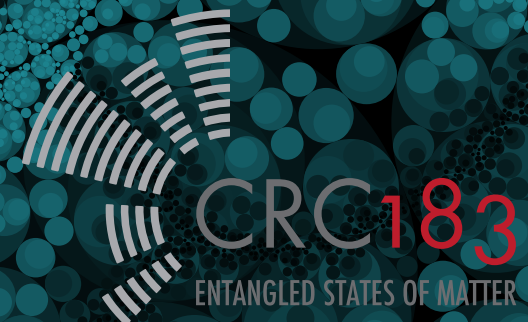


Macroscopic entanglement, spin liquids & Kitaev materials

Entanglement in Strongly Correlated Systems
Benasque, February 2020

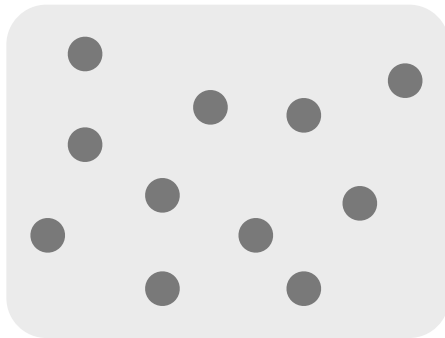
Simon Trebst
University of Cologne

trebst@thp.uni-koeln.de

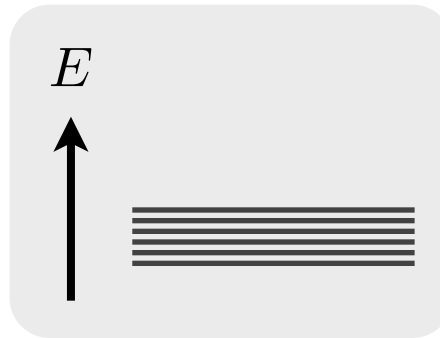


When do interesting things happen?

Some of the most intriguing phenomena in condensed matter physics arise from the splitting of **'accidental' degeneracies**.



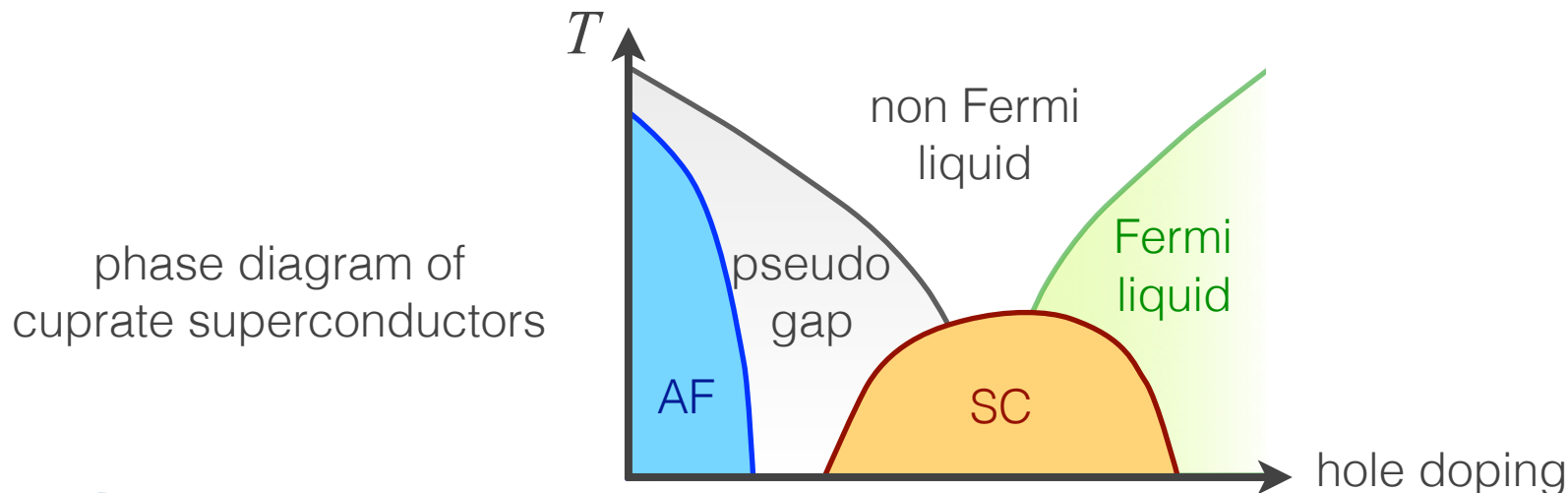
interacting
many-body system



'accidental'
degeneracy

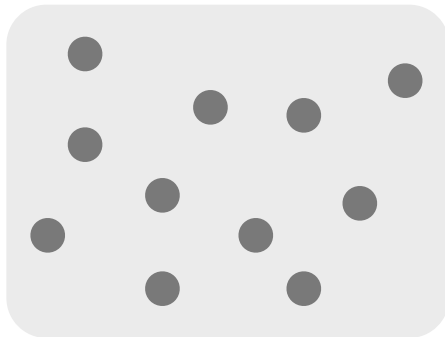


residual effects
select ground state

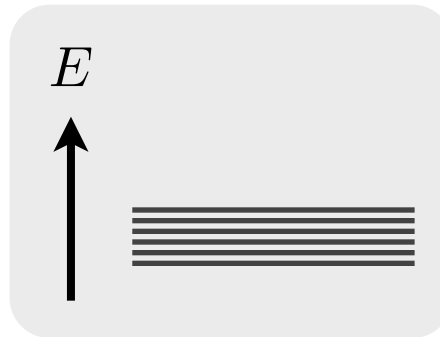


When do interesting things happen?

Some of the most intriguing phenomena in condensed matter physics arise from the splitting of **'accidental' degeneracies**.



interacting
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'accidental'
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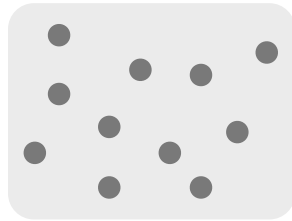


residual effects
select ground state

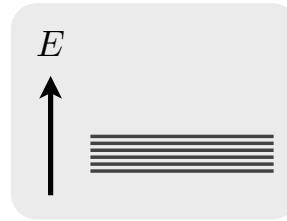
But they are also **notoriously difficult** to handle, due to

- multiple energy scales
- complex energy landscapes / slow equilibration
- strong coupling
- macroscopic entanglement

Example – frustrated magnets



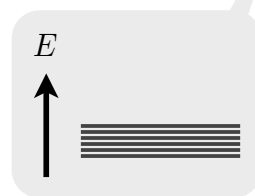
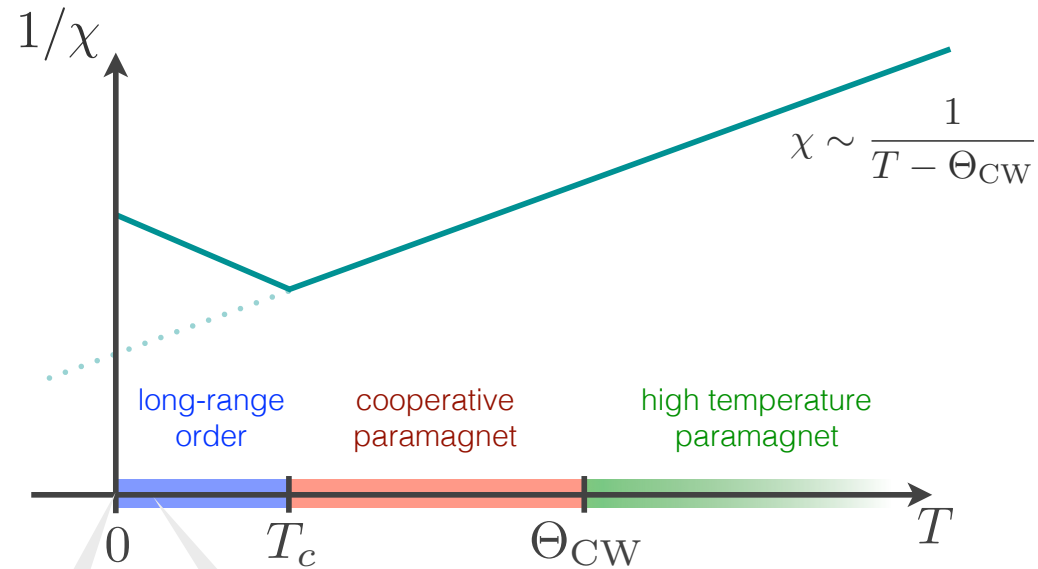
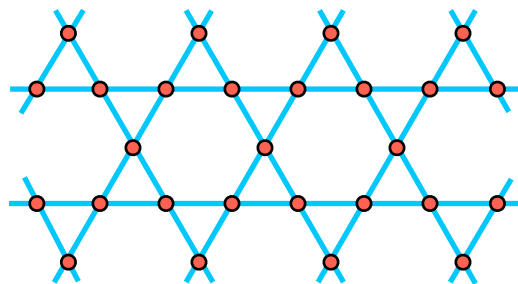
interacting many-body system



