

Explicit solvent effects in the electrostatics of Newton Black Films: evidence from computer simulations

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Acknowledgements:

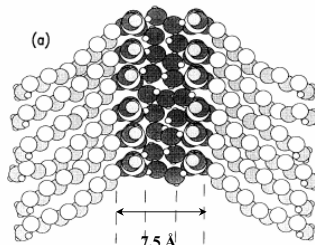
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Simulations were performed at the HPCx Supercomputing facility (Edinburgh, UK) under the Materials Chemistry Consortium

Motivation: to perform a test of basic hypothesis of theories on forces

- **Electrostatics of aqueous systems** usually described by “**primitive**” theories (solvent characterized by a constant ϵ_r)
- This hypothesis is **true at large length scales**, but it is **not clear to be OK** at colloidal scale. Several groups claim that the breakdown of this hypothesis can explain results related to **hydration forces** (see for example the classical review by Leikin, Parsegian and Rau (1993)).
- Some groups (Berkowitz 1995, Berendsen 1996) have tried to settle this issue by Molecular Dynamics simulations (of DPPC or DPPS bilayers, for example) but their results were not clear.
- We decided to perform simulations of **SDS/water/SDS Newton Black films**, taking advantage of the new UK national Supercomputing facility HPCx (Edinburgh).
- The model and potentials were relatively “easy” to construct for this system due to the **availability of experimental data**.



SDS [Na⁺CH₃(CH₂)₁₁OSO₃⁻]

Conclusion: explicit solvent effects are very important

● The analysis of our simulations and re-analysis of other authors simulations clearly show that **explicit solvent effects strongly affect electrostatics** in aqueous media

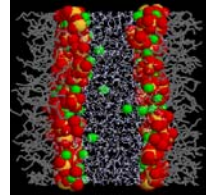
➤ Our Molecular Dynamics simulations of thin SDS/water/SDS films show that:

✓ There is a strong polarization of water at the interfaces, which decays in 10 Å. This is **not in agreement** with the expected dielectric constant of water.

✓ It is not possible to define a **dielectric constant** for the solvent (water) inside these films: there is no local (constitutive) relation between applied field and response of the solvent.

✓ The solvent (water) has an important contribution to the **electrostatic potential profile** inside the film, which is **inconsistent with that expected from a dielectric medium**.

✓ The **strong polarization of water** has an important **repulsive contribution to the electrostatic interaction** between surfactant layers (**hydration force**).



➤ Molecular Dynamics simulations of DPPC bilayers show a similar dielectric response of water.

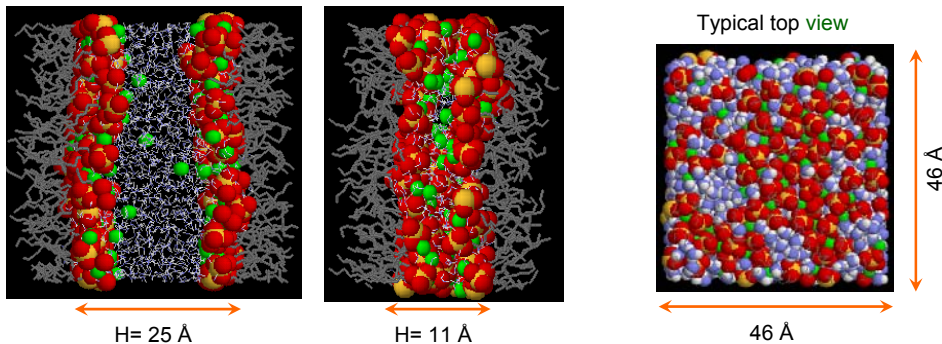
➤ Molecular Dynamics simulations of DPPS bilayers show a dielectric response of water consistent with a small dielectric constant ($\epsilon_r \approx 2.5$)

First, let us briefly describe our Molecular Dynamics Simulations

(this is a little boring part, but I need to do it)

Simulations predict correct results for SDS/water/SDS films

• We have simulated films of different sizes (from $H=32 \text{ \AA}$ to $H=6 \text{ \AA}$) with 33 \AA^2 per surfactant (experimental value observed by B elorgey and Benattar, PRL 1991) at 298 K.



