



Producing smooth flow in racetrack atom circuits by stirring at zero and non-zero temperatures

Mark Edwards

Atomtronics 2019

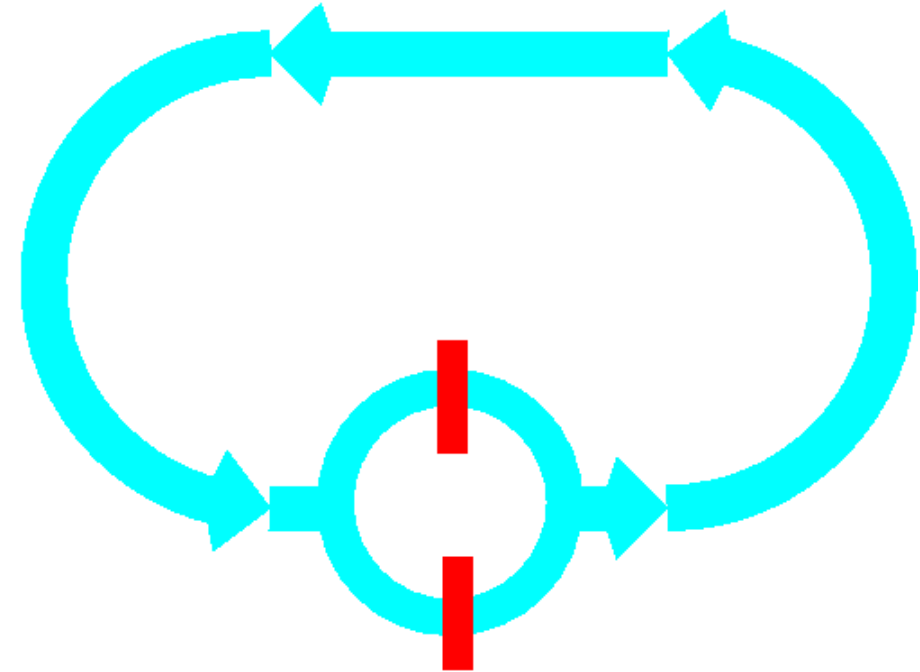
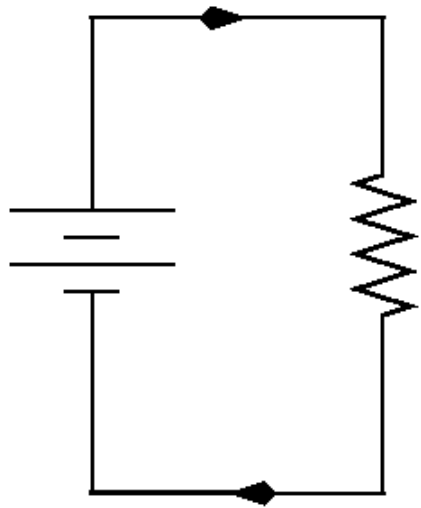
Centro de Ciencias de Benasque Pedro Pascual

14 May 2019



\$\$ = NSF, NIST

Electronic circuits need electron current flow for applications

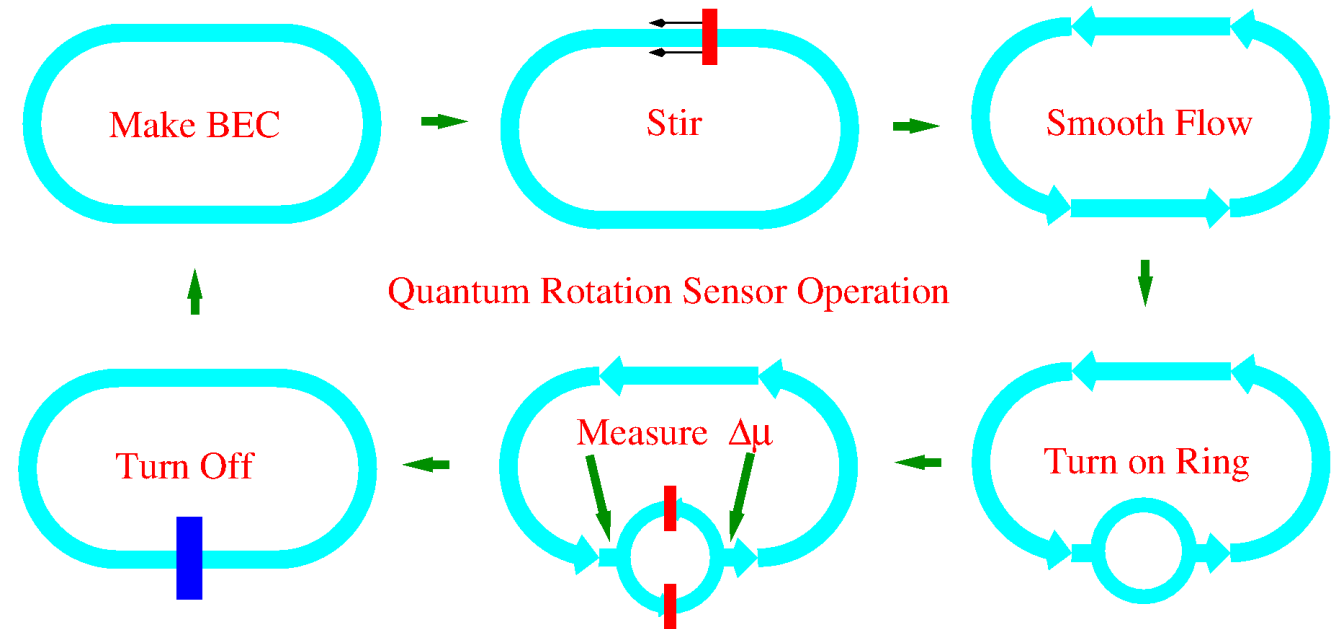


smooth neutral-atom current flow is necessary for atomtronic system applications

Can on-demand smooth flow be produced by stirring?

Racetrack Rotation Sensor Idea

1. Create BEC in racetrack potential
2. Stir with barrier to create smooth flow.
3. Morph on ring channel.
4. Morph on barriers in the ring.
5. Measure the chemical potential difference produced by system rotation.
6. Turn off flow to reset for another measurement.