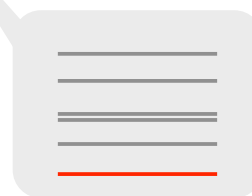
'accidental' degeneracy



residual effects
select ground state

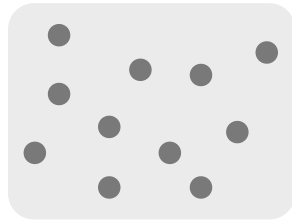


T=0 residual entropy

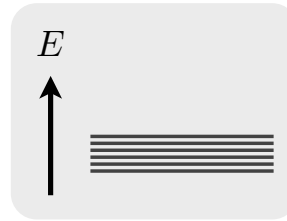


long-range order

Example – quantum Hall liquids



interacting
many-body system

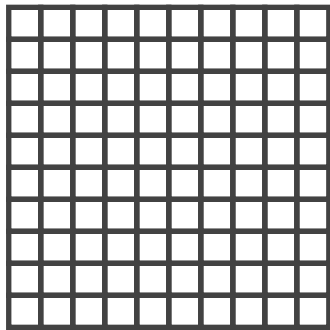


'accidental'
degeneracy



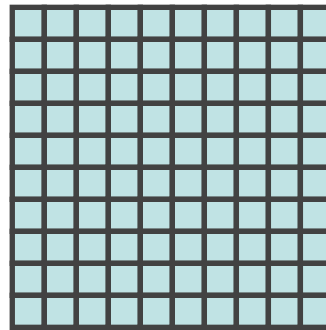
residual effects
select ground state

Landau level **degeneracy**



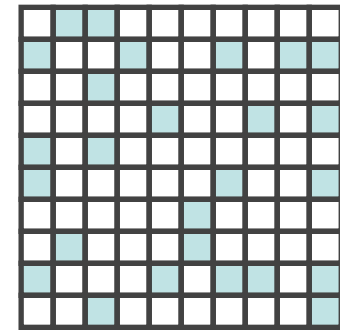
$2\Phi/\Phi_0$
orbital states

integer quantum Hall



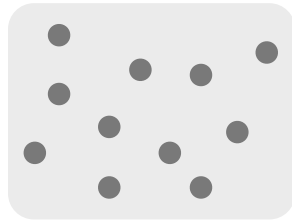
filled level
↓
incompressible liquid

fractional quantum Hall

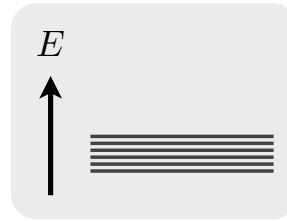


partially filled level
↓ Coulomb repulsion
incompressible liquid

Example – twisted bilayer graphene



interacting
many-body system

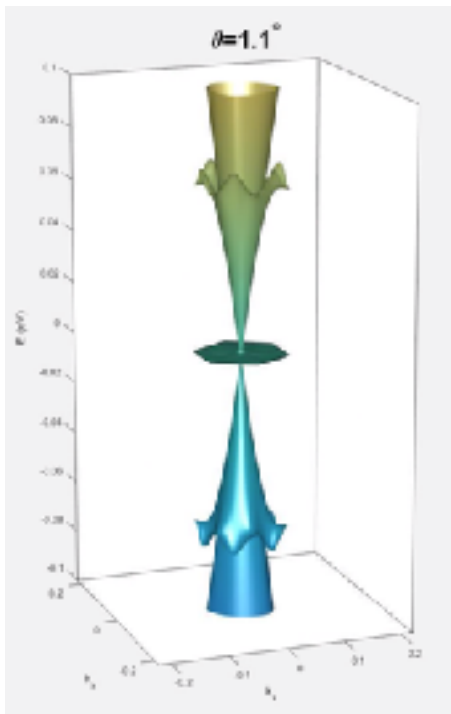


'accidental'
degeneracy

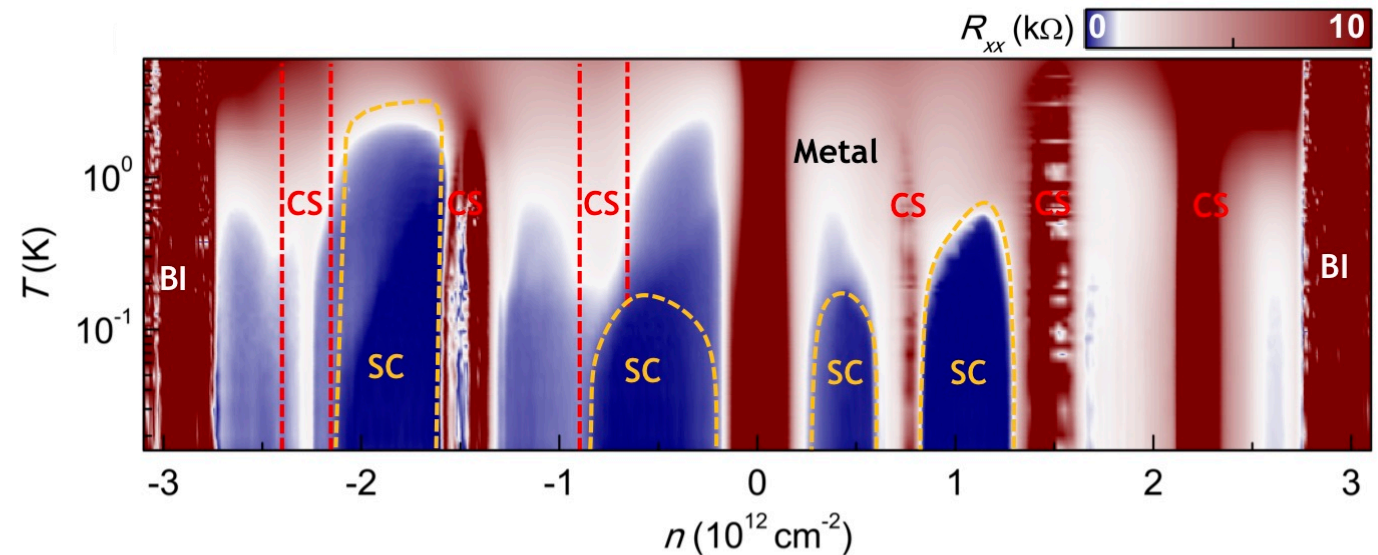


residual effects
select ground state

electronic band structure with a **flat band** in twisted bilayer graphene



Jarillo-Herrero group, Nature **556**, 43 (April 2018)

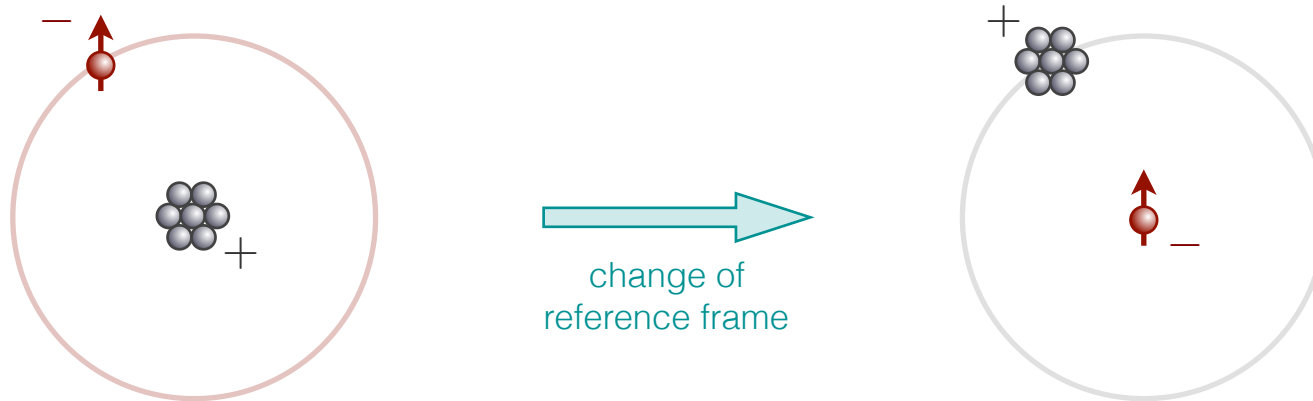


Efetov group, Nature **574**, 653(2019)

spin-orbit materials

Spin-orbit coupling

Spin-orbit coupling 101 – quantum mechanics lecture



relativistic correction

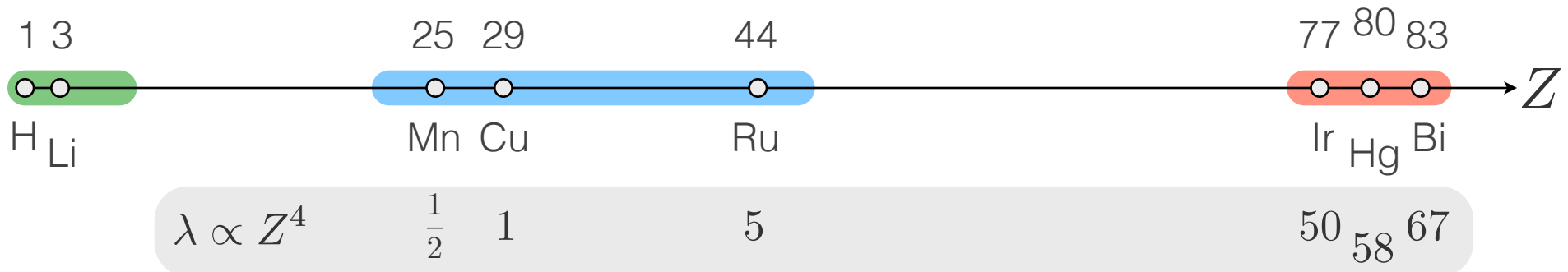
$$\Delta E = \frac{\lambda}{\hbar^2} \vec{l} \cdot \vec{s} = \frac{\lambda}{2} [j(j-1) - l(l-1) - s(s-1)]$$

$$\lambda = \frac{Ze^2\mu_0\hbar^2}{8\pi m_e^2 r^3}$$

$$r \propto 1/Z$$

$$\lambda \propto Z^4$$

Spin-orbit coupling in condensed matter



weak SOC

atomic
fine structure

moderate SOC

multiferroics unconventional
superconductor

TbMnO₃

Sr₂RuO₄

SOC induced DM interaction
competes with **magnetic exchange**

strong SOC

SO-assisted topological
Mott physics insulators

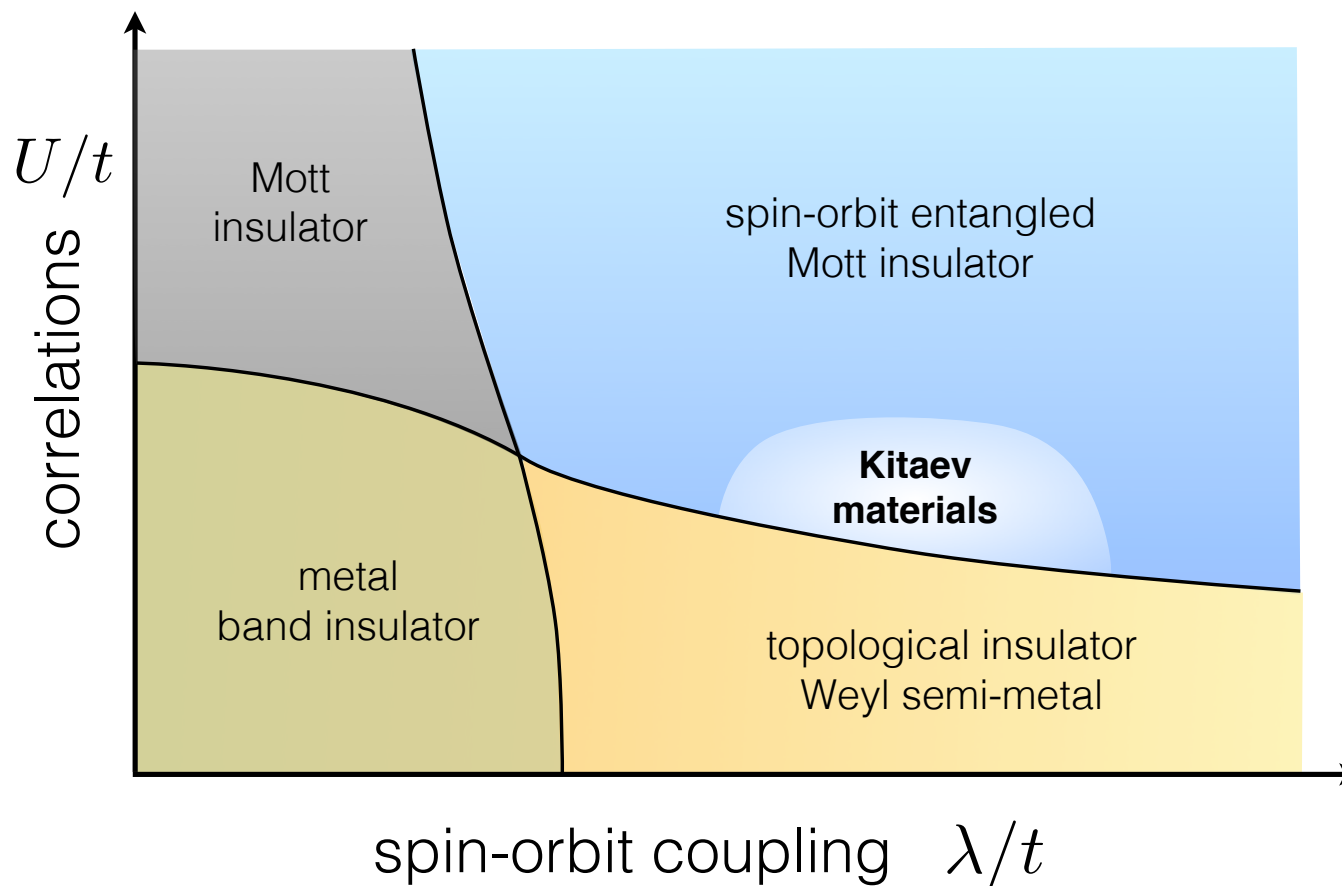
Sr₂IrO₄
Na₄Ir₃O₈
(Na,Li)₂IrO₃

HgTe
Bi₂Se₃

SOC competes directly
with **Hubbard physics**

4d/5d transition metal compounds

Transition metal oxides with **partially filled 4d/5d shells** exhibit an intricate interplay of **spin-orbit coupling**, **electronic correlations**, and **crystal field effects** resulting in a **broad variety of metallic and insulating states**.