- Roughness **OK!** sims $\zeta(R \rightarrow \infty) \rightarrow 2.5 \text{ \AA}$ experiments $\zeta = 2.7 \pm 0.1 \text{ \AA}$
- Surface tension **OK!** sims $\gamma \approx 58 \text{ mN/m}$ experiments $\gamma \approx 63 \text{ mN/m}$
- Other quantities (diffusion coefficients, tilt angle, electron density,...) **OK!**

The SDS and water models are standard and work very well

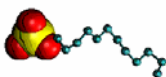
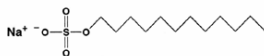
• The force field is based in AMBER 94 with some minor modifications, and reproduces thermodynamic data of SDS and water at ambient conditions. (FB & JF, Langmuir (2004))

- SPC/E model for water, standard in biomolecular simulations



- Low computational cost (tetrahedral water model with 0.1 nm OH distance and partial charges $q_{\text{ox}} = -0.8476e$, $q_{\text{hyd}} = +0.4238e$).
- Correct density, $g(r)$ and diffusion coefficient of liquid water under ordinary conditions.
- Dielectric constant $\epsilon_r \approx 70$ at 1 atm and 300 K

- Surfactant and Na^+ model/effective potentials include description of:



- Bonds and bonds angles
- Bending and torsional potentials
- Nonbonding interactions (Lennard-Jones)
- Electrostatics (partial charges at S, O and CH_2 bonded to oxygen)

Now, let us see what happens inside the film

(this will be the funny part...)

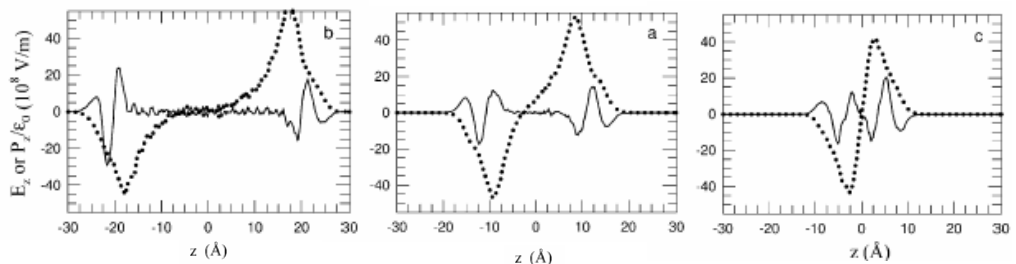
Dielectric constant is not well defined inside SDS/water/SDS films

• The comparison between the Electric field and the Electrostatic polarization shows: $\vec{P} \equiv \vec{D} - \epsilon_0 \vec{E}$

• No dielectric constant near charged interfaces $\vec{P} = \epsilon_0 (\epsilon_r - 1) \vec{E}$

• In fact, it is not possible to define a local (constitutive) relation between applied field and response of the solvent.

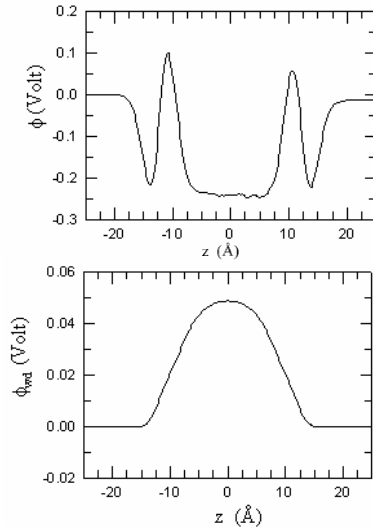
JF & FB, PRL **92**, 236102 (2004)



Polarization maximum at the interface ($P/\epsilon_0 \sim 10^9$ V/m) and decays completely in 10 \AA

Water strongly affects electrostatic potential inside films

- Electrostatic potential inside the aqueous core of the films is **completely different (in value and shape)** from that expected in a film containing a dielectric solvent (characterized with ϵ_r)



JF & FB, PRL **92**, 236102 (2004).

