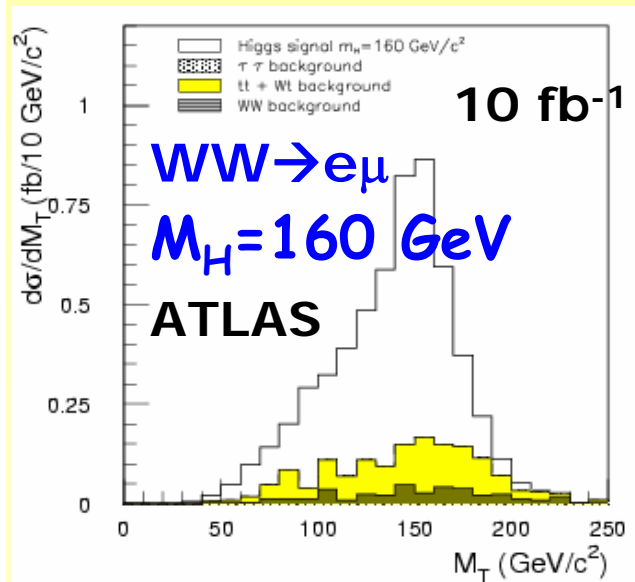
# VBF $H \rightarrow WW$

The best results are obtained if the leptonic decay is considered for both the Ws.

Dominant backgrounds come from  $tt$  production (reduced by the central jet veto and b-jet veto),  $WW$  production

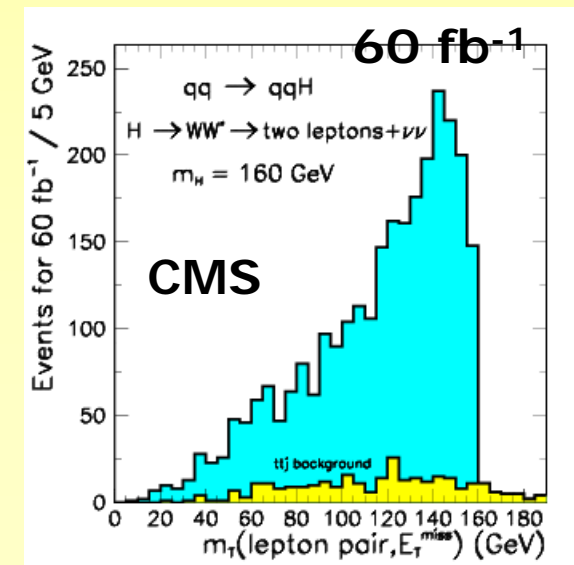
Selection includes fwd tagging, central and b jet veto, spin correlations between leptons

Careful study on the background systematics due to generator uncertainties  $\rightarrow$  most of the multijet final states generated with ME generators. There are preliminary results for NLO  $tt$  calculation.



The normalization of the background value can be estimated at 10% level from data. Background shape taken from MC

The channel is one of the most promising for  $135 \text{ GeV} < M_H < 190 \text{ GeV}$



# VBF $H \rightarrow \tau\tau$

Investigated both in the  $\ell\ell E_T^{\text{miss}}$  and in the  $\ell j E_T^{\text{miss}}$  channel

ATLAS: it provides the cleanest signal for the discovery in the low mass range  
( $M_H < 140\text{-}150$  GeV)

CMS: useful complementary tool to the  $H \rightarrow \gamma\gamma$  channel

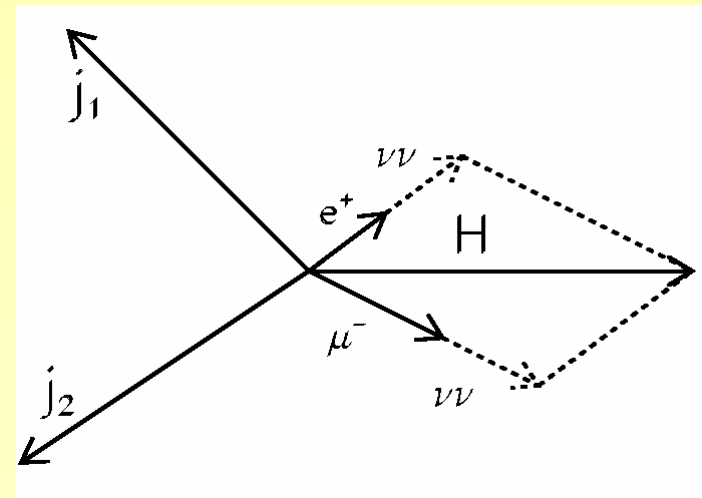
The difference relies on the difference in calorimeter and muon spectrometer design

Main backgrounds:  $Z+nj$ ,  $tt$ ,  $W+nj$  (hadronic decay of one of the two  $\tau$ )

The Higgs mass can be completely reconstructed using the collinear approximation for the  $\tau$  decay.

Defining: 
$$x_i = \frac{P_{l_i}}{P_{l_i} + P_i^{\text{miss}}}$$

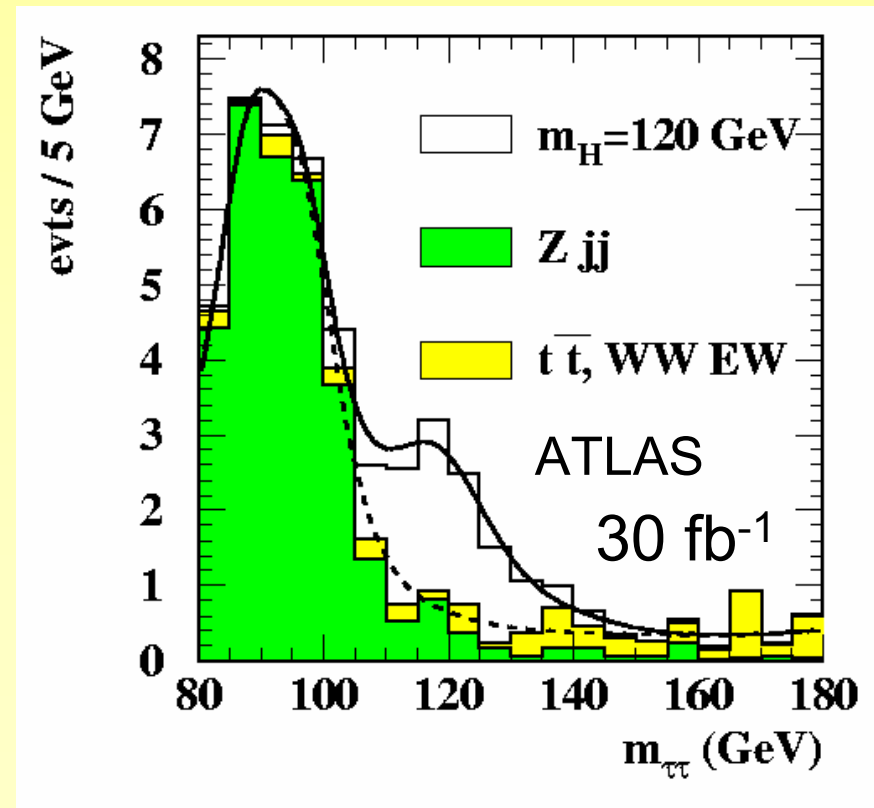
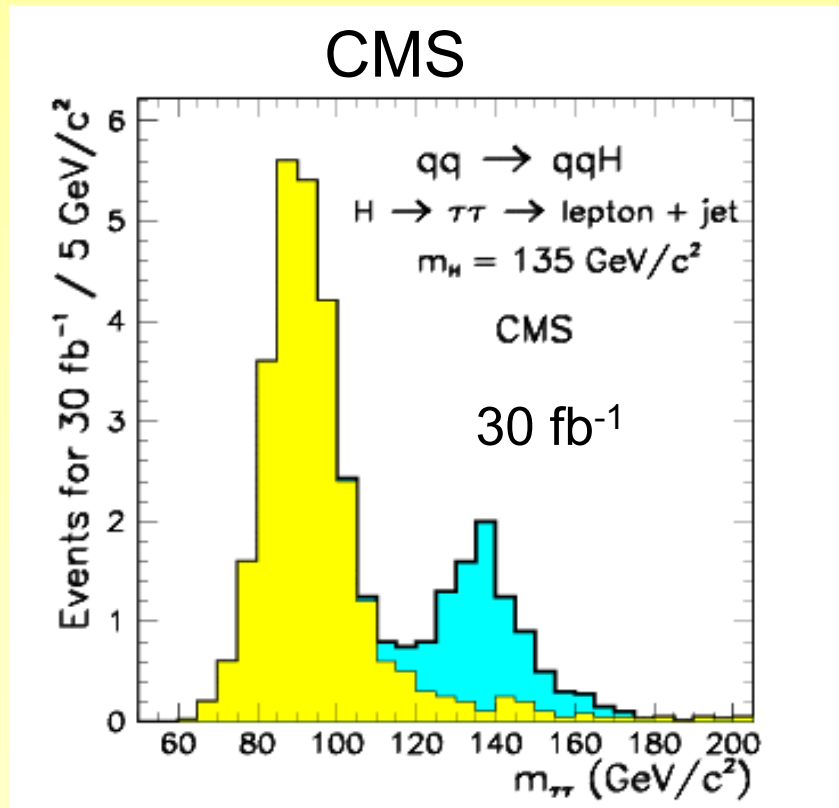
It can be shown that 
$$m_{\tau\tau} = \frac{m_{ll}}{\sqrt{x_1 x_2}}$$