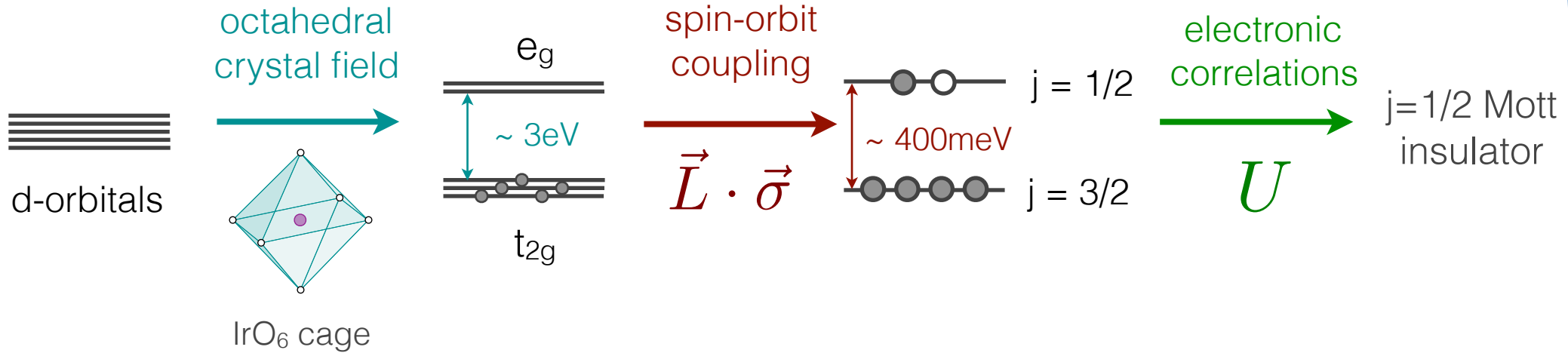


W. Witczak-Krempa, G. Chen, Y. B. Kim, and L. Balents,
Annual Review of Condensed Matter Physics 5, 57 (2014).

$j=1/2$ Mott insulators

most common
Iridium valence

Ir^{4+} ($5d^5$)

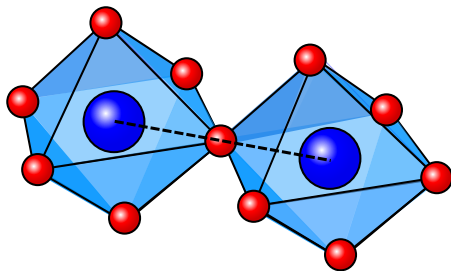


Why are these spin-orbit entangled $j=1/2$ Mott insulators **interesting**?

spin-orbit entangled Mott insulators

Why are these spin-orbit entangled $j=1/2$ Mott insulators **interesting?**

corner-sharing

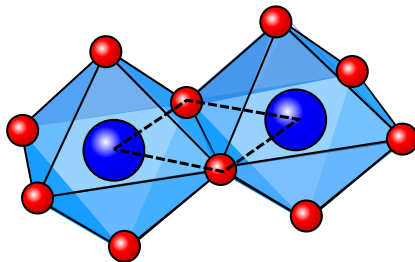


exhibits cuprate-like magnetism
superconductivity?

B.J. Kim et al. PRL 101, 076402 (2008)

B.J. Kim et al. Science 323, 1329 (2009)

edge-sharing



...

exhibit Kitaev-like magnetism
spin liquids?

G. Jackeli and G. Khaliullin, PRL 102, 017205 (2009)

J. Chaloupka, G. Jackeli, and G. Khaliullin, PRL 105, 027204 (2010)

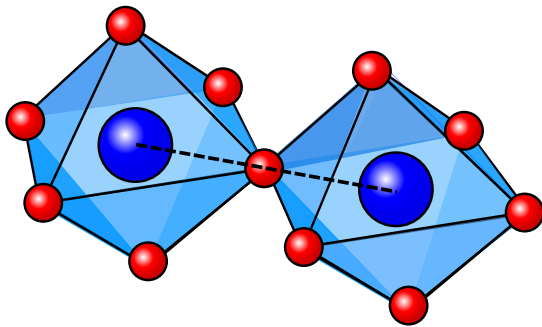


bond-directional exchange



G. Jackeli and G. Khaliullin, PRL 102, 017205 (2009)
J. Chaloupka, G. Jackeli, and G. Khaliullin, PRL 105, 027204 (2010)

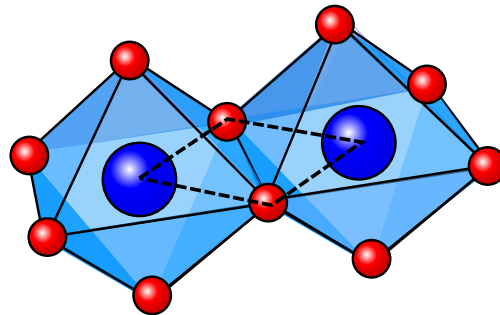
corner-sharing



Sr_2IrO_4

Heisenberg exchange

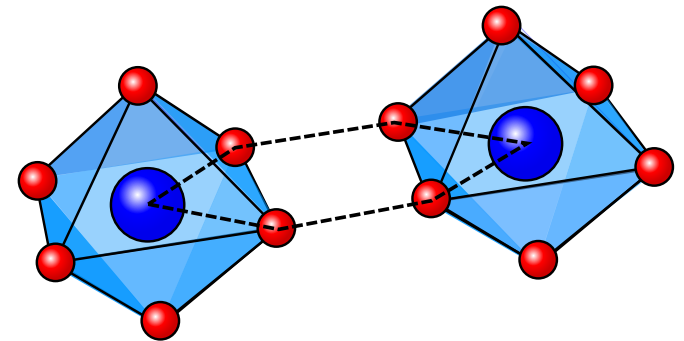
edge-sharing



$(\text{Na, Li})_2\text{IrO}_3$
 RuCl_3

Heisenberg-Kitaev exchange

“parallel edge”-sharing



$\text{Ba}_3\text{IrTi}_2\text{O}_9$

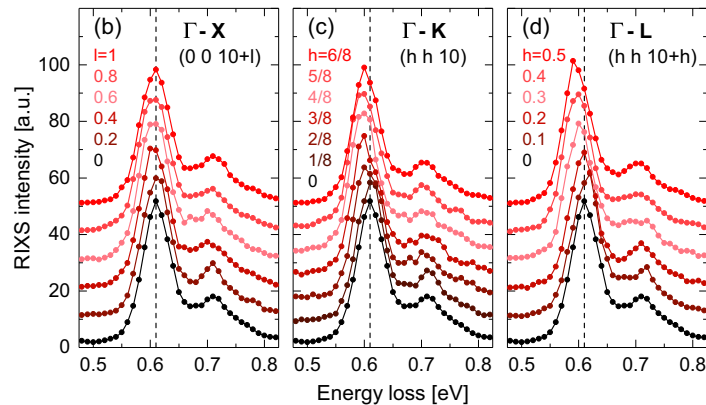
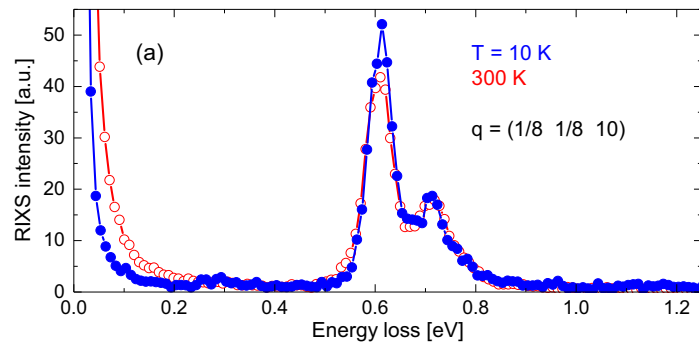
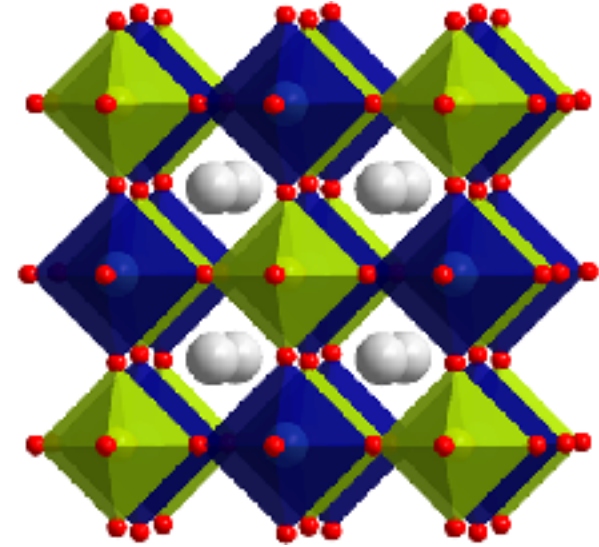
$$H = - \sum_{\gamma\text{-bonds}} J \mathbf{S}_i \mathbf{S}_j + K S_i^\gamma S_j^\gamma + \Gamma \left(S_i^\alpha S_j^\beta + S_i^\beta S_j^\alpha \right)$$

spin-orbit entangled moments

Ba₂CeIrO₆

Revelli *et al.*, PRB **100**, 085139 (2019)

The double perovskite Ba₂CeIrO₆ is the **best $j=1/2$ system** we have ever seen, but not really a “Kitaev material”.



RIXS experiments

- pristine $j=1/2$ physics

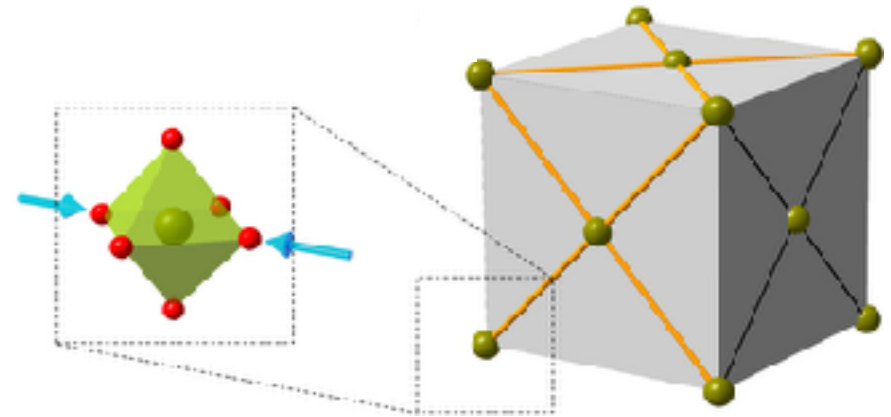
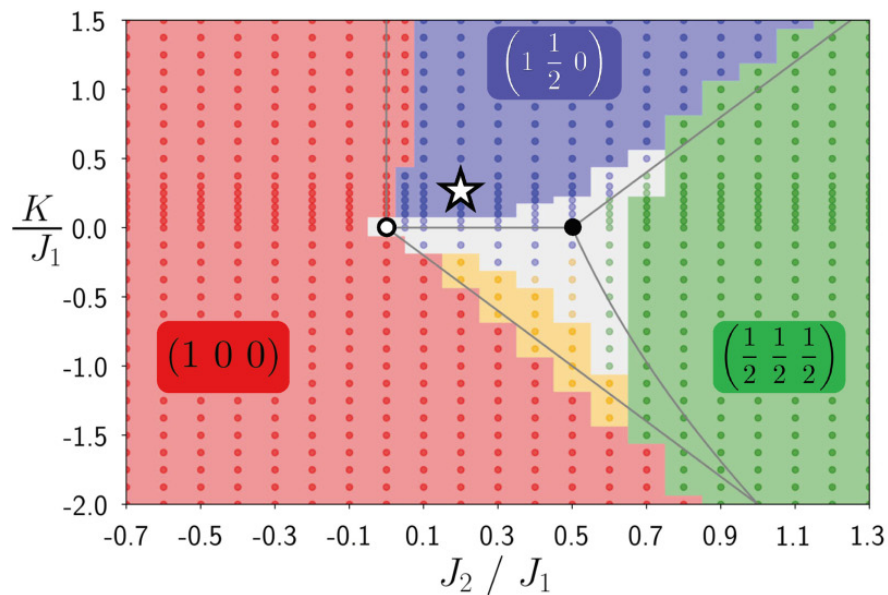
$$|0\rangle = 0.991 \left| \frac{1}{2}, \frac{1}{2} \right\rangle - 0.130 \left| \frac{3}{2}, \frac{1}{2} \right\rangle$$

- frustrated FCC magnetism, but Kitaev interaction relieves frustration
- strong magneto-elastic effect

Ba₂CeIrO₆

Revelli *et al.*, PRB **100**, 085139 (2019)

