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# Extracting universal data from 2d chiral topological states: a CFT perspective

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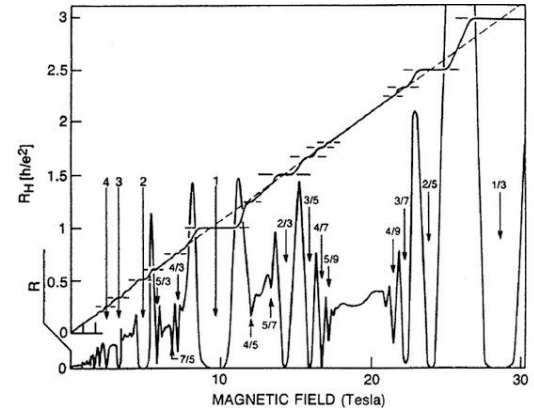
Entanglement in Strongly Correlated Systems  
Benasque, February 13<sup>th</sup>, 2017

# Outline

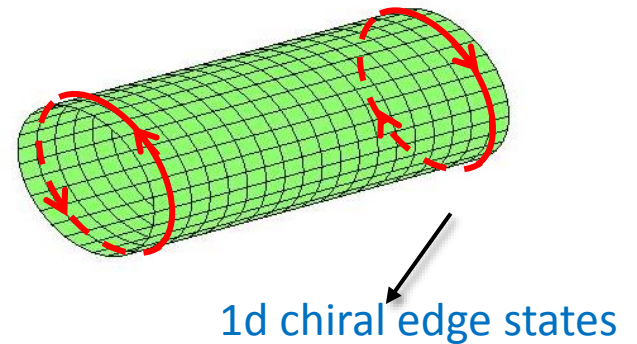
- Quick sketch of chiral conformal field theory (CFT)
- Topological entanglement entropy
- Entanglement spectrum
- Momentum polarization (and perhaps also Klein twist)  
=> tomorrow's talk
- Summary

# 2d chiral topological states

- 2d topological states which break time-reversal symmetry
  - Integer and fractional quantum Hall states
  - p+ip superconductor
  - Chiral spin liquids
  - ...



- 2d chiral topological states have deep relation with CFT
  - Gapless edge excitations
  - Trial wave functions
  - Fusion rules of anyons
  - ...



- Extracting topological data is basically extracting information of the chiral CFT.

# Quick sketch of chiral CFT

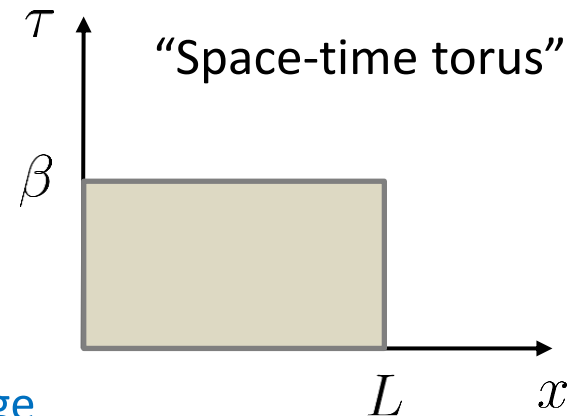
- 1d chiral CFT: gapless one-way movers with linear dispersion

$$H = P$$



$$H = P = \frac{2\pi}{L} \left( L_0 - \frac{c}{24} \right)$$

central charge



Eigenvectors:  $L_0|\alpha\rangle = h_\alpha|\alpha\rangle$

Vacuum:  $L_0|0\rangle = 0$

conformal dimension

- Non-chiral CFTs, describing 1d critical chains, are recovered by combining with a counter-propagating branch.