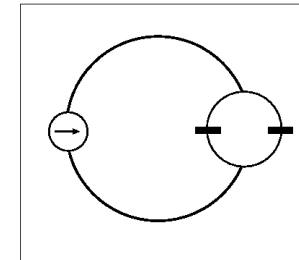


Need the ability to make on-demand smooth flow for this idea to work...we studied smooth-flow production by stirring a BEC in a racetrack potential

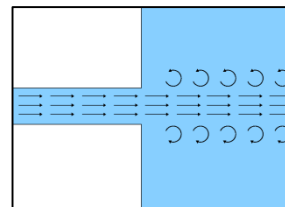
Can we make an atom circuit with smooth flow?

$$\partial \mathbf{v}(\mathbf{r}, t) / \partial t \approx 0 \quad \mathbf{v}(\mathbf{r}, t) \neq 0$$

- Ways to create flow:
 - reshaping and tilting the path of the atom circuit,
 - OAM transfer (Ryu, et al., PRL **99**, 260401 (2007))
 - phase imprint (Kumar, et al., PRA **97**, 043615 (2018))
 - **stirring** (Wright, et al., PRL **110**, 025302 (2013))
- Applications of smooth flow:
 - precision sensors for navigation and metrology



- fundamental physics



- Can then design elements for applications.

Zero-Temperature (GPE) and Non-Zero Temperature (ZNG) Model

ZNG (Zaremba-Nikuni-Griffin) Model

- The ZNG theory describes the behavior of a **thermal equilibrium system that is subject to a weak, external perturbation**.
- Assumes that the ultracold atom **system consists of a Bose-Einstein condensate plus a non-condensate part that is modeled as a classical gas**.
- Accounts for time-dependent condensate and non-condensate behavior and allows for both time-dependent atom clouds to affect each other.
- Quantities that describe the system are the **condensate wave function**, called $\Phi(\mathbf{r}, t)$ and the **single-particle distribution function**, called $f(\mathbf{p}, \mathbf{r}, t)$.

The single-particle distribution function can be understood as:

$$dN(\mathbf{p}, \mathbf{r}, t) \equiv f(\mathbf{p}, \mathbf{r}, t) \frac{d^3p d^3r}{(2\pi\hbar)^3} =$$

the number of non-condensate atoms having momentum in a volume of d^3p around \mathbf{p} and located within a volume d^3r of \mathbf{r} .

Equation of motion for the condensate

Generalized Gross-Pitaevskii Equation

$$i\hbar \frac{\partial}{\partial t} \Phi(\mathbf{r}, t) = \left(\frac{-\hbar^2}{2m} \nabla^2 + V_{\text{trap}}(\mathbf{r}, t) + gn_c(\mathbf{r}, t) + 2g\tilde{n}(\mathbf{r}, t) - iR(\mathbf{r}, t) \right) \Phi(\mathbf{r}, t)$$

$$\tilde{n}(\mathbf{r}, t) = \int \frac{d^3p}{(2\pi\hbar)^3} f(\mathbf{p}, \mathbf{r}, t) = \text{non-condensate density}$$

$R(\mathbf{r}, t)$ = particle exchange rate between BEC and non condensate

Quantum Boltzmann equation

The single-particle distribution function satisfies the Quantum Boltzmann equation (QBE). We can understand this equation in the following way:

$$\frac{\partial f(\mathbf{p}, \mathbf{r}, t)}{\partial t} = \left(\frac{\partial f}{\partial t} \right)_{\text{diff}} + \left(\frac{\partial f}{\partial t} \right)_{\text{force}} + \left(\frac{\partial f}{\partial t} \right)_{\text{coll}}$$

This equation says that the rate of change of the number of particle at phase space point (\mathbf{r}, \mathbf{p}) at time t is affected in three ways between times t and $t + dt$:

1. **Diffusion** – particles are already moving and so either change their position from \mathbf{r} to something else or from something else to \mathbf{r} .
2. **Force** – An external force can change a particle's momentum from \mathbf{p} to something else or from something else to \mathbf{p} .
3. **Collisions** – collisions between particles at \mathbf{r} can have their momenta changed.

ZNG Equation of Motion for non-condensate

Quantum Boltzmann Equation (QBE)

$$\frac{\partial f}{\partial t} - \nabla_r U_{\text{eff}} \cdot \nabla_p f + \frac{\mathbf{p}}{m} \cdot \nabla_r f = C_{12}[f, \Phi] + C_{22}[f]$$

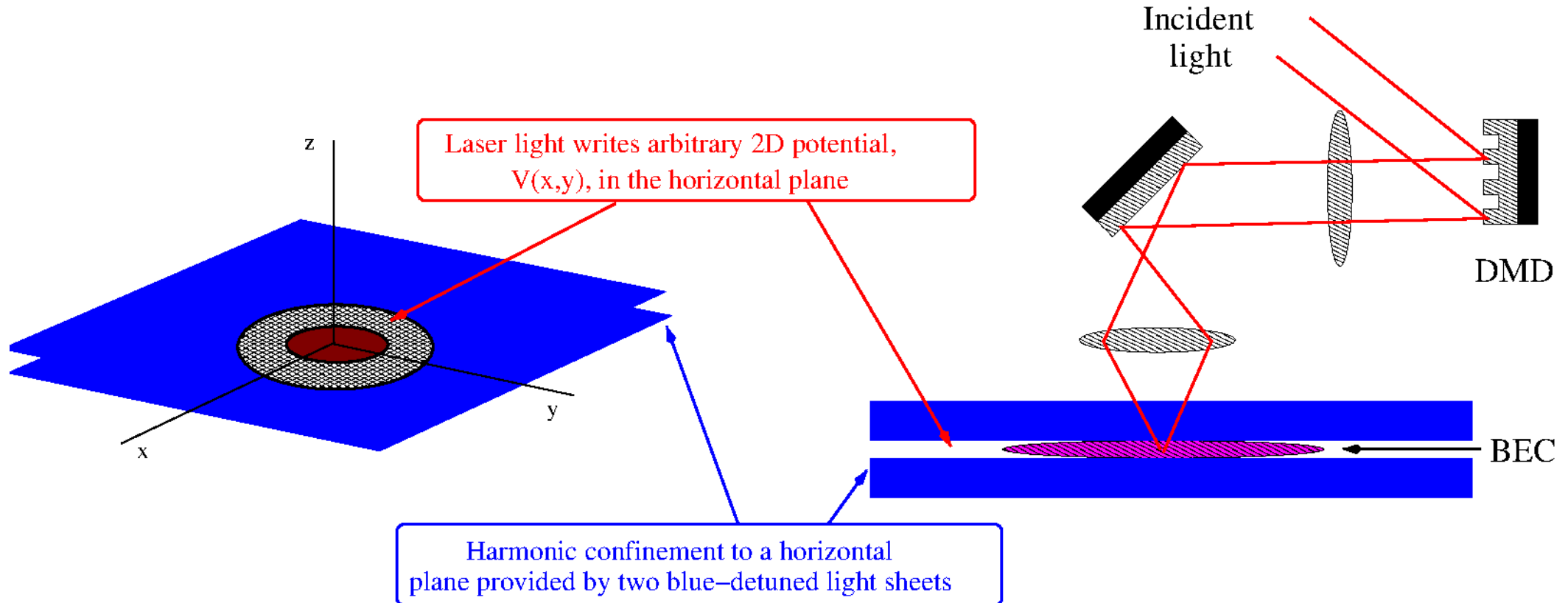
$U_{\text{eff}}(\mathbf{r}, t) = V_{\text{trap}}(\mathbf{r}, t) + 2g(n_c + \tilde{n}) =$ effective potential felt by non-condensate atoms.