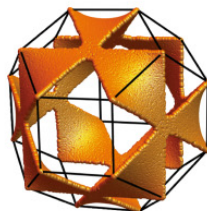
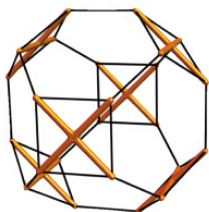
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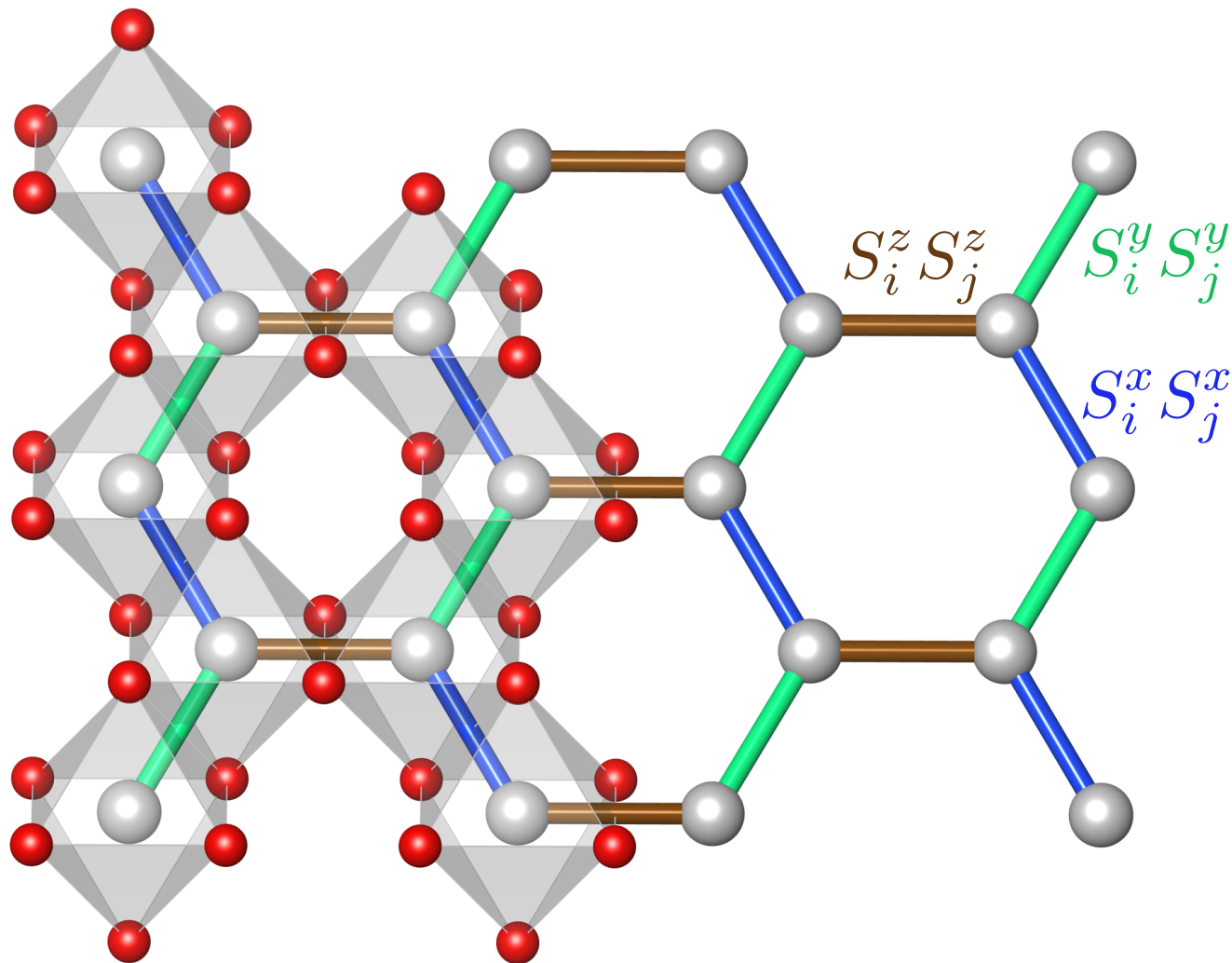


honeycomb Kitaev materials

proximate spin liquids

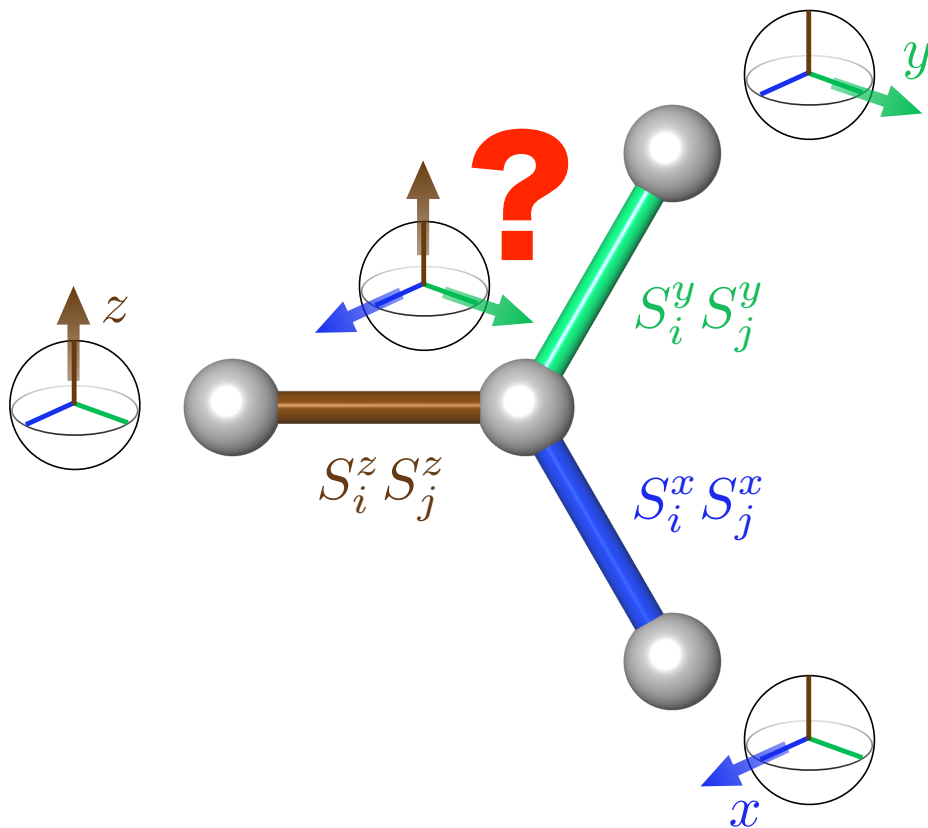
honeycomb Kitaev materials

Na_2IrO_3 , $\alpha\text{-Li}_2\text{IrO}_3$, RuCl_3 , $\text{H}_2\text{LiIr}_2\text{O}_6$

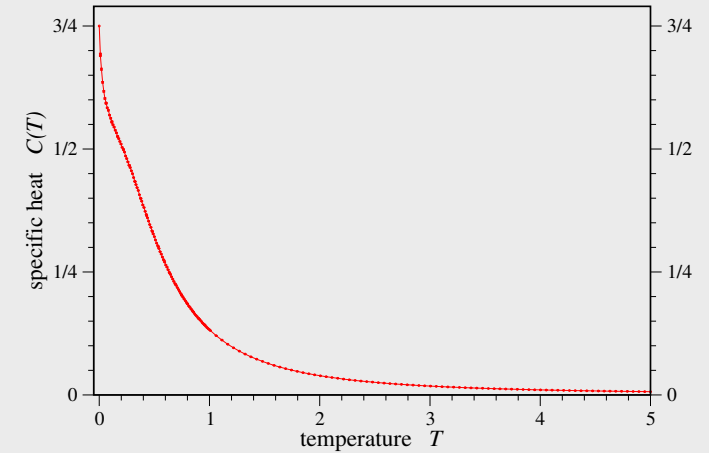


exchange frustration

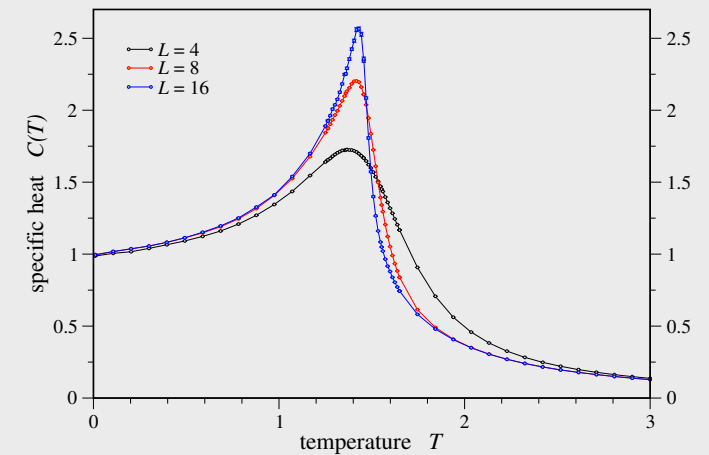
$$H = - \sum_{\gamma\text{-bonds}} J \mathbf{S}_i \mathbf{S}_j + K S_i^\gamma S_j^\gamma$$



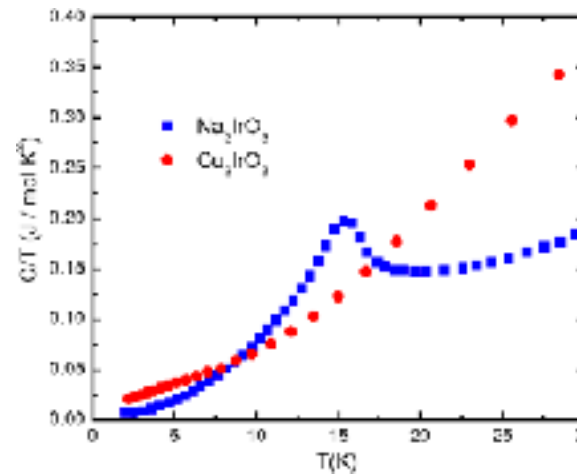
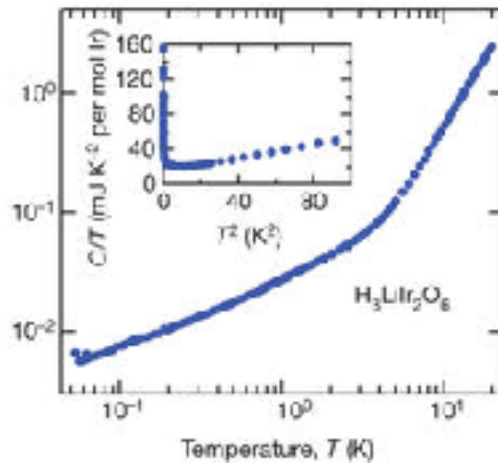
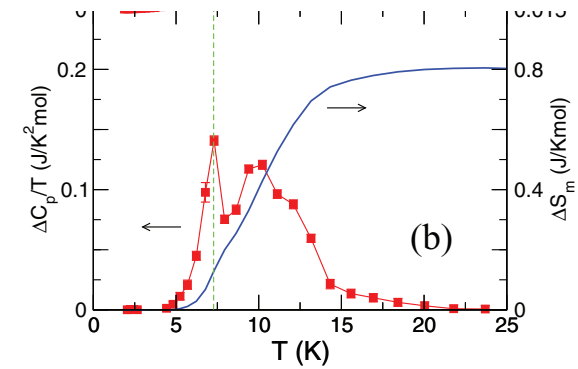
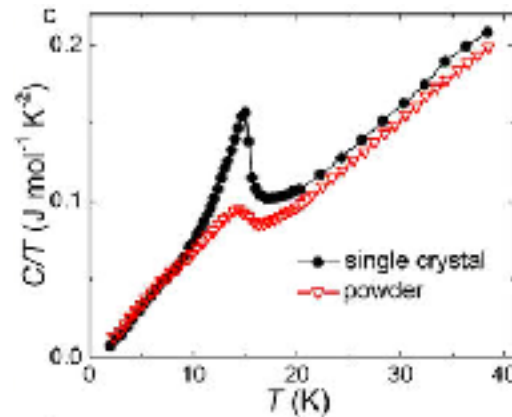
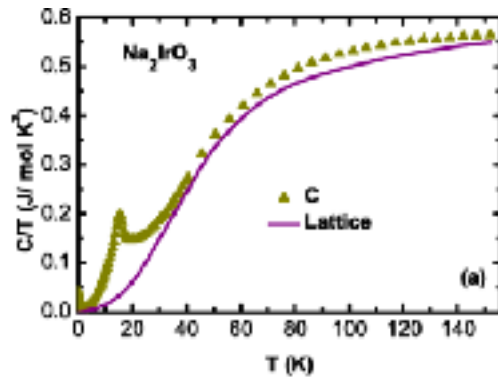
classical Kitaev



classical Heisenberg



Kitaev materials – really?



Candidate materials tend to exhibit **magnetic ordering** at low T .

