Exercise 1: Covariant derivative

Prove that the term $\overline{\Psi}\not{D}\Psi$ where the covariant derivative is given by:

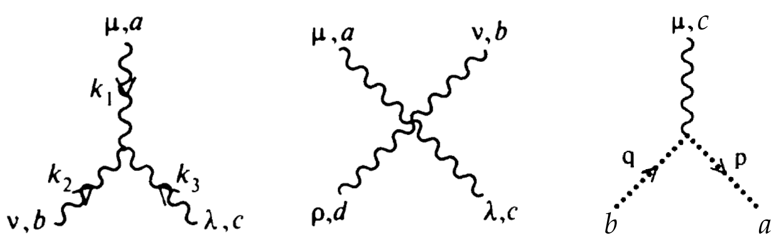
$$D_\mu = \partial_\mu - \mathrm{i}g\widetilde{W}_\mu\,, \quad \widetilde{W}_\mu = T_a W^a_\mu$$

is invariant under gauge transformations:

$$\Psi \mapsto \Psi' = U\Psi\,, \quad U = \exp\{-\mathrm{i}T_a\theta^a(x)\}$$
$$\widetilde{W}_\mu \mapsto \widetilde{W}'_\mu = U\widetilde{W}_\mu U^\dagger - \frac{\mathrm{i}}{g}(\partial_\mu U)U^\dagger$$

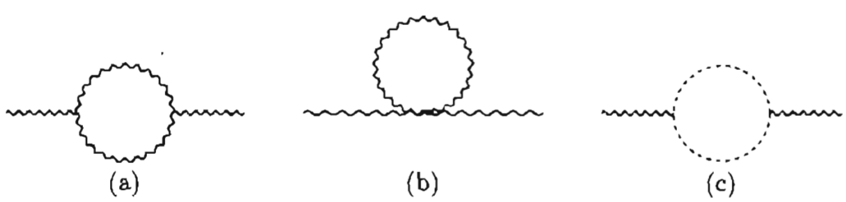
Exercise 2: Feynman rules of general non-Abelian gauge theories

Obtain the Feynman rules for cubic and quartic self-interactions among gauge fields in a general non-Abelian gauge theory, as well as those for the interactions of Faddeev-Popov ghosts with gauge fields:



Exercise 3: Faddeev-Popov ghosts and gauge invariance

Consider the 1-loop self-energy diagrams for non-Abelian gauge theories in the figure. Calculate the diagrams in the 't Hooft-Feynman gauge and show that the sum does not have the tensor structure $g_{\mu\nu}k^2 - k_\mu k_\nu$ required by the gauge invariance of the theory unless diagram (c) involving ghost fields is included.



Hint: Take Feynman rules from previous excercise and use dimensional regularization. It is convenient to use the Passarino-Veltman tensor decomposition of loop integrals:

$$\frac{\mathrm{i}}{16\pi^2}\{B_0,\, B_\mu,\, B_{\mu\nu}\} = \mu^\epsilon \int \frac{\mathrm{d}^D q}{(2\pi)^D}\frac{\{1,\, q_\mu,\, q_\mu q_\nu\}}{q^2(q+k)^2}$$

$$\text{where} \quad B_0 = \Delta_\epsilon + \text{finite}$$
$$B_\mu = k_\mu B_1\,, \quad B_1 = -\frac{\Delta_\epsilon}{2} + \text{finite}$$

$$B_{\mu\nu} = g_{\mu\nu} B_{00} + k_\mu k_\nu B_{11} \ , \quad B_{00} = -\frac{k^2}{12}\Delta_\epsilon + \text{finite} \ , \quad B_{11} = \frac{\Delta_\epsilon}{3} + \text{finite}$$

with $\Delta_\epsilon = 2/\epsilon - \gamma + \ln 4\pi$ and $D = 4 - \epsilon$. You may check that the ultraviolet divergent part has the expected structure or find the final result in terms of scalar integrals, that for massless fields read:

$$B_1 = -\frac{1}{2} B_0 \ , \quad B_{00} = -\frac{k^2}{4(D-1)} B_0 \ , \quad B_{11} = \frac{D}{4(D-1)} B_0 \ .$$

Do not forget a symmetry factor (1/2) in front of (a) and (b), and a factor (-1) in (c).