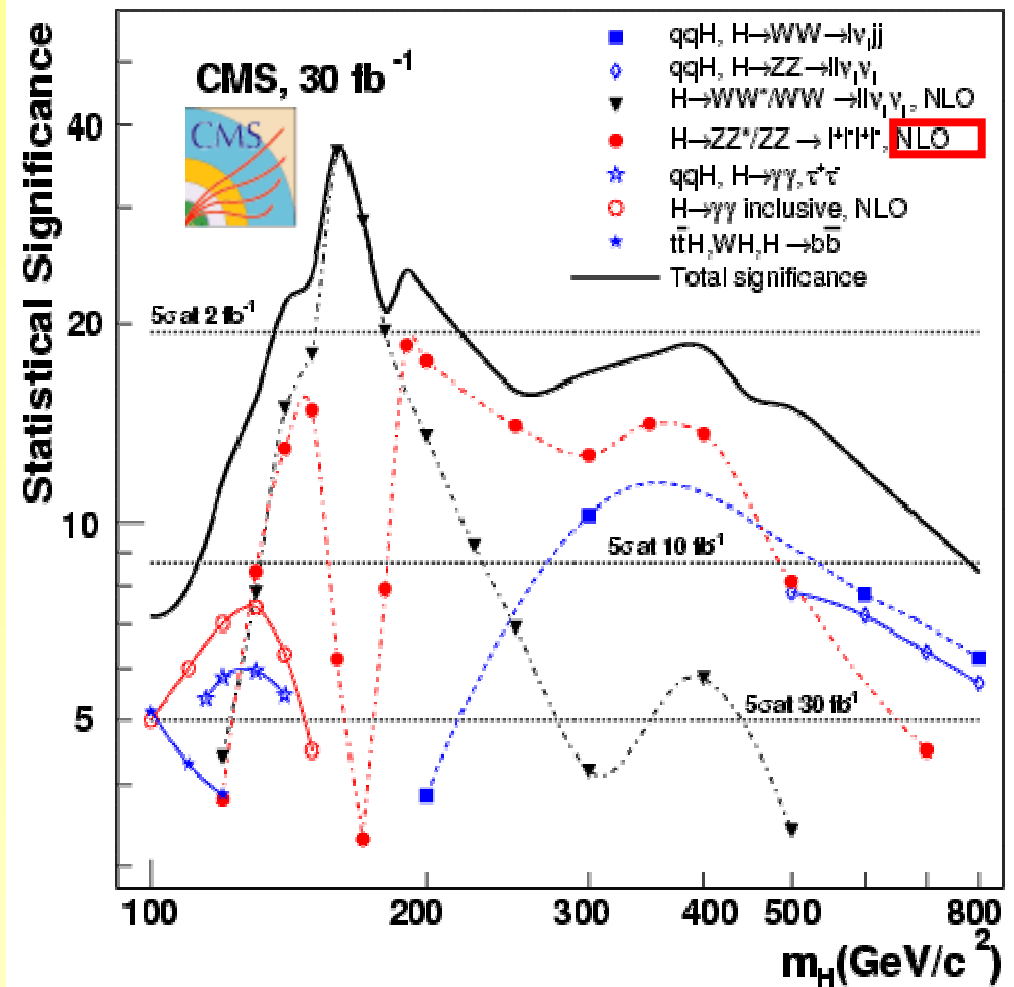
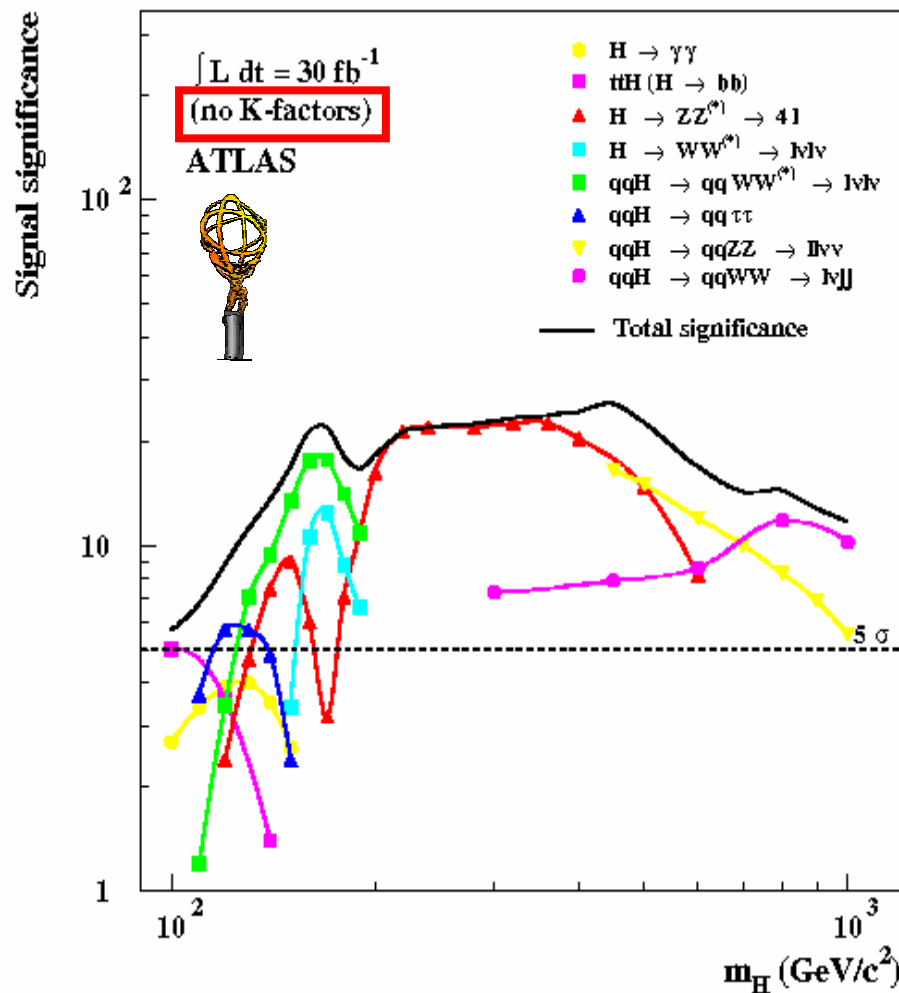
# VBF $H \rightarrow \tau\tau(2)$

Statistical significance (Poisson statistics) above 5 for  $30 \text{ fb}^{-1}$  if  $110 \text{ GeV} < M_H < 140 \text{ GeV}$

Background normalization control at 10% level from the expected  $Z \rightarrow \tau\tau$  peak and from the high sideband



# SM Discovery potential



Almost all allowed mass range explored in 1<sup>st</sup> year ( $10 \text{ fb}^{-1}$ ) for ATLAS-CMS

With  $30 \text{ fb}^{-1}$ , more than  $7 \sigma$  for the whole range (provided systematics on the background are under control)

Low mass region is the most difficult: combination of different channels, all channels put strong requirements on detector

# Determination of Higgs parameters

Once the Higgs boson would be discovered, we want to measure its parameters (mass, spin, partial widths, coupling constants)