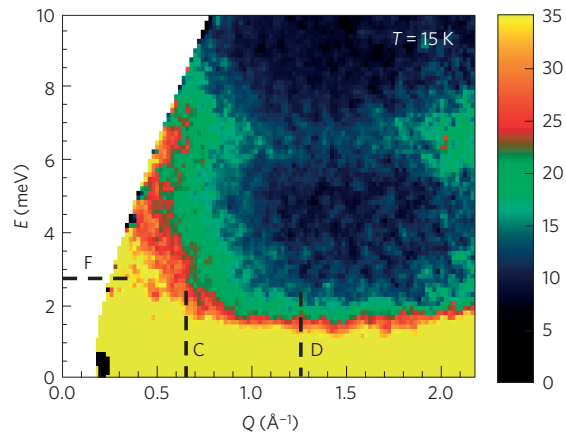
honeycomb Kitaev materials

	magnetic moment $\mu_{\text{eff}} / \mu_{\text{B}}$	ordering temperature T_{N}	Curie-Weiss temperature Θ_{CW}
Na_2IrO_3	1.79(2)	15 K zig-zag order	-125 K
$\alpha\text{-Li}_2\text{IrO}_3$	1.83(5)	15 K counterrotating spirals	-33 K
RuCl_3	2.2	7 K zig-zag order	-150 K
$\text{H}_2\text{LiIr}_2\text{O}_6$?	—	?

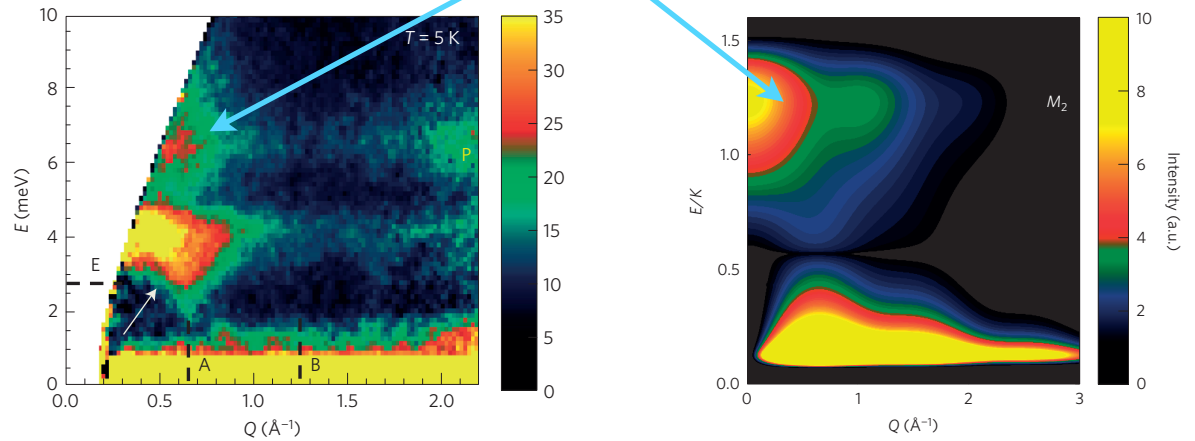
RuCl₃

neutron scattering

Banerjee *et al.*, Nature Materials 4604 (2016)



broad scattering continuum

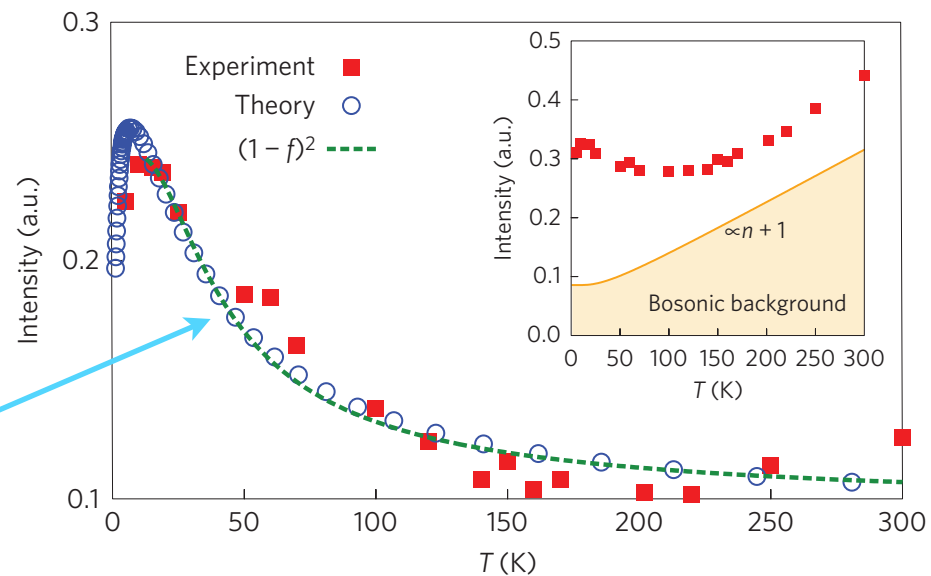


Raman scattering

Nasu *et al.*, Nature Physics 12, 912 (2016)

Sandilands *et al.*, PRL 114, 147201 (2015)

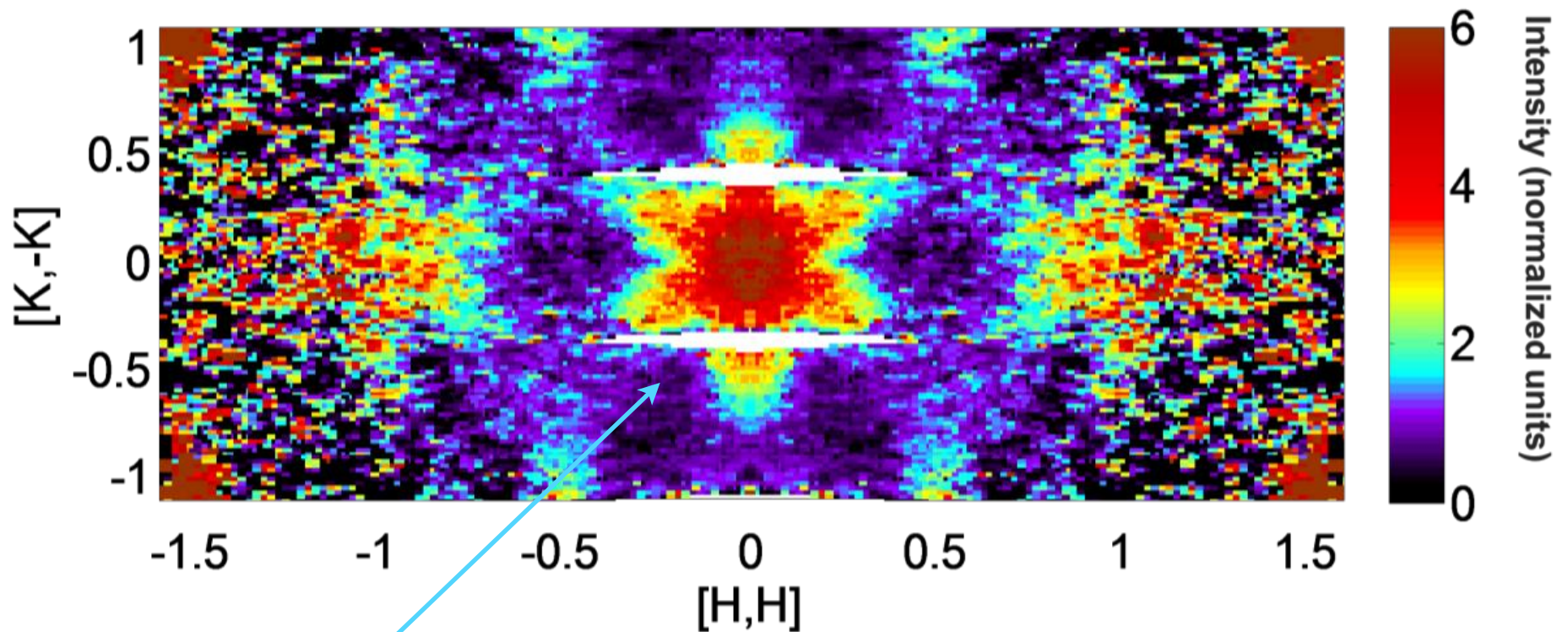
fermionic contribution



RuCl₃

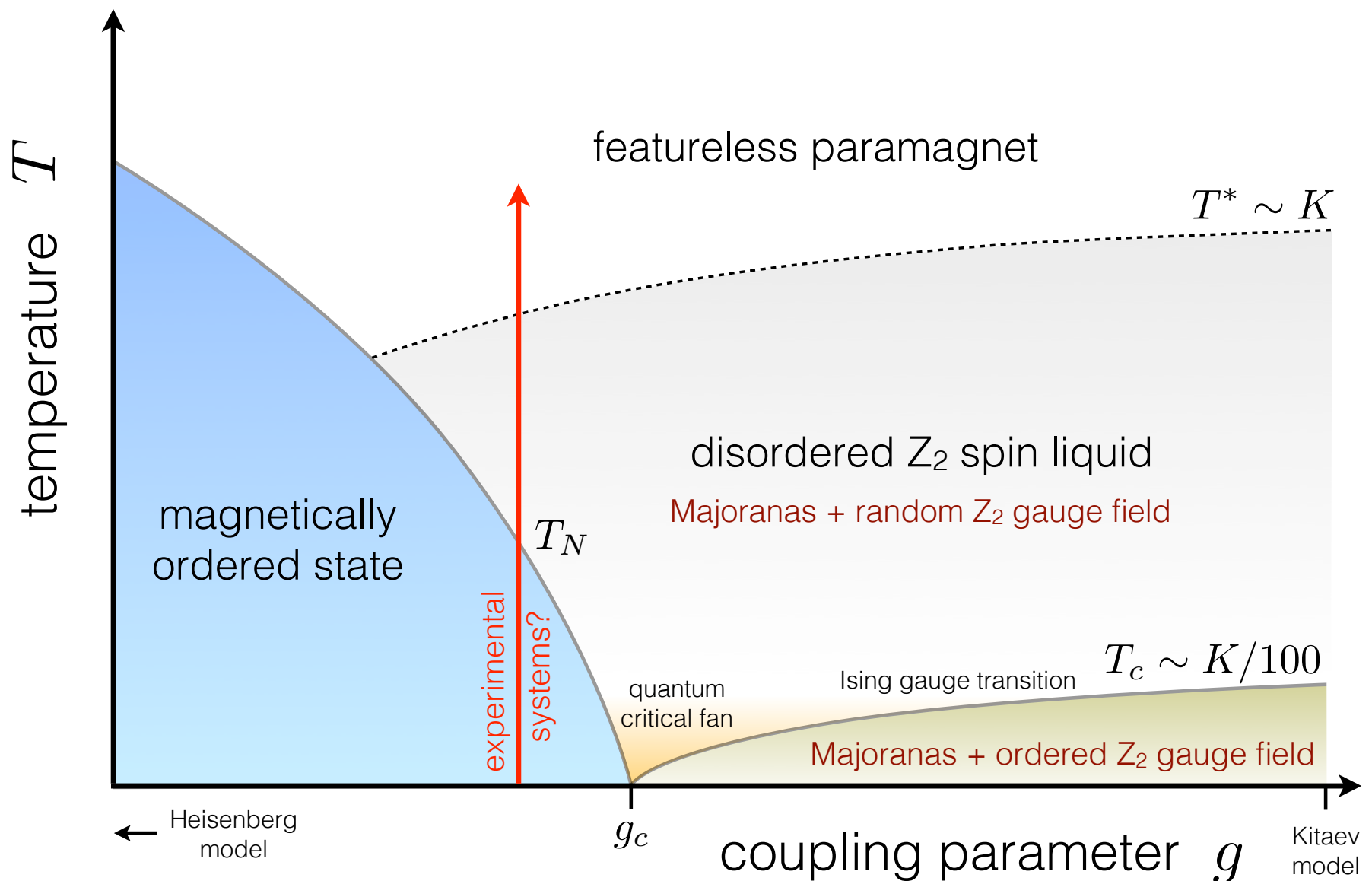
neutron scattering

Banerjee *et al.*, Nature Materials 4604 (2016)



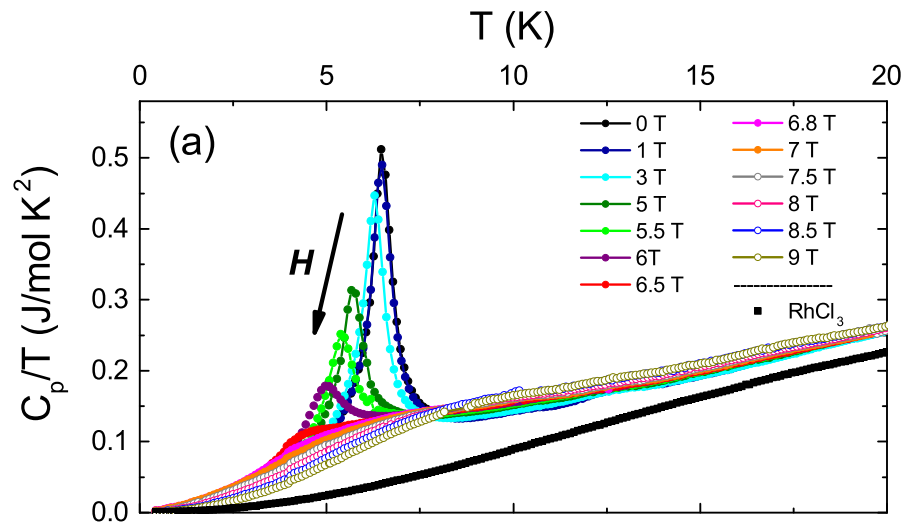
star-like feature arises from interplay of spin-wave and spin liquids physics at intermediate energy scales.

