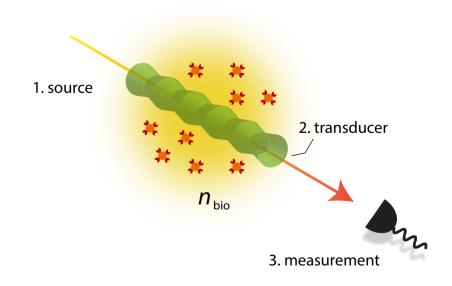
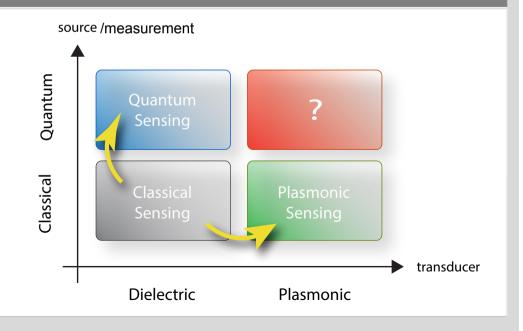


Quantum plasmonic sensing

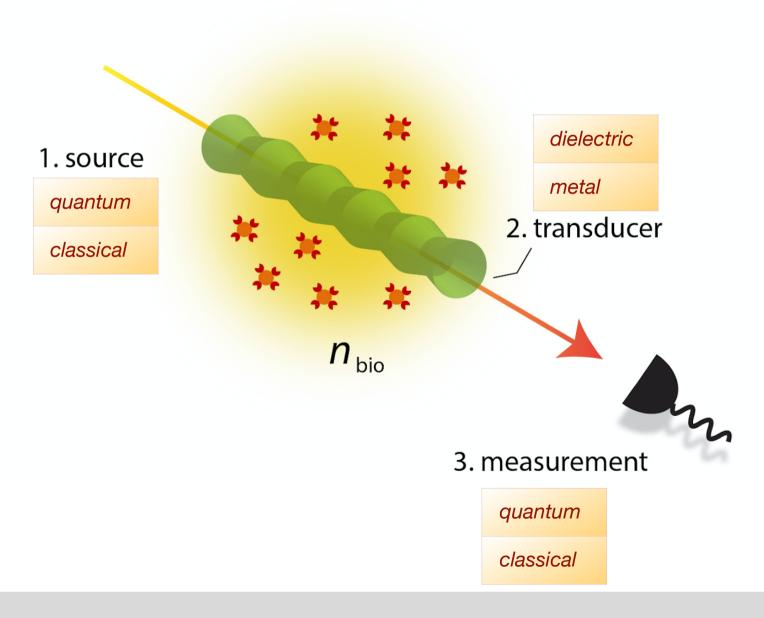
Changhyoup Lee

Institute of Theoretical Solid State Physics

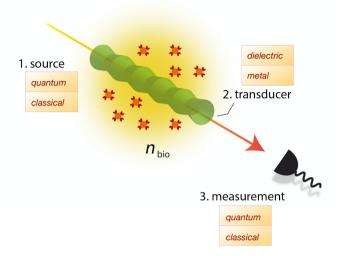


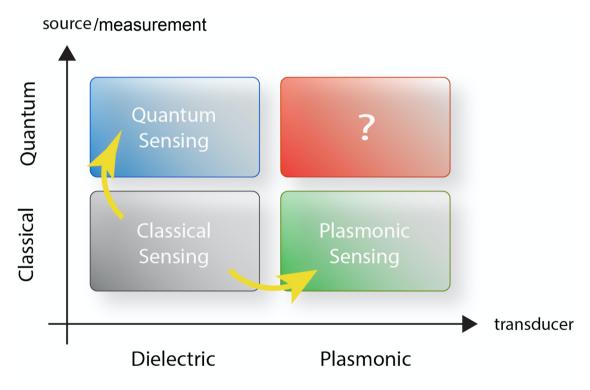




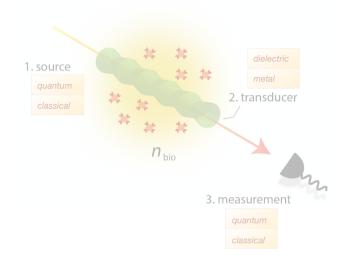


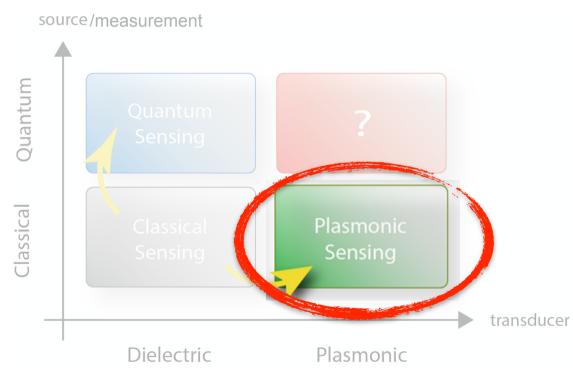








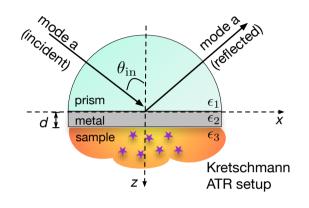


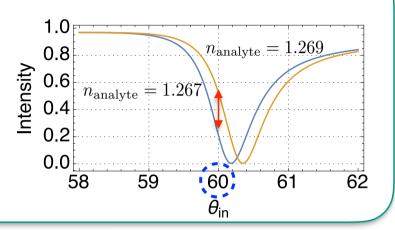


Plasmonic sensing

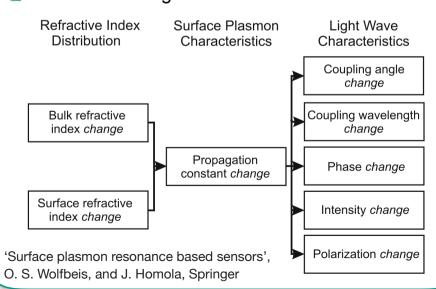


Example: intensity-sensitive ATR sensing





Various sensing mechanisms



