



Spectroscopy from experimentalist's viewpoint

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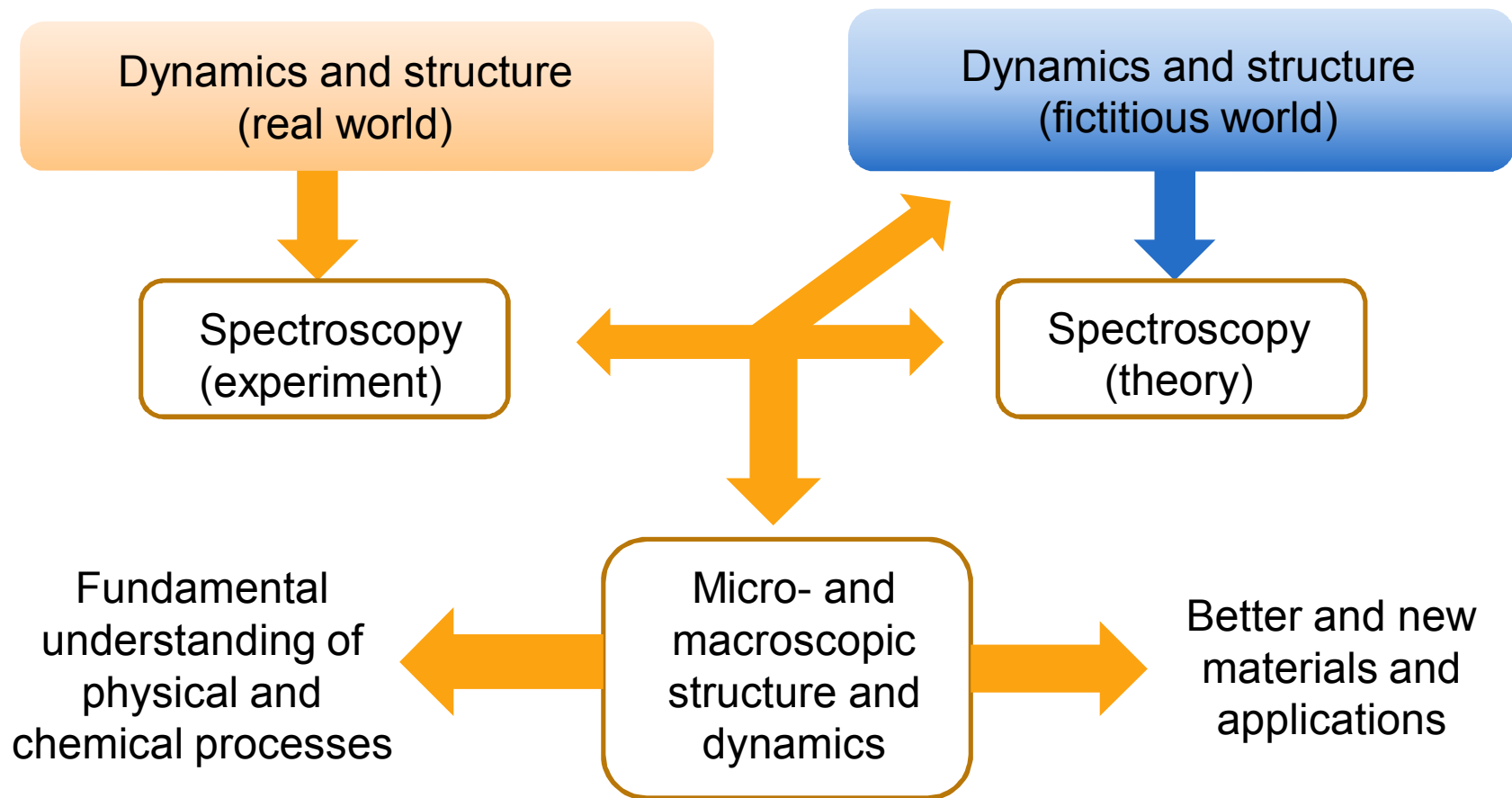
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TDDFT school, Benasque, Spain, January 2014