$C_{12}[f, \Phi]$, $C_{22}[f]$ are collision integrals.

Model System and Simulation Characteristics

Our atom-circuit setup

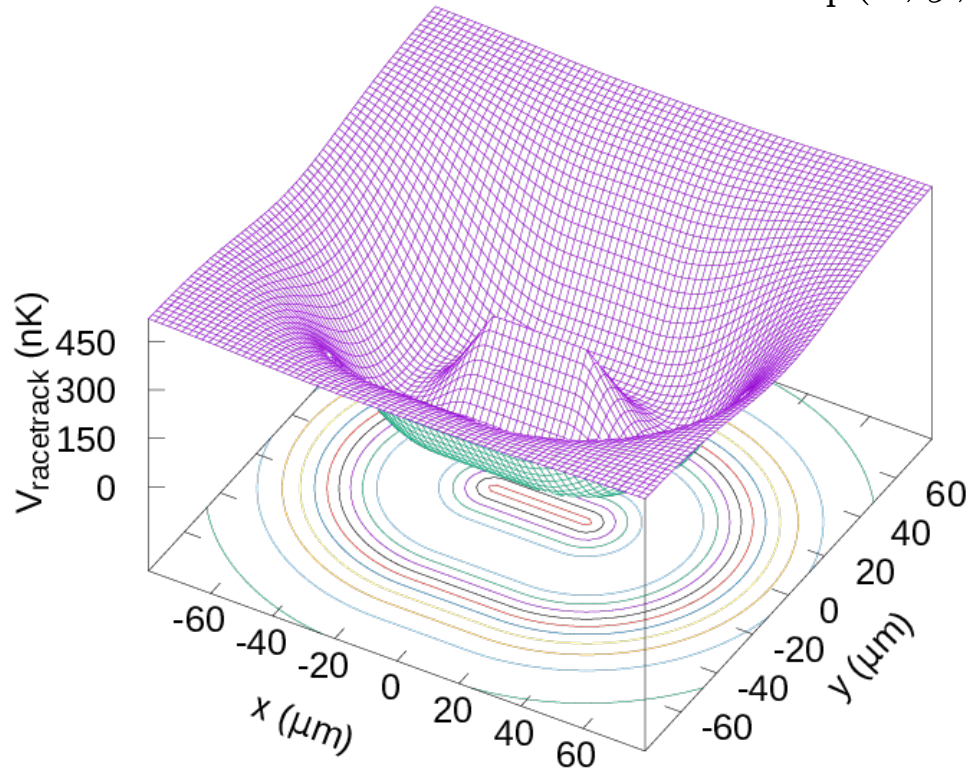
$$V_{\text{trap}}(x, y, z) = \frac{1}{2}M\omega_z^2 z^2 + V(x, y)$$



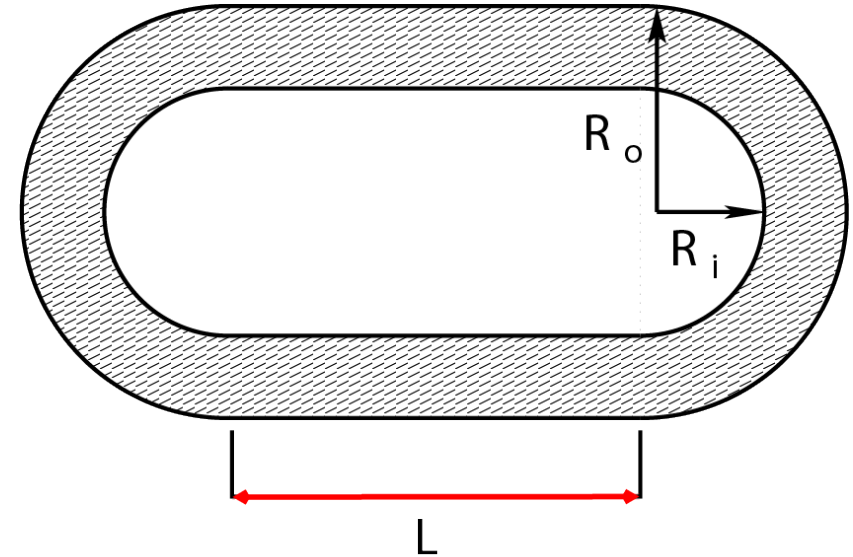
The full potential felt by the condensate atoms is a vertical harmonic potential plus an arbitrary 2D potential, $V(x, y)$, in a horizontal plane:

Our atom-circuit setup

$$V_{\text{trap}}(x, y, z) = \frac{1}{2} M \omega_z^2 z^2 + V(x, y)$$



Racetrack potential



$$V_{\text{racetrack}}(x, y) = V_{\text{rt}} \left\{ \frac{1}{2} \left[\tanh \left(\frac{\rho(x, y) - R_o}{s} \right) + \tanh \left(\frac{R_i - \rho(x, y)}{s} \right) \right] + \tanh \left(\frac{R_o - R_i}{2s} \right) \right\}$$

Systematic zero- and non-zero smooth-flow production studies

- We studied how much flow was produced by stirring a BEC in a racetrack potential using the GPE model for $T=0$.
- The ZNG model was used to study flow production at non-zero T .
- The goal is to develop a simple model to predict the final flow.
- Smooth-flow production simulations were performed where a racetrack condensate was stirred with a fixed-shape paddle to see how much flow could be produced by stirring.
- We varied the following parameters:
 - V_{pmax} , paddle strength
 - TR, paddle stir speed expressed as the number of “Total Revolutions” completed around the $L=0$ racetrack in 4 seconds
 - L , the length of the racetrack straightaway
 - T , the temperature of the initial equilibrium state.

