

# Reducing phase fluctuations in a layered superconductor via driving

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# Reducing phase fluctuations via driving

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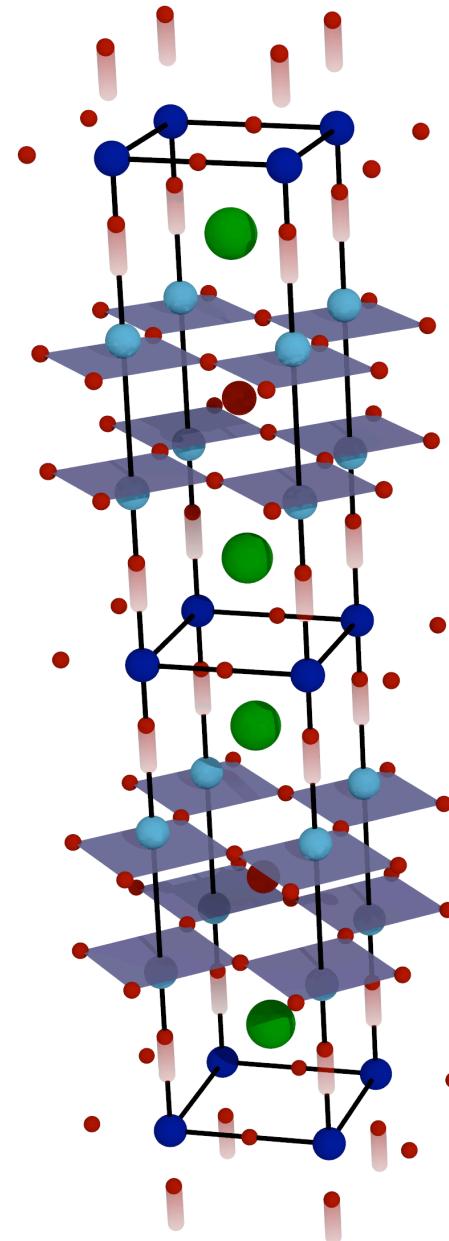
Pump-probe experiments in YBCO

Dynamic enhancement of superconductivity

Modeled via coupled Josephson junctions

Reduction of inter-layer phase fluctuations

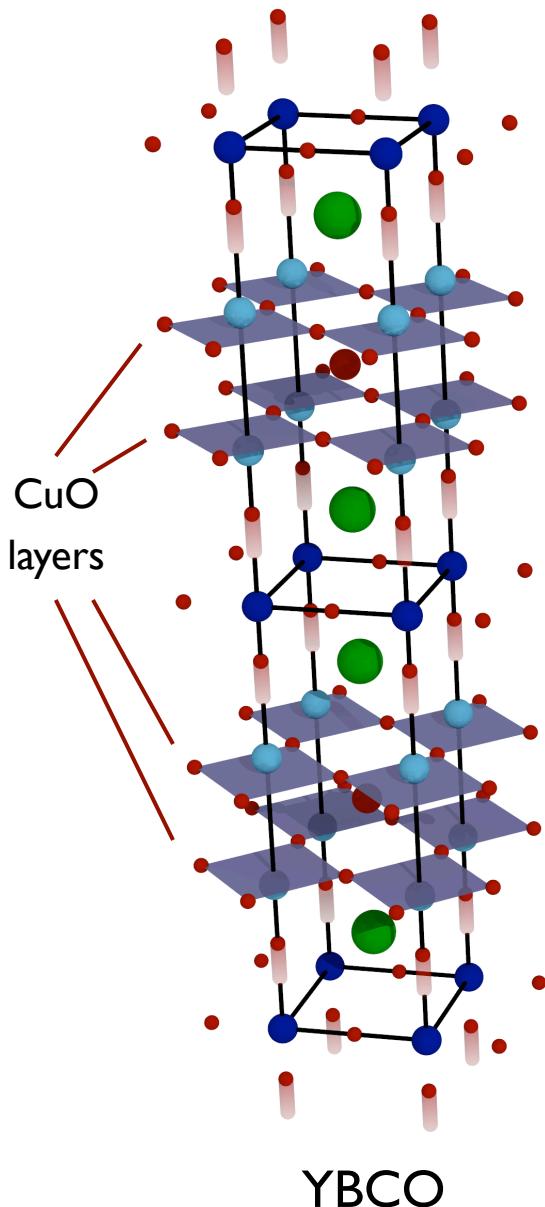
c-m/1406.3609



# Overview

- Motivation and setup of the model
  - YBCO, plasmons, pump-probe experiments, Josephson junctions, fluctuations...
- Toy models: two coupled junctions, single driven junctions
  - Dynamical suppression of fluctuations, power spectrum...
- Full model of the bulk system
  - Comparison to toy models, in-plane dynamics...
- Conclusions

# Layered superconductors



Pairing field

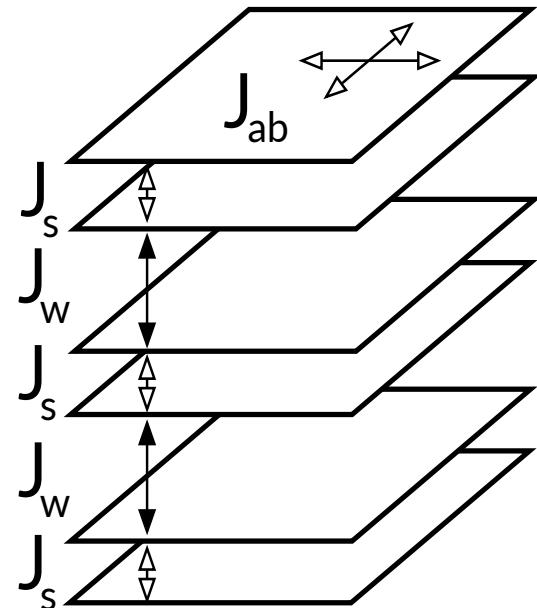
$$\psi_i \approx \sqrt{n_0 + \delta n_i} \exp(i\theta_i)$$