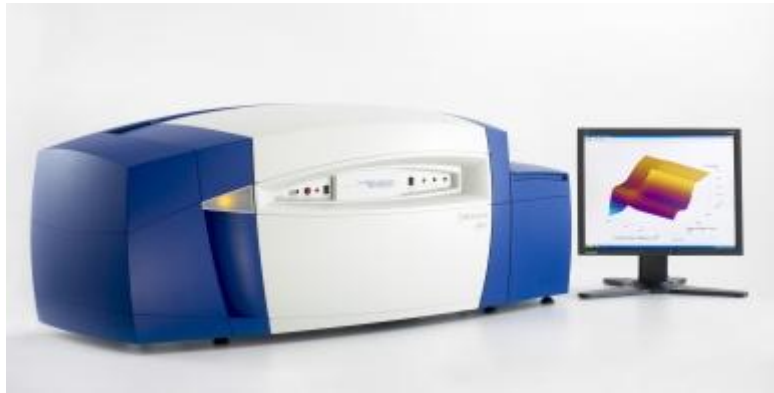
Motivation: why theory



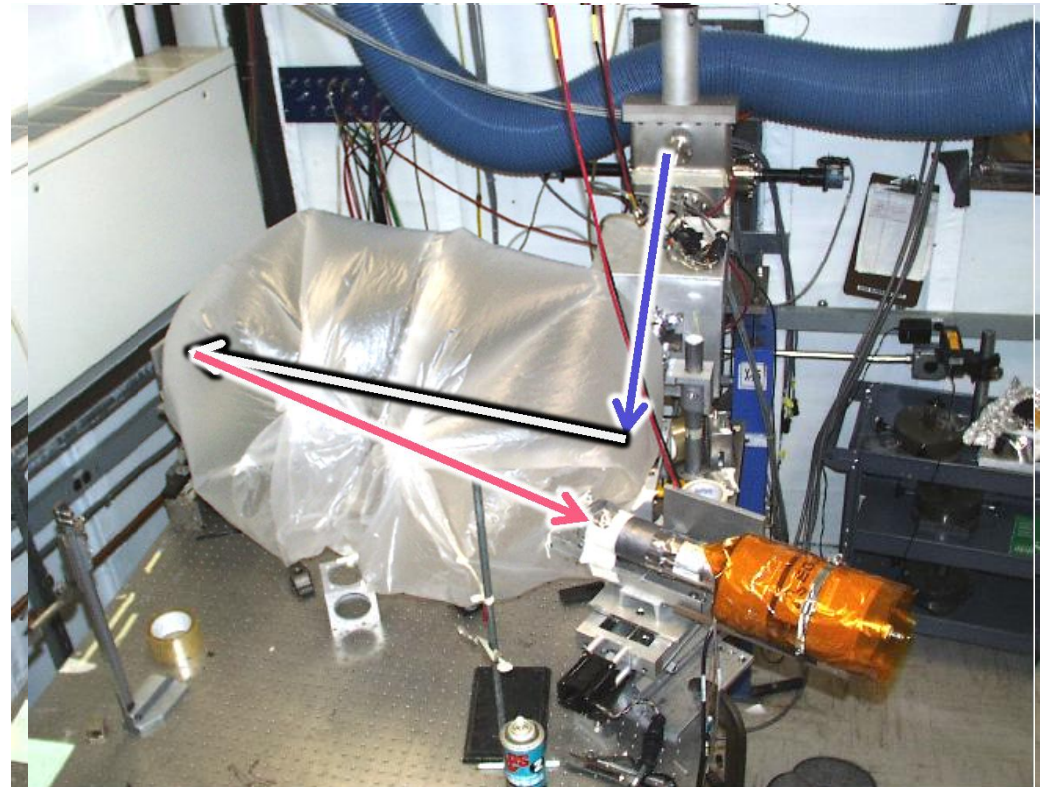


Measurement vs. experiment

Equipment for measurements

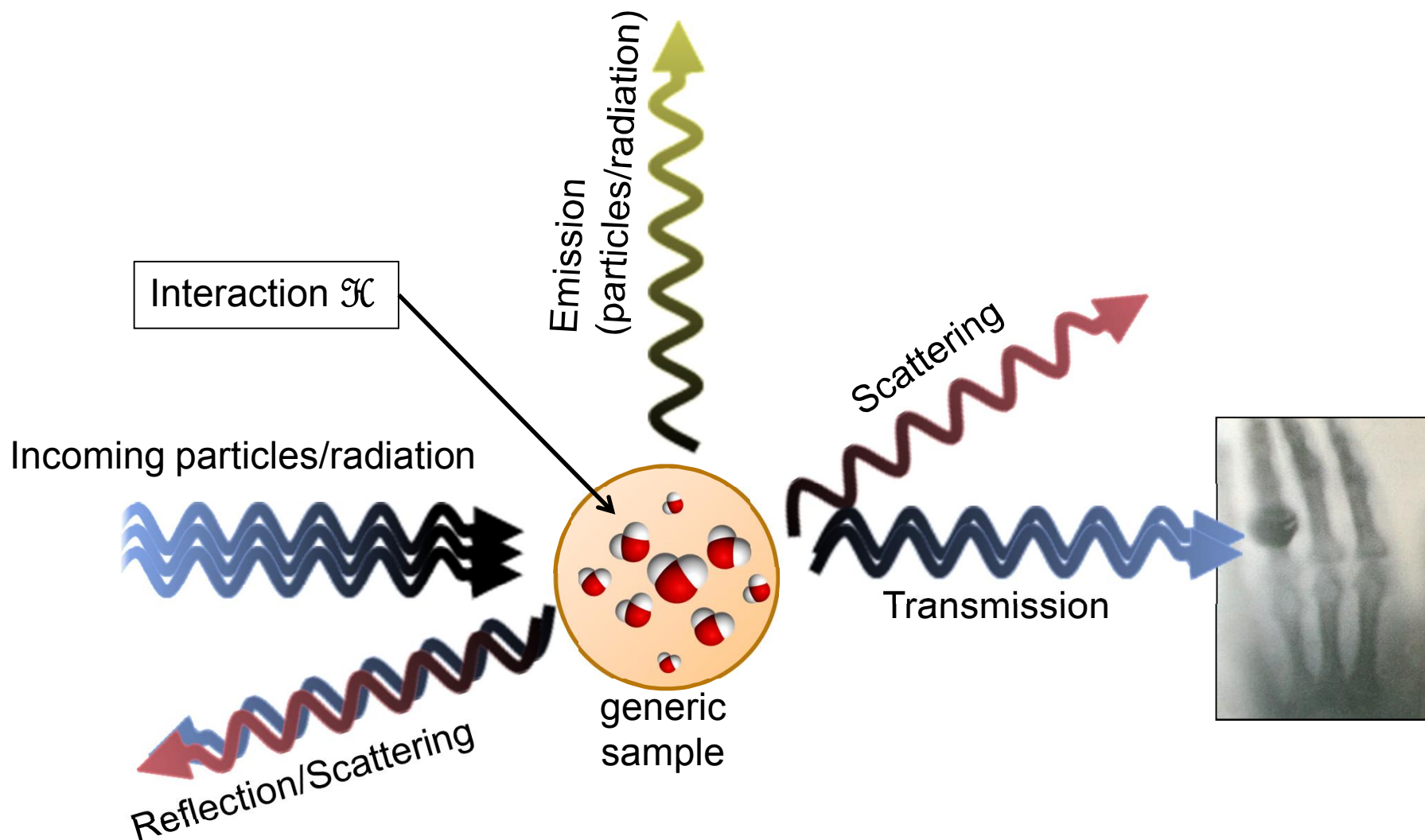


Equipment for experiments





Probes and messengers





Properties of particles

Polarization: linear, circular

- photons
- electrons (how?)
- linear and circular dichroism become possible

Probing depth

- surface vs. bulk sensitivity
- e.g., what is surface enhanced Raman scattering, SERS?

Element or site selectivity

- what methods can selectively probe properties of selected atom in a large system?

Access of dynamic ranges

- energy, momentum, time, space
- resolving power / applicable range



(non-exhaustive) Classification of excitations

	Collective	"Single particle"
Vibrational	Sound waves (phonons)	Molecular bending, stretching...
Spin	Magnons, spin waves	Spin flip
Charge	Plasmons, polarons	Excitons, momentum density, electron removal
Orbital	Orbiton	Crystal-field excitation

- + multiples (bimagnons, double plasmons, ...)
- + coupling (important!)



Electron response to electromagnetic field

Macroscopic **dielectric function** $\epsilon_M(Q, E) = \epsilon_1(Q, E) + i \epsilon_2(Q, E)$

Complex refractive index $\tilde{n} = \sqrt{\epsilon} = n + i\kappa$

Reflectivity $R = \frac{(n-1)^2 + \kappa^2}{(n+1)^2 + \kappa^2}$

Reflectivity measurements

Absorption coefficient $\alpha = 4\pi\kappa/\lambda$

UV/Vis absorption,
x-ray absorption (XAS), ...

Loss function $-\text{Im}[\epsilon_M^{-1}(\mathbf{Q}, E)]$

Electron energy loss,
inelastic x-ray scattering

Dynamic structure factor $S(\mathbf{Q}, E) = -\left(\frac{Q^2}{4\pi^2 e^2 n}\right) \text{Im}[\epsilon_M^{-1}(\mathbf{Q}, E)]$