# Virasoro algebra, primary field, and Verma module

Virasoro algebra:  $[L_n, L_m] = (n - m)L_{n+m} + \frac{c}{12}n(n^2 - 1)\delta_{n+m,0}$

Primary state:  $L_n|\alpha\rangle = 0, \forall n > 0$



Primary field

$$|\alpha\rangle = \lim_{z \rightarrow 0} \mathcal{V}_\alpha(z)|0\rangle$$

Verma module:

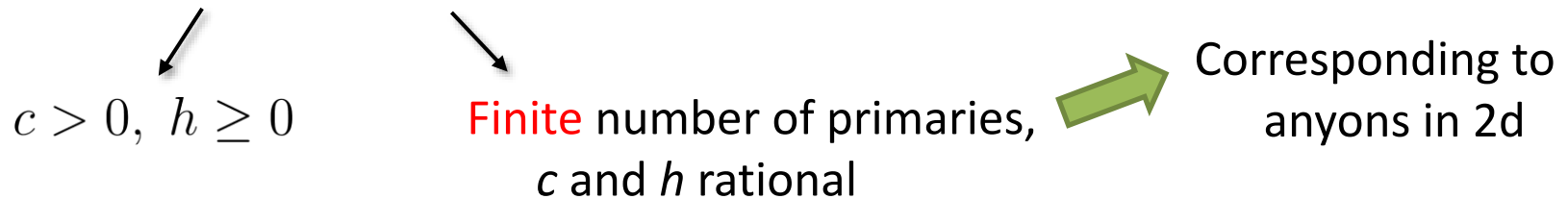
$h_\alpha$	$ \alpha\rangle$				
$h_\alpha + 1$	$L_{-1} \alpha\rangle$				
$h_\alpha + 2$	$L_{-1}^2 \alpha\rangle$	$L_{-2} \alpha\rangle$			
$\vdots$	$L_{-1}^3 \alpha\rangle$	$L_{-1}L_{-2} \alpha\rangle$	$L_{-3} \alpha\rangle$		
	$L_{-1}^4 \alpha\rangle$	$L_{-1}^2L_{-2} \alpha\rangle$	$L_{-1}L_{-3} \alpha\rangle$	$L_{-2}^2 \alpha\rangle$	$L_{-4} \alpha\rangle$
	$\vdots$				

➤ Null vectors with vanishing norm should be removed, e.g.  $L_{-1}|0\rangle = 0$

# Virasoro algebra, primary field, and Verma module

## Further remarks:

- Additional symmetry could further unify a number of Virasoro towers into a single one [e.g.  $SU(2)_1$  CFT with Kac-Moody algebra organizes them into two towers => two (Kac-Moody) primaries].
- Those appearing as edge theories of 2d chiral topological states are **unitary** and **rational** CFTs.



Example: **Ising CFT**     $c = 1/2$

**primaries**     $I, \sigma, \psi$      $\longrightarrow$      $h_I = 0, h_\sigma = \frac{1}{16}, h_\psi = \frac{1}{2}$

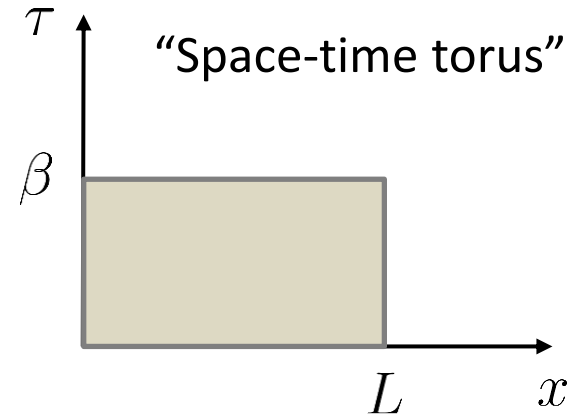
# Chiral character of CFT

- Chiral character is like a partition function (restricted to each tower of states)

$$\begin{aligned}\chi_a(q) &= \text{Tr}_a(e^{-\beta H}) \\ &= \text{Tr}_a[e^{-2\pi\frac{\beta}{L}(L_0 - c/24)}] \\ &= \text{Tr}_a(q^{L_0 - c/24})\end{aligned}$$

↙

$$q = e^{-2\pi\frac{\beta}{L}}$$



Example: **Ising CFT**

$$\begin{aligned}\chi_I(q) &= q^{-1/48}(1 + q^2 + q^3 + 2q^4 + 2q^5 + 3q^6 + \dots) \\ \chi_\sigma(q) &= q^{1/24}(1 + q + q^2 + q^3 + 2q^4 + 2q^5 + 3q^6 + \dots) \\ \chi_\psi(q) &= q^{23/48}(1 + q + q^2 + 2q^3 + 2q^4 + 3q^5 + 4q^6 + \dots)\end{aligned}$$