**EXACT
Poisson eq.**

$$-\epsilon_0 \frac{d\phi}{dz} = \rho_f(z) + \rho_w(z)$$

← charge distribution due to water computed in sims

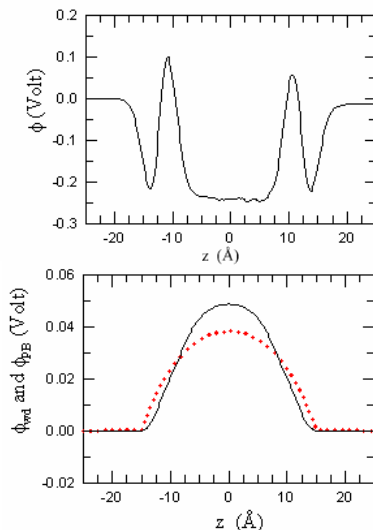
charge distribution due to surfactant, Na⁺ (computed in sims)

Poisson eq. in a dielectric medium

$$-\epsilon_0 \frac{d\phi_{wd}}{dz} = \frac{\rho_f(z)}{\epsilon_r}$$

← water replaced by a dielectric medium

PB equation fails inside these films due to water response



JF & FB, PRL **92**, 236102 (2004).

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Poisson eq.**

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charge distribution computed from sims

Poisson eq. in a dielectric medium

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water replaced by a dielectric medium

Poisson-Boltzmann

$$\begin{cases} -\epsilon_0 \frac{d\phi_{PB}}{dz} = \sum_i q_i n_i(z) \\ n_i(z) = n_{oi} \exp\left[-\frac{q_i \phi(z)}{k_B T}\right] \end{cases}$$

water replaced by a dielectric medium, ion distribution replaced by Mean-Field Boltzmann distribution

Water strongly increases the stabilizing electrostatic force

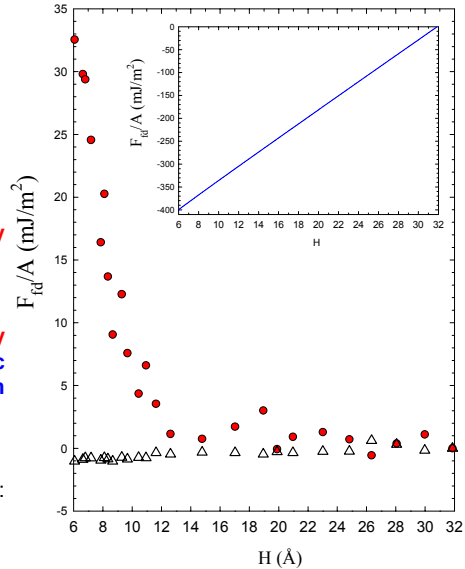
• Calculation of the field Free Energy shows an important repulsive contribution due to the polarization induced near the interfaces.

$$F_{field}(H) = F^{(0)} + \frac{1}{2} \int \rho_f \phi dV$$

• It is **strongly repulsive** using **SDS charge density and potential** computed from simulations (circles).

• It is **weakly attractive** using **SDS charge density computed from simulations** and the **electrostatic potential ϕ_{wd} derived assumed a dielectric medium** (triangles).

• It is **strongly attractive** using Poisson-Boltzmann: charge density and electrostatic potential (solid line).



Water strongly increases the stabilizing electrostatic force

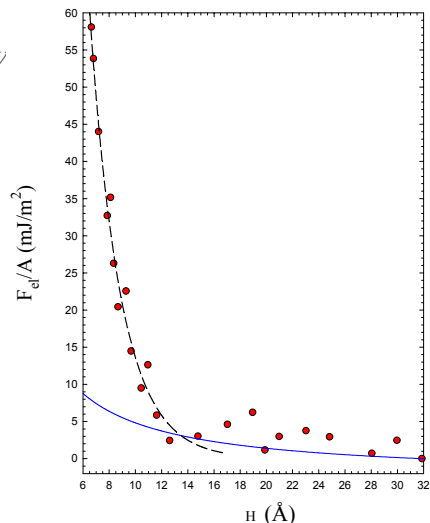
• Approximate calculation of the electrostatic Free Energy shows:

$$F_{el} = F_{field} + F_{mix} \begin{cases} F_{field}(H) = F^{(0)} + \frac{1}{2} \int \rho_f \phi dV \\ F_{mix} = k_B T \sum_i \int n_i \ln X_i dV \end{cases}$$

• A **strong-short distance repulsion** due to **strong water polarization near interfaces**.