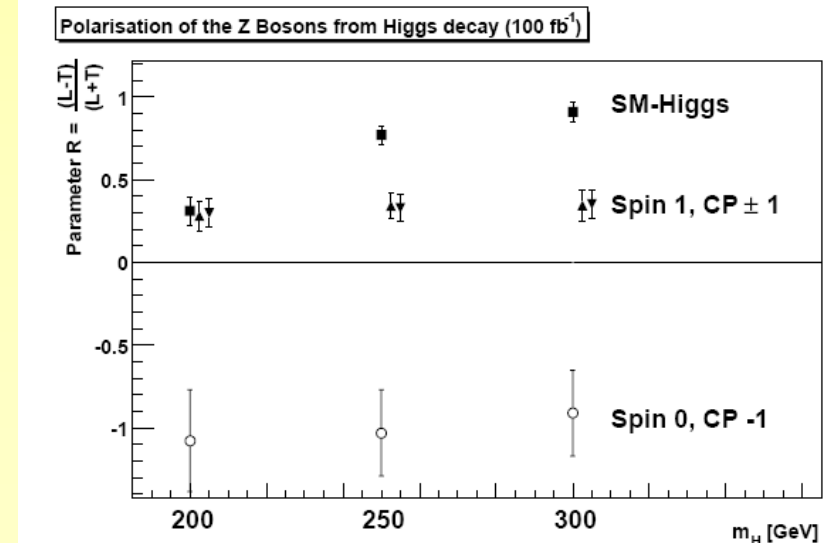
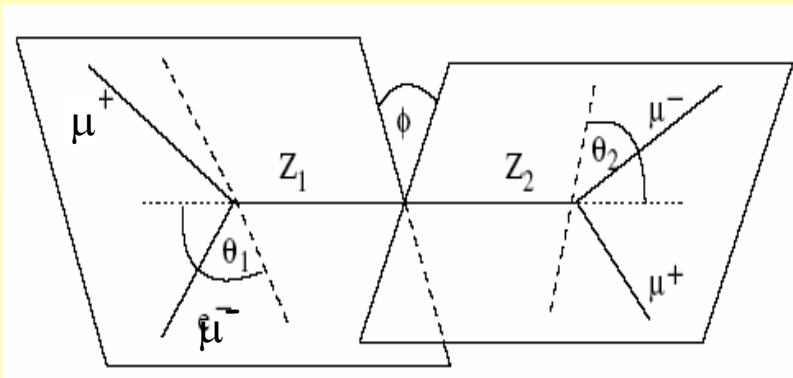
Most of the coupling measurements will be possible because in the whole Higgs mass range more than one decay is accessible at one time

Spin determination examples:

- spin 1 ruled out if  $H \rightarrow \gamma\gamma$  or  $gg \rightarrow H$  is observed
- $H \rightarrow ZZ$ : angular distributions provide information on the spin

$$G(\theta) = T(1 + \cos^2 \theta) + L \sin^2 \theta$$

$$R = \frac{L - T}{L + T}$$

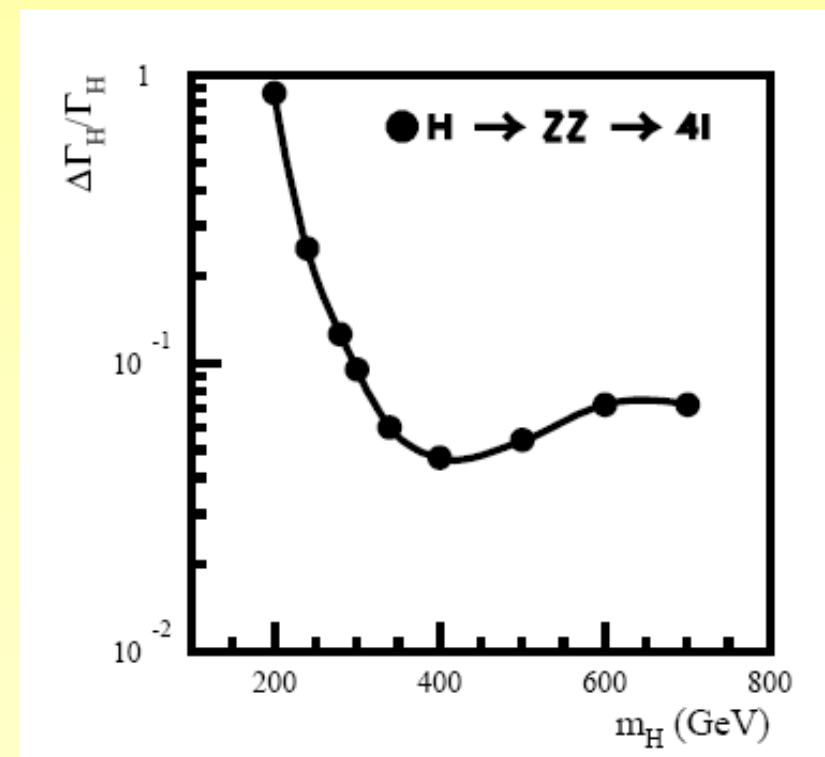
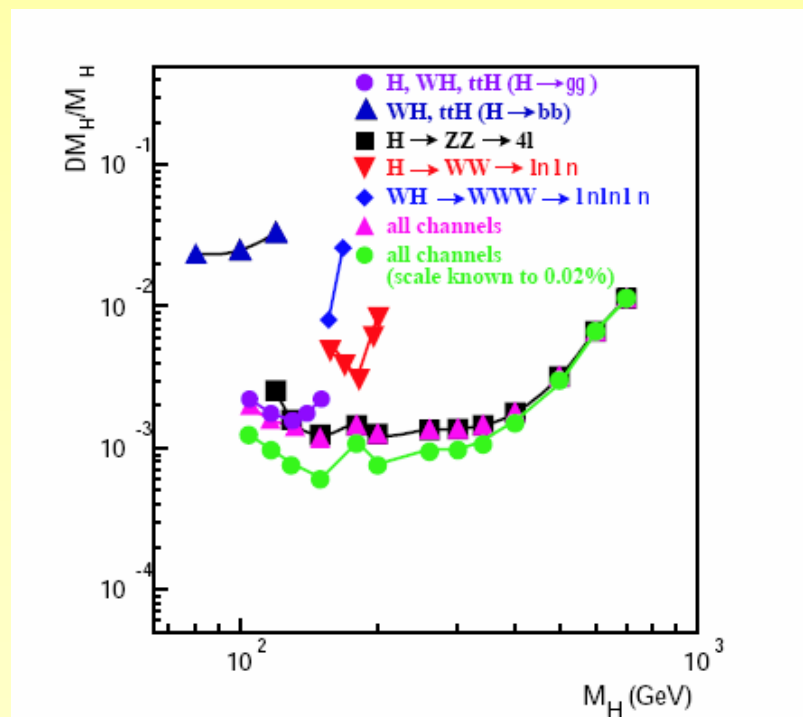




## Determination of Higgs parameter(2)

Precision on the mass: it assumes 0.1% precision on lepton measurement, 1% precision on jet measurement (probably optimistic....)

Width determination: detector dominated is  $M_H < 200$  GeV. If  $M_H > 200$  GeV the  $H \rightarrow 4l$  is used for direct determination.



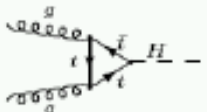

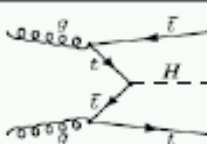
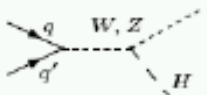

# Measuring Higgs couplings

- Use all the observable signals to perform a likelihood fit.

Take into account:

- Statistical fluctuations
- Cross talk between signal channels.
- Systematic uncertainties (luminosity, efficiencies).
- Statistical and systematic uncertainties on the backgrounds.

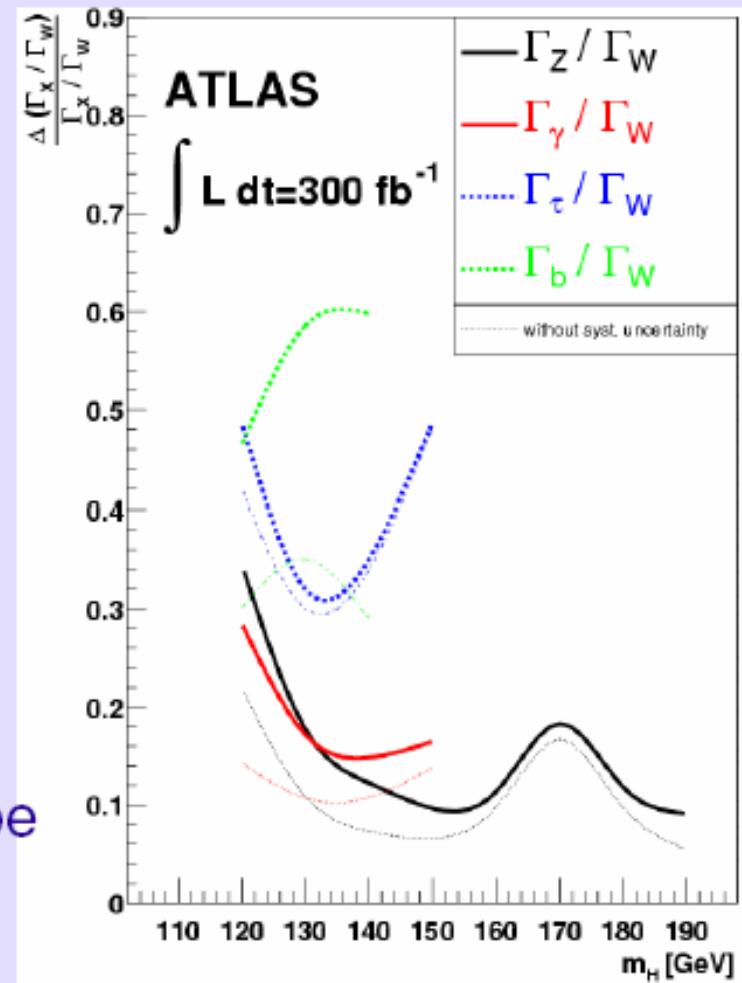
Making different assumptions we can extract different quantities.