Various figures of merit for sensing performance

Sensitivity
$$S = \frac{\delta Y}{S} \frac{\delta n_{ef}}{S}$$

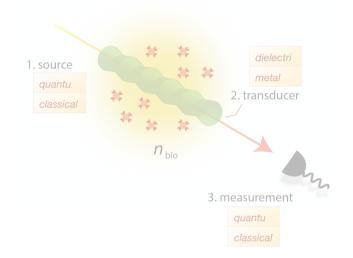
- Resolution $\delta n_{e\!f}$ δn
- Accuracy
- Precision
- ► etc..

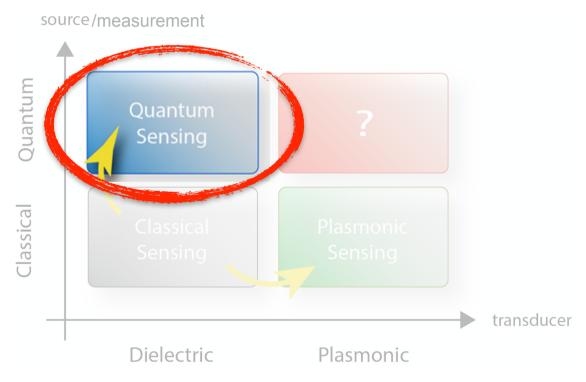
$$\sigma_{RI} = rac{\sigma_{so}}{S}$$
 standard deviation



- ✓ beyond the diffraction limit
- ✓ enhanced sensitivity





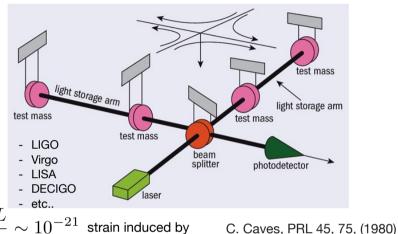


Quantum sensing (quantum metrology)

C. Caves, PRD 23, 1693 (1981)

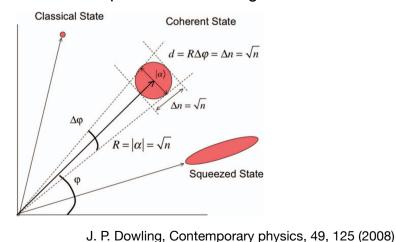


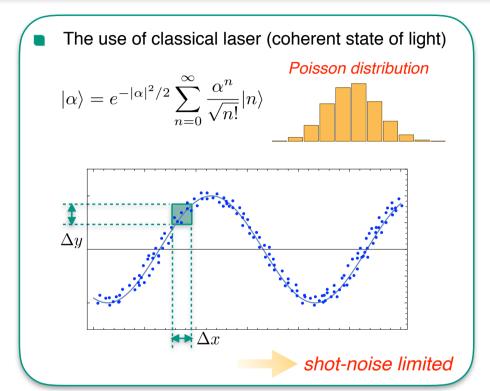
Original idea: Gravitational wave detector