# Outline


- Quick sketch of chiral conformal field theory (CFT)
- **Topological entanglement entropy**
- Entanglement spectrum
- Momentum polarization (and perhaps also Klein twist)  
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# Topological entanglement entropy

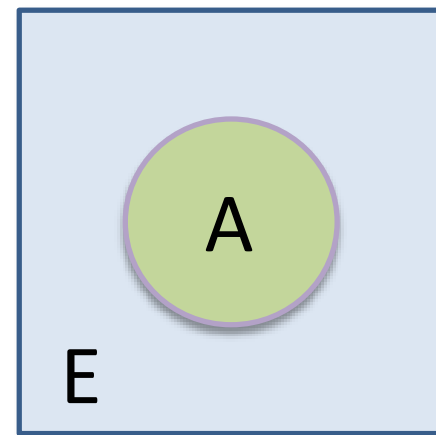
Entanglement entropy:

$$S = \alpha L - \gamma + \dots$$



$$\gamma = \ln(D) \quad \text{D: total quantum dimensions of anyons}$$

- For 2d chiral topological states, Kitaev & Preskill also provided a heuristic “proof” based on CFT...



$$\rho = \text{Tr}_E(|\Psi\rangle\langle\Psi|)$$

$$S = -\text{Tr}(\rho \ln \rho)$$

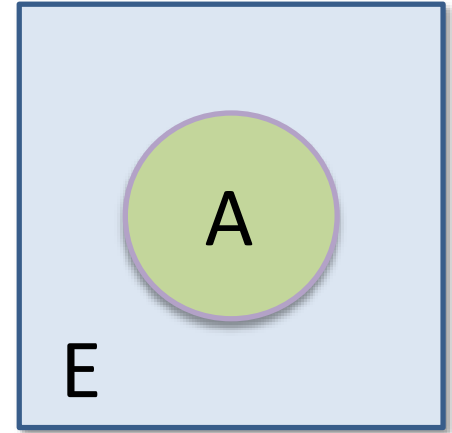
Kitaev & Preskill, '06

Levin & Wen, '06

# Topological entropy from chiral CFT

Kitaev & Preskill's  
assumption:

$$\rho \propto e^{-\beta H} |I\rangle$$

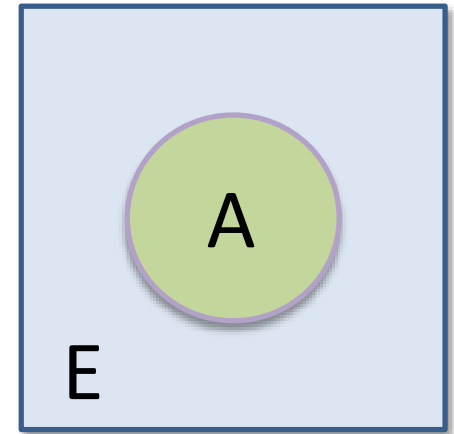
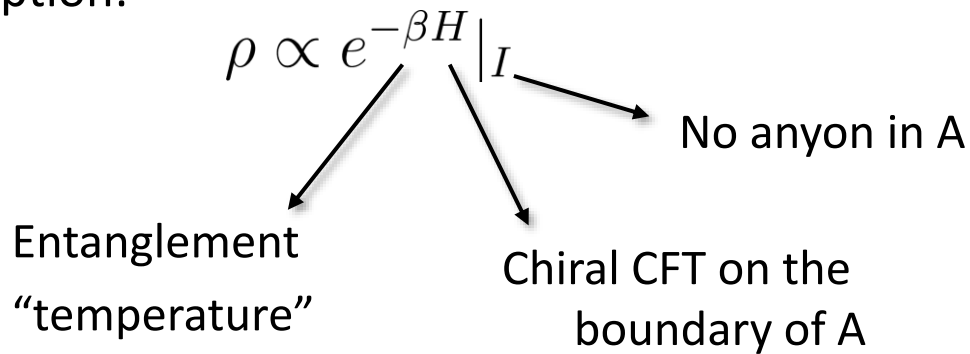