Stirring schedule for all simulations

Barrier Characteristics:

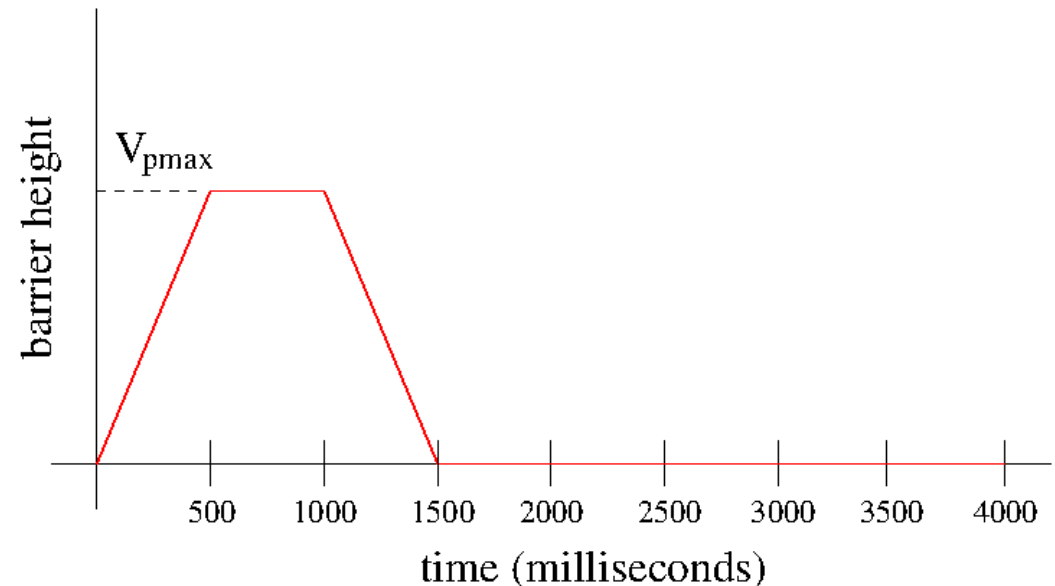
- Rectangular and perpendicular to the racetrack midline
- Perpendicular width is twice that of the racetrack
- Always moving at constant speed
- Energy height varies with time

At 40 ms intervals of the system evolution we calculated:

- Optical density
- **Phase distribution**
- Velocity distribution x and y components
- Vorticity z component

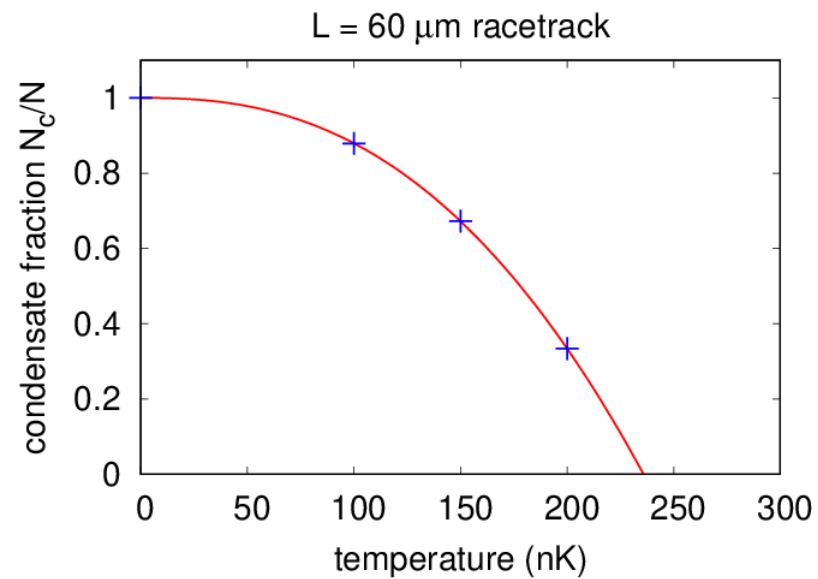
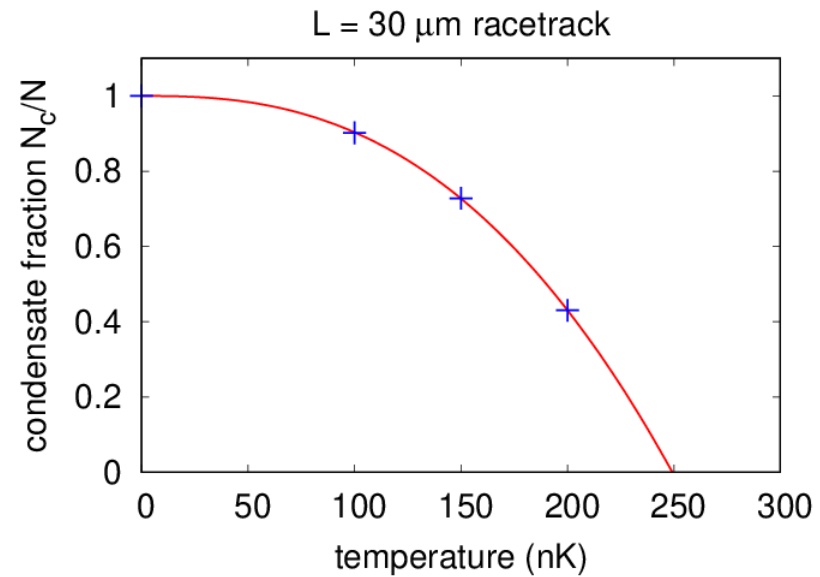
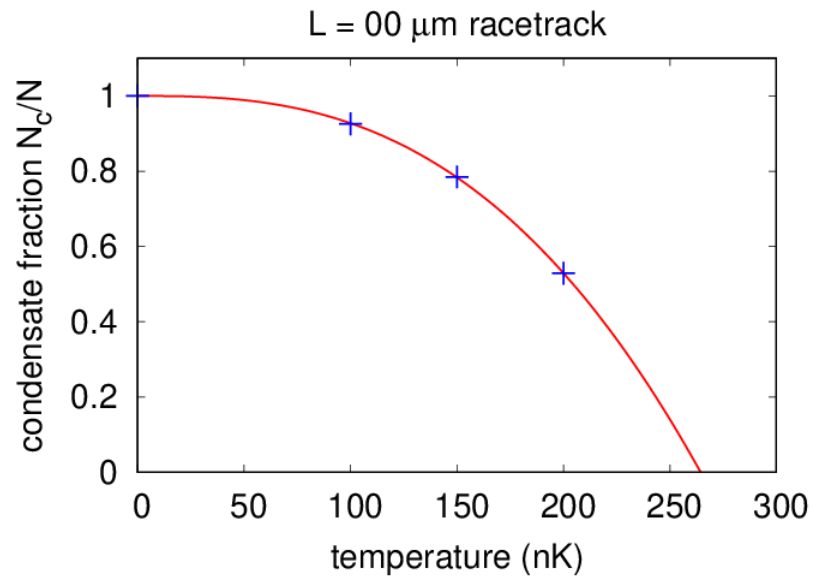
Barrier energy height schedule :

- Ramp up to V_{pmax} over 500 ms
- Remain at V_{pmax} for 500 ms
- Ramp down to zero over 500 ms
- **Simulation runs for a total of 4000 ms for $T=0$ and 2000 ms for non-zero T .**



condensate fractions vs racetrack length

(Total # of atoms fixed at 500,000)

