

Recall from first let we

Defn A spectral triple (A, \mathcal{H}, D) is given by

A : \ast -algebra of bdd operators on

\mathcal{H} : Hilbert space

D : ess. self-adj. operator in \mathcal{H}

st. - $[D, a]$ extends to bdd commutator

- $(1 + D^2)^{-1}$ is cpt. operator.

grading: γ

real structure: J

$$\gamma D = -D \gamma$$

$$J^2 = \varepsilon, J D = \varepsilon' D, J \gamma = \varepsilon'' \gamma$$

$$\varepsilon, \varepsilon', \varepsilon'' \in \{\pm 1\}.$$

N.B. Distance formula makes sense on $S(A)$:

$$d(\varphi, \psi) = \sup_{a \in A} \{ |\varphi(a) - \psi(a)| : \|[D, a]\| \leq 1 \}$$

for all $\varphi, \psi \in S(A)$. (positive, normed lin. functionals)

Basic example: $(M, g) \quad \mathcal{J}_M, \quad D_M, \quad \gamma_M, \quad J_M, \dots$

Finite spaces $X = \begin{matrix} \bullet & \bullet & \dots & \bullet \\ 1 & 2 & & N \end{matrix}$

$X \mapsto (\mathbb{C}^N, \mathbb{C}^N, D = D^\dagger)$: diagonal matrices $\begin{pmatrix} \phi_1 & & \\ & \dots & \\ & & \phi_N \end{pmatrix}$
 hermitian matrix.

Reconstruction: any metric on $\{1, 2, \dots, N\}$ can be obtained from such a sp. tr.

Ex. $N=2$ $(\mathbb{C}^2, \mathbb{C}^2, D = \begin{pmatrix} 0 & c \\ \bar{c} & 0 \end{pmatrix})$
 $d(1, 2) = \frac{1}{|c|}$

Finite noncommutative spaces

$\begin{pmatrix} a_1 & & \phi \\ & a_2 & \\ \phi & & \ddots \\ & & & a_N \end{pmatrix}$ block diagonal matrices

$A_F = M_{n_1}(\mathbb{C}) \oplus \dots \oplus M_{n_N}(\mathbb{C})$

$H_F = V$ \mathbb{C} vector space.

$D_F = D_F^\dagger$ herm. matrix

Ex (YM) $\left\{ \begin{array}{l} A = M_n(\mathbb{C}) \\ H = \mathbb{C}^n \end{array} \right. , \quad D = D^\dagger$

Ex (SM)

$$A = \mathbb{C} \oplus \mathbb{H} \oplus M_3(\mathbb{C})$$

$$\mathcal{H} = \mathbb{C}^{96}$$

$$= \left(\mathbb{C}^2 \otimes \mathbb{C}^2 \otimes \left(\mathbb{C}^2 \oplus \mathbb{C}^2 \otimes \mathbb{C}^3 \right) \right)^{\oplus 3}$$

$\gamma: f \otimes \bar{f}$
 $\gamma: L \otimes R$

$$D = \begin{pmatrix} J & T \\ T^* & \bar{J} \end{pmatrix}$$

$$S = \begin{pmatrix} 0 & \gamma_{L,R} \\ \gamma_{L,R} & 0 \end{pmatrix}$$

$$T = \begin{pmatrix} 0 & \gamma_R \\ \gamma_R & 0 \end{pmatrix}$$

gen
 \downarrow
 $\oplus 3$

Ex (PS)

$$A = \mathbb{H}_R \oplus \mathbb{H}_L \oplus M_4(\mathbb{C})$$

$$\mathcal{H} = \mathbb{C}^2 \otimes \mathbb{C}^2 \otimes \left(\mathbb{C}^2 \otimes \mathbb{C}^4 \right)^{\oplus 3}$$

Ex. (AC manifolds): product $M \times F$

$$\mathbb{C}^\infty(M \otimes A_F, \mathcal{L}(\mathbb{J}_M) \otimes \mathcal{H}_F, D_M \otimes 1 + \gamma \otimes D_F)$$

$$M \text{ even-dimensional } \left. \begin{array}{l} \gamma_M \otimes 1 + 1 \otimes \gamma_F \\ \mathbb{J}_M \otimes 1 + 1 \otimes \mathbb{J}_F \end{array} \right\}$$

Symmetries of spectral triples.

(A, \mathcal{H}, D)

Most natural : unitary equivalence.

Defn. $(A_1, \mathcal{H}_1, D_1)$ and $(A_2, \mathcal{H}_2, D_2)$ are called unitarily equivalent if there exists isomorphism $\alpha: A_1 \xrightarrow{\sim} A_2$ and unitary operator $U: \mathcal{H}_1 \rightarrow \mathcal{H}_2$

s.t.

$$\alpha(a) = U a U^*$$

$$D_2 = U D_1 U^*$$

Special case: $u \in U(A) : \begin{cases} A_1 = A_2 = A \\ \mathcal{H}_1 = \mathcal{H}_2 = \mathcal{H} \end{cases}$

$$\alpha(u) = u a u^\dagger = \alpha_u(a)$$

inner automorphisms $\rightarrow \alpha: U(A) \rightarrow \text{Inn}(A)$.

$$u D u^\dagger = D + u [D, u^\dagger]$$

pure gauge fields

Ex (manifold): $u [D, u^\dagger] = i \gamma^\mu (u \partial_\mu u^\dagger)$

$$u = e^{iX} \quad = \gamma^\mu \partial_\mu X$$

Ex (AC) $(C^\infty(M) \otimes M_N(\mathbb{C}), C^\infty(M) \otimes \mathbb{C}^N, D_N \otimes 1 + \gamma_\mu \otimes D)$

Then $U(A) = C^\infty(M, u(n))$

and $u [D_M, u^\dagger]$ is a $u(n)$ -pure gauge field

Can we extend this mathematical description of pure gauge fields to capture arbitrary gauge fields?