The use of quantum state of light

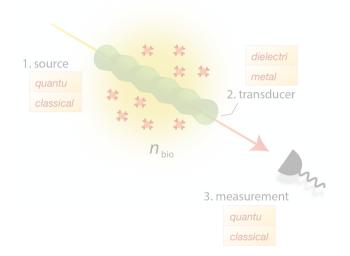
neutron star binary

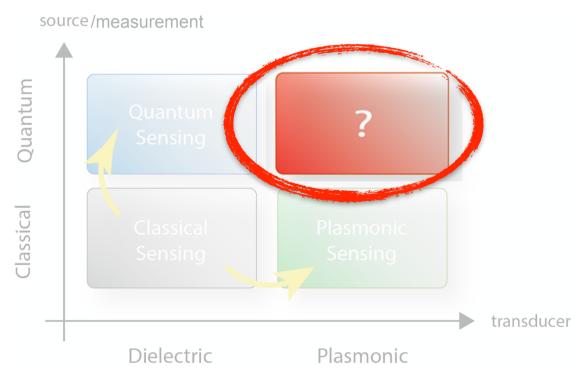




✓ beats the shot-noise limit

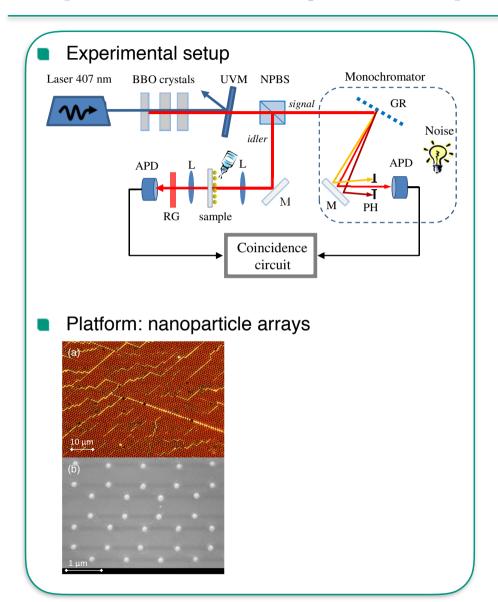


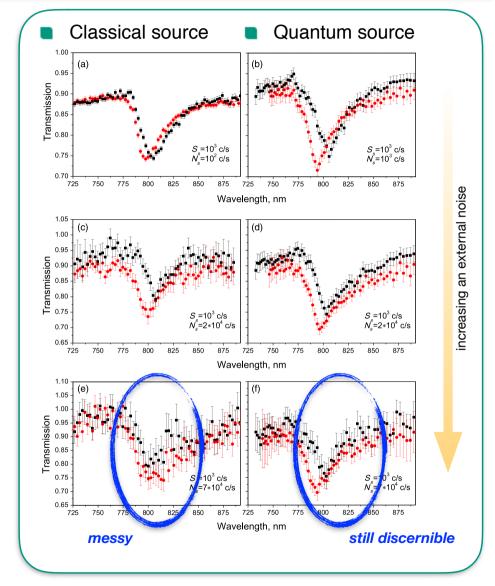




Experiment I on quantum plasmonic sensing



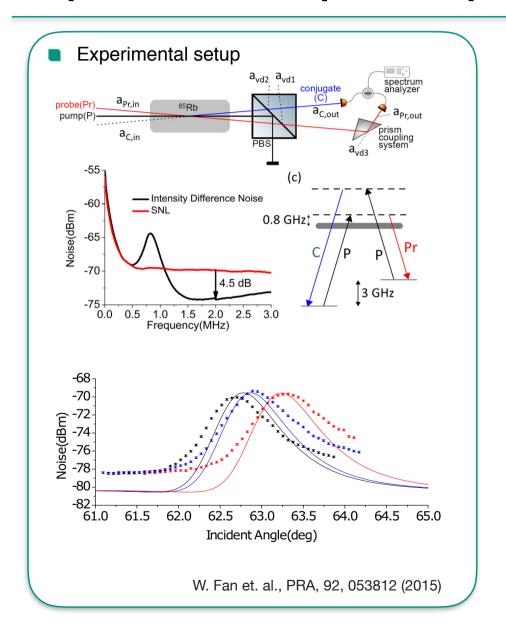


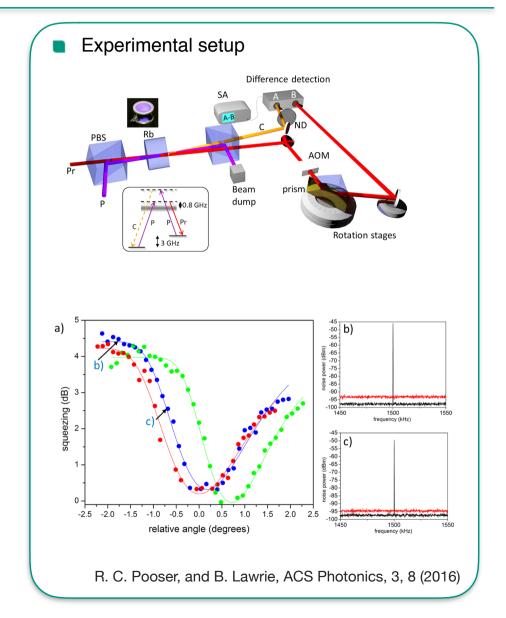


D. A. Kalashnikov et. al., PRX, 4, 011049 (2014)

Experiment II on quantum plasmonic sensing

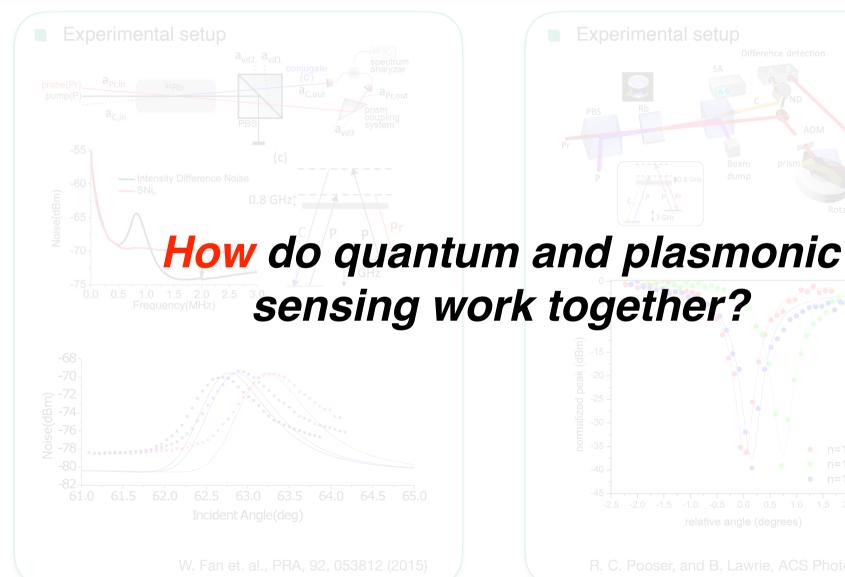


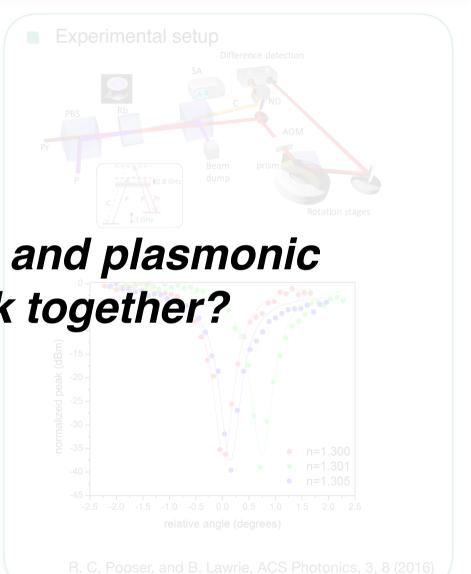




Experiment II on quantum plasmonic sensing

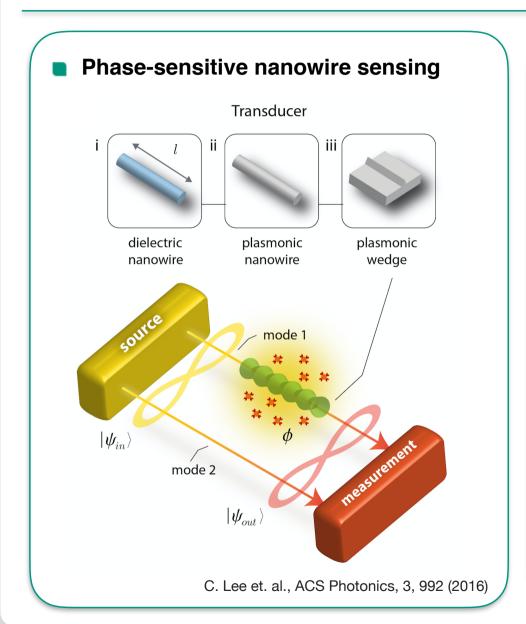


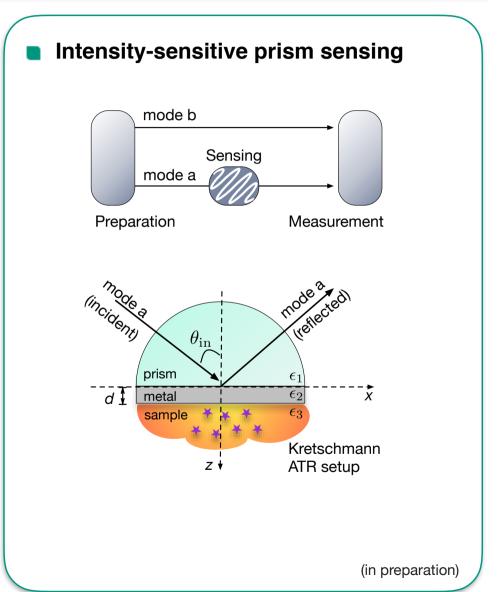




Quantum plasmonic sensing

