

Phase relaxation in a superfluid atomic ring

arXiv:1705.02650

Jérôme Beugnon

Benasque, May 2017

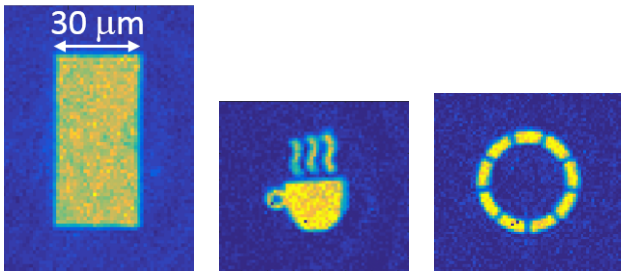
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UPMC-Sorbonnes Universités
Paris, France

Our approach to atomtronics

Confine Bose gases in a single plane and design arbitrary optical potentials

Our approach to atomtronics

Confine Bose gases in a single plane and design arbitrary optical potentials



- ▶ Custom-shaped potentials
- ▶ Time-dependent potentials
- ▶ Coupling between clouds
- ▶ Manipulate spin degrees of freedom

Phase relaxation in atomtronic devices

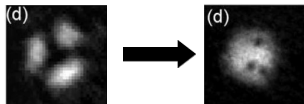
When coupling different subsystems, we usually consider the role of ΔT , $\Delta\mu$.

With superfluids, phase difference $\Delta\phi$ is crucial (Josephson physics).

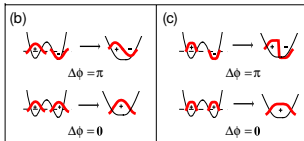
- ▶ What happens when connecting independent superfluids ?
- ▶ Does the phase homogenize ?
- ▶ How much time does it take to uniformize the phase ?
- ▶ What are the underlying microscopic mechanisms ?

Related works

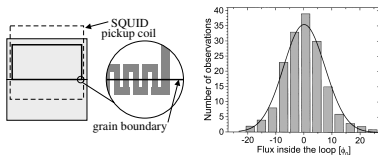
- ▶ Merging 3 BECs \rightarrow vortices (Scherer et al. PRL **98**, 110402 (2007))



- ▶ Merging 2 BECs \rightarrow heating (Jo et al. PRL **98** 180401 (2007))



- ▶ Connecting Josephson junctions (Carmi et al. PRL **84**, 4966 (2000))

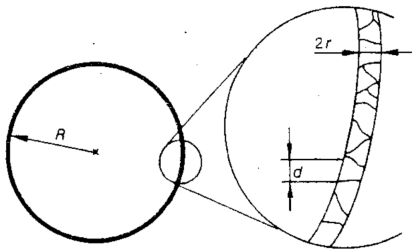


Link to Kibble-Zurek physics

Kibble-Zurek mechanism predicts the formation of topological defects when quench cooling a system across a phase transition.

Zurek's gedanken experiment:

Zurek Nature 307, 505 (1985)

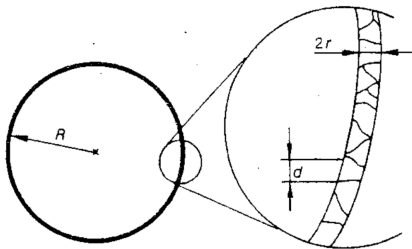


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Independent phase domains form because of the divergence of the thermalization time around T_c .

Domain size:

$$\hat{\xi} \propto \tau_Q^{\nu/(1+\nu z)}$$

ν, z : critical exponents

Winding number:

$$\langle n_w^2 \rangle \propto \hat{\xi}^\alpha$$

α depends on dimensionality.

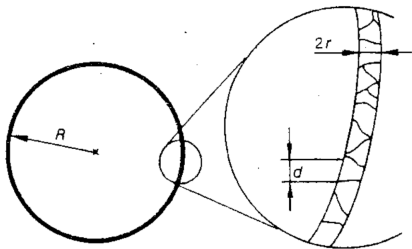
Corman et al. PRL 113, 135302 (2014)