Proximate spin liquids

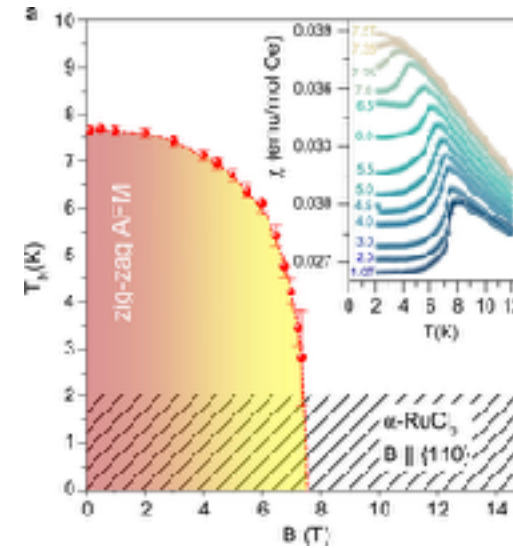


Spin liquids?!

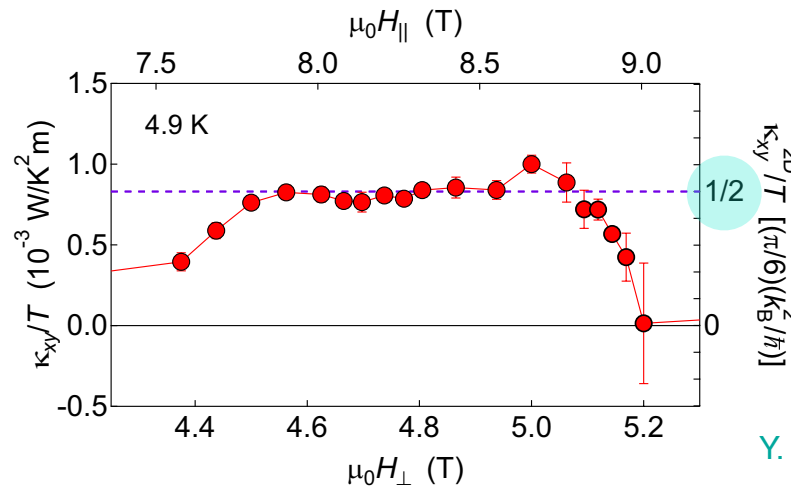
Something interesting happens for **RuCl₃** in a magnetic field.



A. U. B. Wolter et al., PRB **96**, 041405(R) (2017)



A. Banerjee et al., npj Quantum Mater. **3**, 8 (2018)



κ_{xy}^{2D}/T [$(\pi/6)(k_B^2/h)$]

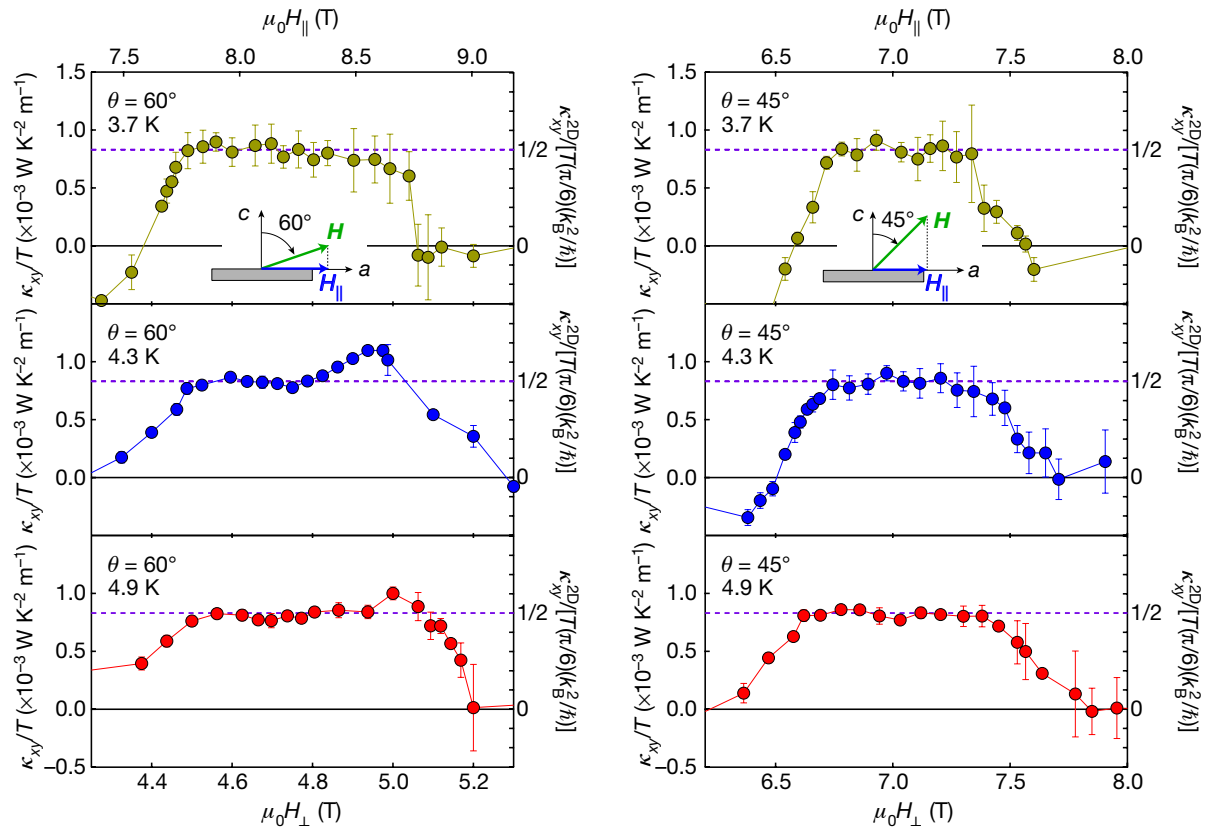
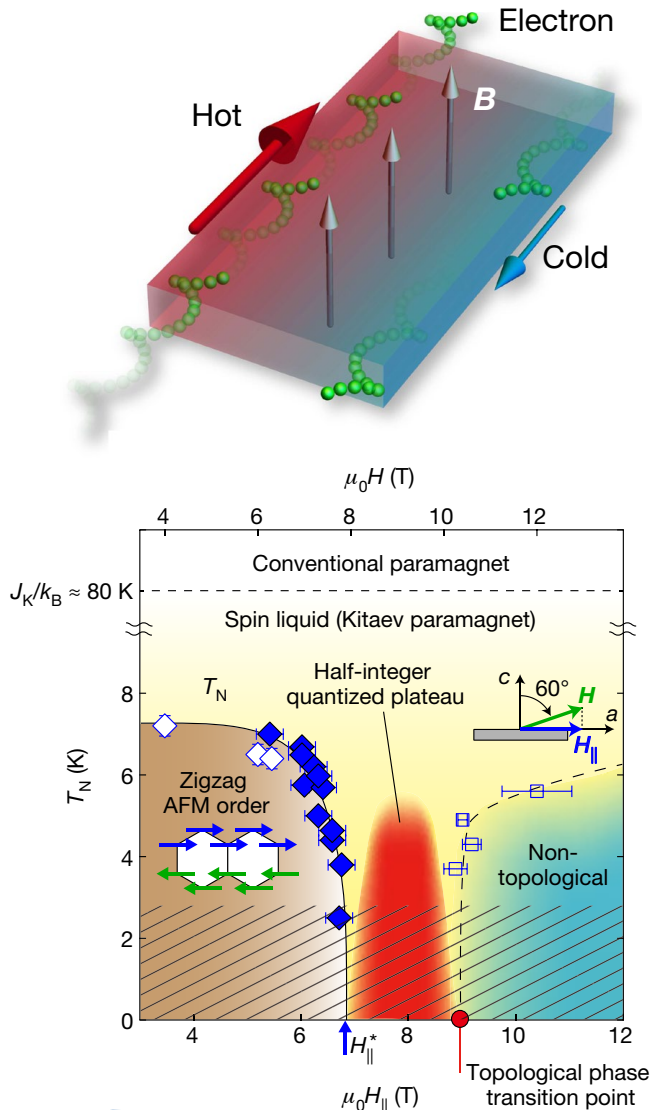
← **half-integer quantized thermal Hall effect**

Y. Kasahara et al., Nature **559**, 227-231 (2018)

a new quantum Hall effect

S. Trebst & A. Rosch,
Physik-Journal (2018)

Something interesting happens for **RuCl₃** in a magnetic field.



Y. Kasahara et al., Nature **559**, 227-231 (2018)

Why does this work in the presence of phonons?

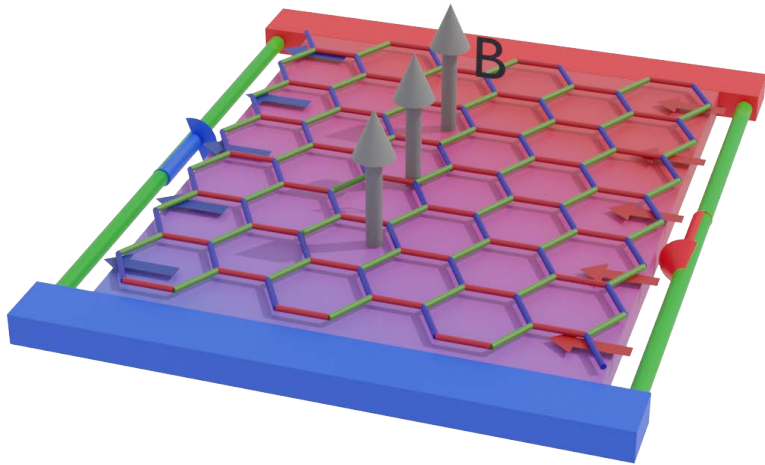
Y. Vinkler-Aviv & A. Rosch, PRX **8**, 031032 (2018)

M. Ye, G. Halász, L. Savary, and L. Balents, PRL **121**, 147201 (2018)

thermal Hall effect in RuCl_3

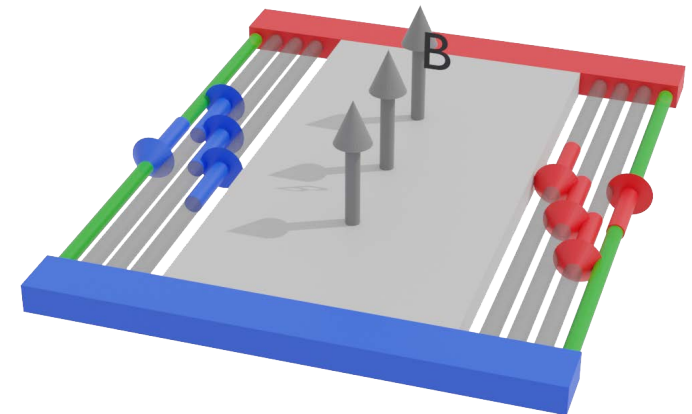
$$\kappa_{xy} = \frac{n}{2} \cdot \frac{\pi}{3} \frac{k_B^2 T}{\hbar}$$

A **half-quantized thermal Hall response** is direct evidence for gapless **Majorana modes**.