Spectral density function $A(k, \omega)$

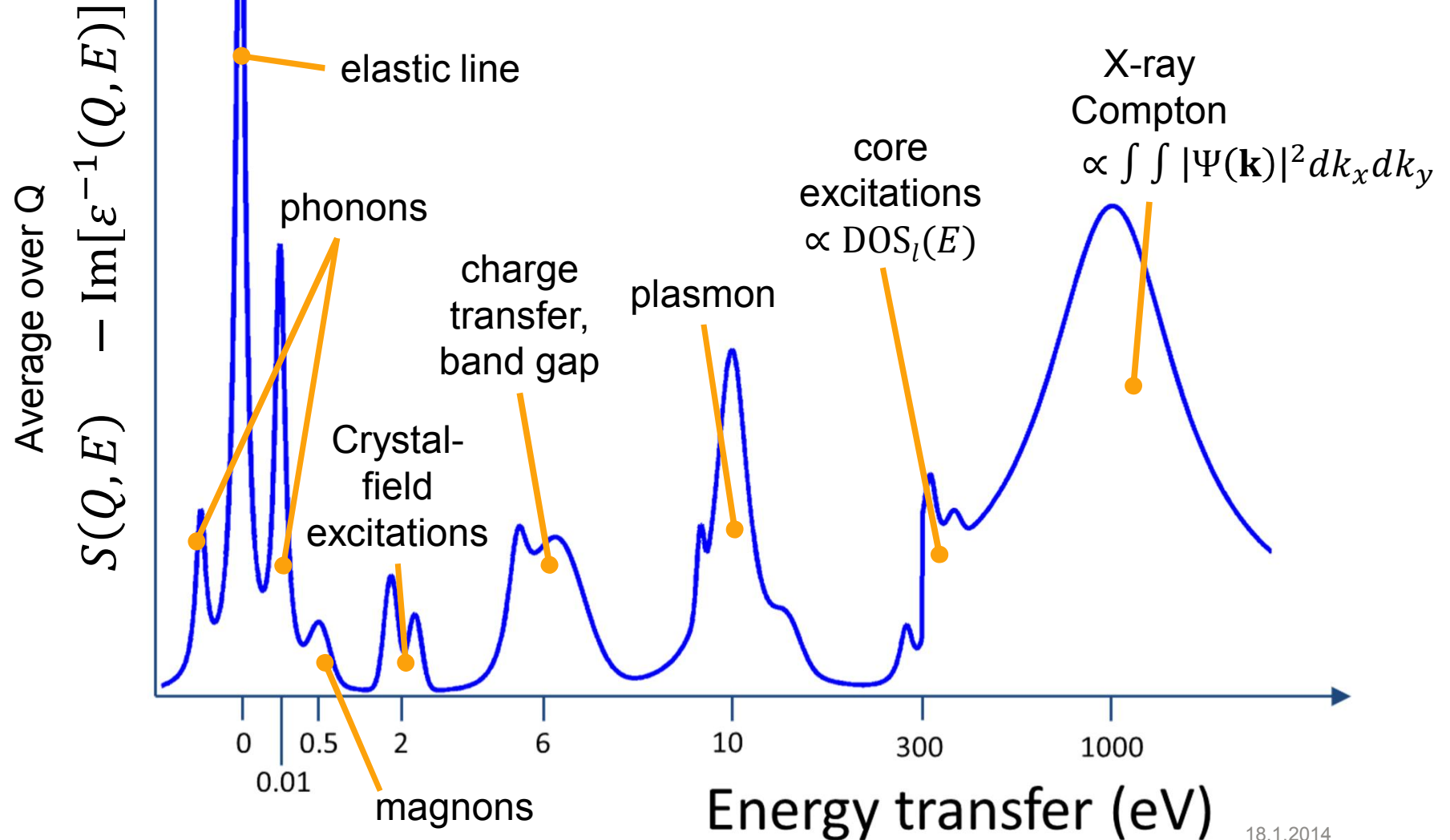
Resonant Raman spectra

Dynamic structure factor is also the Fourier transform of the **density correlation function**





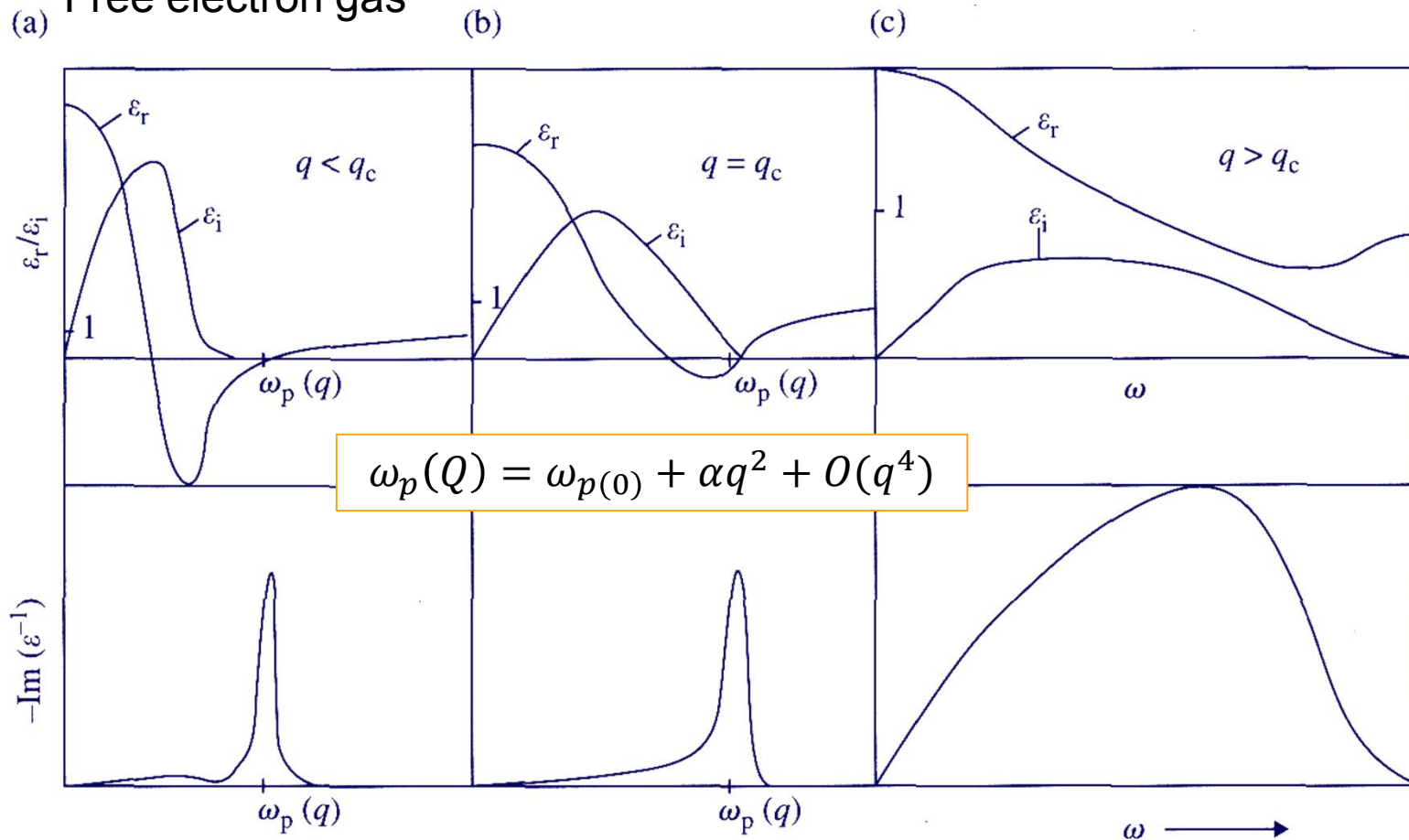
Energy-loss spectrum by inelastic x-ray scattering





Plasmons

Dielectric function from the Lindhard polarization function (RPA)
Free electron gas