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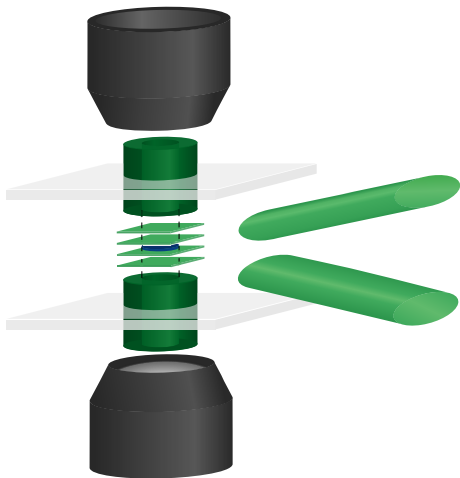
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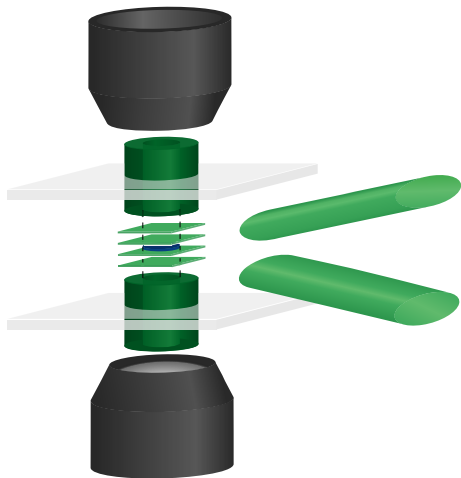
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Zurek assumes a \sqrt{N} -scaling of $\sqrt{\langle n_w^2 \rangle}$. We can test it !

A single layer of ultracold atoms



A single layer of ultracold atoms



- ▶ $\omega_z = 2\pi \times 1.6 \text{ kHz}$
($\Rightarrow \Delta z \approx 300 \text{ nm}$)
- ▶ $T \approx 10 \text{ nK}$
- ▶ $n_{2D} \approx 30 \mu\text{m}^{-2}$
- ▶ Atom number $\approx 10^5$
- ▶ Spatial light modulator (DMD) to shape the potential.
(10 kHz refresh rate)

Other physics to be explored

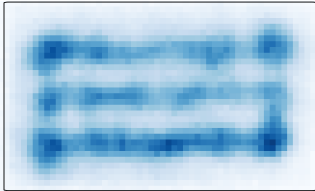
- ▶ Demixing dynamics in a two-component cloud



Non-miscible components evolve into well-defined patterns

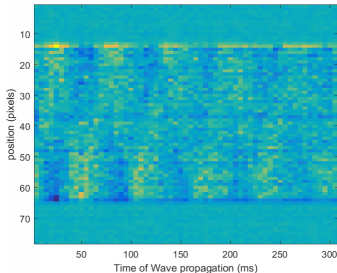
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- ▶ Demixing dynamics in a two-component cloud



Non-miscible components evolve into well-defined patterns

- ▶ Sound wave propagations

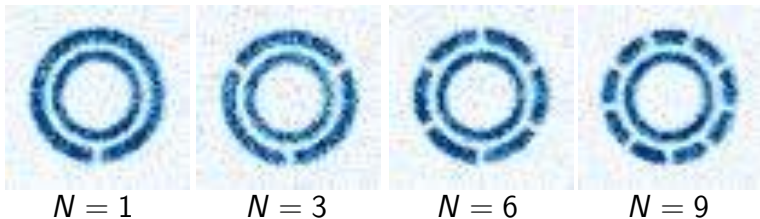


Observation of traveling waves in a box

Merging N condensates with independent phases in a ring geometry and monitor relaxation

Preparing independent BECs

We load a hot cloud in a N -segments configuration and perform evaporative cooling down to $T \approx 10$ nK.

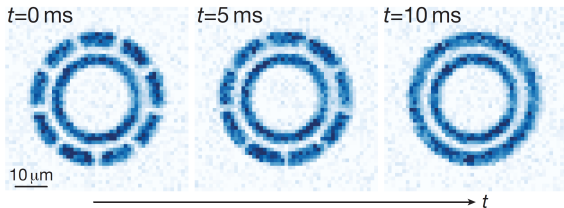


We checked that:

- ▶ The two rings have independent phases
- ▶ Our results are independent of the separation between domains in the range $2-3 \mu\text{m}$. (Here $2.5 \mu\text{m}$)

Merging

We merge the BECs by decreasing the barrier width in 10 ms.



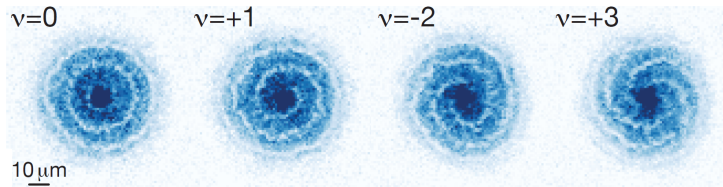
We empirically chose 10 ms because

- ▶ Faster merging leads to additional excitations.
- ▶ Slower merging leads to asynchronous merging.