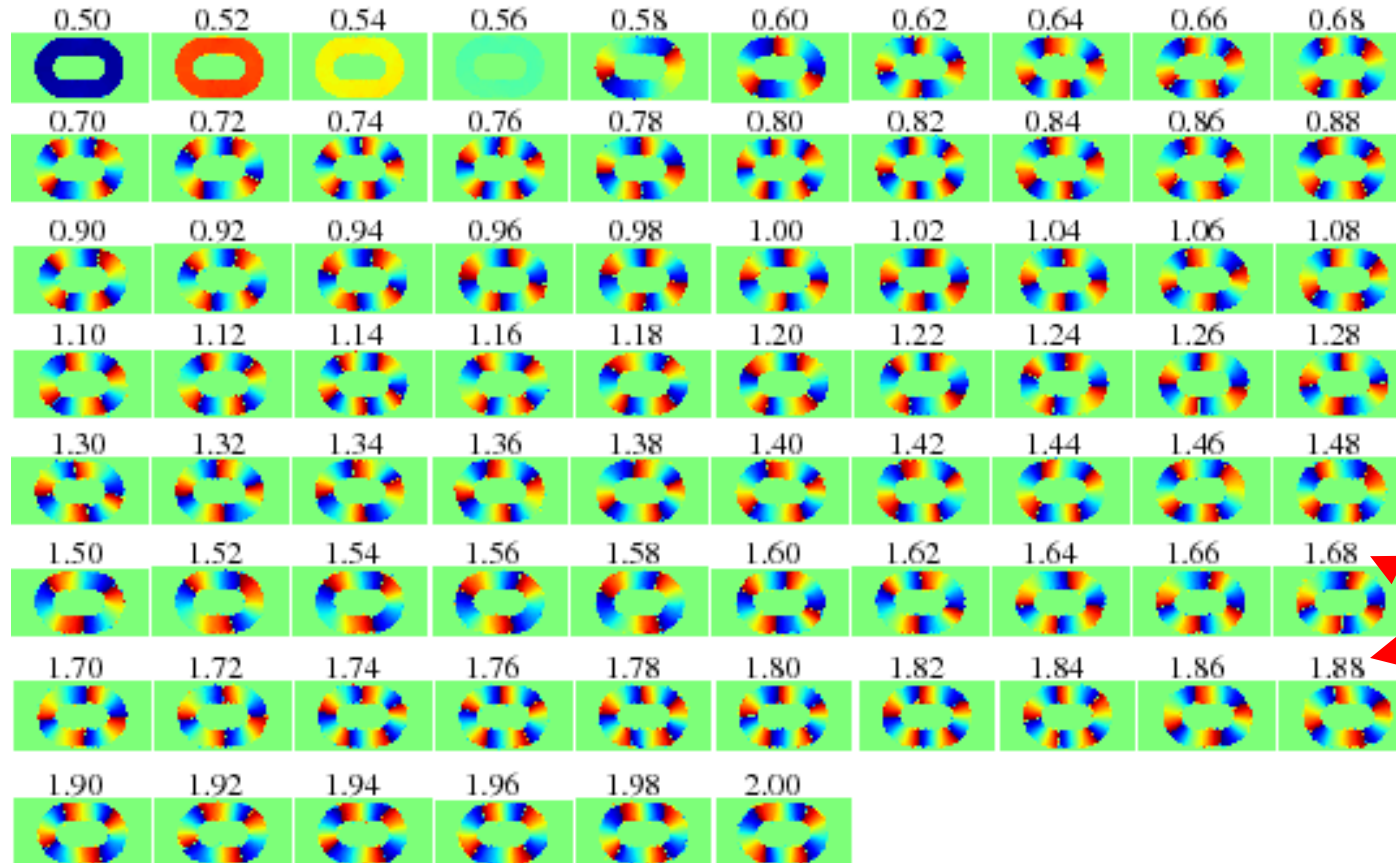


Overview of zero-T systematics study results and extension to non-zero-T

Final phase plots show how much flow is produced

Length (L) = 30 mm, Stir speed (TR) = $9 \times 37 = 333$ mm/s, Temperature (T) = 0 nK

L_30_TR_09



L = racetrack length (μm)
TR = stir speed (1 TR = $37 \mu\text{m/s}$)
T = temp (nK)

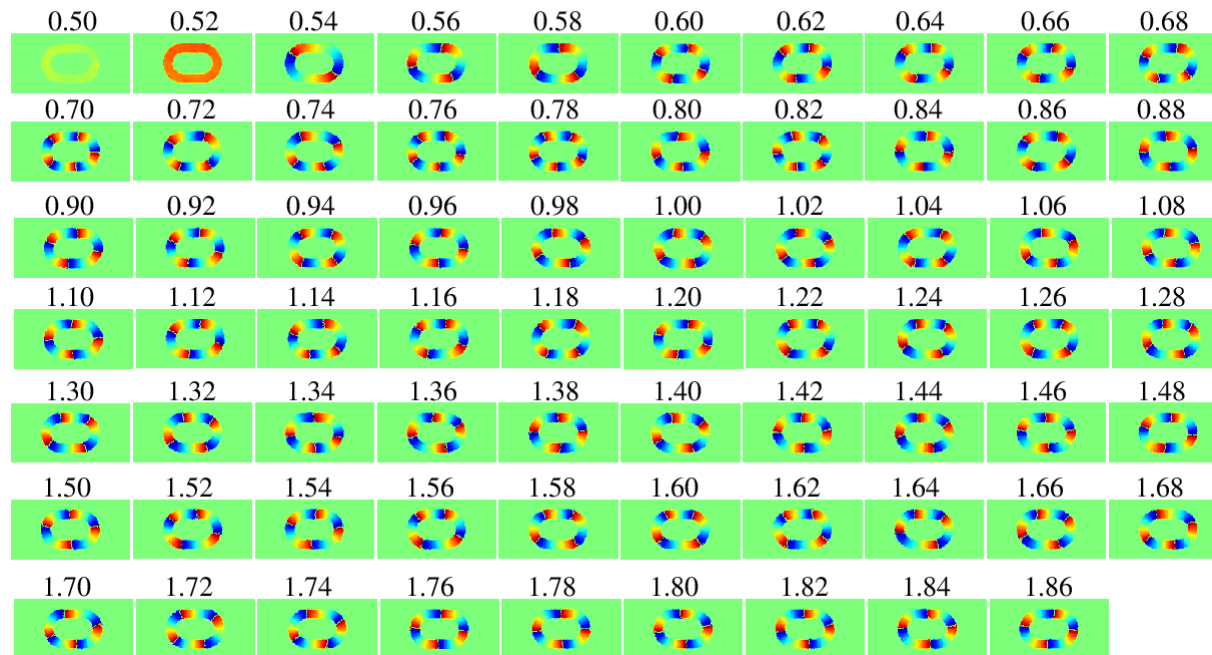
These numbers give the maximum barrier height, $V_{p\text{max}}$, in units of the chemical potential, μ , of the initial condensate.