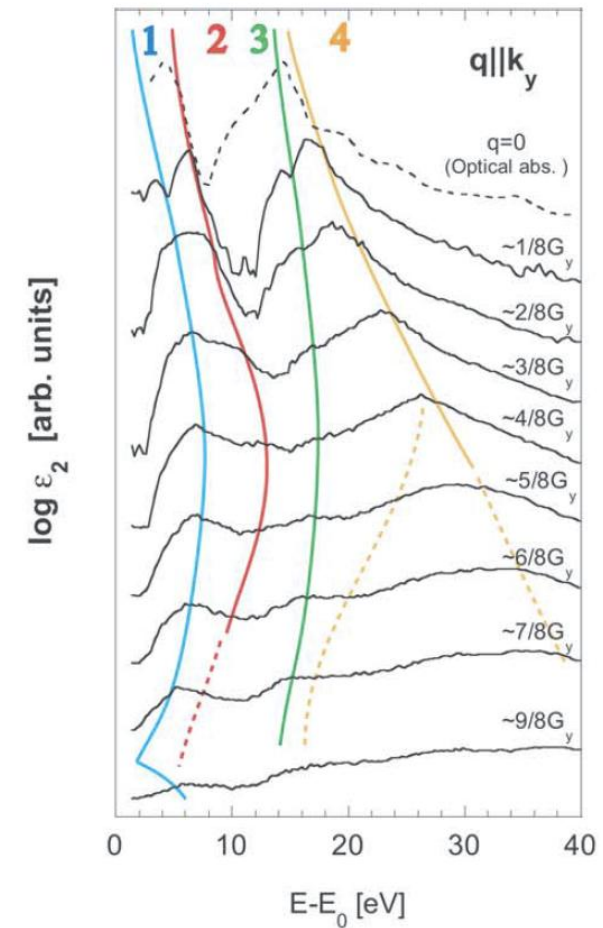
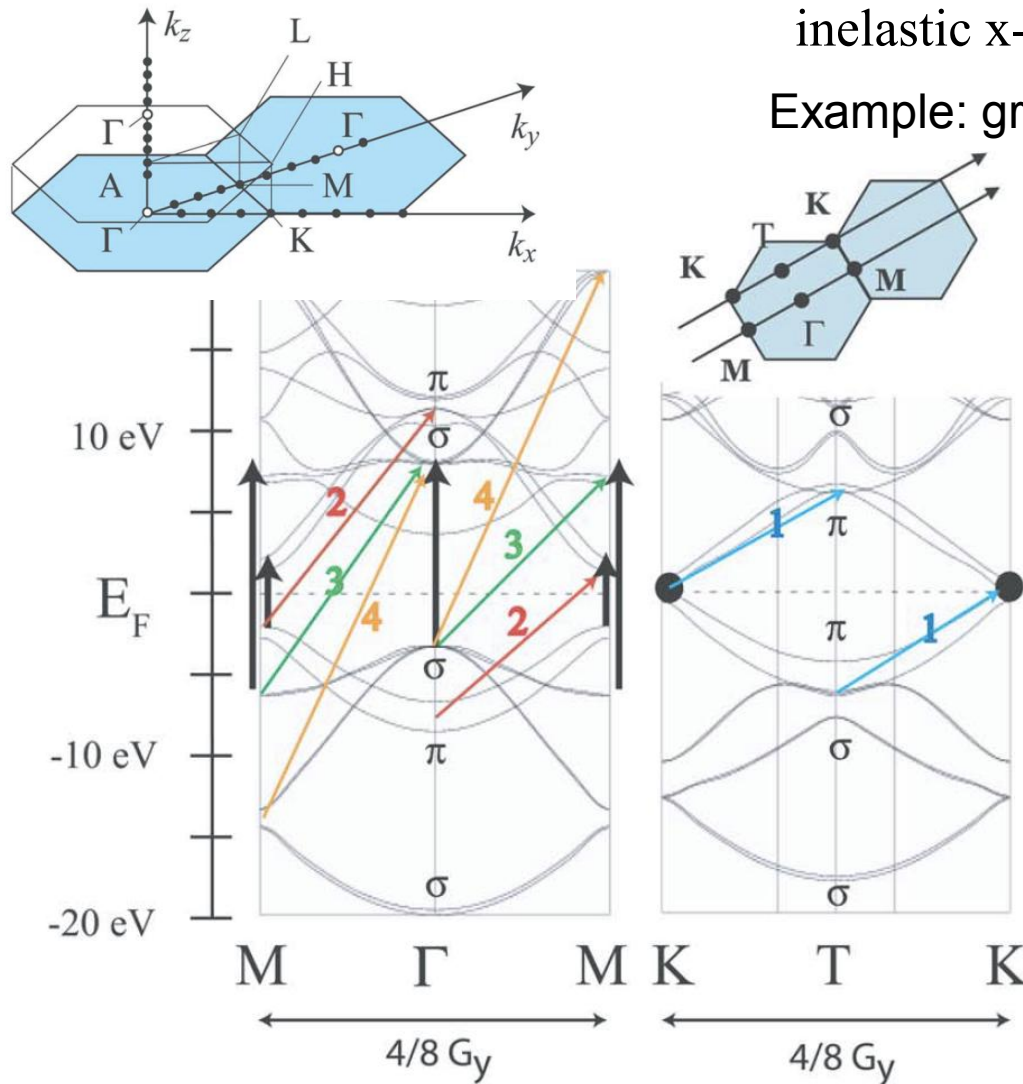
W. Schülke: Electron dynamics by inelastic x-ray scattering (Oxford Univ. Press)



Inter- and intraband excitations

Both *energy* and *momentum* can be controlled in inelastic x-ray scattering spectroscopy

Example: graphite





Outgoing particle

Incoming particle

	Nothing	Photon	Electron
Nothing	Skiing at the Pyrenees	Light emission, photoluminescence	Electron emission
Photon	Absorption (IR, UV/vis, vacuum UV, x-ray..) CD; XMCD	Ellipsometry, reflectometry, Raman, Brillouin, x-ray scattering, Compton scattering	Photoemission Auger spectroscopy
Electron	Electron absorption	Inverse photoemission Cathodo-luminescence	Electron energy loss Tunneling

- + Time-of-flights
- + Non-linear phenomena

Differences in:

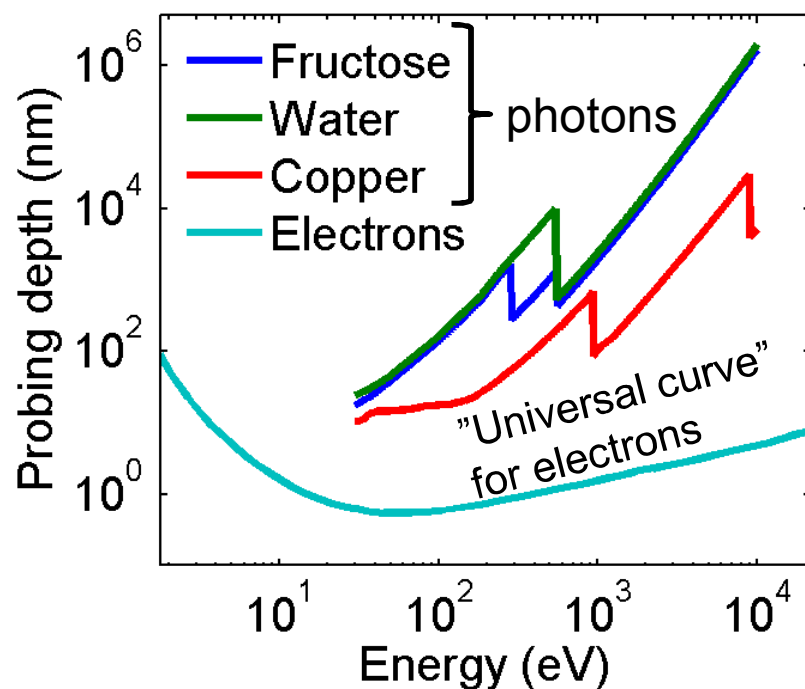
- Dynamic range of Q and E, coupling to spin, nuclei, ...
- Bulk or surface sensitivity
- Resolving power in energy, momentum, time, space...
- Element specificity (useful for complicated systems)



Particle probing depth

Large probability for interaction \Rightarrow higher surface sensitivity

Dependence on atomic number \Rightarrow elemental sensitivity



Question: why it is useful to do transmission microscopy of organic matter using photon energies of 400 eV?