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# Topological entropy from chiral CFT

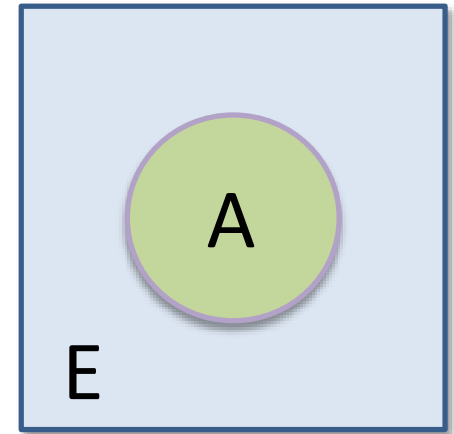
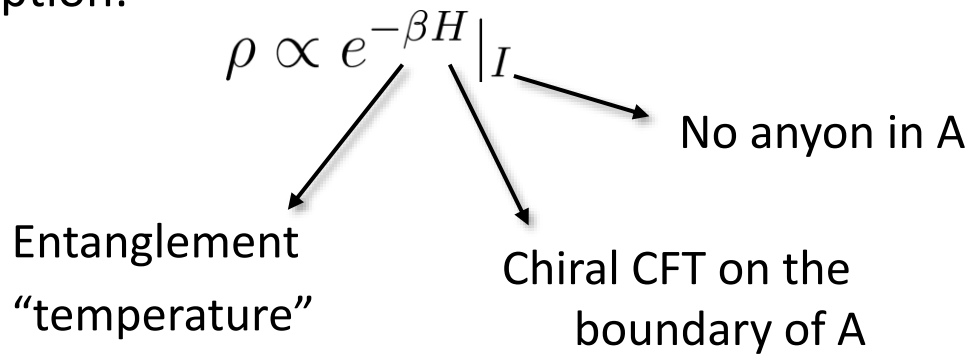
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# Topological entropy from chiral CFT

Kitaev & Preskill's  
assumption:



$$\rho = \text{Tr}_E(|\Psi\rangle\langle\Psi|)$$

$$S = -\text{Tr}(\rho \ln \rho)$$

Renyi entropy:

$$S_n = \frac{1}{1-n} \ln \text{Tr}(\rho^n)$$

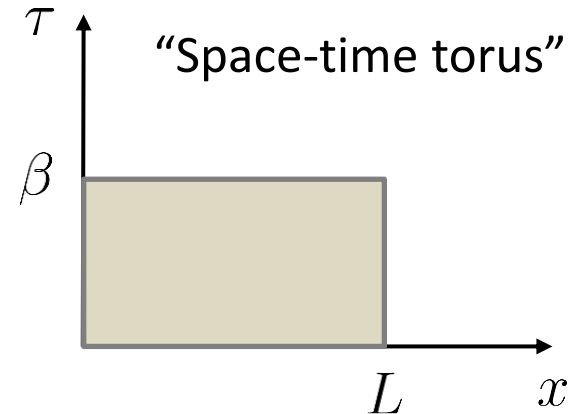
$$\text{Tr}(\rho^n) = \frac{\text{Tr}_I(e^{-n\beta H})}{[\text{Tr}_I(e^{-\beta H})]^n} = \frac{\chi_I(q^n)}{[\chi_I(q)]^n}$$

Recall  $q = e^{-2\pi\frac{\beta}{L}}$

# Calculating chiral character

$$\begin{aligned}\chi_I(q) &= \text{Tr}_I(q^{L_0 - c/24}) \\ &= q^{-c/24}(1 + d_1q + d_2q^2 + \dots)\end{aligned}$$

$$q = e^{-2\pi\frac{\beta}{L}} \rightarrow 1 \quad \text{for } L \gg \beta$$



Modular transformation:

$$\begin{aligned}\chi_I(q) &= \sum_a S_{I,a} \chi_a(q') \\ &\simeq S_{I,I} \chi_I(q' \rightarrow 0) \\ &\simeq S_{I,I} (q')^{-c/24} \\ &= \boxed{S_{I,I} e^{\frac{\pi c}{12} \frac{L}{\beta}}}\end{aligned}$$