• The **repulsive pressure** has the typical exponential behaviour of the **so-called hydration force** (see review by Leikin, Parsegian & Rau (1993)).

$$\Pi_{el} = - \frac{1}{A} \frac{\partial F_{el}}{\partial H} \Big|_T \quad \Pi = \Pi_0 e^{-H/\lambda} \quad \begin{array}{l} \lambda = 2.35 \text{ \AA} \\ \Pi_0 = 4.09 \times 10^9 \text{ Pa} \end{array}$$



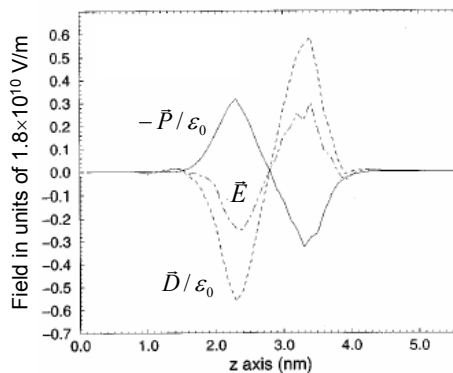
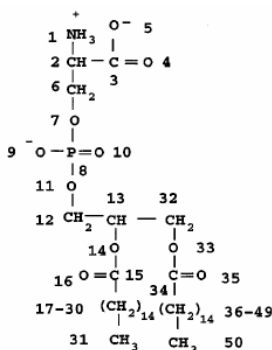
... and some more surprises appear by looking to other systems...

(this is the last part of the talk!!!)

Water inside DPPS films has a small dielectric constant

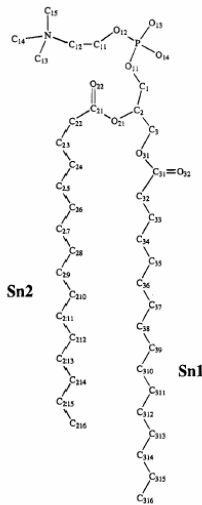
• Simulations of DPPS/water+Na⁺/DPPS⁻ by Berendsen group (1996) at 350 K and 1 atm, with an area per lipid 54 Å² and aqueous core of thickness ≈ 25 Å.

- They do not analyze the electric field but it can be obtained from their published data.
- The results are consistent with a relation between the electric field and water polarization, $\vec{P} = \epsilon_0(\epsilon_r - 1)\vec{E}$ with a **water dielectric constant $\epsilon_r \approx 2.5$** !! .



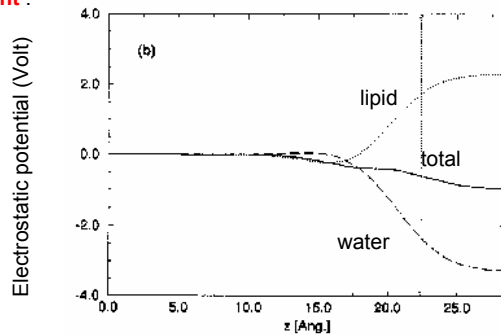
Re-examination of sims in other systems shows similar behaviour

• Simulations of **DPPC bilayers in water** by Berkowitz group (1995) at 333 K, with an area per lipid 65.8 Å² and aqueous core of thickness ≈ 12 Å.



4. Nomenclature of the atoms in the DPPC molecule.

- The electrostatic potential is monotonic (no oscillations!!).
- **Overcompensation**: The electrostatic potential created by water is larger (and opposite) than that created by the lipids.
- Using available data, it is not possible to introduce a well-defined **dielectric constant**.

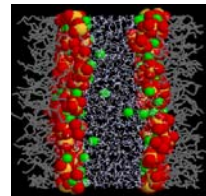


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