Charge fluctuations

(extended) Lawrence-  
Doniach model

$$H_\theta = - \sum_{\langle ij \rangle} J_{ij} \cos(\theta_i - \theta_j)$$

phase

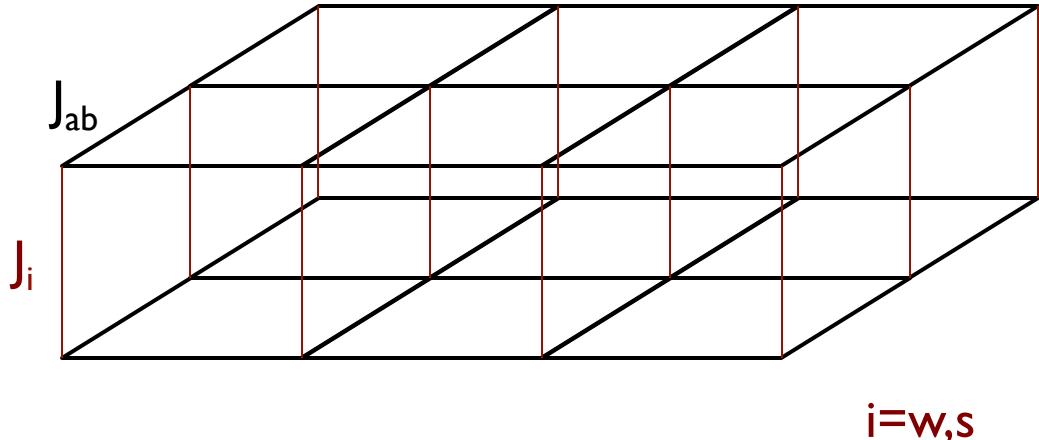


Layered system of coupled  
Josephson junctions

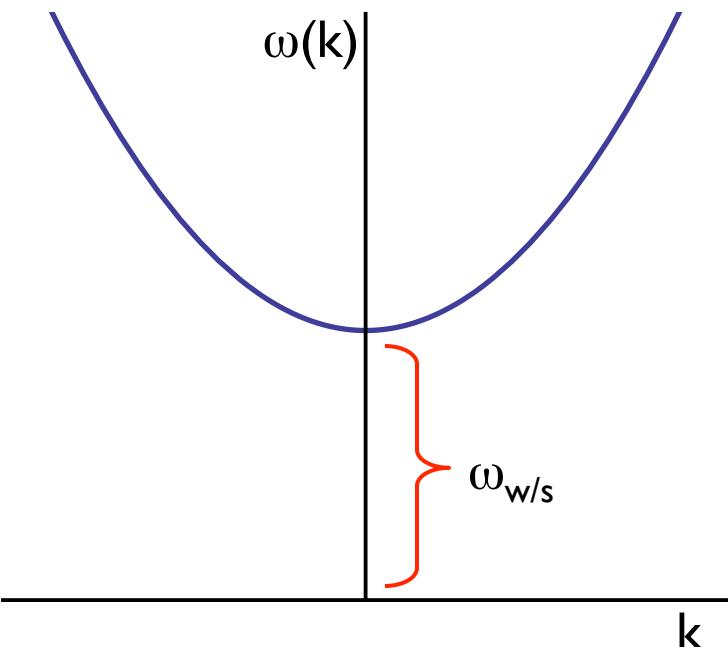
# Plasmon modes

Charging energy

$$H_n = \frac{E_c}{2} \sum_i \delta n_i^2$$