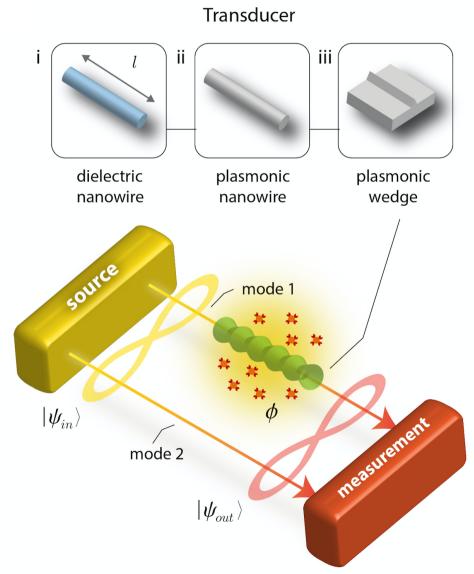






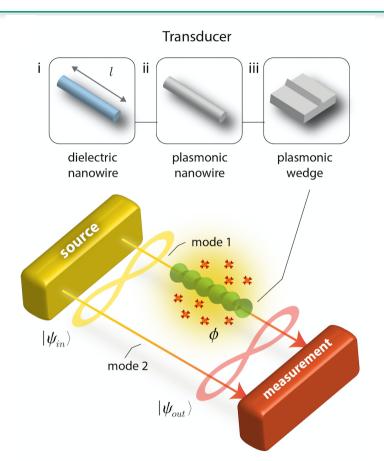
Two-mode interferometric nanowire sensing





Two-mode interferometric nanowire sensing





Classical reference (shot-noise limited)

$$|\alpha\rangle = e^{-|\alpha|^2/2} \sum_{m=0}^{\infty} \frac{\alpha^m}{\sqrt{m!}} |m\rangle$$

 $\hat{M} = \hat{I}_a - \hat{I}_b$

Quantum example (Heisenberg limited)

$$|\psi_{\rm in}\rangle = (|N0\rangle_{12} + |0N\rangle_{12})/\sqrt{2}$$
$$\hat{A} = |0, N\rangle\langle N, 0| + |N, 0\rangle\langle 0, N|$$

1.15

plasmonic (C)

1.10

1.20

1.25

 $n_{\rm bio}$

Monitored output: expectation values $\frac{\langle \hat{M} \rangle = M_0 \cos(\phi(n_{\rm bio}))}{\langle \hat{A} \rangle = A_0 \cos(N\phi(n_{\rm bio}))}$