$q' = e^{-2\pi\frac{L}{\beta}} \rightarrow 0 \quad \text{for } L \gg \beta$   
 $\chi_a(q') = (q')^{h_a - c/24} (1 + d_1q' + \dots)$   
 $a = I \quad \text{dominant!}$

# Topological entropy from chiral CFT

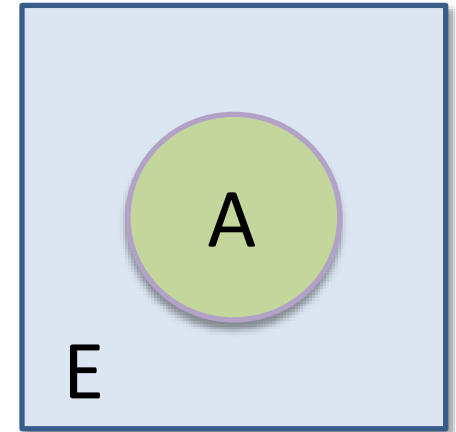
$$\begin{aligned} \text{Tr}(\rho^n) &= \frac{\chi_I(q^n)}{[\chi_I(q)]^n} = \frac{S_{I,I} \exp(\frac{\pi c}{12} \frac{L}{n\beta})}{(S_{I,I})^n \exp(\frac{\pi c}{12} \frac{nL}{\beta})} \\ &= (S_{I,I})^{1-n} \exp\left[\frac{\pi c}{12\beta} \left(\frac{1}{n} - n\right)L\right] \end{aligned}$$

←  $\chi_I(q) \simeq S_{I,I} e^{\frac{\pi c}{12} \frac{L}{\beta}}$

Renyi entropy:

$$\begin{aligned} S_n &= \frac{1}{1-n} \ln \text{Tr}(\rho^n) \\ &= \left(1 + \frac{1}{n}\right) \frac{\pi c}{12\beta} L + \ln(S_{I,I}) \\ &= \boxed{\alpha_n L - \ln(D)} \end{aligned}$$

↙  $S_{I,I} = \frac{1}{D}$



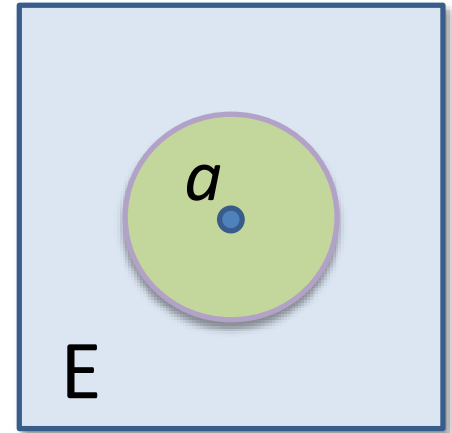
- This CFT-based derivation is consistent with the TQFT approach.

# Topological entropy in the presence of anyon

Kitaev & Preskill's  
assumption:

$$\rho \propto e^{-\beta H} |_a \longrightarrow a \text{ sector}$$

$$\text{Tr}(\rho^n) = \frac{\chi_a(q^n)}{[\chi_a(q)]^n} = (S_{a,I})^{1-n} \exp \left[ \frac{\pi c}{12\beta} \left( \frac{1}{n} - n \right) L \right]$$



$$\chi_a(q) \simeq S_{a,I} e^{\frac{\pi c}{12} \frac{L}{\beta}}$$

Renyi entropy:

$$\begin{aligned} S_n &= \frac{1}{1-n} \ln \text{Tr}(\rho^n) \\ &= \left(1 + \frac{1}{n}\right) \frac{\pi c}{12\beta} L + \ln(S_{a,I}) \\ &= \boxed{\alpha_n L - \ln(D/d_a)} \end{aligned}$$

$$\searrow S_{a,I} = \frac{d_a}{D}$$

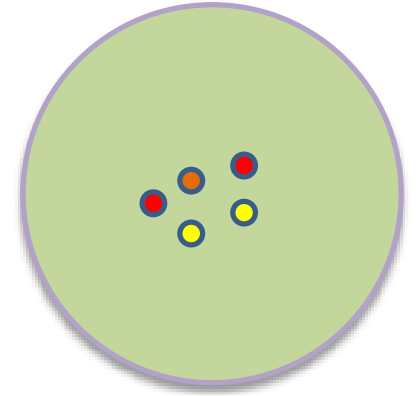
$d_a$ : quantum dimension  
of anyon  $a$