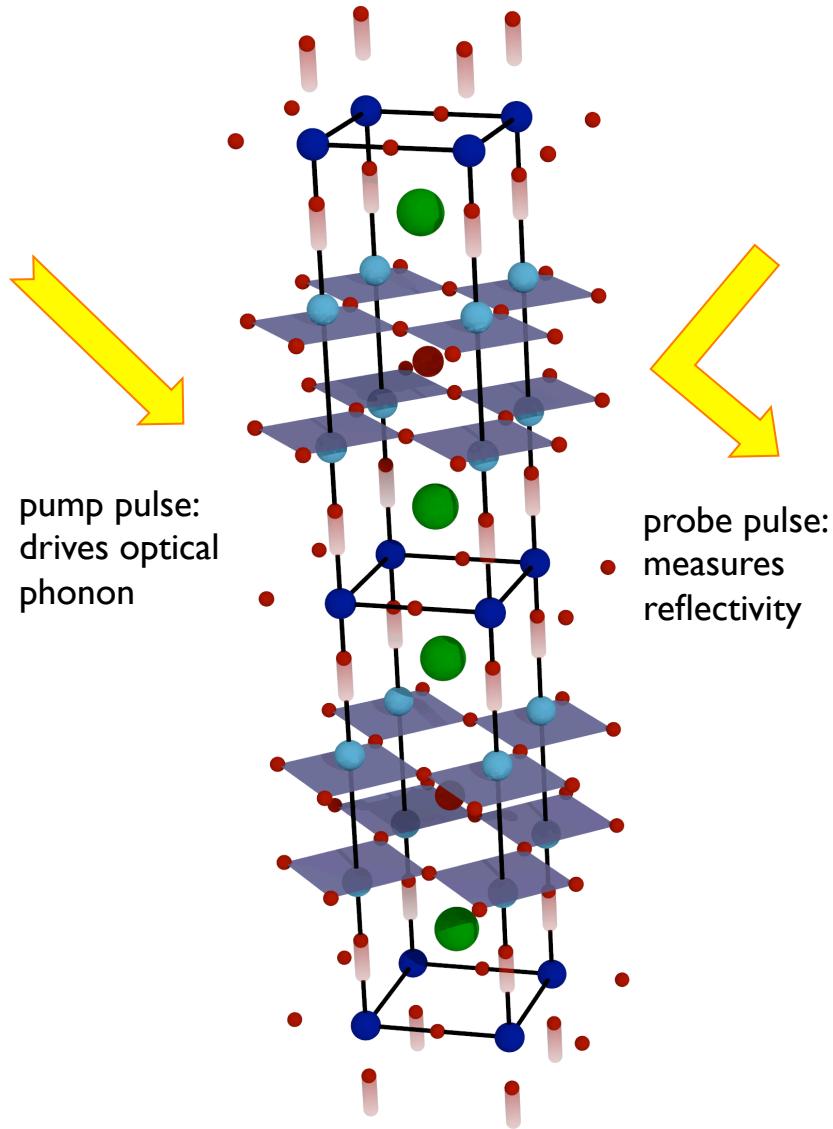
$i=w,s$



Plasmon frequencies

$$\omega_{w/s} = \sqrt{2J_{w/s}E_c}$$

# Pump-probe experiments

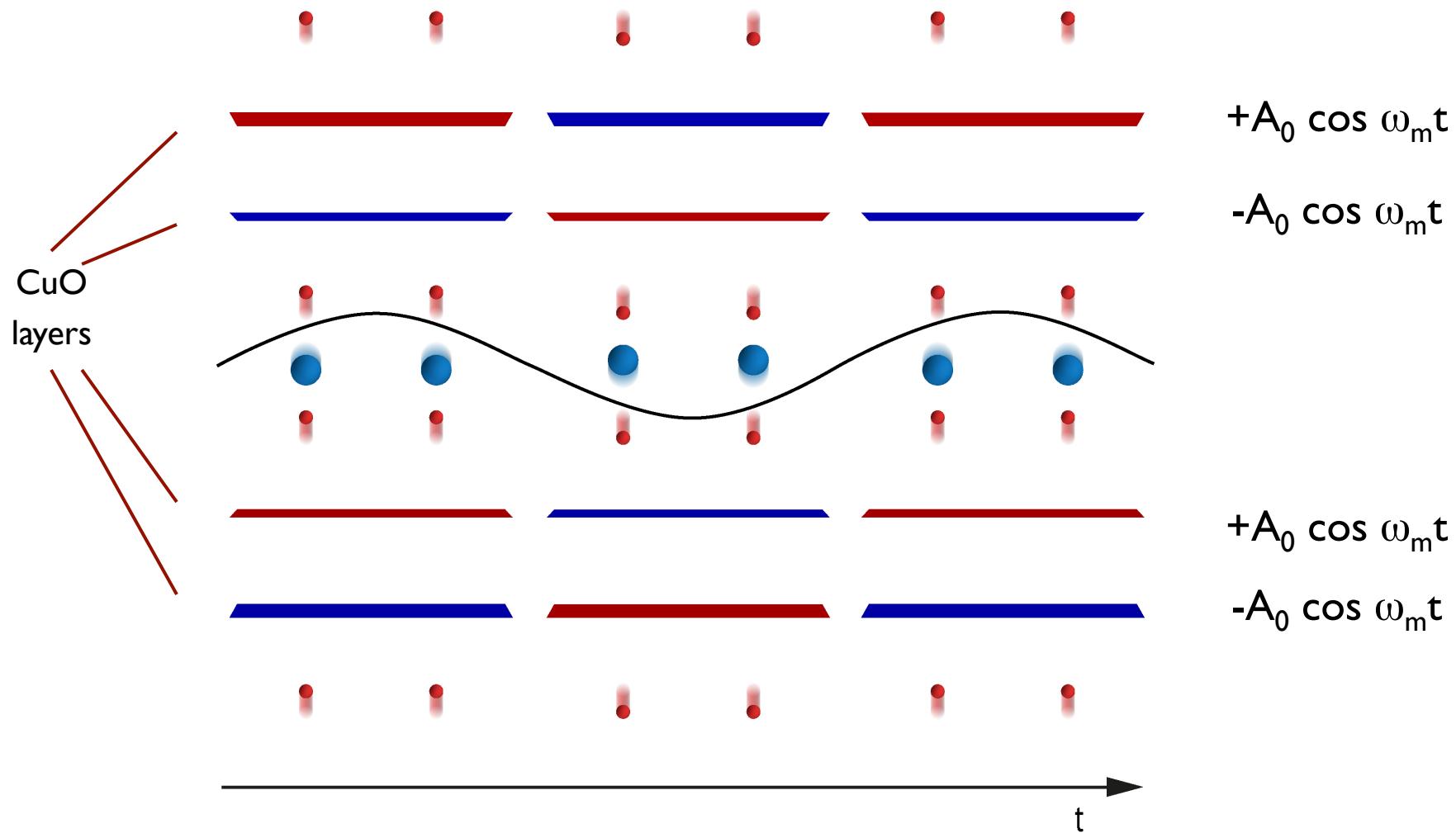


Observation:

Pump pulse boosts the optical conductivity at low frequencies

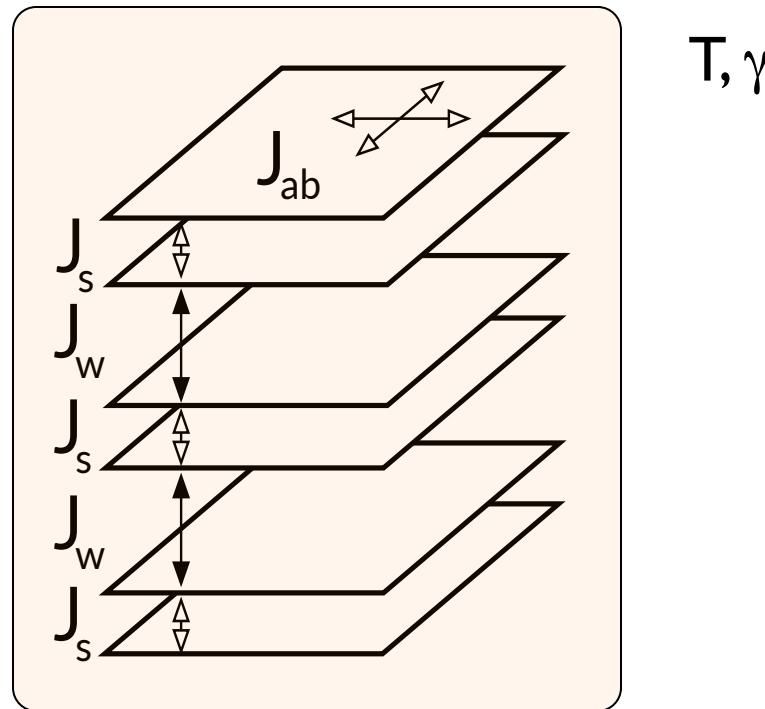
Superfluid density  $\omega \sigma_2(\omega, t)$  is enhanced.