Production	Decay
 GF: Gluon Fusion ( $gg \rightarrow H$ )	$H \rightarrow ZZ^{(*)} \rightarrow 4l$ $H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$ $H \rightarrow \gamma\gamma$
 WBF: Weak Boson Fusion ( $qq \rightarrow H$ )	$H \rightarrow ZZ^{(*)} \rightarrow 4l$ $H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$ $H \rightarrow \tau\tau \rightarrow l\nu l\nu$ $H \rightarrow \tau\tau \rightarrow l\nu \text{ had}\nu$ $H \rightarrow \gamma\gamma$
 $t\bar{t}H$	$H \rightarrow WW^{(*)} \rightarrow l\nu l\nu (l\nu)$ $H \rightarrow b\bar{b}$ $H \rightarrow \tau\tau$ (not included) $H \rightarrow \gamma\gamma$
 $WH$	$H \rightarrow WW^{(*)} \rightarrow l\nu l\nu (l\nu)$ $H \rightarrow \gamma\gamma$
 $ZH$	$H \rightarrow \gamma\gamma$

# Ratio of partial decay widths

Likelihood method assuming:  
Spin=0, CP-even and only one Higgs boson.

- $H \rightarrow bb$ : low S/B, large systematics.
  - Uncertainty on the background dominates.
- For all other channels:
  - Ratios of partial widths can be measured with accuracy better than 50%



# Higgs boson couplings

Same method but assuming:

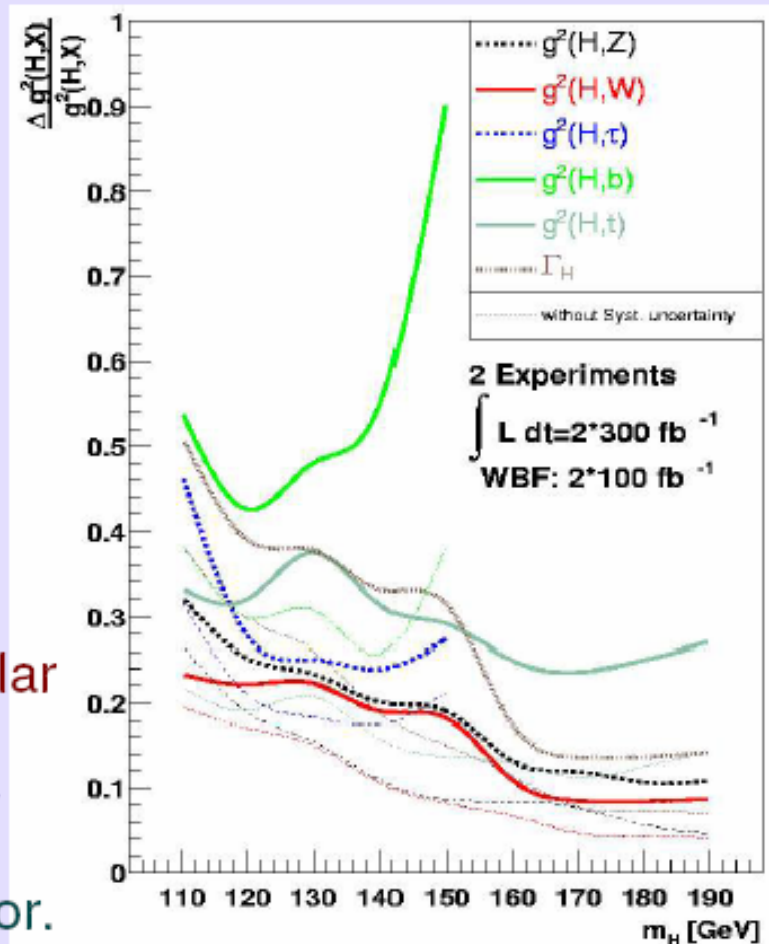
- $\Gamma_V \leq \Gamma_V^{\text{SM}}$ ,  $V = W, Z$ .
- Allow for undetected Higgs decays and contributions to loops from non SM particles.

Assumptions valid for an arbitrary number of Higgs doublets (with or without Higgs singlet) and in particular for MSSM.

Precisions of 20-50% for  $m_H < 150$  GeV

Systematics account for half of the error.

Dominant systematic for  $m_H < 150$  GeV: background.



# Higgs boson self coupling

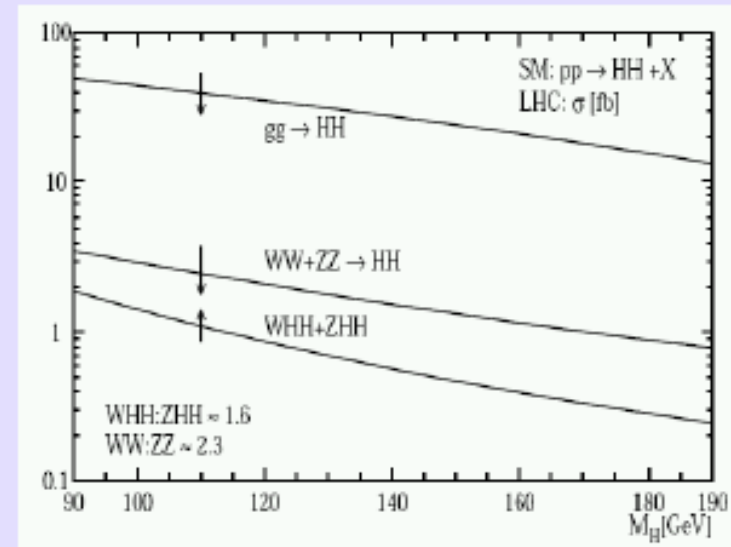
Higgs potencial:  $V_H = \frac{1}{2}m_H^2 H^2 + \lambda v H^3 + \frac{1}{4}\lambda' H^4$

$v = (\sqrt{2}G_F)^{-\frac{1}{2}}$  : vacuum expectation value.

SM:  $\lambda = \lambda' = \frac{m_H^2}{2v^2}$

HHHH not accesible at LHC

- Studied gg fusion for  $HH \rightarrow WWWW \rightarrow \nu l q q \nu l q q$ .
- Main backgrounds:
  - $WWWjj$
  - $ttW$
  - $ttj$  + lepton from b decay.
- S/B ratio ~ 10%.



- Signal not observable at LHC
- Some sensitivity could be possible at Super-LHC.
- Bounds to be investigated.

## SM Higgs boson - conclusion

- ATLAS will be able to detect a Higgs signal with significance  $> 5\sigma$  for all  $m_H$ , for  $30 \text{ fb}^{-1}$  of data.
- After the discovery, precise measurements of its properties have to be done:
  - The Higgs boson mass can be measured with good resolution for all  $m_H$  (1‰ for  $m_H < 400 \text{ GeV}$ ).
  - Non SM spin/CP hypothesis can be ruled out with  $100 \text{ fb}^{-1}$  for  $m_H > 230 \text{ GeV}$ .
  - Coupling constants could be measured combining all available signals, with a precision of 10-50%, with  $300 \text{ fb}^{-1}$  of data.
  - Higgs self couplings may be accessible at SLHC.

# MSSM Higgs Sector

- MSSM: 2 Higgs doublets → 5 physical bosons:  $h, H, A, H^+, H^-$
- phenomenology at **Born level** described by  $\tan\beta, m_A$
- mass prediction:  $M_h < M_Z$
- couplings:  $g_{\text{MSSM}} = \xi \cdot g_{\text{SM}}$ 
  - no coupling of  $A$  to  $W/Z$
  - large  $\tan\beta$ : large  $\text{BR}(h, H, A \rightarrow \tau\tau, bb)$

$\xi$	$t$	$b/\tau$	$W/Z$
$h$	$\cos\alpha/\sin\beta$	$-\sin\alpha/\cos\beta$	$\sin(\alpha-\beta)$
$H$	$\sin\alpha/\sin\beta$	$\cos\alpha/\cos\beta$	$\cos(\alpha-\beta)$
$A$	$\cot\beta$	$\tan\beta$	-----

$\alpha$ : mixing angle between CP even Higgs bosons  
(calculable from  $\tan\beta$  and  $M_A$ )