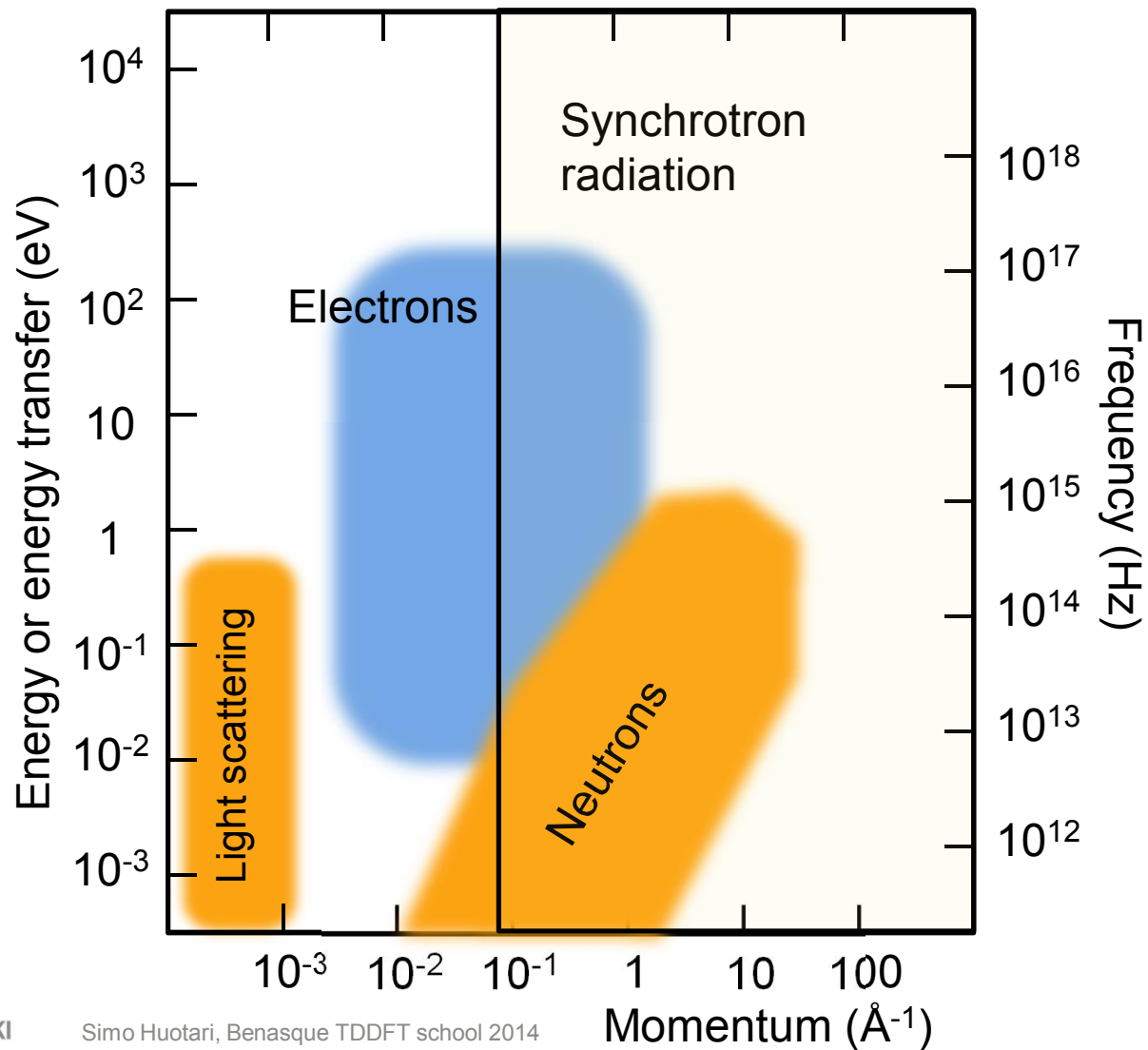
Advanced question: Which of the techniques are sensitive to bulk and which are sensitive to surface?



Dynamic ranges for scattering

Characteristic length scale (Å) $\zeta = 2\pi/Q$

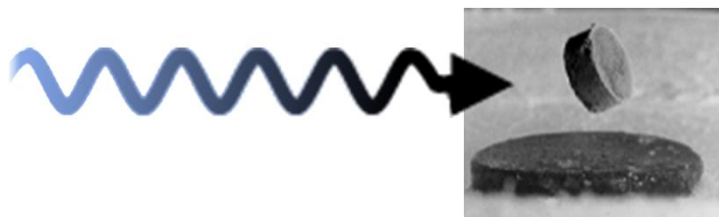
10^4 10^3 10^2 10 1 0.1





Spectroscopy - interaction

Photon-electron $H = \frac{e^2}{2mc^2} \mathbf{A}^2 + \frac{e}{mc} \mathbf{A} \cdot \mathbf{p}$



Scattering

Absorption and emission in first order, resonant scattering in second order

Electron-electron $V = e^2 / |\mathbf{r} - \mathbf{r}_i|$

Transition rate from state $|A\rangle$ to state $|B\rangle$ is given by the Fermi's Golden Rule:

$$W_{BA} = \frac{2\pi}{\hbar} \left| \langle B | H | A \rangle + \sum_I \frac{\langle B | H | I \rangle \langle I | H | A \rangle}{E_A - E_I + i\Gamma/2} \right|^2 \delta(E_B - E_A) + O(H^3)$$



Kramers-Heisenberg formula

From $\mathbf{A} \cdot \mathbf{A}$ - Non-resonant scattering (Raman, inelastic x-ray,)

$$\frac{d^2\sigma}{d\Omega d\omega_2} = r_0^2 \left(\frac{\omega_2}{\omega_1} \right) \left| \langle \Psi_f | e^{i\mathbf{Q} \cdot \mathbf{r}} | \Psi_i \rangle (\boldsymbol{\epsilon}_1 \cdot \boldsymbol{\epsilon}_2) \right.$$

$$+ \frac{1}{m} \sum_n \left[\frac{\langle \Psi_f | \boldsymbol{\epsilon}_2 \cdot \mathbf{p} e^{-i\mathbf{k}_2 \cdot \mathbf{r}} | \Psi_n \rangle \langle \Psi_n | \boldsymbol{\epsilon}_1 \cdot \mathbf{p} e^{i\mathbf{k}_1 \cdot \mathbf{r}} | \Psi_i \rangle}{E_i - E_n + \omega_1 + i\Gamma_n} \right.$$

$$\left. + \frac{\langle \Psi_f | \boldsymbol{\epsilon}_1 \cdot \mathbf{p} e^{i\mathbf{k}_1 \cdot \mathbf{r}} | \Psi_n \rangle \langle \Psi_n | \boldsymbol{\epsilon}_2 \cdot \mathbf{p} e^{-i\mathbf{k}_2 \cdot \mathbf{r}} | \Psi_i \rangle}{E_i - E_n - \omega_2} \right] \left| \delta(E_i - E_f + \omega) \right|^2$$

From $\mathbf{p} \cdot \mathbf{A}$ - Resonant scattering (Resonant Raman, resonant inelastic x-ray scattering)



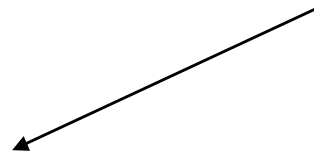
Dynamic structure factor

Average density-density correlation function

$$G(\mathbf{r}, t) = \frac{1}{N} \int \langle \rho(\mathbf{r}' - \mathbf{r}, t) \rho(\mathbf{r}', 0) \rangle d\mathbf{r}'$$

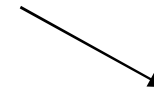
Dynamic structure factor

$$S(\mathbf{Q}, E) = \int \int G(\mathbf{r}, t) e^{i\mathbf{Q}\cdot\mathbf{r}} e^{-iEt}$$



Small Q

interference effects are important
(collective excitations)