# Effect of driven phonons



$$H_{dr} = \sum_i A_i(t) \delta n_i$$

# Coupled Josephson junctions, with a thermal bath



Example: single Josephson junction coupled to a thermal bath

$$\ddot{\theta} = -\omega_w^2 \sin \theta + F_0 \sin(\omega_m t) - \gamma \dot{\theta} + \xi$$

deterministic      damping      noise

## Parameters

Plasmon frequencies

$$\omega_w = \sqrt{2J_w E_c} \quad \sim 1 \text{ Thz}$$

$$\omega_s = \sqrt{2J_s E_c} \quad \sim 10 \text{ Thz}$$

Kosterlitz-Thouless scale

$$J_{ab} \sim 10^2 \text{ K}$$

Ratios

$$J_{ab} : J_s : J_w \sim 1000 : 100 : 1$$

Driving energy

$$A_0 \sim 50 \text{ K} \sim 1 \text{ THz}$$

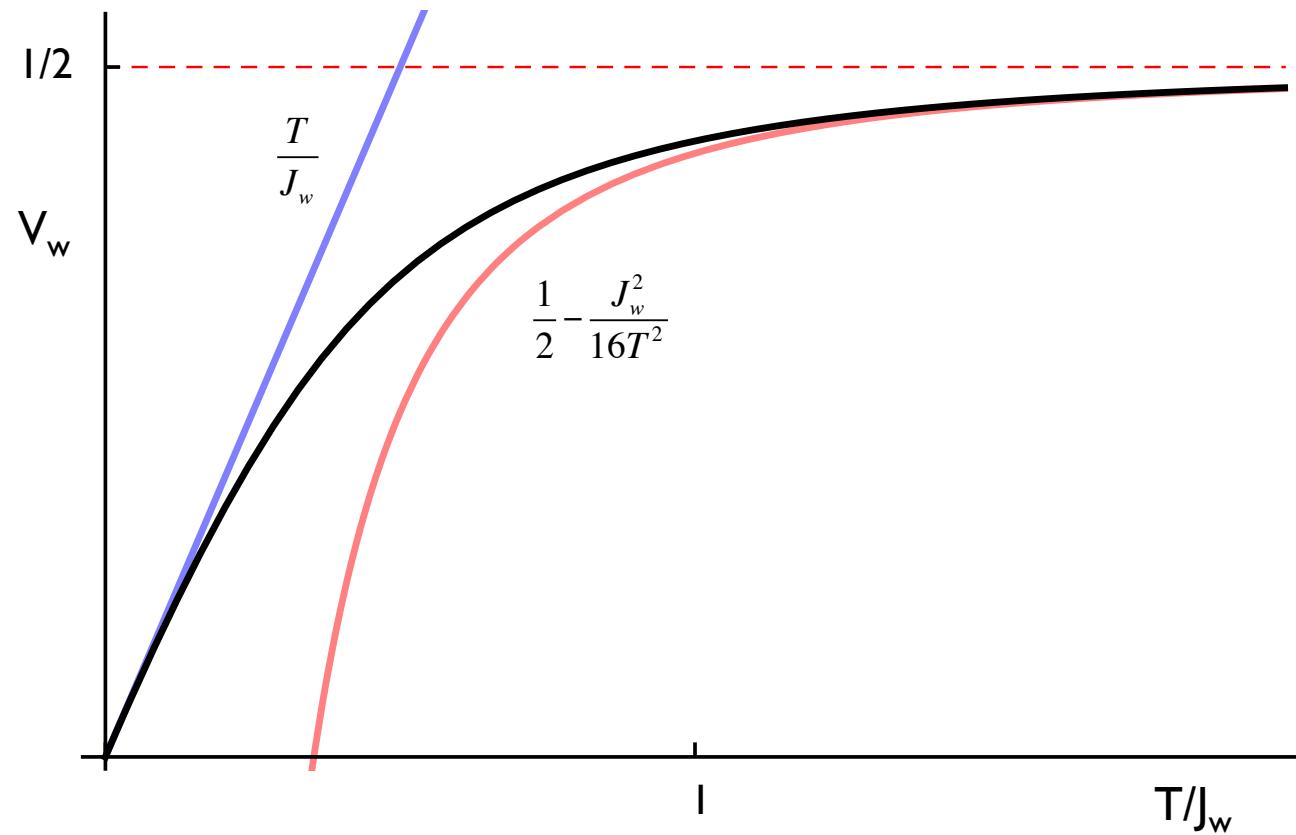
Damping

$$\omega_s > \gamma > \omega_w \quad \gamma \sim 10^1 - 10^2 \text{ K}$$

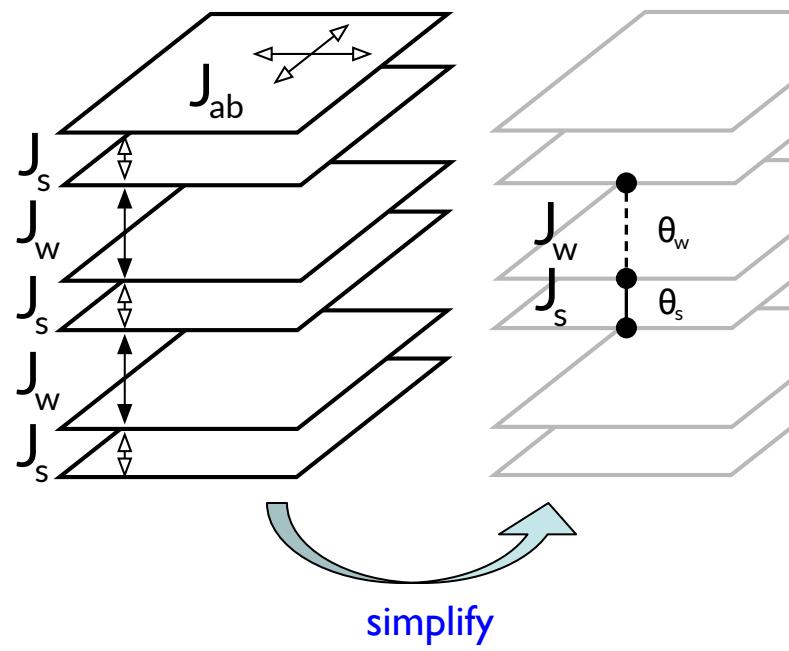
## Phase and current fluctuations

$$V_w = \langle \sin^2 \theta_w \rangle - \langle \sin \theta_w \rangle^2 \sim \langle j_w^2 \rangle - \langle j_w \rangle^2 \quad j_w = 2J_w \sin \theta_w$$