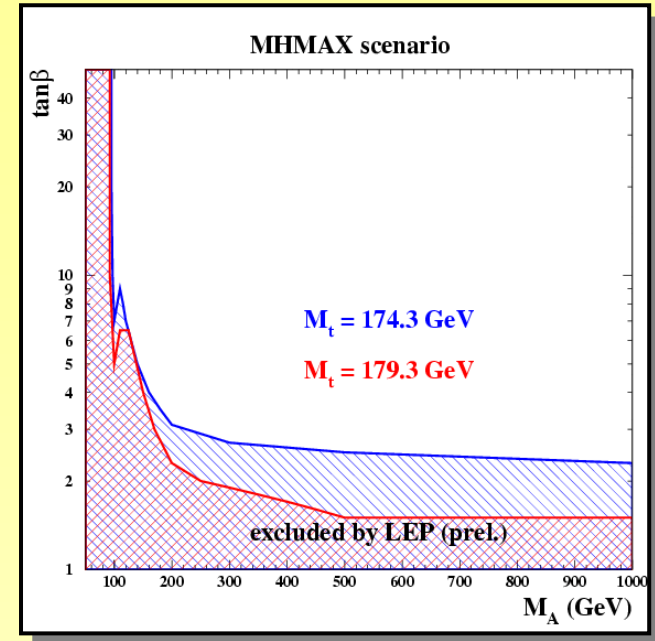
- **large loop corrections** to masses and couplings
- mainly dependent on  **$t/\tilde{t}$  sector**
- parameters:  
 $M_{\text{top}}$  and  $X_t, M_{\text{SUSY}}, M_2, \mu, M_{\text{gluino}}$
- mass prediction  $M_h < 133 \text{ GeV}$   
(for  $M_t = 175 \text{ GeV}$ )

for exclusion bounds and discovery potential:  
**fix the 5 parameters in benchmark scenarios and scan  $(\tan\beta, M_A)$ - plane**

# The $(\tan\beta, M_A)$ -Plane

## current exclusion in $(\tan\beta, M_A)$ -plane:

- LEP excludes low  $\tan\beta$  and low  $M_A$  region
- note: no exclusion from LEP for  $M_t$  larger  $\sim 183$  GeV



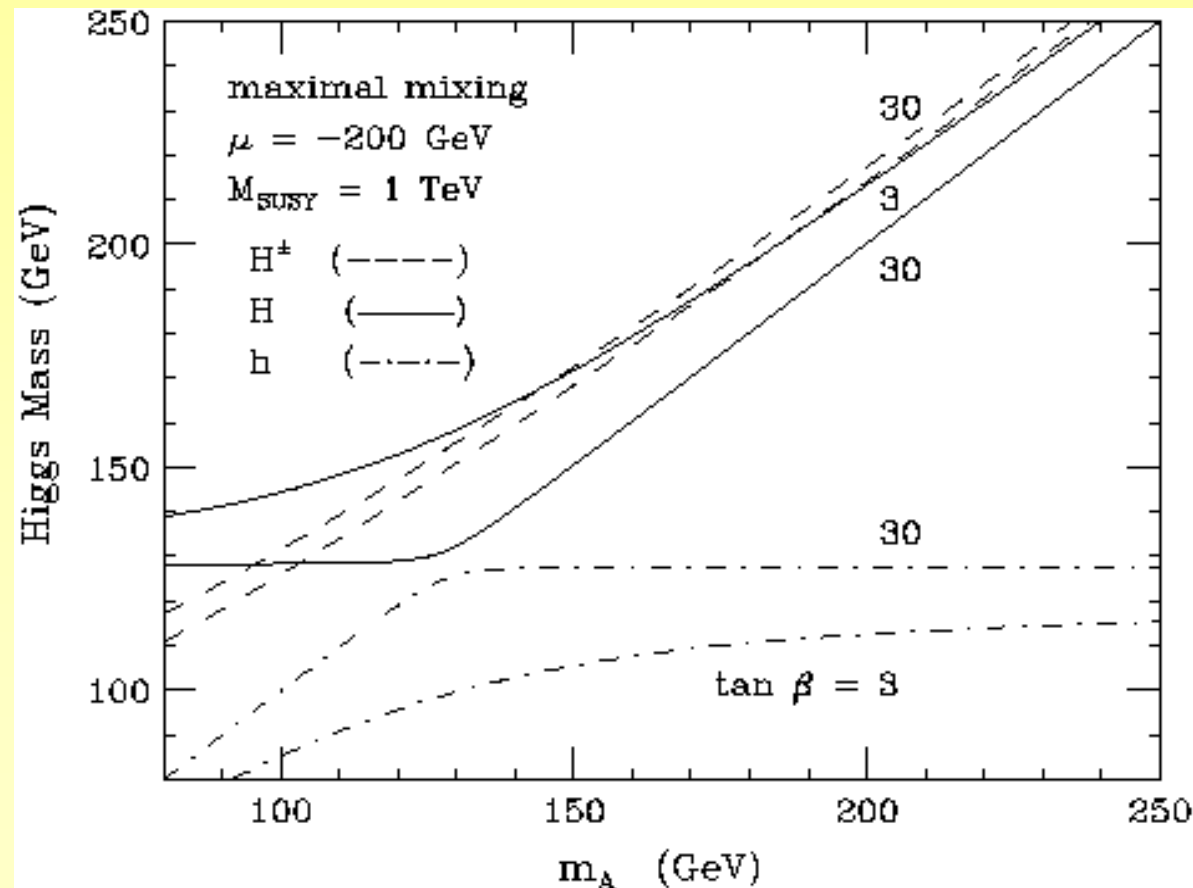
## main questions for LHC/ ATLAS:

- Can at least 1 Higgs be discovered in the allowed parameter space?
- How many Higgs bosons can be observed ?
- Can the SM be discriminated from models with extended Higgs sectors (like MSSM) ?



# MSSM

At high  $M_A$  the heavy bosons degenerate in mass while the  $h$  saturate at a limit value (around 130 GeV)

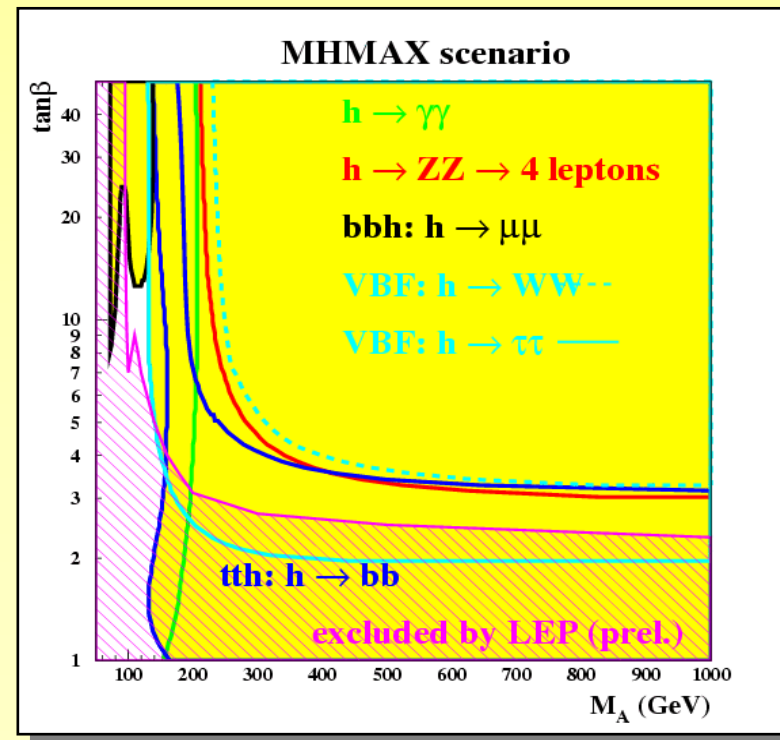


# h search at LHC

Most of the studies done in the LEP **maximal mixing scenario** ( $\mu = -200$  GeV,  $M_{\text{SUSY}} = 1$  TeV,  $M_2 = 200$  GeV,  $M_g = 800$  GeV)

In the decoupling limit (high  $M_A$ ), the  $h$  behaves like a SM Higgs boson. Decay into  $bb$  enhanced at high  $\tan\beta$

- Most of the parameter plane covered by  $h \rightarrow bb$
- VBF channels also important
- large area covered by several channels
- stable discovery and parameter determination possible



Discovery contours

# Heavy Neutral Higgs Bosons

$bbH/A \rightarrow bb\mu\mu$ : cross section higher than in the SM at high  $\tan\beta$ .

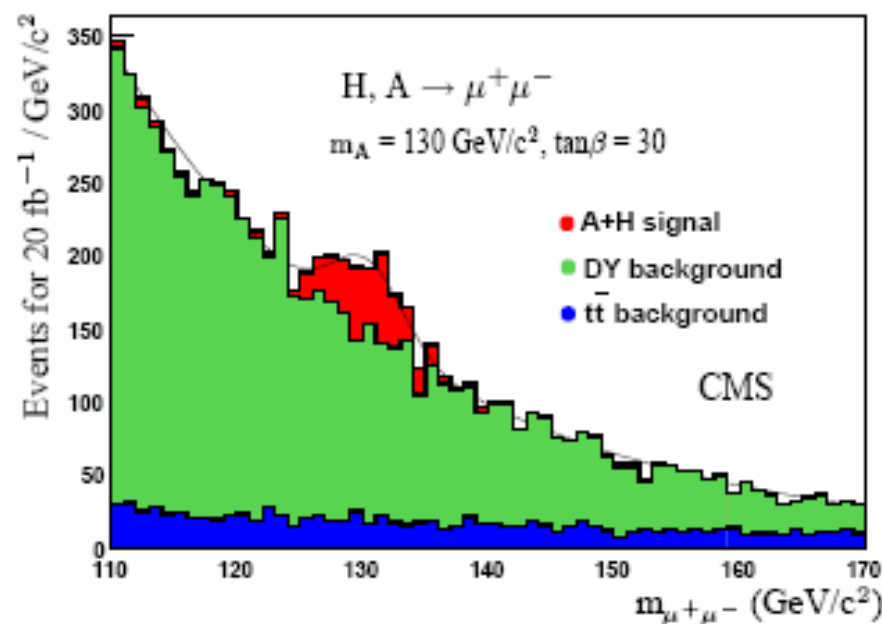
Studied performed in full simulation by CMS.