Large Q

independent particle excitations



Absorption and scattering

Absorption $P(\omega) \propto \sum_f |\langle \Psi_f | \mathbf{p} \cdot \mathbf{A} | \Psi_i \rangle|^2 \delta(E_i - E_f + \omega)$ **Dipole rule**

Scattering $S(\mathbf{Q}, \omega) \propto \sum_f |\langle \Psi_f | e^{i\mathbf{Q}\cdot\mathbf{r}} | \Psi_i \rangle|^2 \delta(E_i - E_f + \omega)$

What happens with the momentum transfer dependence of the scattering matrix element?

$$e^{i\mathbf{Q}\cdot\mathbf{r}} = 1 + i\mathbf{Q} \cdot \mathbf{r} - (\mathbf{Q} \cdot \mathbf{r})^2/2 + \dots$$



dipole

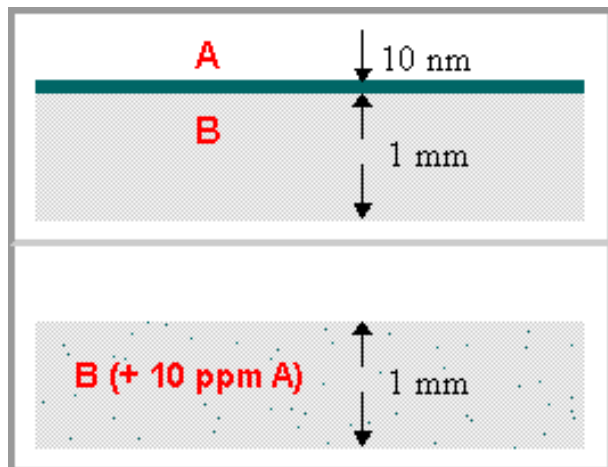


higher order multipoles

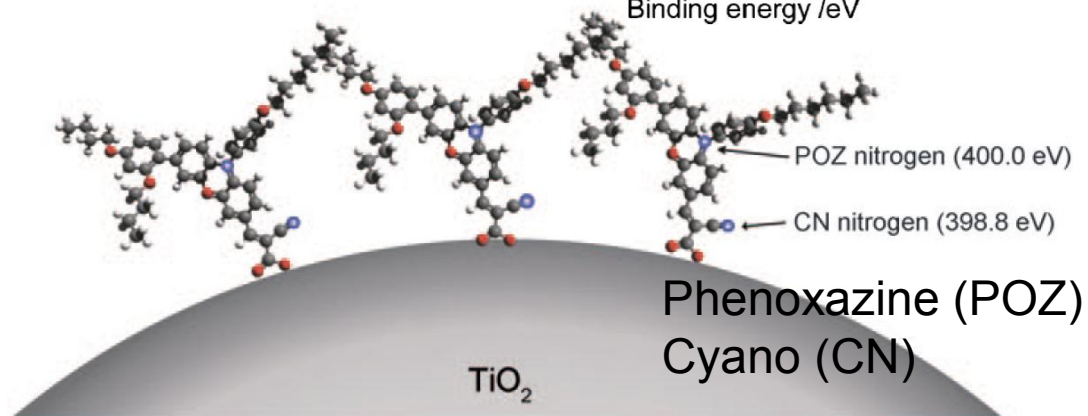
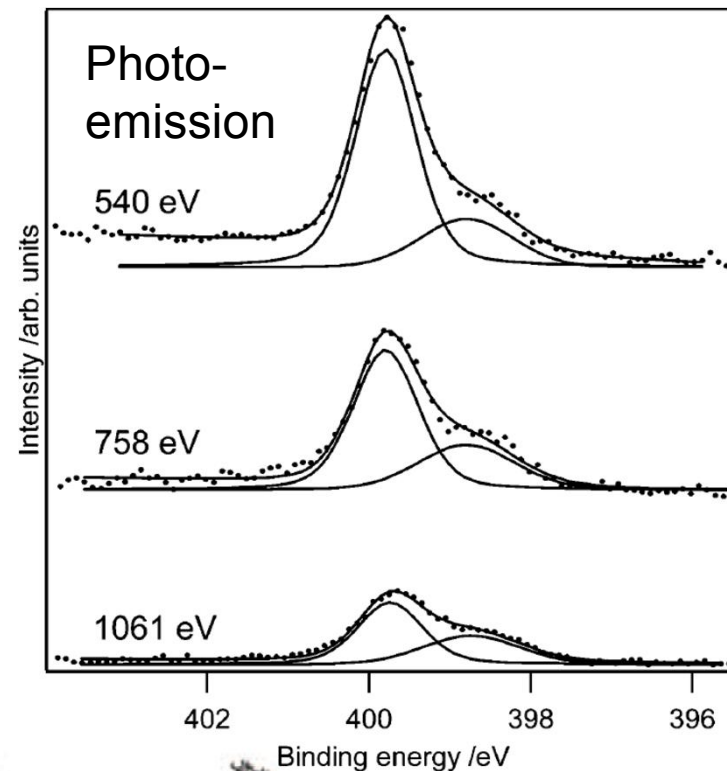
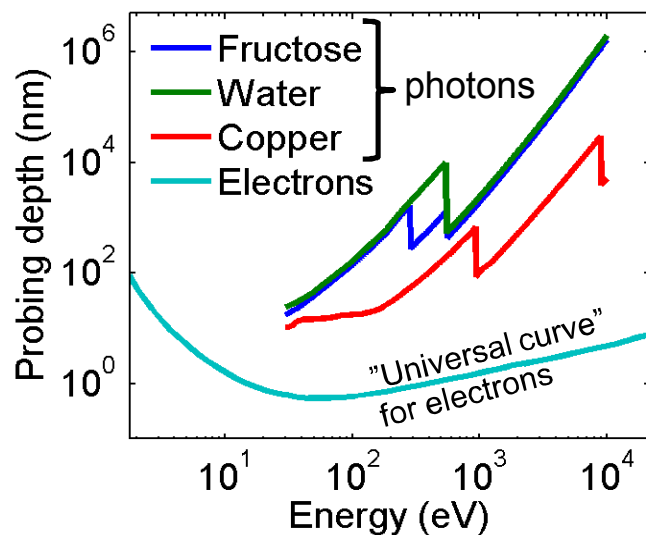


Surface sensitivity: molecular orientation

E.M.J. Johansson et al.
J. Phys. Chem. C 111, 8580 (2007)