1.30

1.35

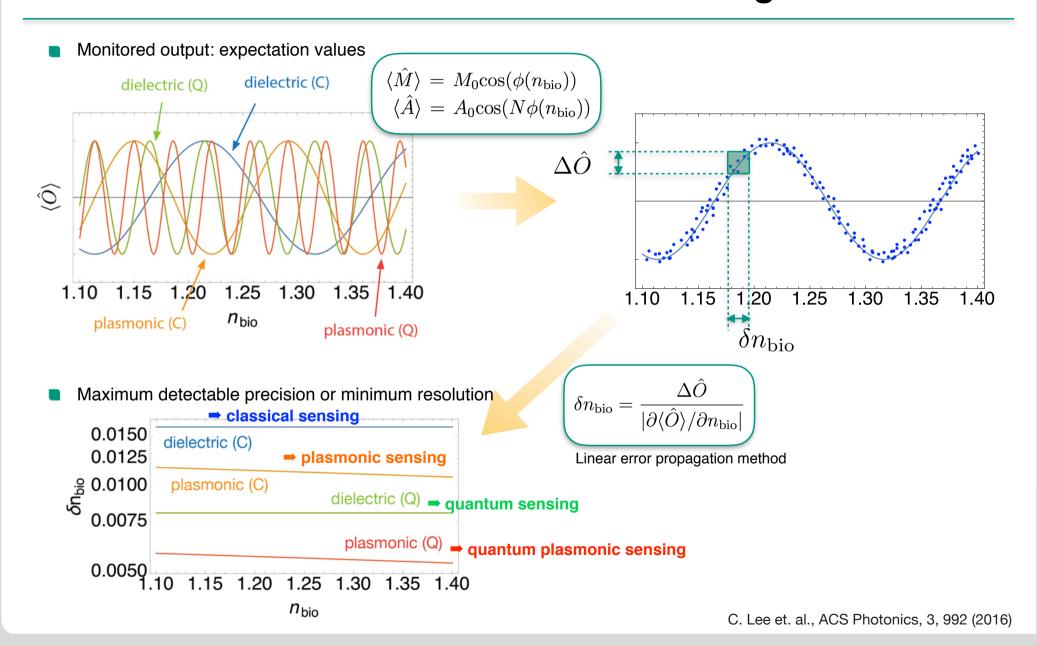
plasmonic (Q)

1.40

I

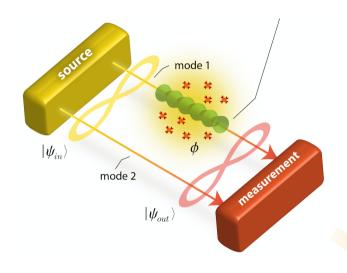
Carlsruhe Institute of Technology

Two-mode interferometric nanowire sensing



Roles of 'quantum' and 'plasmonic' sensing



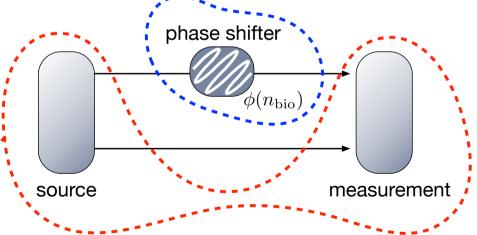


Plasmonic sensing

: how sensitively the transducer can induce the change in optical phase when an environment is altered.



: how sensitively the source and measurement can identify the change of optical phase when it occurs.



Roles of 'quantum' and 'plasmonic' sensing

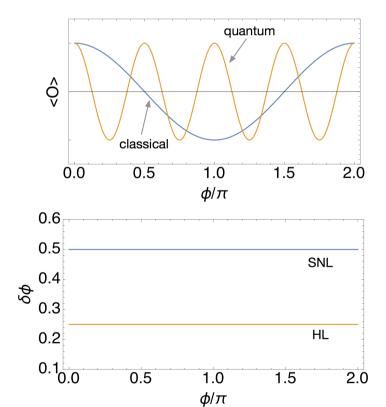


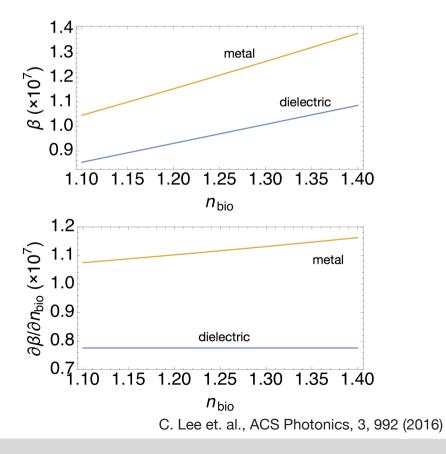
Quantum sensing

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Plasmonic sensing

: how sensitively the transducer can induce the change in optical phase when an environment is altered.