Kitaev material RuCl_3

Y. Kasahara et al., *Nature* **559**, 227-231 (2018)



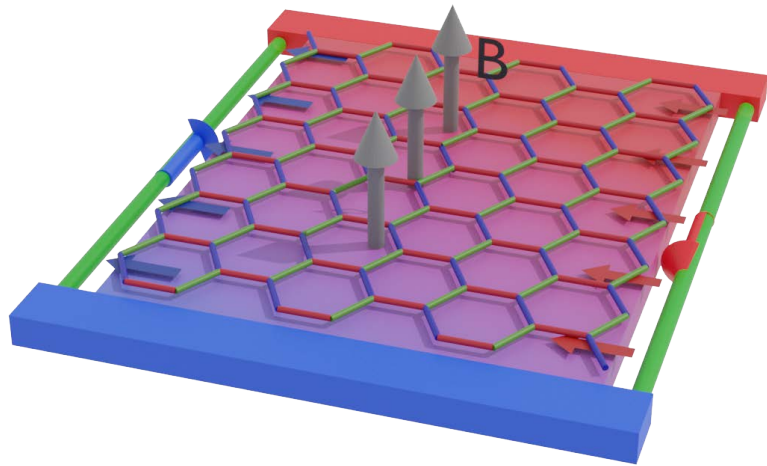
$\nu=5/2$ FQH state in a 2DEG

M. Banerjee et al., *Nature* **559**, 205-210 (2018)

thermal Hall effect in RuCl_3

A **half-quantized thermal Hall response** is direct evidence for gapless **Majorana modes**.

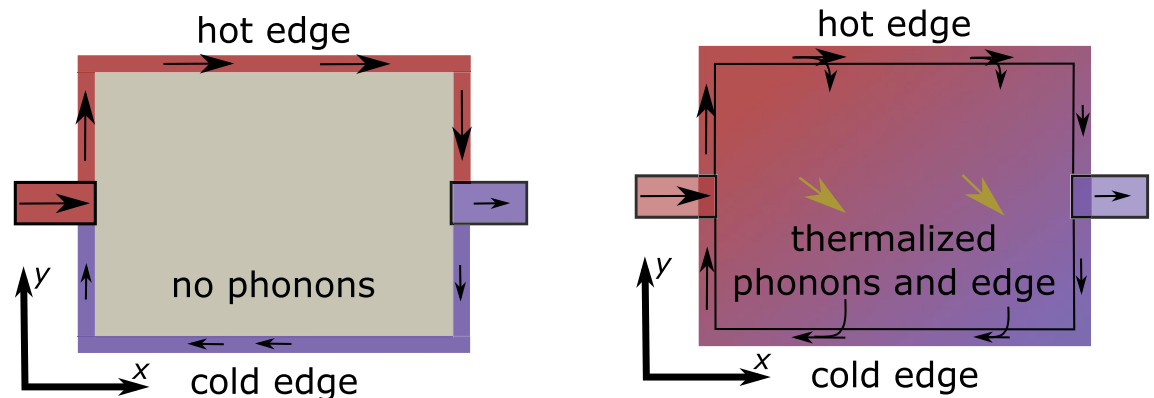
Y. Vinkler-Aviv & A. Rosch, PRX **8**, 031032 (2018)
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Kitaev material RuCl_3

Y. Kasahara et al., Nature **559**, 227-231 (2018)

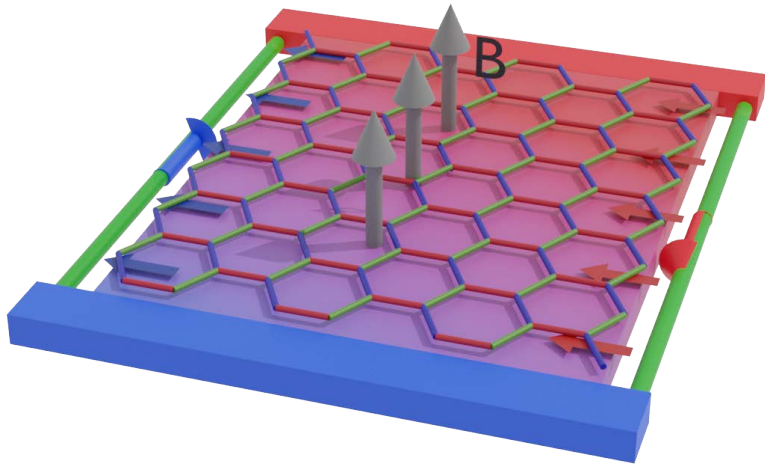
But why is it quantized in the first place?
Why is there no leakage into the bulk,
via gapless acoustic phonons?



thermal Hall effect in RuCl_3

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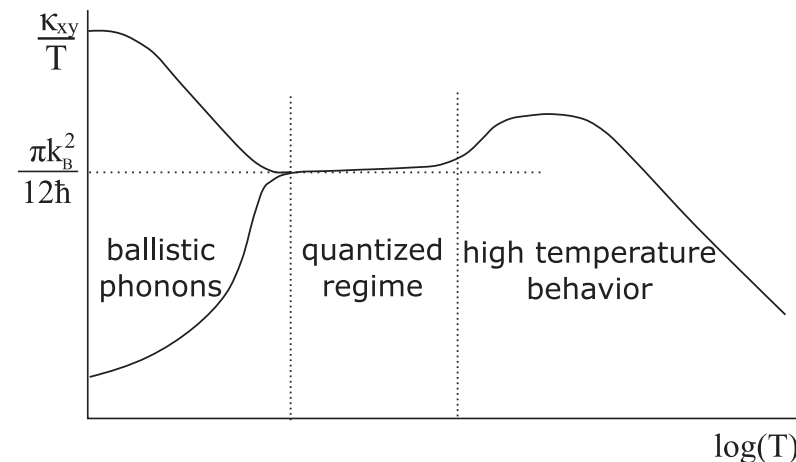
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Y. Kasahara et al., Nature **559**, 227-231 (2018)

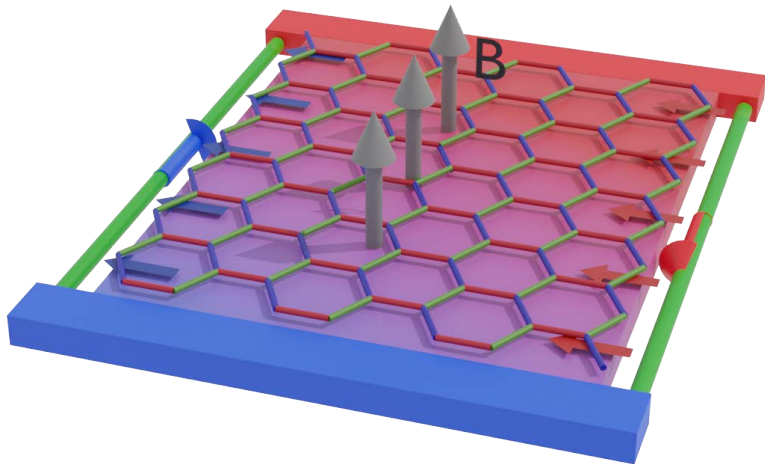
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thermal Hall effect in RuCl_3

A **half-quantized thermal Hall response** is direct evidence for gapless **Majorana modes**.

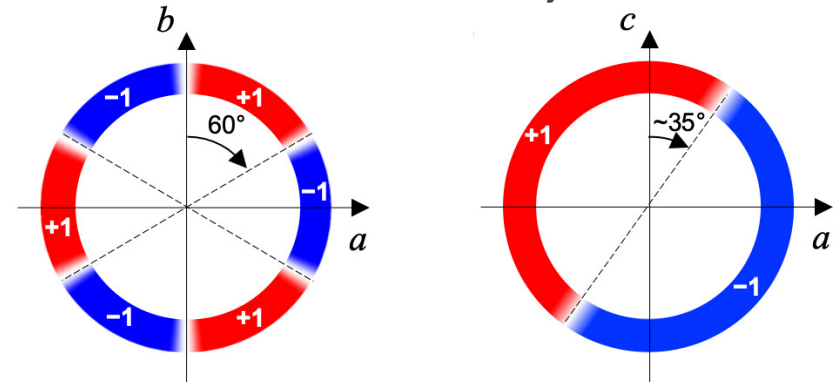
How can we distinguish whether the quantized thermal Hall effect arises from the formation of **Landau levels** or a non-trivial **Chern insulator**?



Kitaev material RuCl_3

Y. Kasahara et al., *Nature* **559**, 227-231 (2018)

The Kitaev spin liquid is a **chiral spin liquid**, a Chern insulator of Majoranas.



Its Hall quantization is **angle-dependent** and occurs even for an in-plane field (**anomalous thermal Hall effect**).

T. Yokoi et al., arXiv:2001.01899

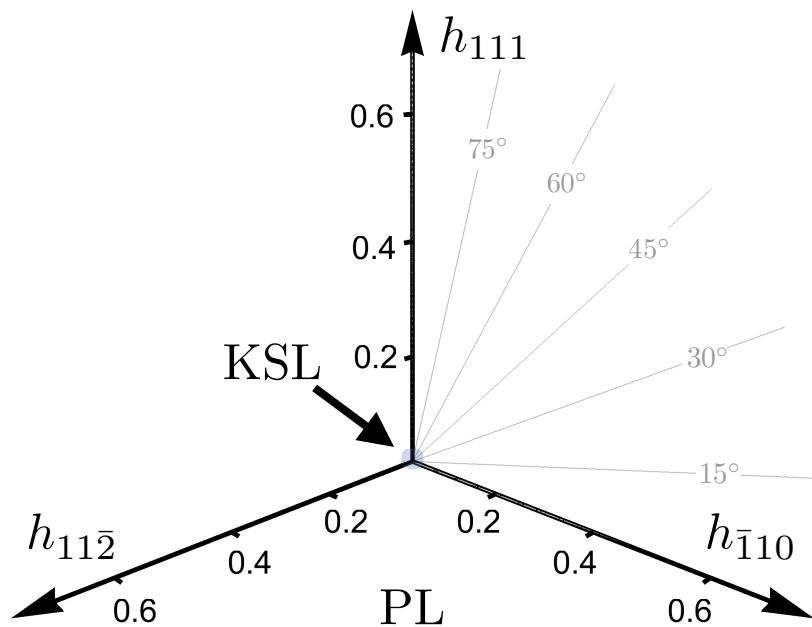
back to the Kitaev model

microscopic relevance

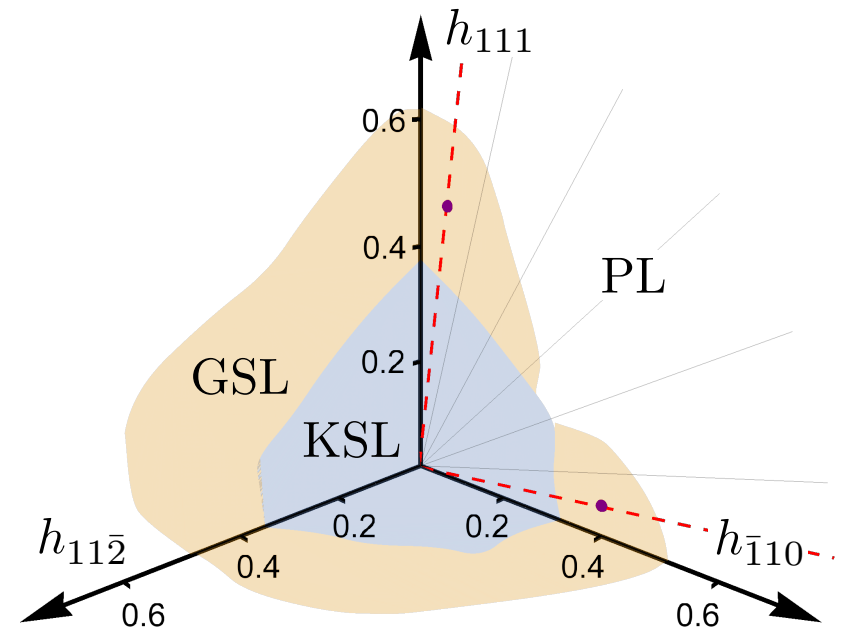
Kitaev model – magnetic field effects

$$\mathcal{H} = - \sum_{\gamma\text{-bonds}} K_{\gamma} S_i^{\gamma} S_j^{\gamma} - \sum_i \mathbf{h} \cdot \mathbf{S}_i$$

FM Kitaev coupling



AFM Kitaev coupling

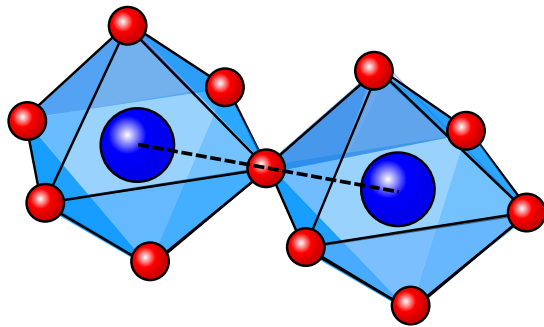


bond-directional exchange



G. Jackeli and G. Khaliullin, PRL 102, 017205 (2009)
 J. Chaloupka, G. Jackeli, and G. Khaliullin, PRL 105, 027204 (2010)