Main backgrounds coming from Drell-Yan production and  $t\bar{t}$ .

At 130 GeV, the resolution (1%) is not enough to resolve the mass difference between H and A

Most of the background rejected with the requirement of b-tag.

Since the associated b are soft ( $E_T < 50$  GeV), a special tagging algorithm (use of secondary vertex and impact parameter without the jet requirement in the calorimeter) has been used to increase the b-jet ID



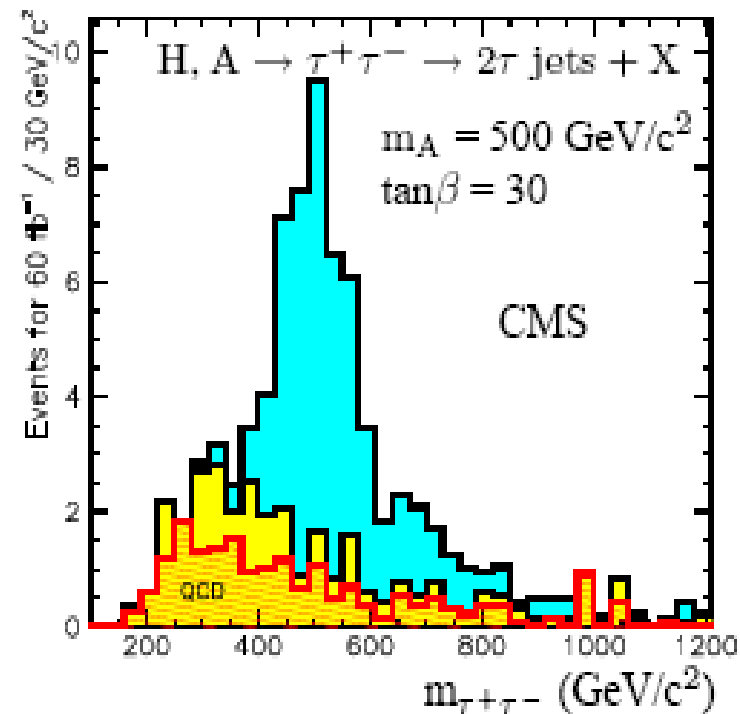
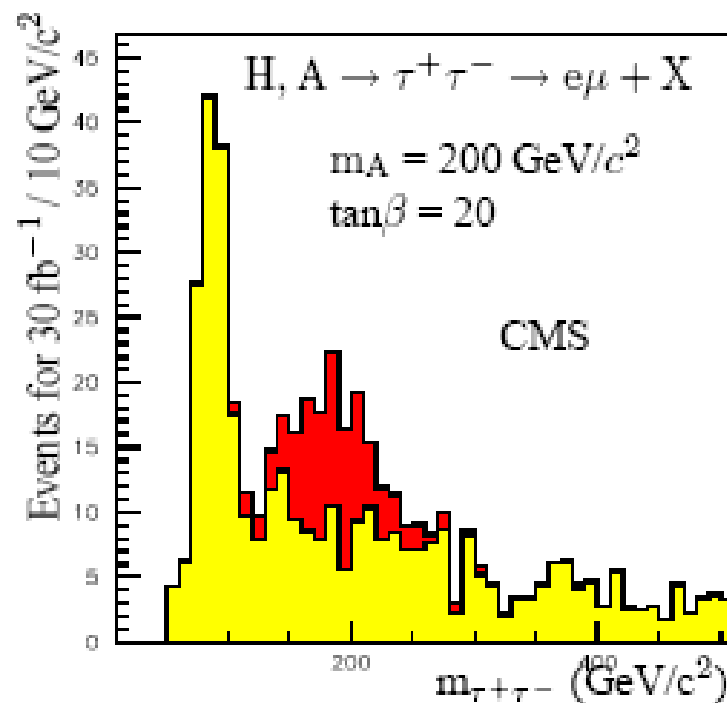
# H/A decay into $\tau$ 's

The tau decays provides the cleanest signature for the heavy Higgs discovery at high mass (and relatively high  $\tan\beta$ )

Investigated all the final states (ll, lj, jj). All of them contribute (at different  $M_A$ ).

The associated production (bbA/H) provides additional rejection against the main backgrounds (Z+jet, WW, QCD (double hadronic tau decay only))

Trigger issues solved w.r.t. the full hadronic final state



## Overall coverage of the $M_A$ - $\tan\beta$ plane for $A/H$

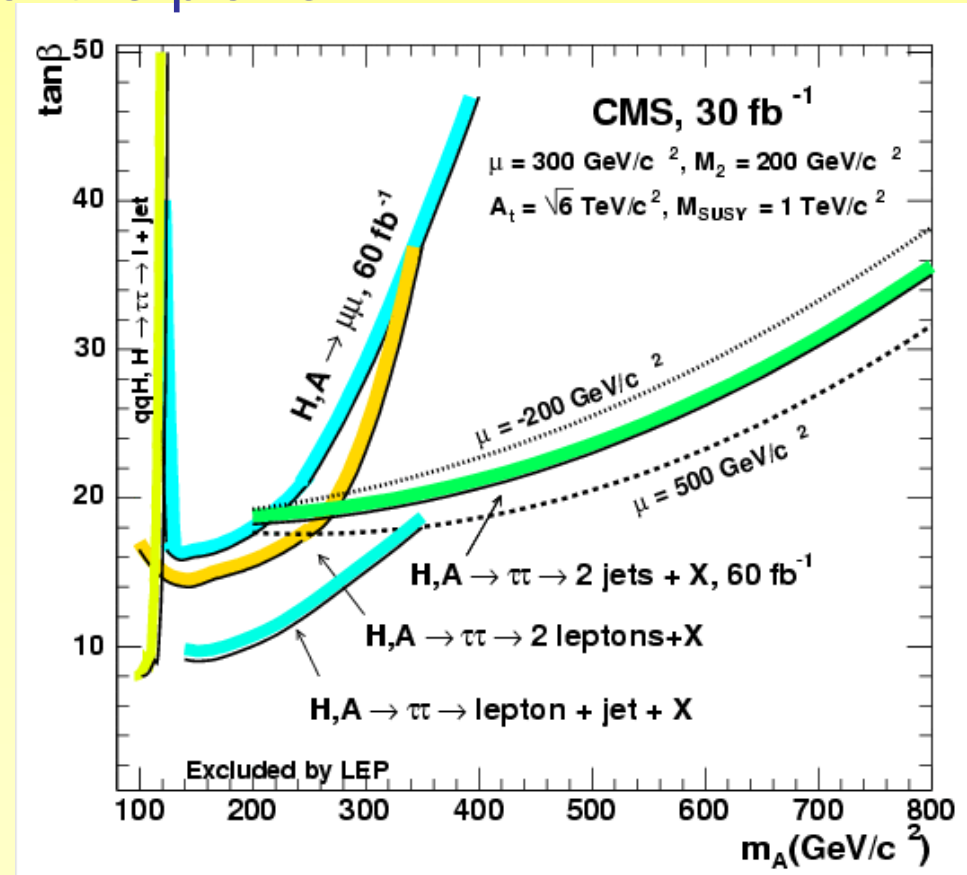
The  $A/H$  Higgs boson can be detected thanks to the  $A/H \rightarrow \mu\mu$  decay in the low mass- high  $\tan\beta$  region

The decay into  $\tau\tau$  allows (with different final states) to cover a large part of the plane

There is no coverage (with  $30 \text{ fb}^{-1}$ ) for the large mass-intermediate  $\tan\beta$  region

The very high mass region can be reached only by the full hadronic final state of the  $\tau\tau$

**Similar plot for ATLAS**



# Charged Higgs decay

Two channels are of particular interest if  $M_H > M_t$ .