Flow produced by stirring with fixed trap geometry and barrier speed:

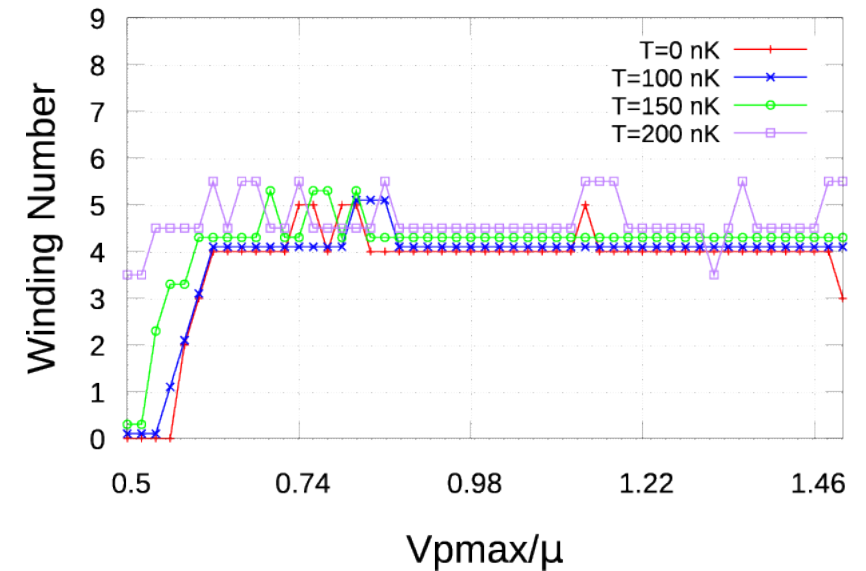
- It is possible to make smooth flow by stirring!
- Flow produced is not monotonic with increasing barrier energy height.
- Flow is quantized but winding number can jump by more than one unit from one barrier height to the next.

can also make flow at finite temperature

L_30_TR_09_T_150

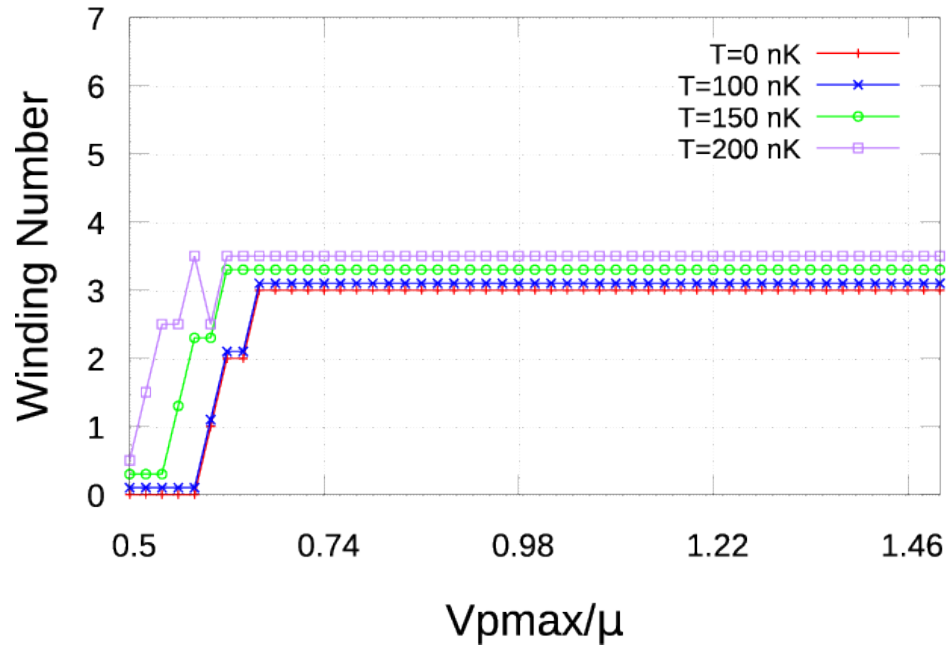


L=30 μm , TR=9, All Temperatures

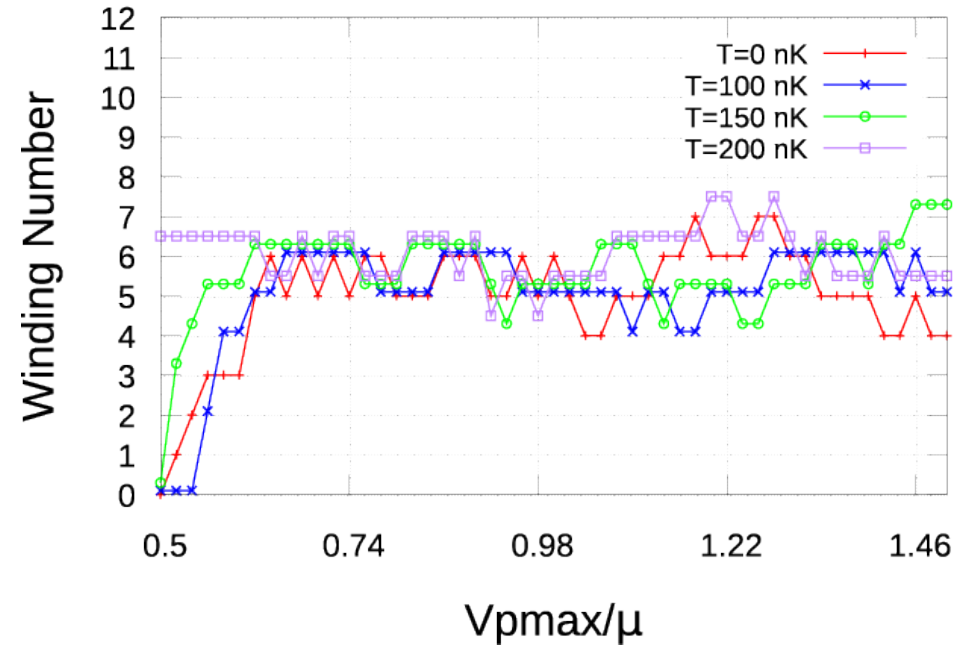


L = racetrack length (μm)
TR = stir speed
(1 TR = 37 $\mu\text{m/s}$)
T = temp (nK)

L=0 μm , TR=9, All Temperatures



L=60 μm , TR=9, All Temperatures



- There are two general features of these plots for all cases considered:
 1. There is a critical max barrier strength, $V_{pmax,c}$, below which no flow is present at the end of the simulation
 2. As V_{pmax} increases, the flow levels off to an average value, w_{avg} , around which it can oscillate.