We wait for 500 ms to let the phase homogenize.

Detection

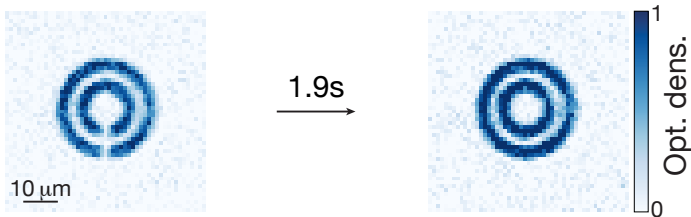
We perform matter-wave interference between the two rings



The number of spiral arms gives the winding number in the outer ring

Phase reference

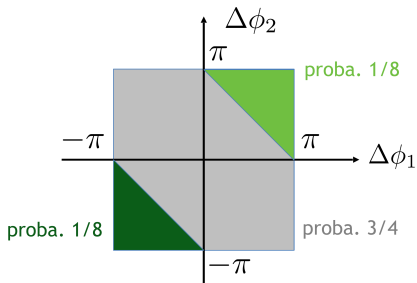
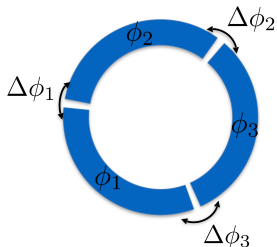
Check that the inner ring as no phase winding



- ▶ Cut the rings during evaporative cooling
- ▶ Close very slowly the rings

We detected 0 spiral pattern over 159 shots.

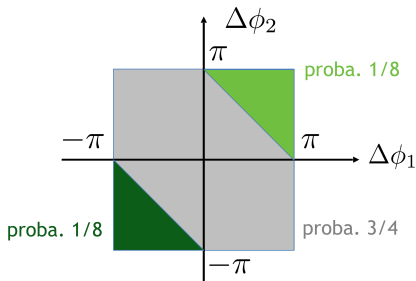
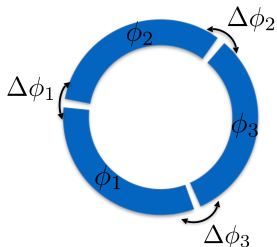
The 3-segments case



$$\Delta\phi_1, \Delta\phi_2 \in (-\pi, \pi]$$

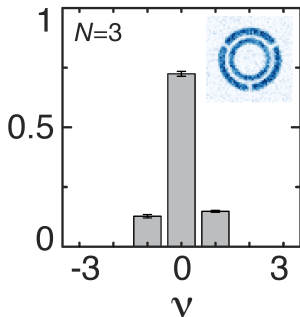
- ▶ $\nu = +1$ if $\Delta\phi_1 + \Delta\phi_2 > \pi$
- ▶ $\nu = -1$ if $\Delta\phi_1 + \Delta\phi_2 < -\pi$
- ▶ $\nu = 0$ otherwise

The 3-segments case

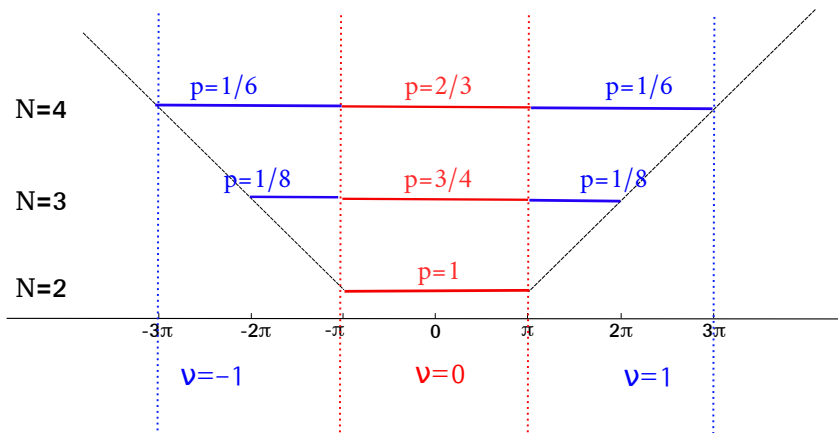


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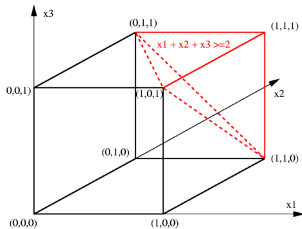