- $H^+ \rightarrow \tau \nu$ ;  $\tau \rightarrow \text{hadr}$

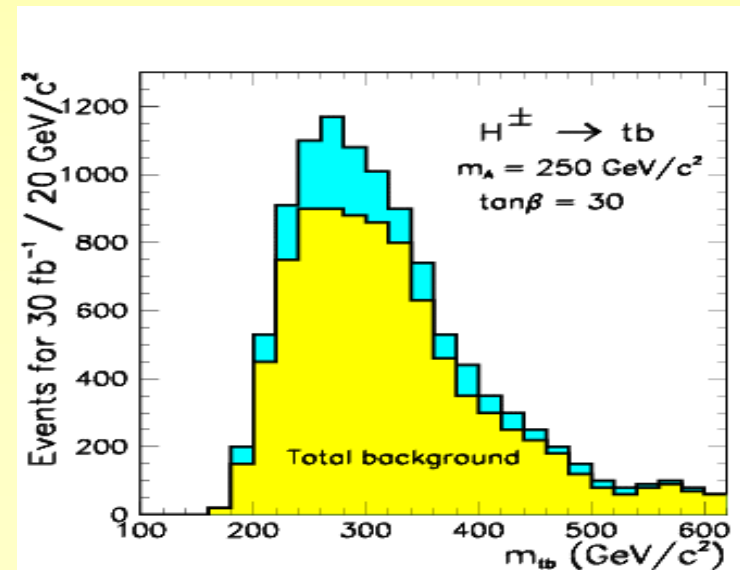
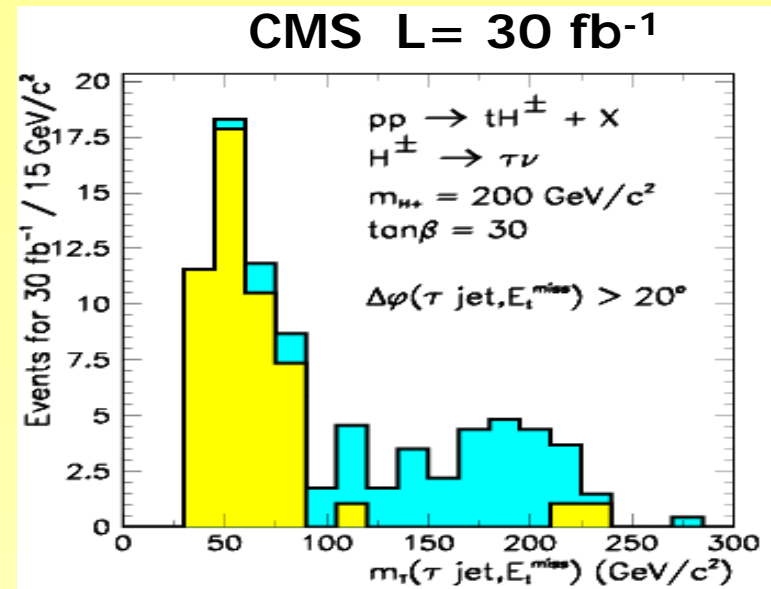
Reconstruction of the top decay and of the hadronic tau.

$t\bar{t}$  and  $Wt$  backgrounds suppressed exploiting spin correlations  $\rightarrow$  One single  $\pi$  has to carry the largest part of the tau energy.

- $H^+ \rightarrow t\bar{b}$   $t \rightarrow bW \rightarrow b\ell\nu q\bar{q}$

large background ( $t\bar{t}$ +jets)

reconstruct tops and Higgs  $\rightarrow M_H$

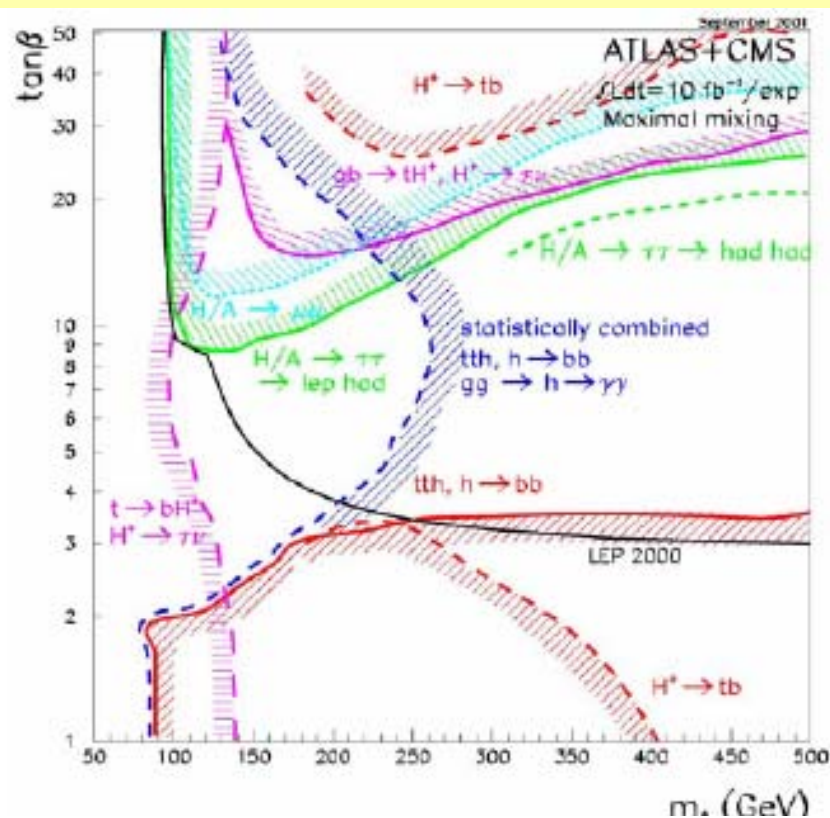


# Overall view ( $10 \text{ fb}^{-1}$ )

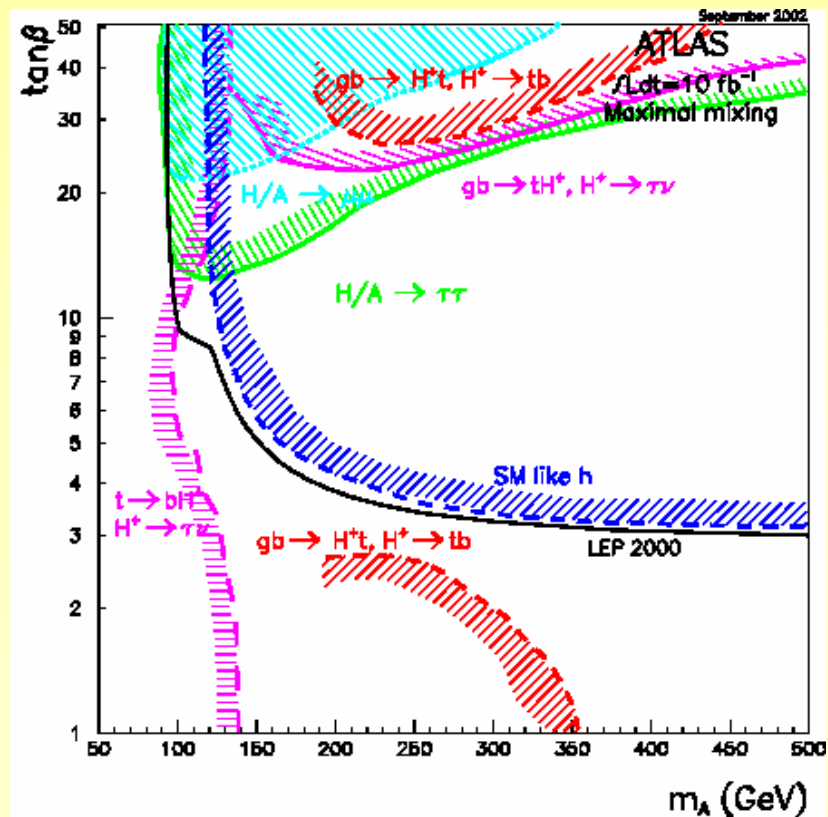
Rather small statistics (of good data) is enough to cover most of the  $M_A$ - $\tan\beta$  plane

But a lot of questions to address before: machine operations, detector understanding, background measurement....