Example: single junction,  $H = J \cos \theta$



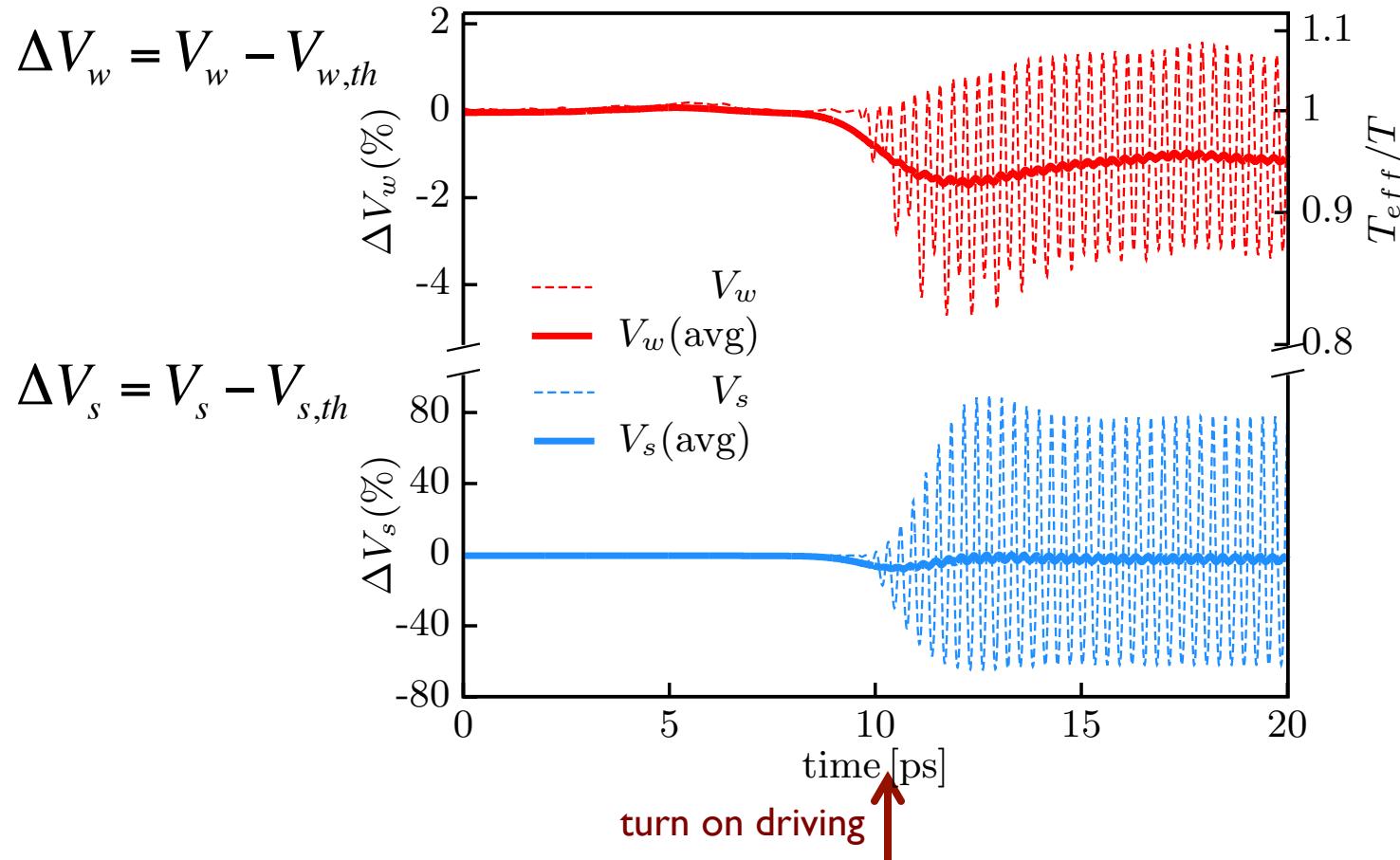
# Toy model



Langevin and Fokker-Planck approach

# Time evolution

$$V_{w,dr} = V_{w,th}(T_{eff})$$

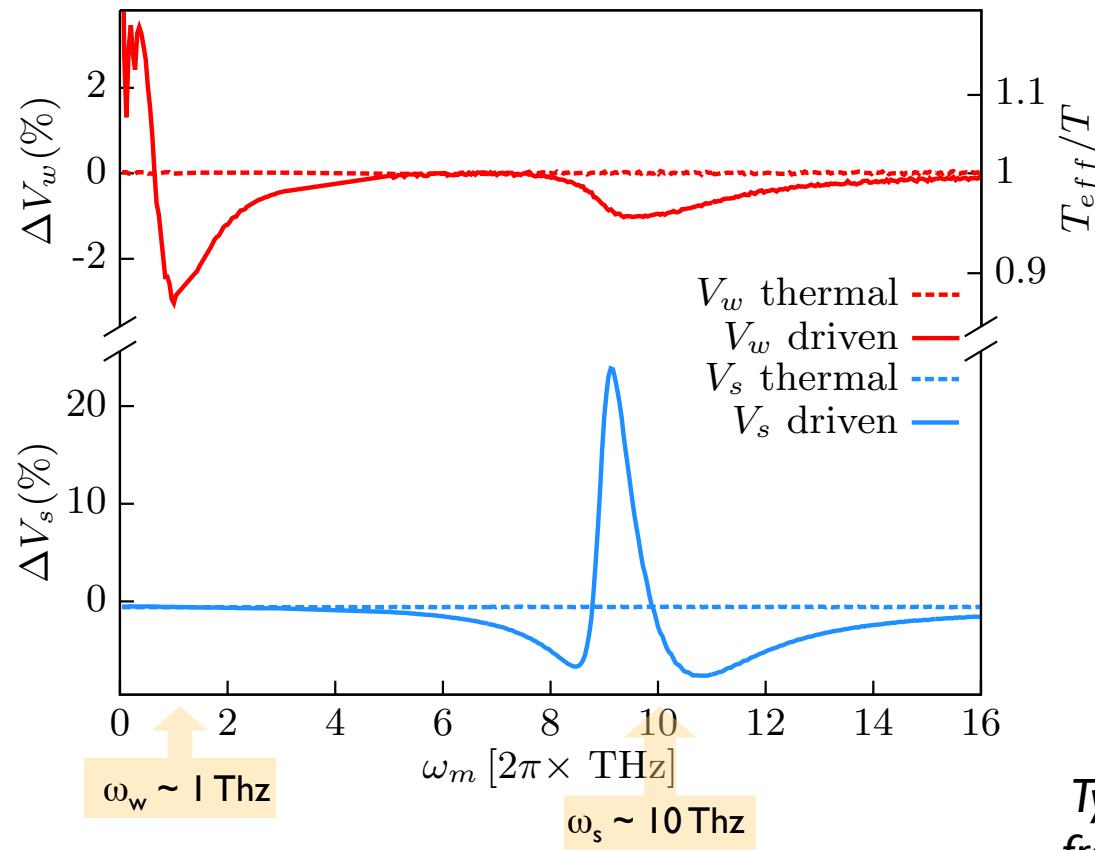


Reduction of fluctuations!

# Frequency scan

Time-average, in steady state

$$\bar{V}_w = \langle V_w(t) \rangle_t$$



Strong junction amplifies external driving