Varying the number of segments



Varying the number of segments

Compute hypervolumes of
an hypercube



Already done by mathematicians !
Euler-Frobenius distribution gives:

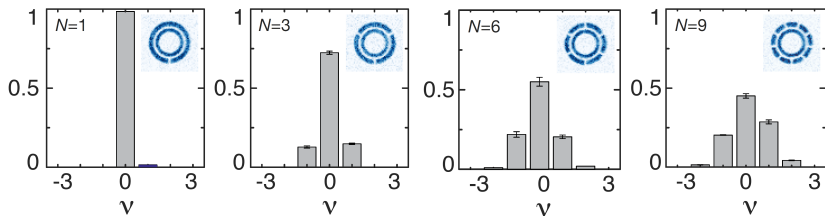
$$\nu_{\text{rms}} = \frac{\sqrt{N}}{2\sqrt{3}} \text{ for } N \geq 3$$

with ν_{rms} the rms width of the
distribution.

S. Jansen, Online J. Anal. Comb. **8** (2013)

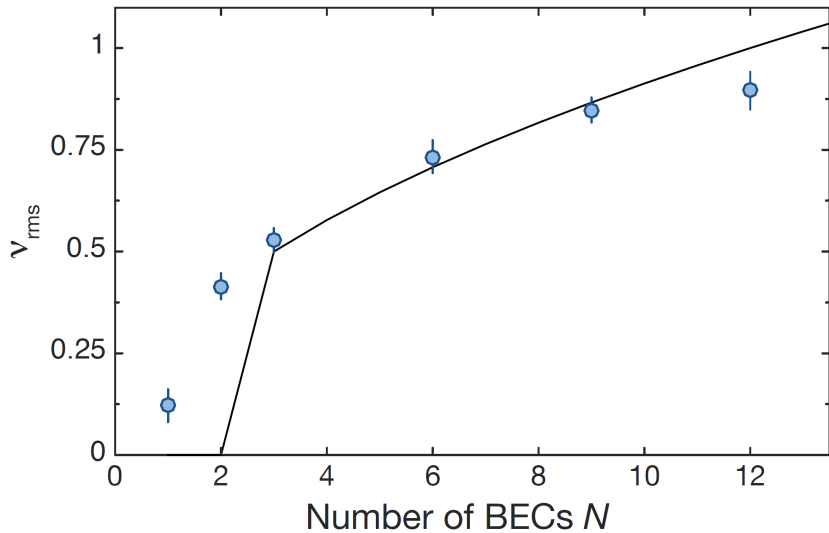
Results N segments

Vary number of segments N from 1 to 12.

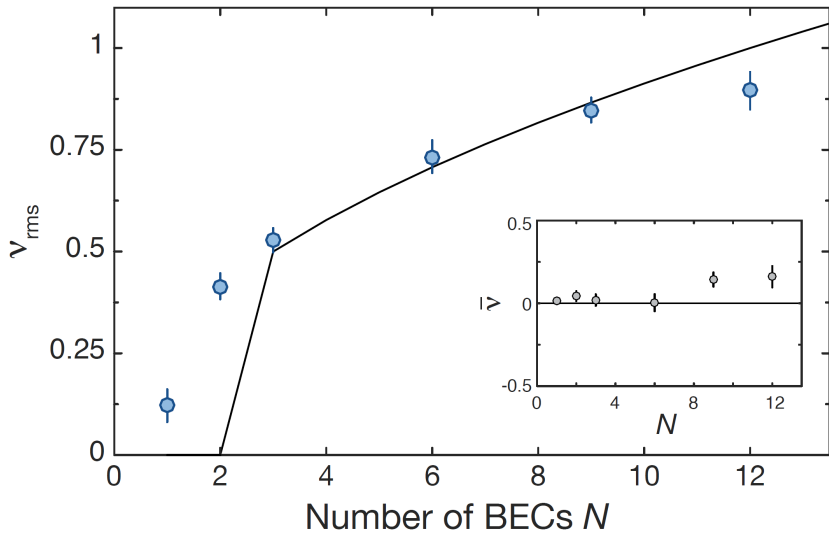


The distribution broadens for increasing N

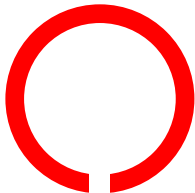
Results for N segments



Results for N segments



Special cases $N = 1$ and $N = 2$

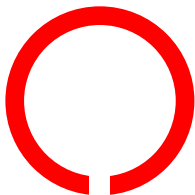


No phase winding expected at zero temperature.