No VBF



with VBF



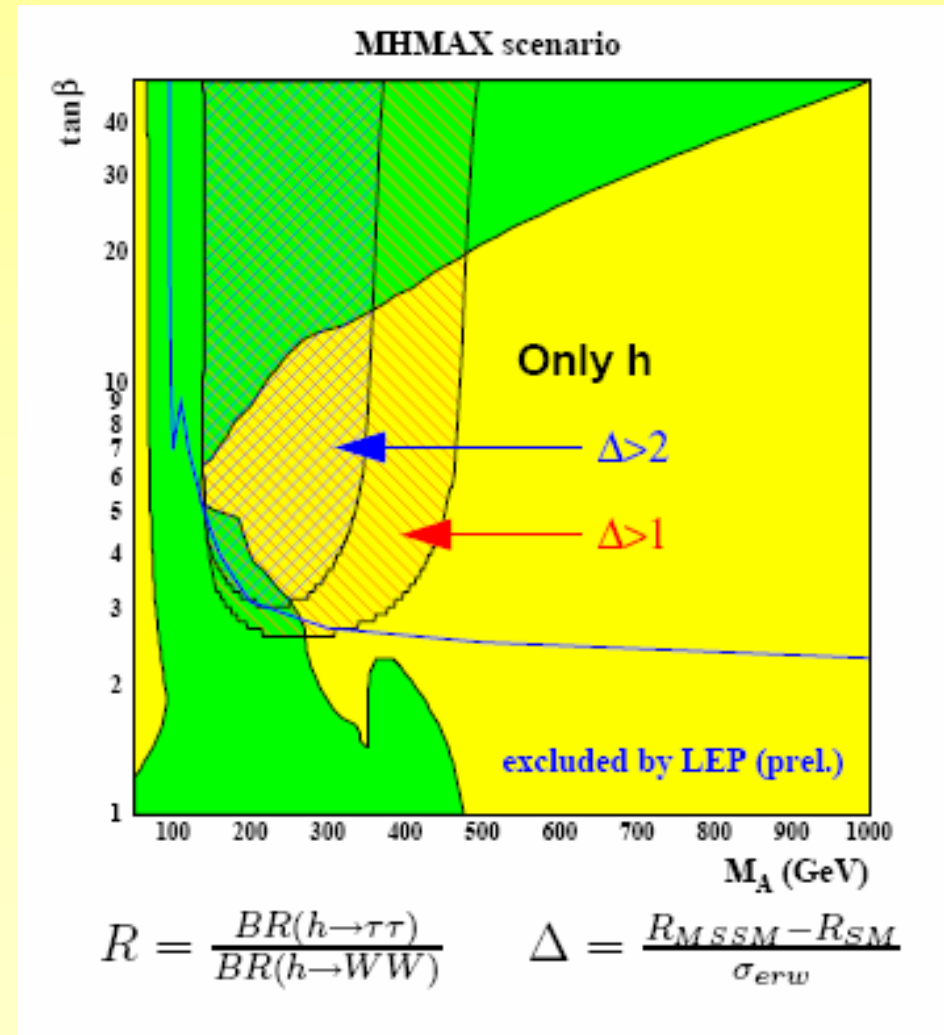
# Overall view ( $300 \text{ fb}^{-1}$ )

Can we disentangle between SM and MSSM just from the Higgs sector?

In most of the parameters space more than 1 Higgs boson is detectable.

In the yellow region just the  $h$  is detectable

Still in a region we can disentangle looking at the ratio between the BR into  $\tau\tau$  and  $WW$ . Anyway, it is difficult....





# Conclusions

-Standard Model Higgs detectable with few tens of  $\text{fb}^{-1}$  over all the mass range. In the most likely region, VBF production plays a relevant role

-Spin measurement feasible in its ZZ decay.

-Coupling measurement will require full luminosity in most of the mass range. Achievable precisions: 15-50% depending on the channel

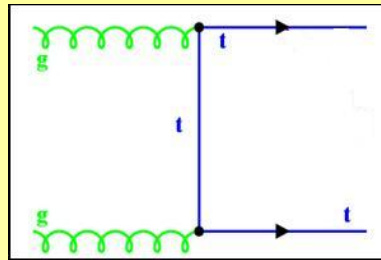
-MSSM plane coverage requires few tens of  $\text{fb}^{-1}$  in a maximal mixing scenario. Higgs decays into third generation fermions play an extremely relevant role. VBF important for the h in the decoupling limit

# Back-up slides

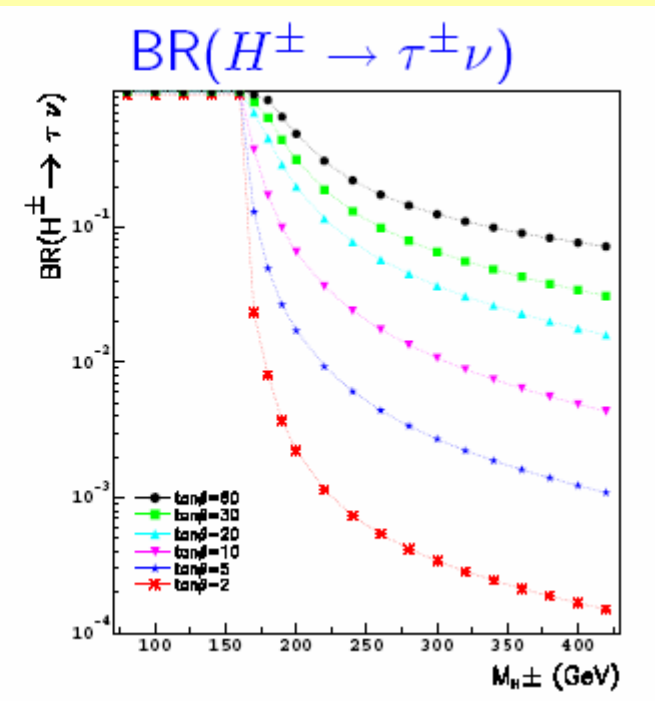
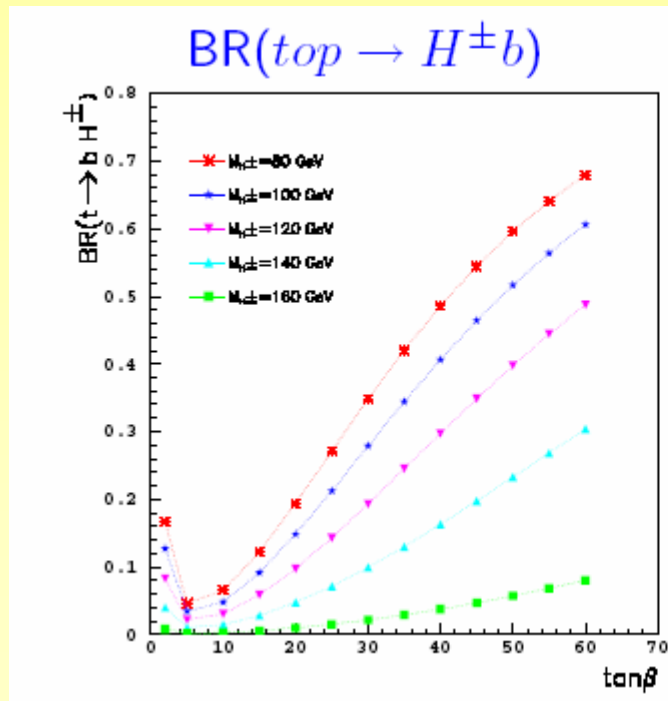
# Higgs search – problem

“Light” charged Higgs search :  $m_{H^\pm} < m_{\text{top}}$

$\sigma_{tt} \sim 1\text{nb}$



large  $t\bar{t}$  sample is available



non excluded region:  $Br t \rightarrow H^\pm b$  20%  $\downarrow$  0%

$Br H^+ \rightarrow \tau \nu \sim 93\%$

# Higgs search – problem

How to observe the signal ?

- Draw  $t\bar{t}$  decay topology in most important decay modes
- $\text{Br}(t \rightarrow Wb)$  complement of  $\text{Br}(t \rightarrow Wb)$
- $\text{Br}(W \rightarrow l\nu) \sim 10\%$        $\text{Br}(W \rightarrow qq) \sim 70\%$
- Check how well the final state is constrained for various decay possibilities
- Think about information that can be obtained from  $p_T^{\text{miss}}$
- Do not forget the trigger

# LHC Phenomenology – problem

## Final states with e or $\mu$ done in '94 in ATLAS

$$\begin{aligned}t &\rightarrow W^+ b \\ W^+ &\rightarrow e \text{ (or } \mu) + X \\ \bar{t} &\rightarrow H^- \bar{b} \\ H^- &\rightarrow \tau^- \nu_\tau \\ \tau^- &\rightarrow \text{had. } \bar{\nu}_\tau\end{aligned}$$

- Trigger on the isolated lepton
- $\text{BR}(W^+ \rightarrow e(\mu) + X)$ : 24.8%
- $M_{H^\pm}$  reconstruction difficult  
(3 or more  $\nu$  in 2 hemispheres)

## Hadronic final state new approach

$$\begin{aligned}t &\rightarrow W^+ b \\ W^+ &\rightarrow q\bar{q}' \\ \bar{t} &\rightarrow H^- \bar{b} \\ H^- &\rightarrow \tau^- \nu_\tau \\ \tau^- &\rightarrow \text{had. } \bar{\nu}_\tau\end{aligned}$$

- Trigger based on jet +  $\cancel{E}_T$
- $\text{BR}(W^+ \rightarrow q\bar{q}')$ : 68.5%
- Simplified  $M_{H^\pm}$  reconstruction  
(2  $\nu$  in the same hemisphere)

Can only look for an excess of taus  
over the SM prediction

Some mass reconstruction is possible:

- show that it is a particle
- measure its mass

# LHC Phenomenology – problem

$$M_T = \sqrt{2 \cdot E_T^{miss} \cdot p_T(\tau - jet) \cdot \cos(\Delta\phi)}$$