# What about many anyons?

First fuse the anyons:

$$a \times b \times c \times \dots = I + e + f + \dots$$

$$\rho \propto x_1 \rho_I + x_e \rho_e + x_f \rho_f + \dots$$



- The reduced density matrix depends on which final state it is...
- If they fuse into a particular anyon sector (e.g.  $f$ ), entanglement entropy is **smaller**



Such basis is called **minimally entangled basis**, convenient for extracting topological data



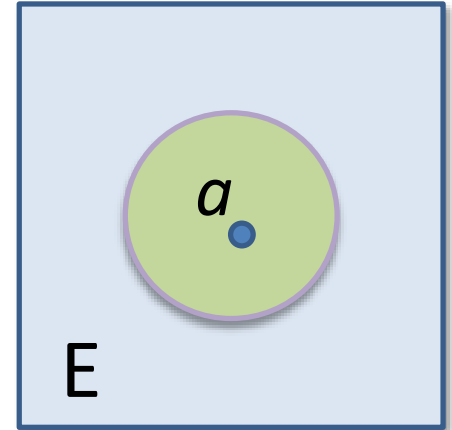
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# Entanglement spectrum

- Li & Haldane: “the low-lying entanglement spectrum can be used as a fingerprint to identify topological order.” [H. Li & Haldane, '08](#)

$$\rho_a = \frac{1}{\chi_a(q)} e^{-\beta H} |_a = \frac{1}{\chi_a(q)} q^{L_0 - c/24} |_a$$



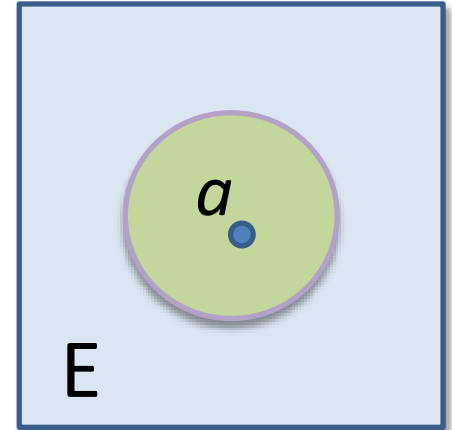
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$$\rho_a = \frac{1}{\chi_a(q)} e^{-\beta H} |_a = \frac{1}{\chi_a(q)} q^{L_0 - c/24} |_a$$

- Largest eigenvalue in sector  $a$ :

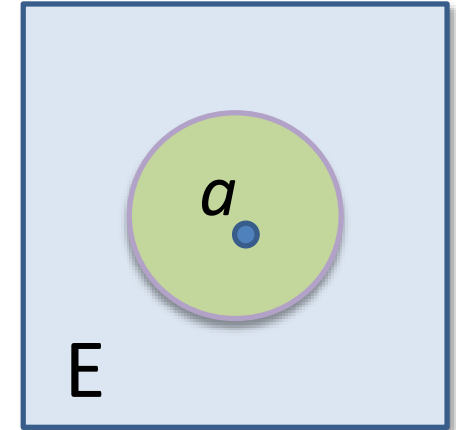
$$\lambda_a = \frac{1}{\chi_a(q)} q^{h_a - c/24}$$



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$$\chi_a(q) \simeq S_{a,I} e^{\frac{\pi c}{12} \frac{L}{\beta}}$$

$$-\ln(\lambda_a) = \frac{\pi c}{12\beta} L - \ln(D/d_a) + 2\pi \frac{\beta}{L} \left( h_a - \frac{c}{24} \right)$$