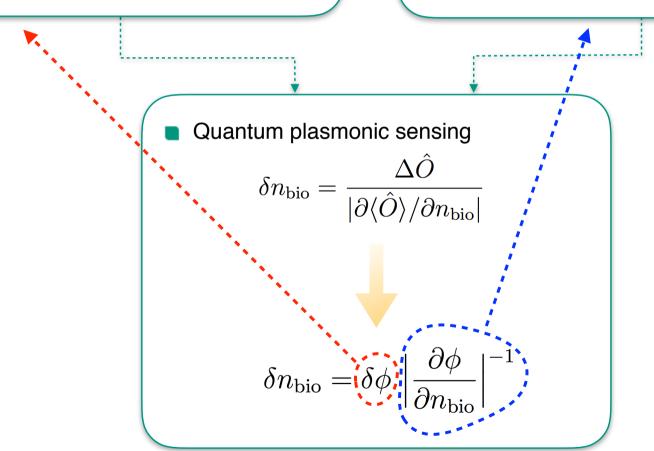
Roles of 'quantum' and 'plasmonic' sensing



- Quantum sensing
 - : how sensitively the source and measurement can identify the change of optical phase when it occurs.

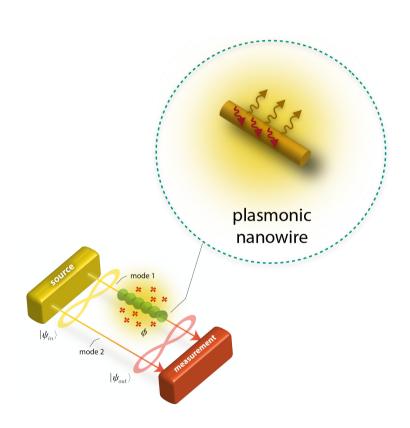
Plasmonic sensing

: how sensitively the transducer can induce the change in optical phase when an environment is altered.



Lossy sensing





Quantum example (Heisenberg limited)

$$|\psi_{\rm in}\rangle = (|N0\rangle_{12} + |0N\rangle_{12})/\sqrt{2}$$



General N-photon state

$$|\psi_{\rm in}\rangle = \sum_{n=0}^{N} c_n |n, N-n\rangle$$

Minimum resolution via Cramer-Rao bound

$$\delta n_{\rm bio} = \delta \phi \left| \frac{\partial \phi}{\partial n_{\rm bio}} \right|^{-1}$$

$$\delta\phi = F_Q^{-1/2}$$

U. Dorner, et al., PRL, 102, 040403 (2009)

Lossy sensing

0.5

0.4

0.3

0.2

0.1

0.0

 $|c_2|^2$

 $|c_0|^2$

1.15

1.20

1.25

 $n_{\rm bio}$

1.30

1.35 1.40



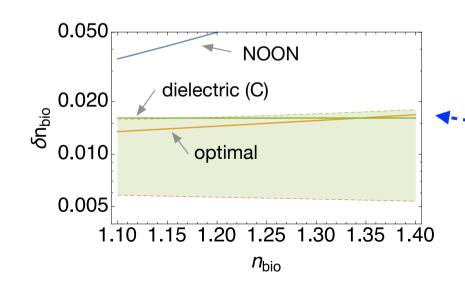
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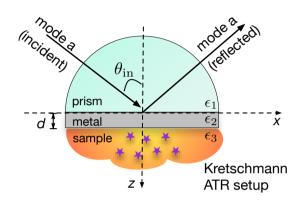
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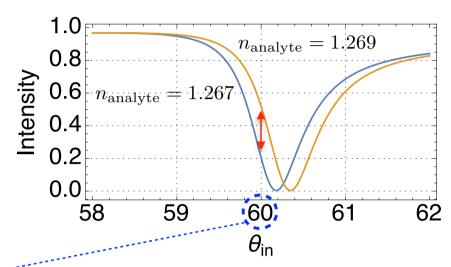
U. Dorner, et al., PRL, 102, 040403 (2009)

Single-mode intensity-sensitive ATR sensing

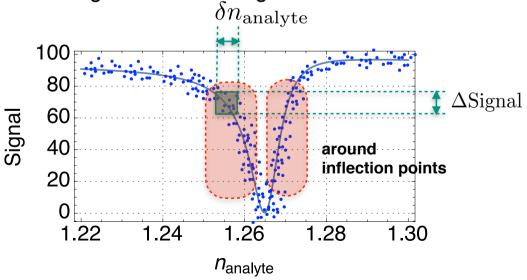


Intensity-sensitive ATR sensing





For a given incident angle



The average photon number of initial probe beam, as an example,

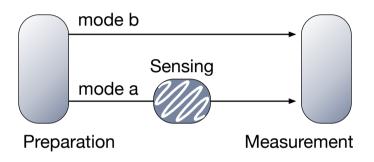
$$ext{Signal} = |r_{ ext{spp}}|^2 N \ \Delta ext{Signal} = \sqrt{|r_{ ext{spp}}|^2 N}$$

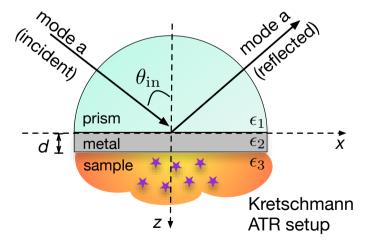
Parameter estimation

$$\delta n_{\rm analyte} = \frac{\Delta \text{Signal}}{\left| \frac{\partial \text{Signal}}{\partial n_{\rm analyte}} \right|}$$



Two-mode SPR sensing with differential intensity measurement





- Example states: twin modes
- * Classical state $|\alpha\rangle|\alpha\rangle$

$$N_a = N_b$$
$$\Delta N_a = \Delta N_b$$

* Two-mode squeezed vacuum state (SPDC)