http://www.chem.qmul.ac.uk/surfaces/scc/scat5_1.htm





Modern 3rd generation synchrotrons

Advanced Photon Source, USA



Super Photon Ring 8 (Spring-8), Japan

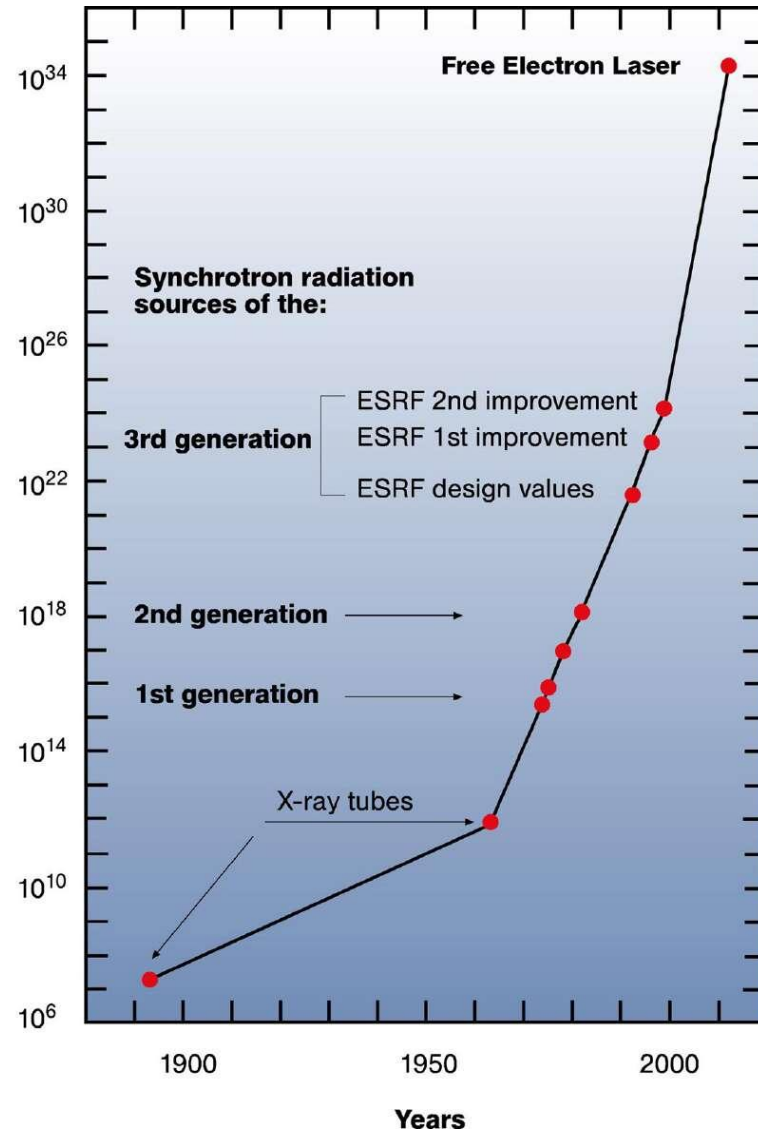
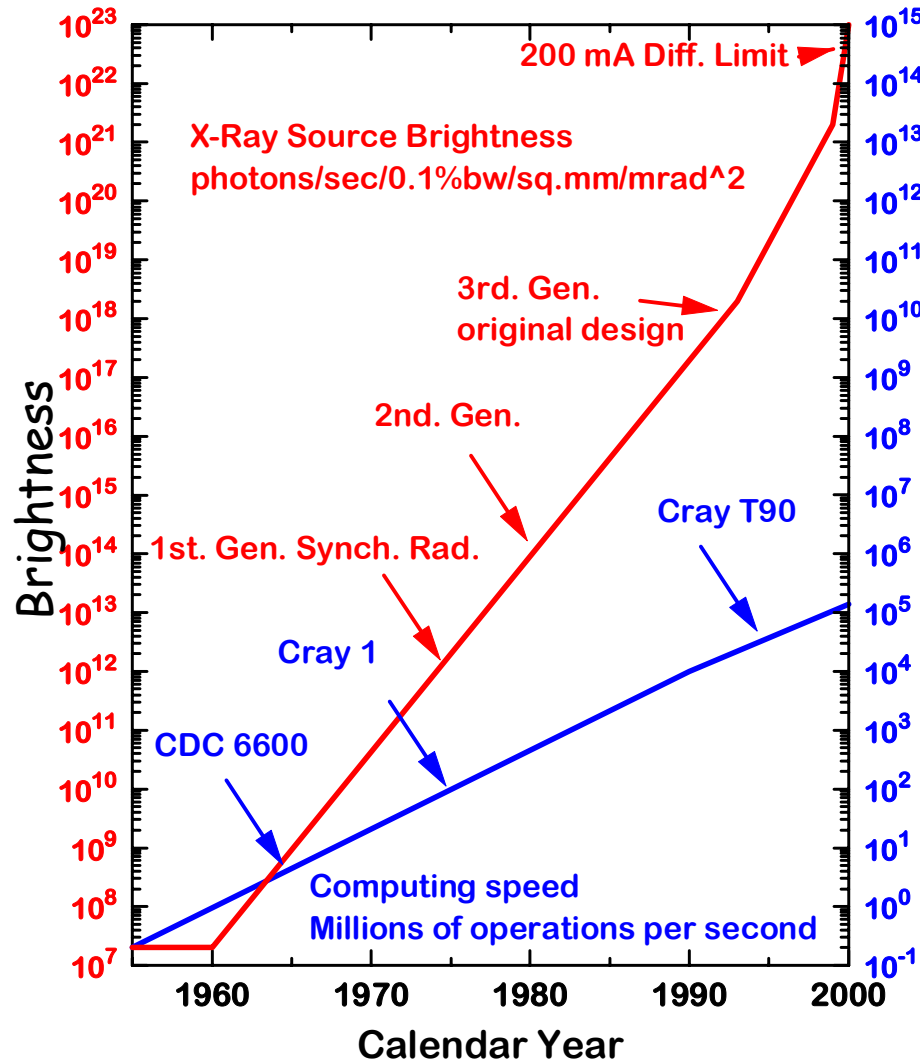