We find :

$$P(0) = 98\% \quad P(\pm 1) = 2\%$$

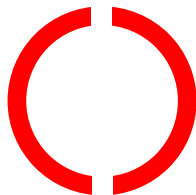
Special cases $N = 1$ and $N = 2$



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Marginal situation at zero temperature

Phase winding if $\Delta\phi = \pi$

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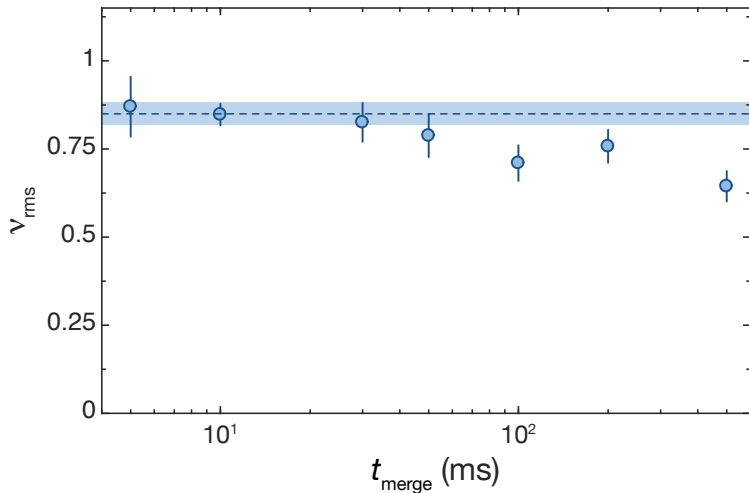
$$P(0) = 84\% \quad P(\pm 1) = 8\%$$

Timescales

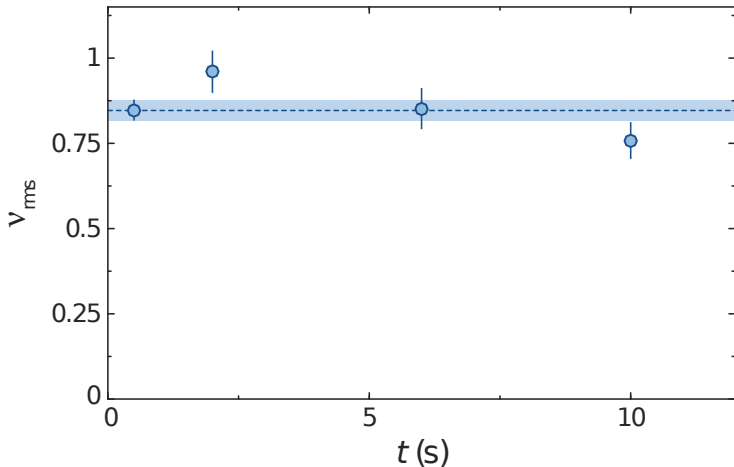
Merging < Sound round-trip < Relaxation < Lifetime

10 ms < 100 ms < 500 ms < 20 s

Influence of the merging time

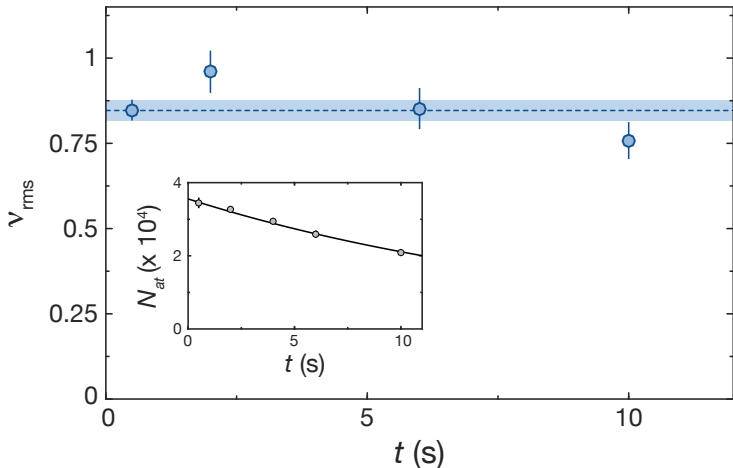


Lifetime



Superfluid current lifetime is very long (>10 s) and larger than the atomic lifetime.