# Power spectrum of the weak junction

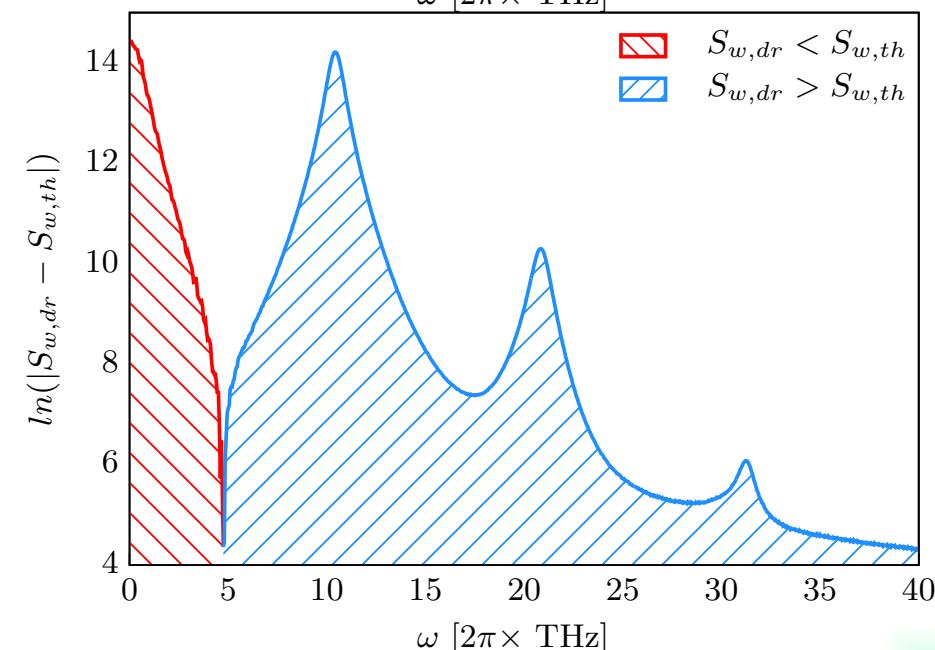
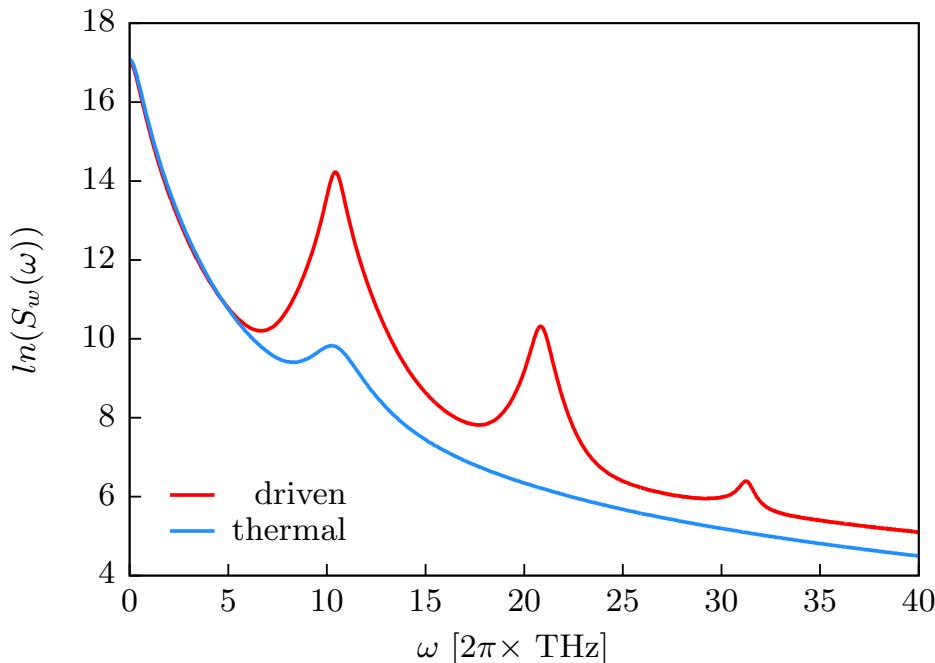
$$S_w(\omega) = \langle j_w(-\omega) j_w(\omega) \rangle$$

with  $j_w = 2J_w \sin \theta_w$

- ✓ Suppression of low-frequency fluctuations
- ✓ Enhancement high-frequency fluctuations

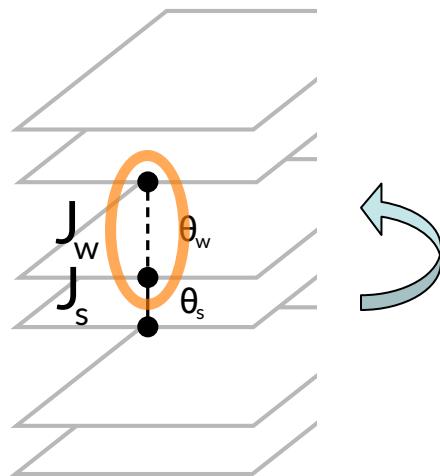


Redistribution of spectral weight!