Summary of Results

- Trends for $V_{pmax,c}$
 - For **fixed length and stirring speed**, with increasing **temperature** $V_{pmax,c}$ gets lower.
 - For **fixed length and temperature**, with increasing **stirring speed** $V_{pmax,c}$ gets lower.
 - For **fixed stirring speed and temperature**, with increasing racetrack **length** $V_{pmax,c}$ gets slightly lower (small effect).
- Trends for w_{avg}
 - For **fixed length and stirring speed**, increasing **temperature** has little effect on w_{avg}
 - For **fixed length and temperature**, with increasing **stirring speed** w_{avg} gets higher.
 - For **fixed stirring speed and temperature**, with increasing racetrack **length** w_{avg} gets higher.

Mechanism of zero-T flow production

Condensate flow dynamics

We found that the final flow depends on the following experimental conditions for fixed stirring schedule:

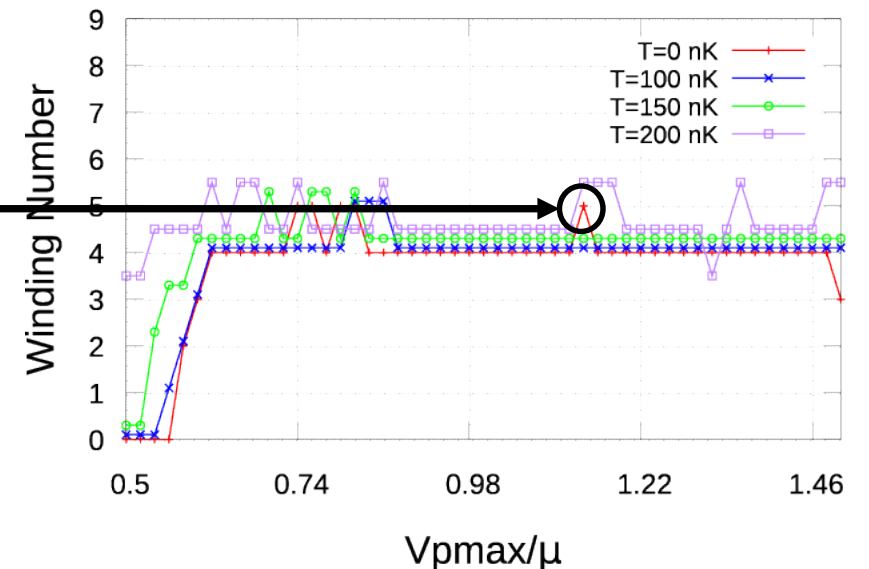
- Speed of stirring, v_s or TR
- Maximum barrier height, V_{pmax}
- Confinement geometry, L
- Temperature, T

What is the mechanism for producing the final flow in the GPE model?

We performed a fine-timescale study of the condensate dynamics during the stirring process for the case with the following characteristics:

- $T = 0$ nK
- $L = 30$ μm
- $V_{pmax} = 54$ nK = 1.14 μ
- $v_s = 339$ $\mu\text{m}/\text{second}$ (TR=09)

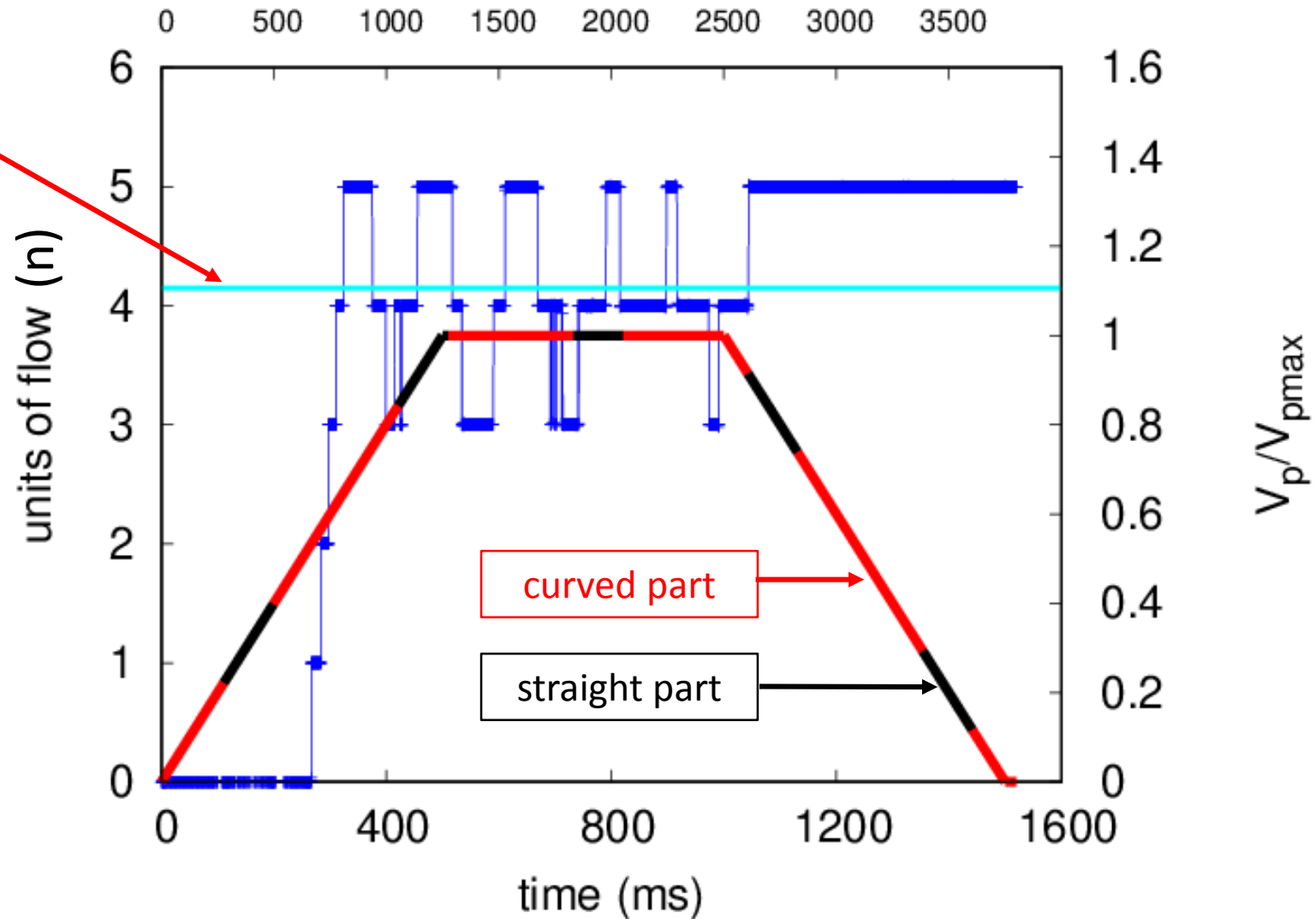
$L=30$ μm , TR=9, All Temperatures



BEC midline circulation during stirring phase (L_30_TR_09_T_000)