Exercise 4: Propagator of a massive vector boson field

Consider the Proca Lagrangian of a massive vector boson field

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} M^2 A_\mu A^\mu \ , \quad \text{with} \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \ .$$

Show that the propagator of A_μ is

$$\widetilde{D}_{\mu\nu}(k) = \frac{\mathrm{i}}{k^2 - M^2 + \mathrm{i}0} \left[-g_{\mu\nu} + \frac{k_\mu k_\nu}{M^2} \right]$$

Exercise 5: Propagator of a massive gauge field

Consider the U(1) gauge invariant Lagrangian $\mathcal{L}$ with gauge fixing $\mathcal{L}_{\text{GF}}$:

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + (D_\mu \phi)^\dagger (D^\mu \phi) - \mu^2 \phi^\dagger \phi - \lambda (\phi^\dagger \phi)^2$$
$$\mathcal{L}_{\text{GF}} = -\frac{1}{2\xi} (\partial_\mu A^\mu - \xi M_A \chi)^2 \ , \quad \text{with} \quad D_\mu = \partial_\mu + \mathrm{i} e A_\mu \ , \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

where $M_A = ev$ after spontaneous symmetry breaking ($\mu^2 < 0$, $\lambda > 0$) when the complex scalar field ϕ acquires a VEV and is parameterized by

$$\phi(x) = \frac{1}{\sqrt{2}} [v + \varphi(x) + \mathrm{i}\chi(x)] \ , \quad \mu^2 = -\lambda v^2.$$

Show that the propagators of φ, χ and the gauge field A_μ are respectively

$$\widetilde{D}^\varphi(k) = \frac{\mathrm{i}}{k^2 - M_\varphi^2 + \mathrm{i}0} \quad \text{with } M_\varphi^2 = -2\mu^2 = 2\lambda v^2$$
$$\widetilde{D}^\chi(k) = \frac{\mathrm{i}}{k^2 - \xi M_A^2 + \mathrm{i}0} \ , \quad \widetilde{D}_{\mu\nu}(k) = \frac{\mathrm{i}}{k^2 - M_A^2 + \mathrm{i}0} \left[-g_{\mu\nu} + (1-\xi) \frac{k_\mu k_\nu}{k^2 - \xi M_A^2} \right]$$

Exercise 6: The conjugate Higgs doublet

Show that $\Phi^c \equiv i\sigma_2 \Phi^* = \begin{pmatrix} \phi^{0*} \\ -\phi^- \end{pmatrix}$ transforms under SU(2) like $\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$, with $\phi^- = (\phi^+)^*$. What are the weak isospins, hypercharges and electric charges of ϕ^0, ϕ^{0*}, ϕ^+, ϕ^-? *Hint*: Use the property of Pauli matrices: $\sigma_i^* = -\sigma_2 \sigma_i \sigma_2$.

Exercise 7: Lagrangian and Feynman rules of the Standard Model

Try to reproduce the Lagrangian and the corresponding Feynman rules of as many Standard Model interactions as you can. Of particular interest/difficulty are [VVV] and [VVVV].

Exercise 8: Z pole observables at tree level

Show that

(a) $\Gamma(f\bar{f}) \equiv \Gamma(Z \to f\bar{f}) = N_c^f \frac{\alpha M_Z}{3}(v_f^2 + a_f^2)\,, \quad N_c^f = 1\ (3)$ for $f =$ lepton (quark)

(b) $\sigma_{\text{had}} = 12\pi \frac{\Gamma(\mathrm{e}^+\mathrm{e}^-)\Gamma(\text{had})}{M_Z^2 \Gamma_Z^2}$

(c) $A_{FB} = \frac{3}{4} A_f\,, \quad$ with $A_f = \frac{2 v_f a_f}{v_f^2 + a_f^2}$

Exercise 9: Higgs partial decay widths at tree level

Show that

(a) $\Gamma(H \to f\bar{f}) = N_c^f \frac{G_F M_H}{4\pi\sqrt{2}} m_f^2 \left(1 - \frac{4m_f^2}{M_H^2}\right)^{3/2}\,, \quad N_c^f = 1\ (3)$ for $f =$ lepton (quark)

(b)
$$\Gamma(H \to W^+W^-) = \frac{G_F M_H^3}{8\pi\sqrt{2}} \sqrt{1 - \frac{4M_W^2}{M_H^2}} \left(1 - \frac{4M_W^2}{M_H^2} + \frac{12 M_W^4}{M_H^4}\right)$$

$$\Gamma(H \to ZZ) = \frac{G_F M_H^3}{16\pi\sqrt{2}} \sqrt{1 - \frac{4M_Z^2}{M_H^2}} \left(1 - \frac{4M_Z^2}{M_H^2} + \frac{12 M_Z^4}{M_H^4}\right)$$