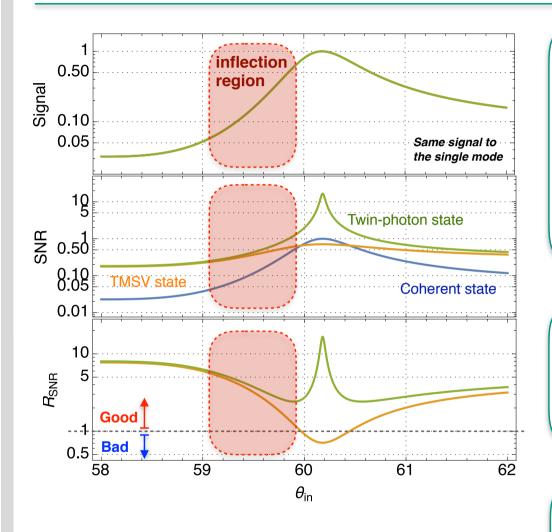
$$|\text{TMSV}\rangle = \frac{1}{\cosh(r)} \sum_{n=0}^{\infty} (-1)^n e^{in\phi} \tanh^n(r) |n, n\rangle$$

- * Twin photons $|n_a\rangle|n_b\rangle$
- Measurement: intensity difference

$$\hat{M} = \hat{b}^{\dagger} \hat{b} - \hat{a}^{\dagger} \hat{a}$$

- * two-mode correlation can be used.
- * the excess noise can be eliminated.





Example states: twin modes

$$N_a = N_b$$

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$$\hat{M} = \hat{b}^{\dagger} \hat{b} - \hat{a}^{\dagger} \hat{a}$$

- * two-mode correlation can be used.
- * the excess noise can be eliminated.

Average input power

$$N_a = N_b = 1$$

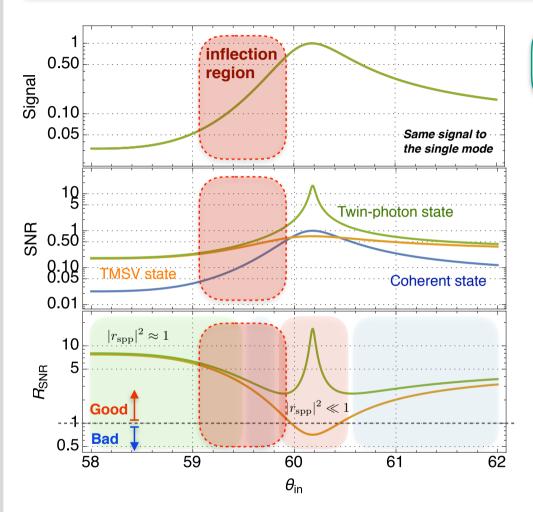
No channel losses at both modes $\eta_a = \eta_b = 1$

$$\eta_a = \eta_b = 1$$

Refractive index of sample

$$n_{\text{analyte}} = 1.267$$





$$R_{\rm SNR} = \frac{SNR}{SNR_c} = \sqrt{\frac{1+|r_{\rm spp}|^2}{(1-|r_{\rm spp}|^2)^2Q_{\rm M}+2\sigma|r_{\rm spp}|^2+(1-|r_{\rm spp}|^2)}}$$
 Mandel Q-factor
$$Q_{\rm M} = \frac{\left\langle \Delta N_a^2\right\rangle}{\left\langle N_a\right\rangle} - 1$$
 Degree of correlation
$$\sigma = \frac{\left\langle \Delta (N_a-N_b)^2\right\rangle}{\left\langle N_a\right\rangle+\left\langle N_b\right\rangle}$$

	$Q_{ m M}$	σ	R
Coherent state	0	1	1
TMSV state	>0	0	$\sqrt{\frac{1 + r_{\rm spp} ^2}{(1 - r_{\rm spp} ^2)^2 Q_{\rm M} + (1 - r_{\rm spp} ^2)}}$
Twin-photon state	-1	0	$\sqrt{rac{1 + r_{ m spp} ^2}{(1 - r_{ m spp} ^2) r_{ m spp} ^2}}$

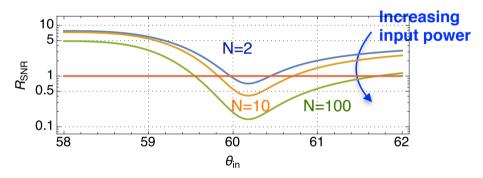


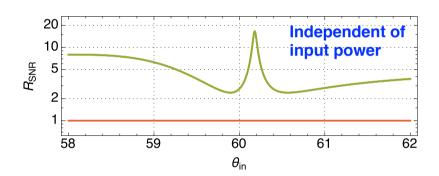
Two-mode squeezed state (SPDC)

$$|\text{TMSV}\rangle = \frac{1}{\cosh(r)} \sum_{n=0}^{\infty} (-1)^n e^{in\phi} \tanh^n(r) |n, n\rangle$$

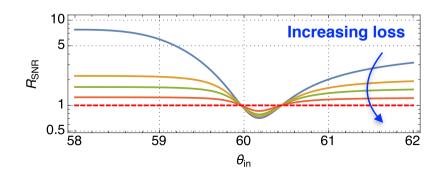
Twin photons $|n_a
angle |n_b
angle$

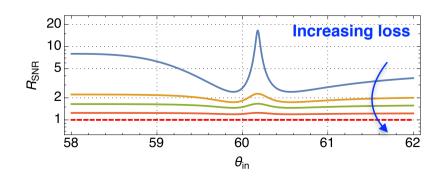
Effect of increasing the input photon number