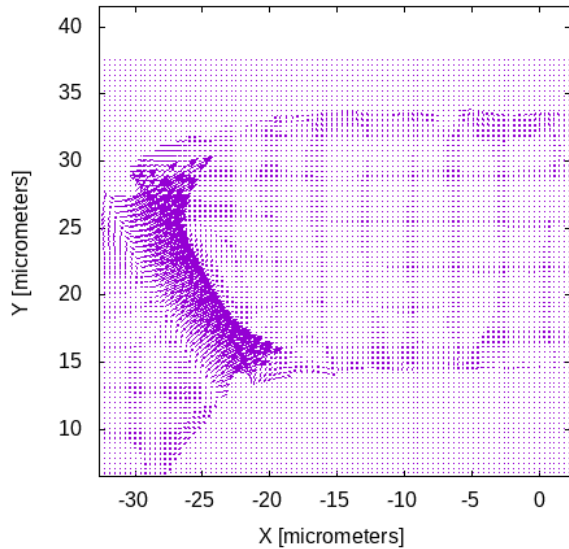
$$\oint_{\mathcal{C}} \mathbf{v} \cdot d\mathbf{l} = \frac{\hbar}{m} \times 2\pi n$$

\mathcal{C} = midline track



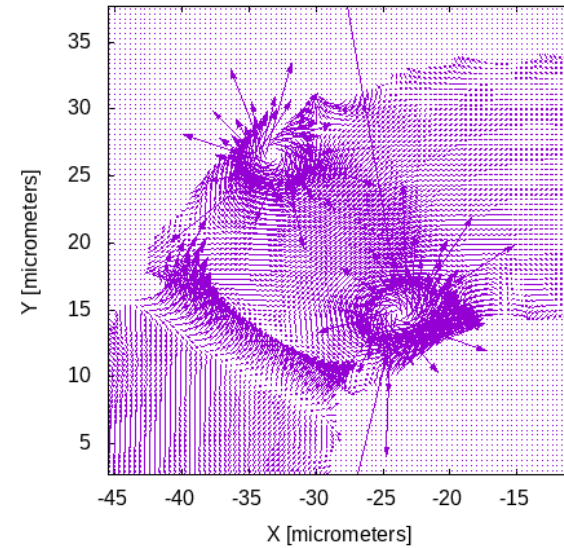
Velocity Distributions of Flow Production: Overview

velocity distribution frame 0590



Backflow develops in the barrier region as the barrier height increases

velocity distribution frame 0680

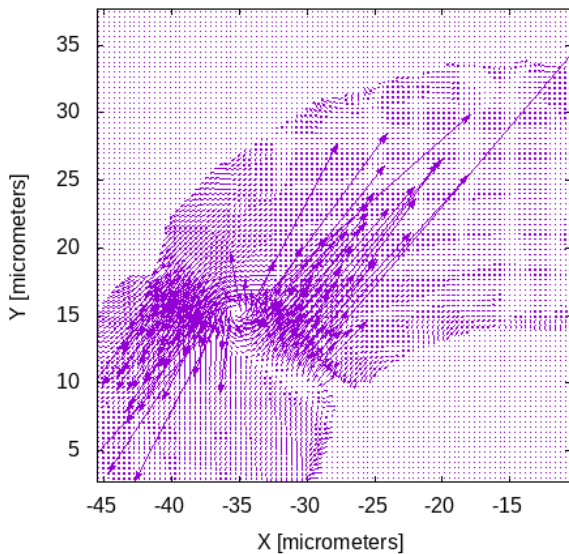


Vortices continue to be formed until the induced flow moves faster than the barrier.

At this point a forward flow develops. Flow is lost/gained afterward by

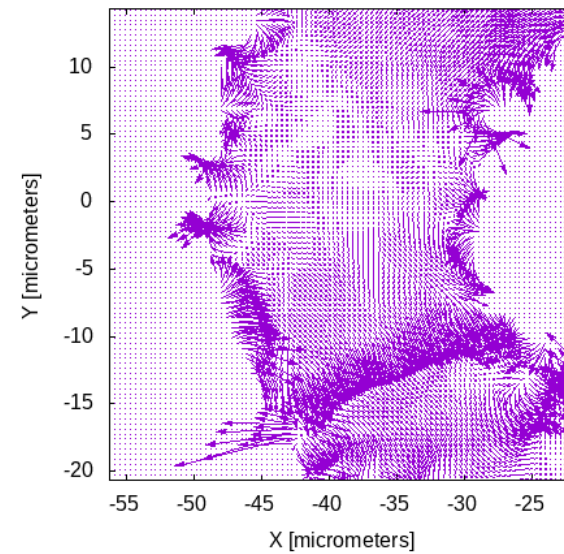
- forward flow in the barrier
- straight-curved or curved to straight transitions

velocity distribution frame 0665



Above some critical flow speed in the barrier region a vortex forms on the outside end of the barrier and migrates into the middle. A vortex/anti-vortex pair is formed.

velocity distribution frame 0864



Research Group and Collaborators



Edwards Research Group

Ben Eller
Olatunde Oladehin
Colson Sapp
Daniel Fogarty
Charles Henry
Elizabeth Ashwood
Brennan Coheleach
Jason Pinto
Anne DeLua

JQI/NIST Theory Collaborators

Charles Clark
Yi-Hsieh Wang
Ted Jacobson

JQI/NIST Experimental Collaborators

Gretchen Campbell
Bill Phillips
Wendell Hill III
Steve Eckel
Fred Jendrzejewski
Chris Lobb

Summary and Future Work

- We investigated whether it is possible to create smooth flow in a “racetrack” atom circuit by stirring with a rectangular-shaped barrier at zero and non-zero temperature.
- We found that it is possible to make smooth flow in this way at all temperatures and other parameters considered and studied the dependence of the final flow produced on stirring speed, max energy height, racetrack length, and the temperature of the initial state.
- The final flow produced was not monotonic with increasing max barrier height for fixed stirring speed, racetrack length, and temperature.
- Our simulations suggest the creation of vortices by flow through the barrier region. Their transfer in and out of the interior of the BEC is responsible for the increase/decrease of overall circulation around the racetrack.
- Future research will focus on more in-depth study (dynamics) of non-zero temperature cases and the effect of the racetrack length on the circulation produced by stirring, as well as the effect of collisions between BEC and non-condensate atoms.