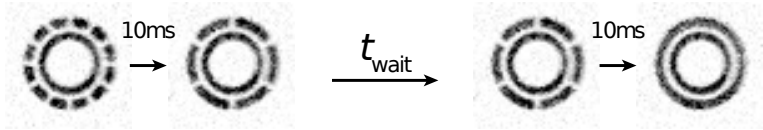
Lifetime



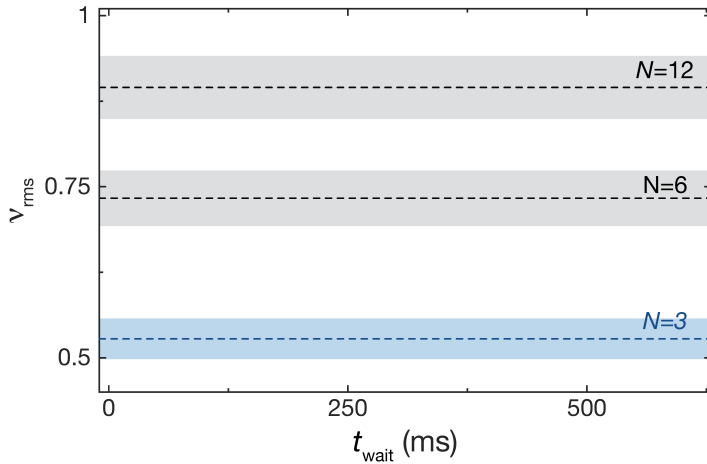
Superfluid current lifetime is very long (>10 s) and larger than the atomic lifetime.

Clear illustration of topological protection !

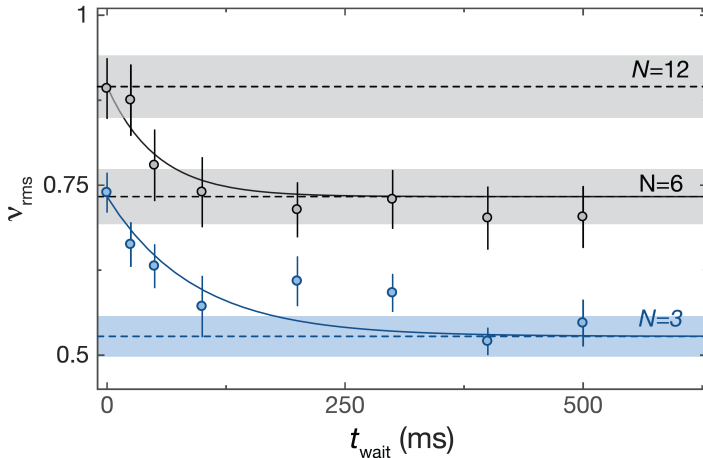
Two-step merging



Two-step merging

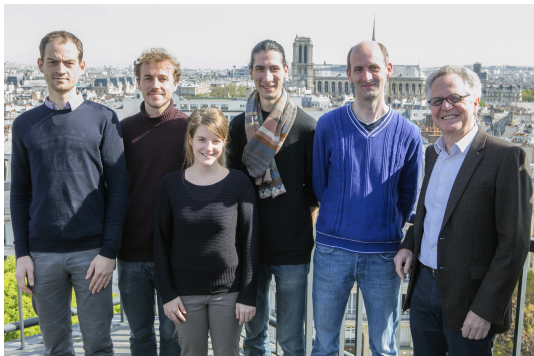


Two-step merging



Exponential fits : $\tau_{12} = 52(17)$ ms, $\tau_6 = 90(30)$ ms
Shorter segments homogenize faster. Microscopic mechanism ?

The team



From left to right:

Sylvain Nascimbene

Jean-Loup Ville (PhD)

Monika Aidelsburger
(Postdoc)

Raphaël Saint-Jalm(PhD)

Jérôme Beugnon

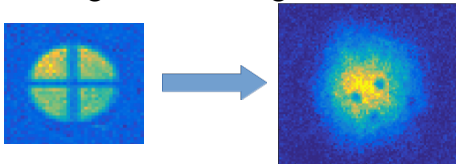
Jean Dalibard

Funding:



Perspectives

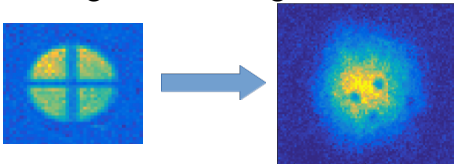
- ▶ Investigate different geometries



- ▶ Control the relative phase
- ▶ Observe microscopic dynamics

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! Post-doc position available !