European Synchrotron Radiation Facility, France





Moore's law

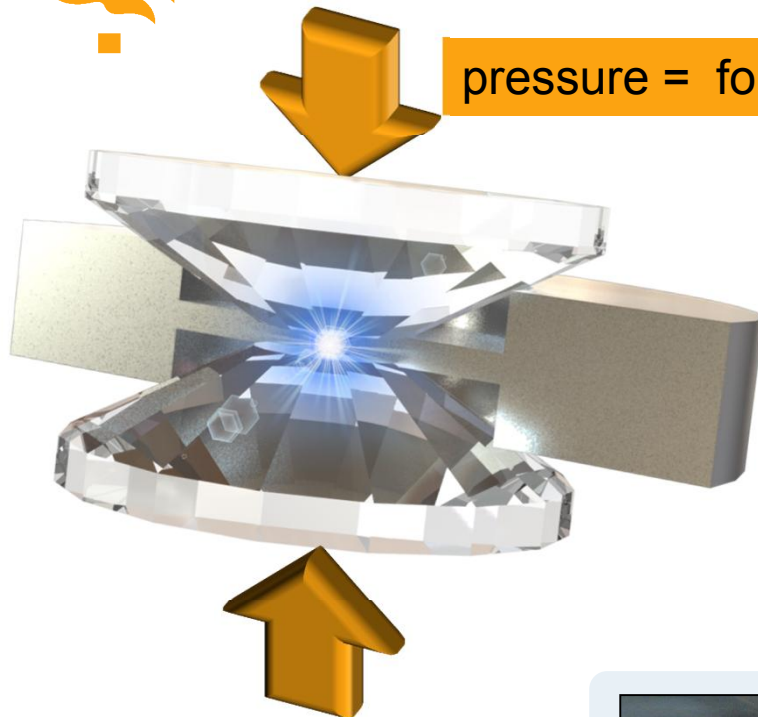




Extremely high pressures

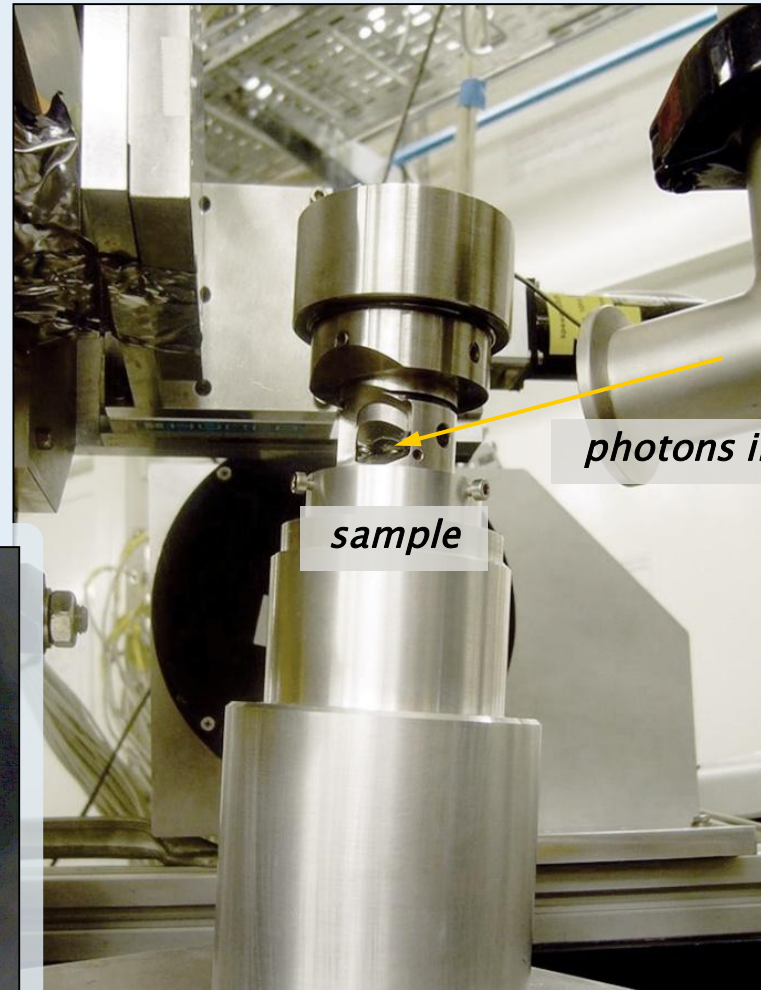
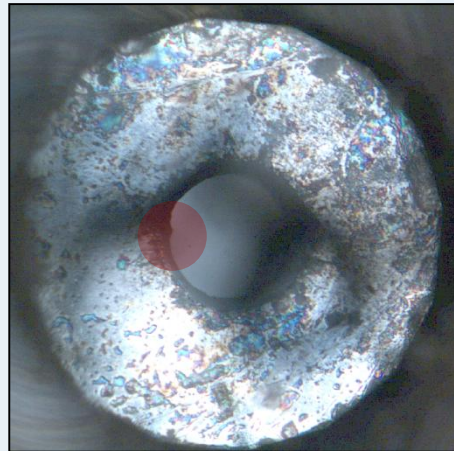
pressure = force / area

V.M.Giordano, T. Pylkkänen et al. ESRF ID16



Experimental details

- 1 mm diamond tip (culet)
- \varnothing 5 mm Be gasket
- \varnothing 350 micron sample size
- X-ray beam $100 \times 50 \mu\text{m}^2$
- Ruby chip for P calibration



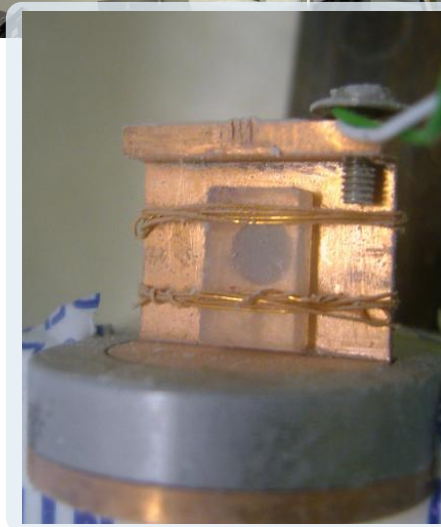


Low temperatures

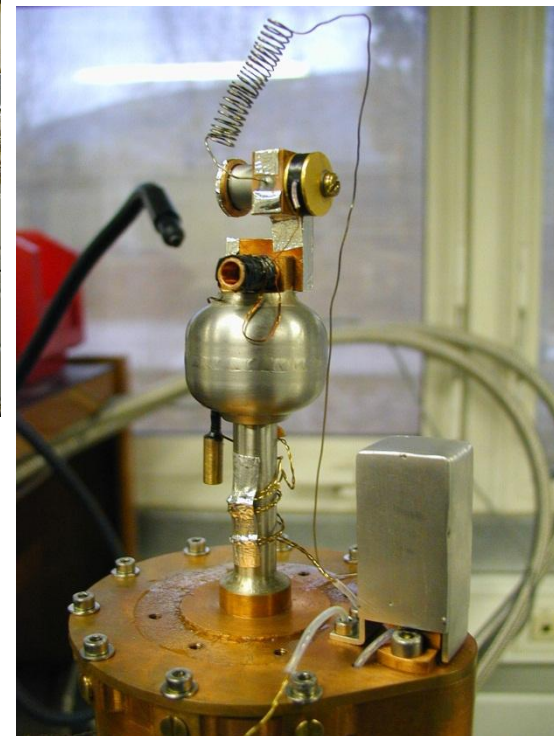
Liquid nitrogen cryo-jet (77 K)



Liquid He Cryostat (4 K)



He sorption pump (1 K)



F. Albergamo et al., ESRF

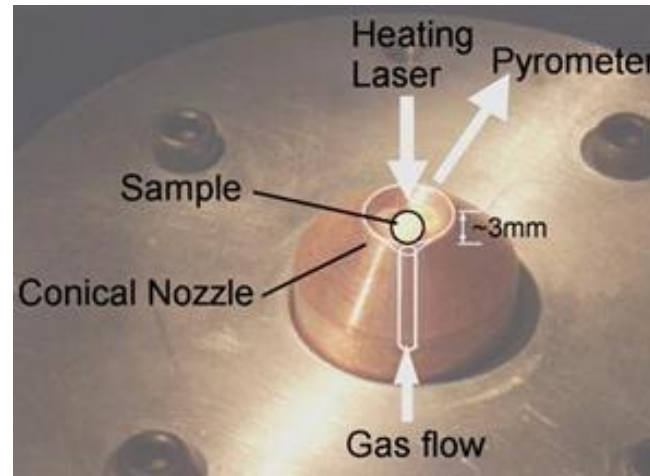


Extremely high temperatures (thousands of deg): Laser heating and aerodynamic levitation

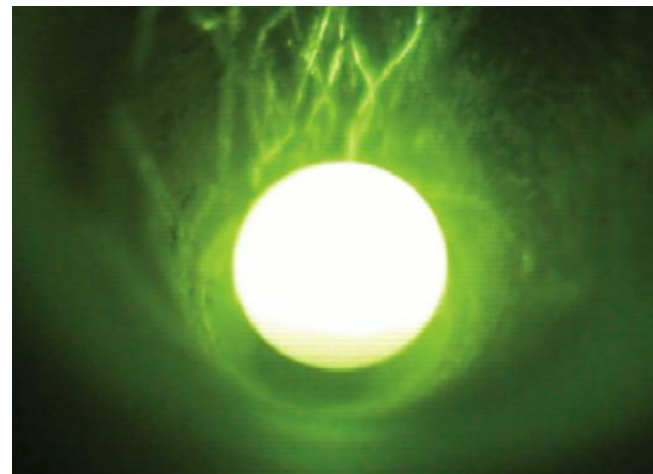
High temperatures
(up to 1000 deg C):
furnaces, hot air guns



ESRF, Sample group



Wikipedia:
Aerodynamic
levitation



Droplet of liquid basalt BCR-2 during levitation. The sample was heated from the top using a CO₂-laser. The diameter of the sphere was ~2 mm. A. Pack et al., Geochemical transactions 2010, 11:4