# Single junction approximation

Focus on weak junction



Treat strong junction as external drive

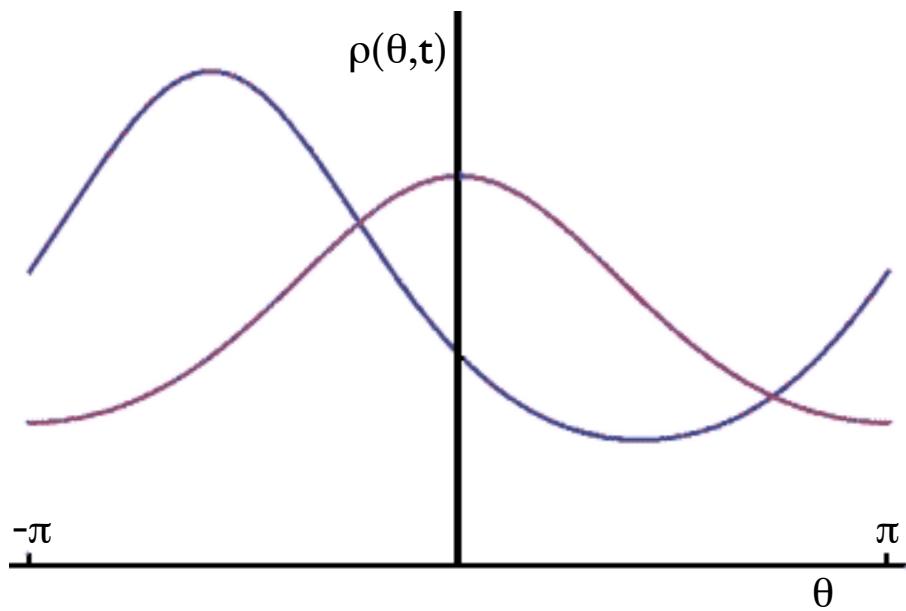
Driven, damped Josephson junction in thermal bath

$$\ddot{\theta} = -\omega_w^2 \sin \theta + F_0 \sin(\omega_m t) - \gamma \dot{\theta} + \xi$$

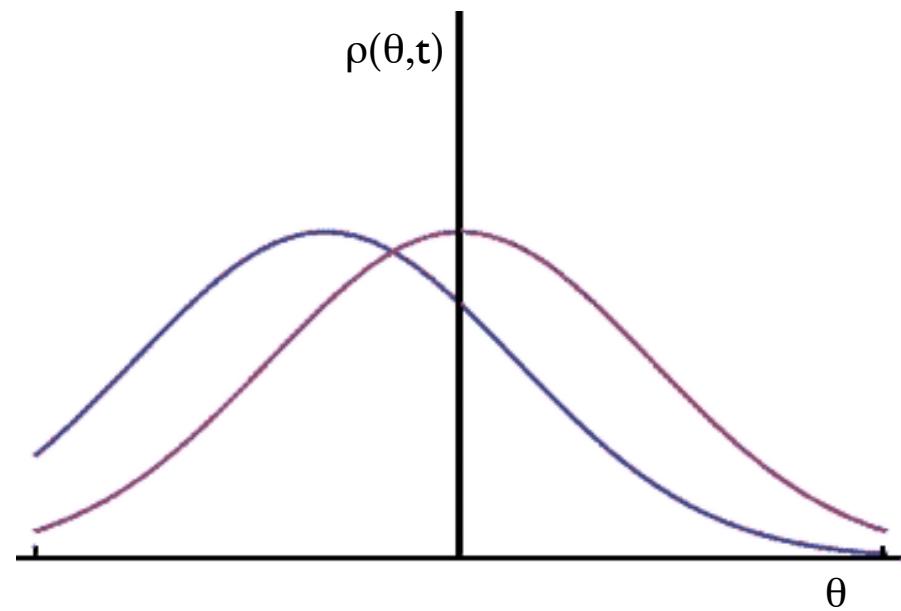
# Overdamped, driven Josephson junction

$$\cancel{\ddot{\theta}} = -\omega_w^2 \sin \theta + F_0 \sin(\omega_m t) - \gamma \dot{\theta} + \xi \quad \rightarrow \quad \dot{\theta} = -\frac{\omega_w^2}{\gamma} \sin \theta + \frac{F_0 \sin(\omega_m t)}{\gamma} + \frac{\xi}{\gamma}$$

Fokker-Planck solution gives  $\rho(\theta, t)$



Josephson junction



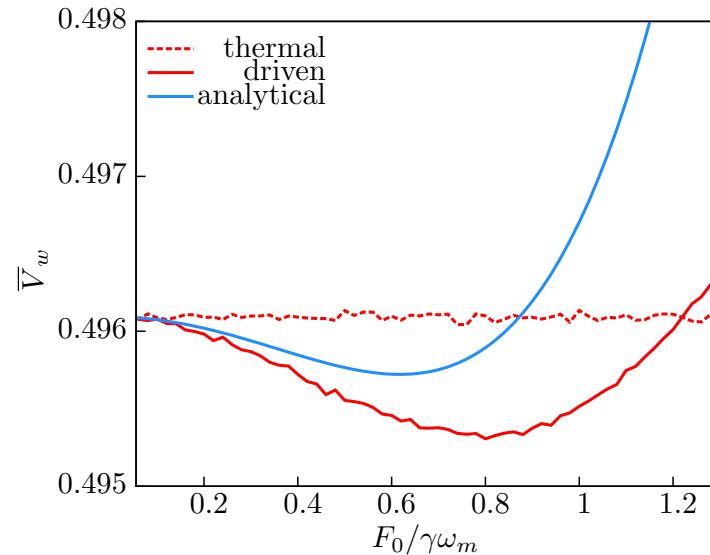
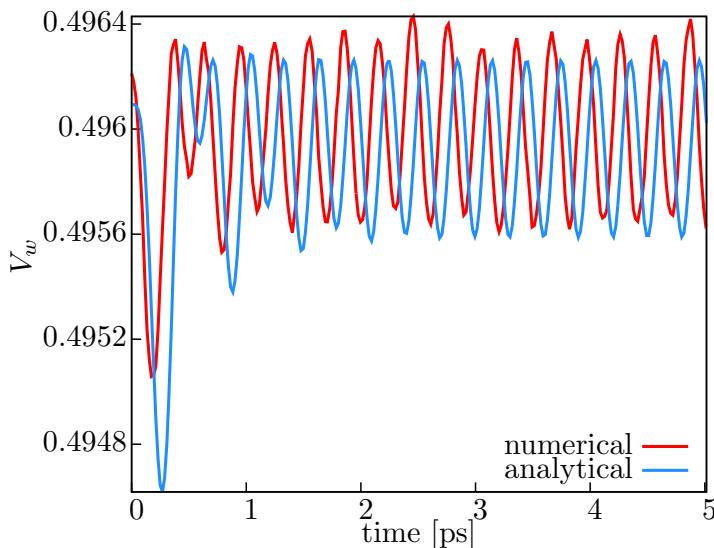
Harmonic oscillator