Topological spin  
of anyons

# Entanglement spectrum

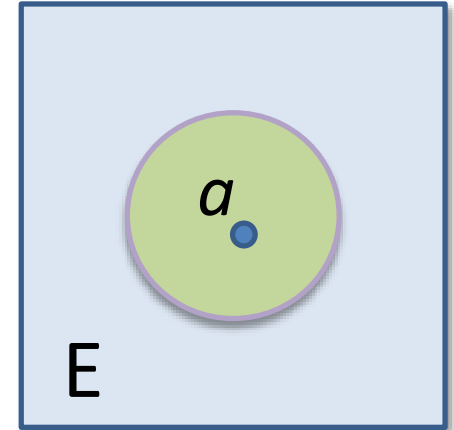
$$-\ln(\lambda_a) = \frac{\pi c}{12\beta} L - \ln(D/d_a) + 2\pi \frac{\beta}{L} (h_a - \frac{c}{24})$$

## Further remarks:

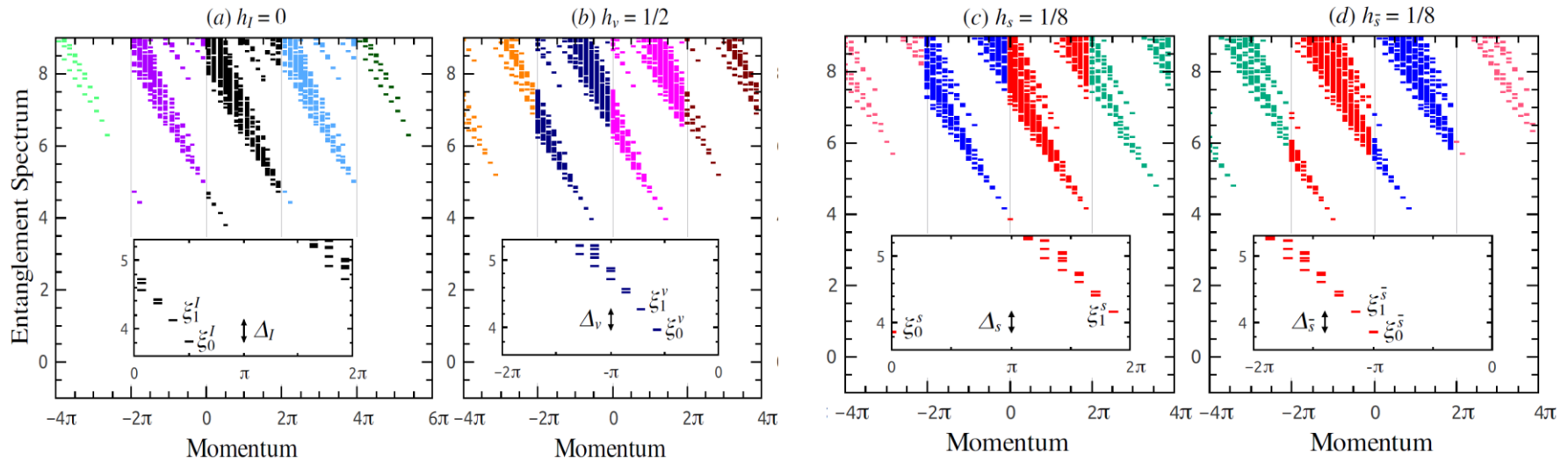
- It's consistent with Renyi entropy result with  $n = \infty$ , but the  $1/L$  correction is not within TQFT framework.
- Non-CFT contributions may exist (presumably satisfy the area law), so better use

$$-\ln(\lambda_a) = \alpha L - \ln(D/d_a) + \frac{\mu}{L} (h_a - \frac{c}{24})$$

- If the entanglement spectrum is close to a chiral CFT, then the formula also works for low-lying entanglement “excited” states.