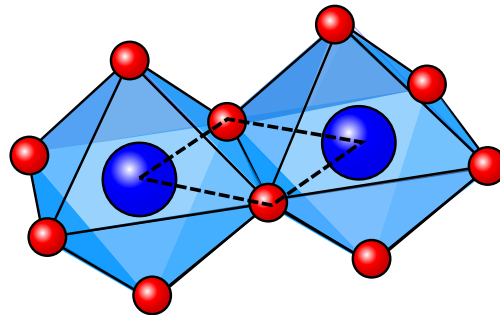
corner-sharing



Sr_2IrO_4

Heisenberg exchange

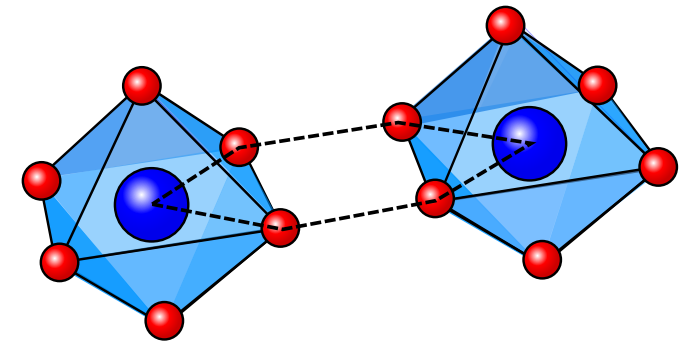
edge-sharing



$(\text{Na,Li})_2\text{IrO}_3$
 RuCl_3

Heisenberg-Kitaev exchange

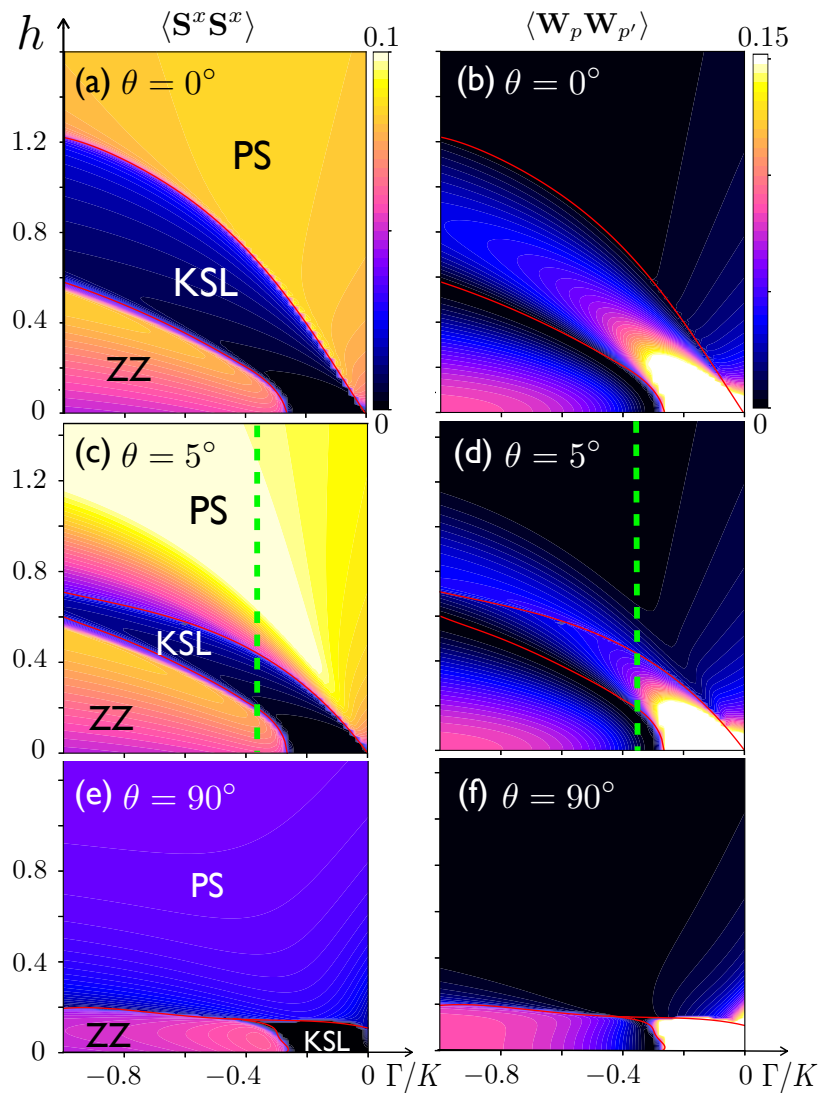
“parallel edge”-sharing



$\text{Ba}_3\text{IrTi}_2\text{O}_9$

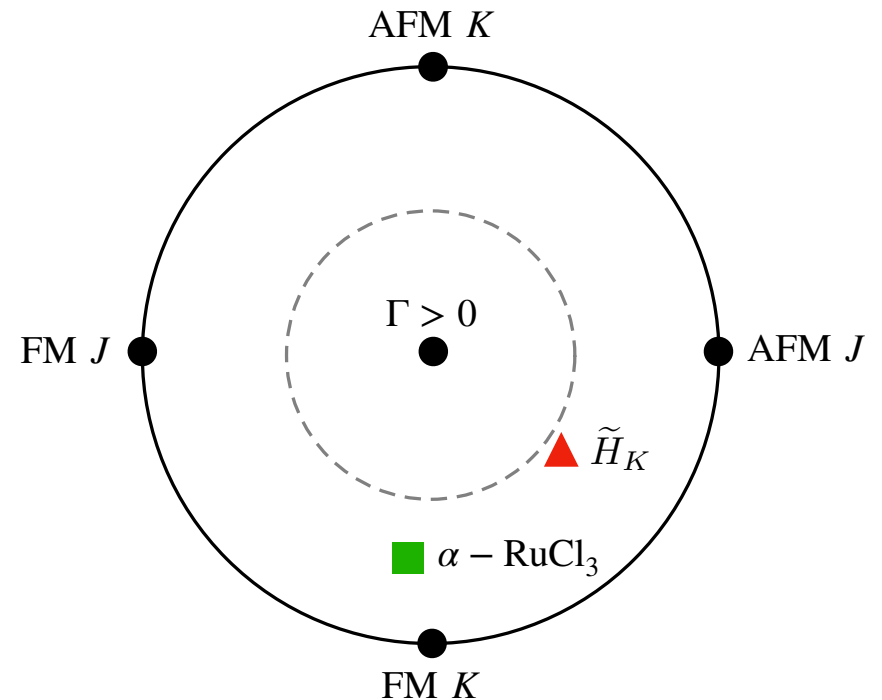
$$H = - \sum_{\gamma\text{-bonds}} J \mathbf{S}_i \mathbf{S}_j + K S_i^\gamma S_j^\gamma + \Gamma \left(S_i^\alpha S_j^\beta + S_i^\beta S_j^\alpha \right)$$

Microscopic relevance to RuCl₃



AFM off-diagonal coupling

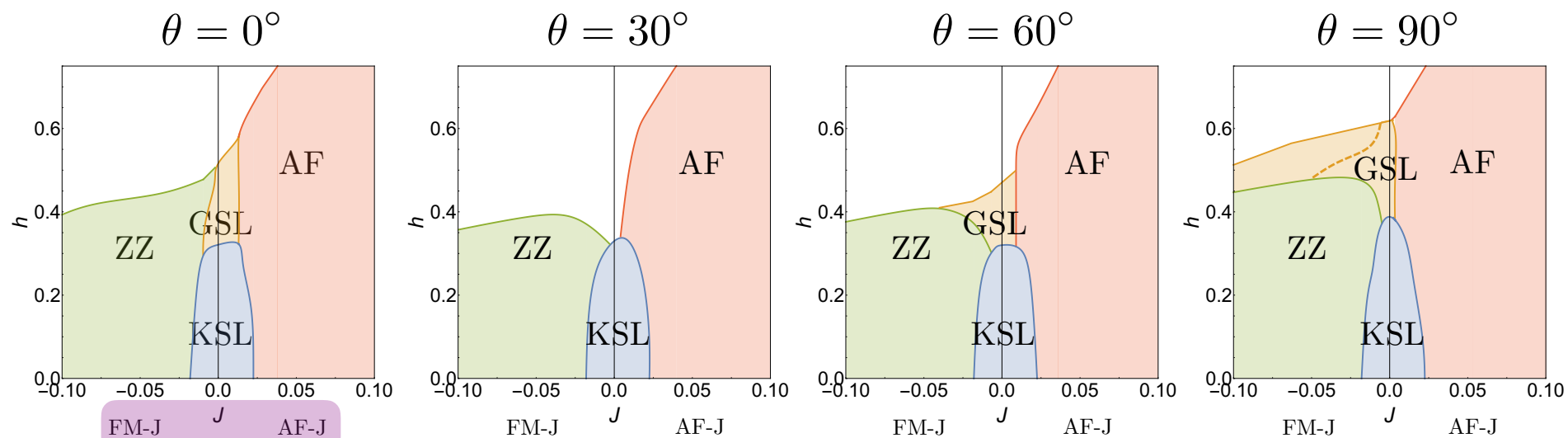
Gordon et al., Nature Comm. 10, 2470 (2019)



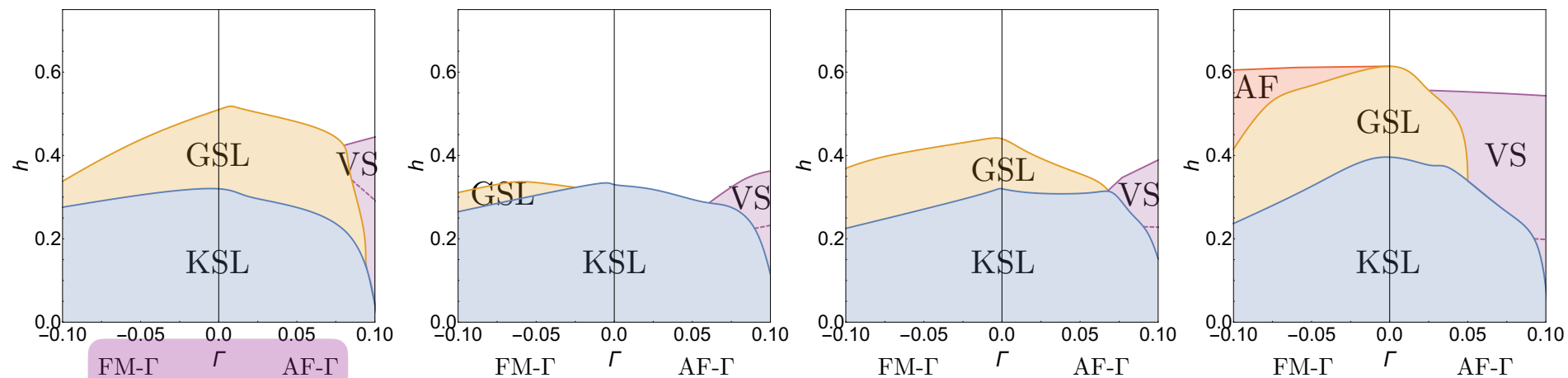
“hidden” AFM Kitaev model
(via Klein duality)

Kaib, Winter & Valenti, PRB 100, 144445 (2019)

Higgsed U(1) spin liquid?



Heisenberg



Gamma

Kitaev materials

- a family of spin-orbit assisted $j=1/2$ Mott insulators
- bond-directional exchange induces frustration
- unconventional forms of magnetism

Bond-directional exchange

- (proximate) spin liquids
- signatures of Majorana fermions and Z_2 gauge field
- spin textures

Family of lattice geometries

- honeycomb – Na_2IrO_3 , $\alpha\text{-Li}_2\text{IrO}_3$, $(\text{H}_{3/4}\text{Li}_{1/4})_2\text{IrO}_3$, RuCl_3
- triangular – $\text{Ba}_3\text{IrTi}_2\text{O}_9$, $\text{Ba}_3\text{Ir}_2\text{TiO}_9$, $\text{Ba}_3\text{Ir}_2\text{InO}_9$
- 3D – $\beta\text{-Li}_2\text{IrO}_3$, $\gamma\text{-Li}_2\text{IrO}_3$, metal-organic compounds

Thanks!