Effect of increasing the channel losses



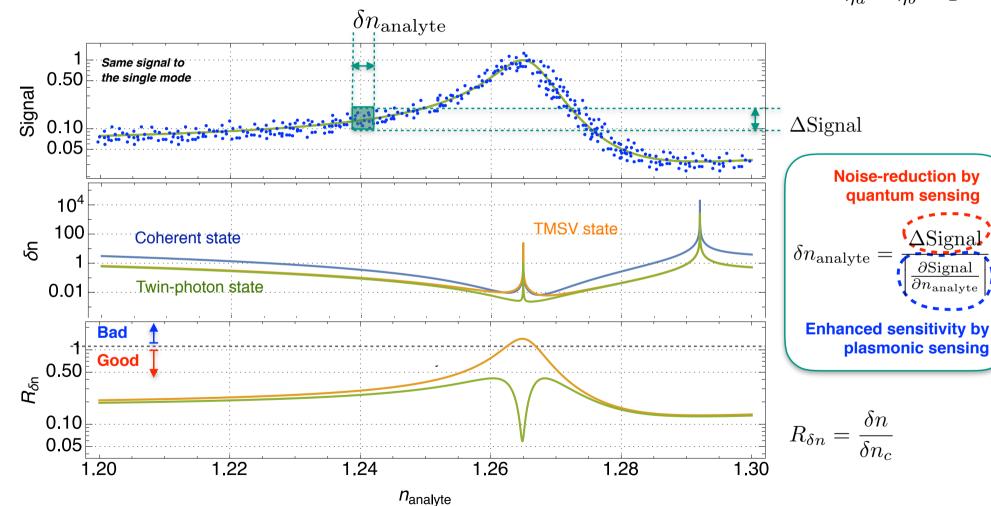


Estimation of the refractive index



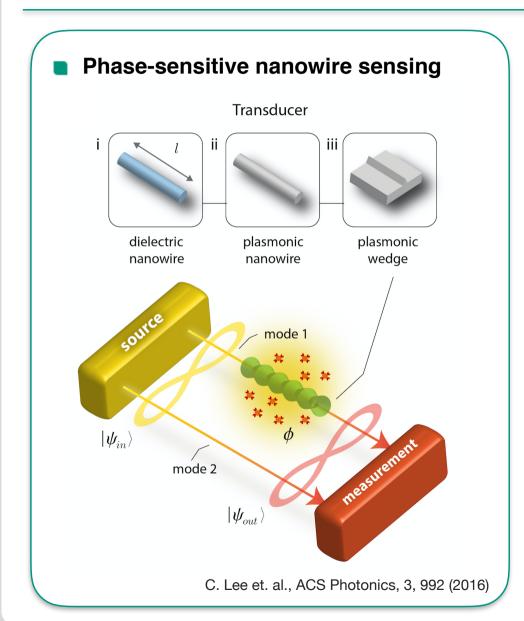
 \blacksquare for a given incident angle $\theta_{\rm in}=60$

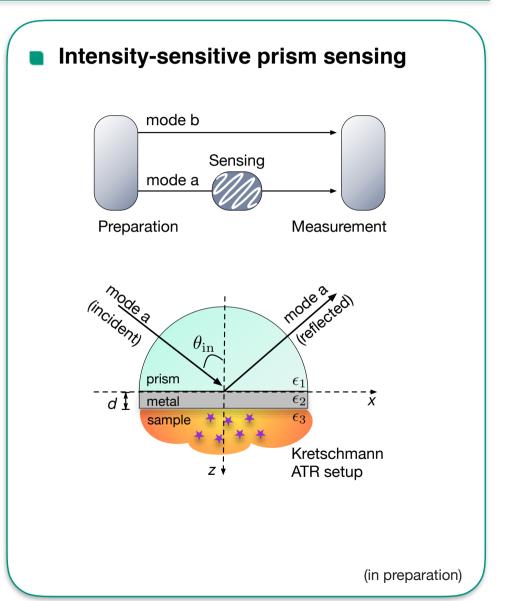
$$N_a = N_b = 1$$
$$\eta_a = \eta_b = 1$$



Conclusion

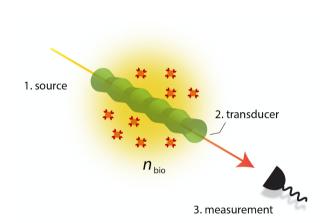


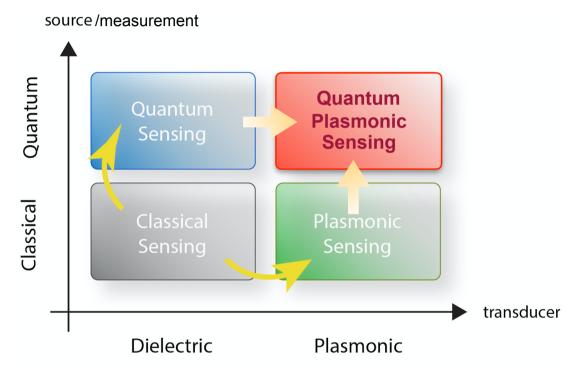




Conclusion









Thank you!

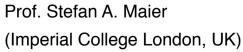
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Prof. Jinhyoung Lee (Hanyang University, South Korea)



Mr. Frederik Dieleman (Imperial College London, UK)