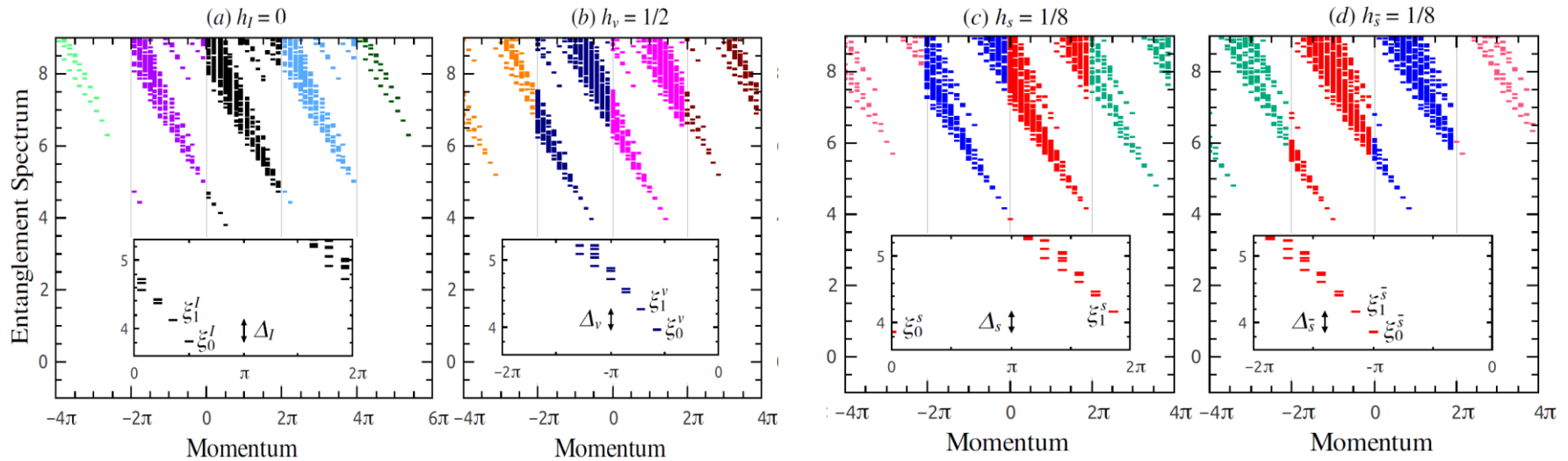


# Example: $\frac{1}{4}$ Laughlin-like states on the lattice

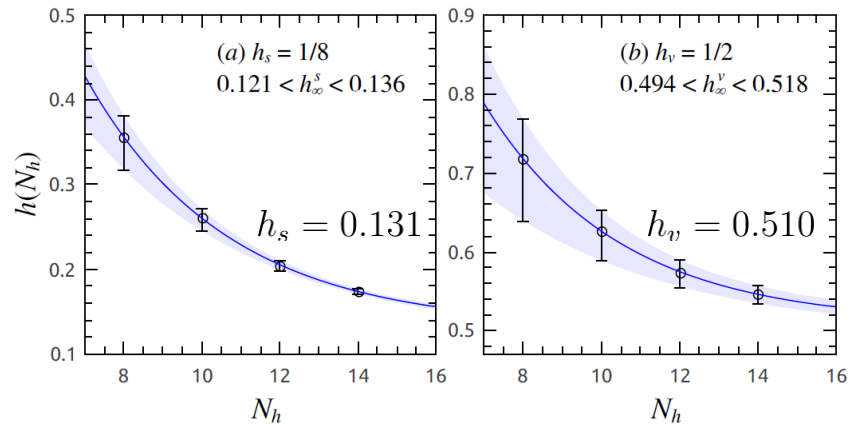


- Expect  $\frac{\xi_0^v - \xi_0^I}{\xi_1^I - \xi_0^I} \rightarrow \frac{1}{2}$

# Example: $\frac{1}{4}$ Laughlin-like states on the lattice



- Expect  $\frac{\xi_0^v - \xi_0^I}{\xi_1^I - \xi_0^I} \rightarrow \frac{1}{2}$
- Increase  $L$  as much as possible



# Summary

- For 2d chiral topological states, CFT provides a very useful tool for understanding their physical properties.
- The universal topological data from entanglement measures, such as topological entanglement entropy and entanglement spectrum, can be derived naturally from the assumption that the reduced density matrix is a thermal state of a chiral CFT.
- Other CFT-based approaches for extracting universal data available  
=> momentum polarization & Klein twist (tomorrow)



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*Thank you for your attention!*