# Reduction in steady state, transient behavior

## High temperature expansion

$$V_w \approx \frac{1}{2} - \frac{J_w^2}{16T^2} - \frac{1}{32} \frac{J_w^2}{T^2} \frac{F_0^2}{\gamma^2 \omega_m^2} + \frac{21}{512} \frac{J_w^2}{T^2} \frac{F_0^4}{\gamma^4 \omega_m^4}$$

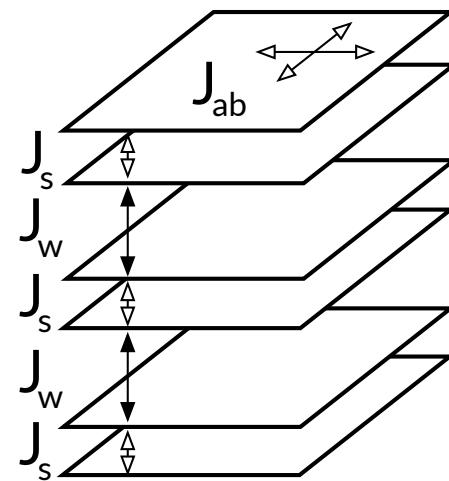
Reduction of  $T_{\text{eff}}$ : 5 – 10%



## Transient behavior

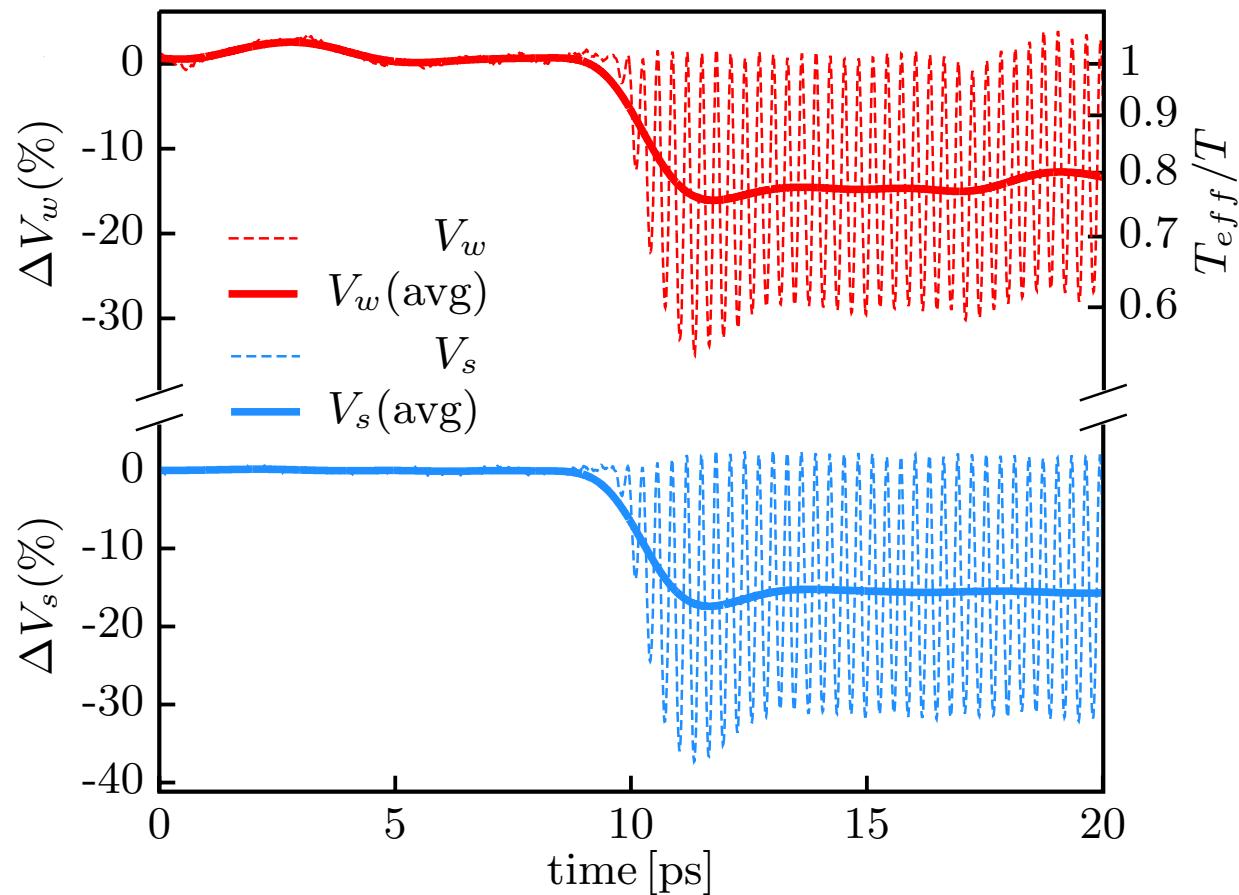
$$t_{tr} = \frac{\gamma}{2TE_c}$$

# Bulk system



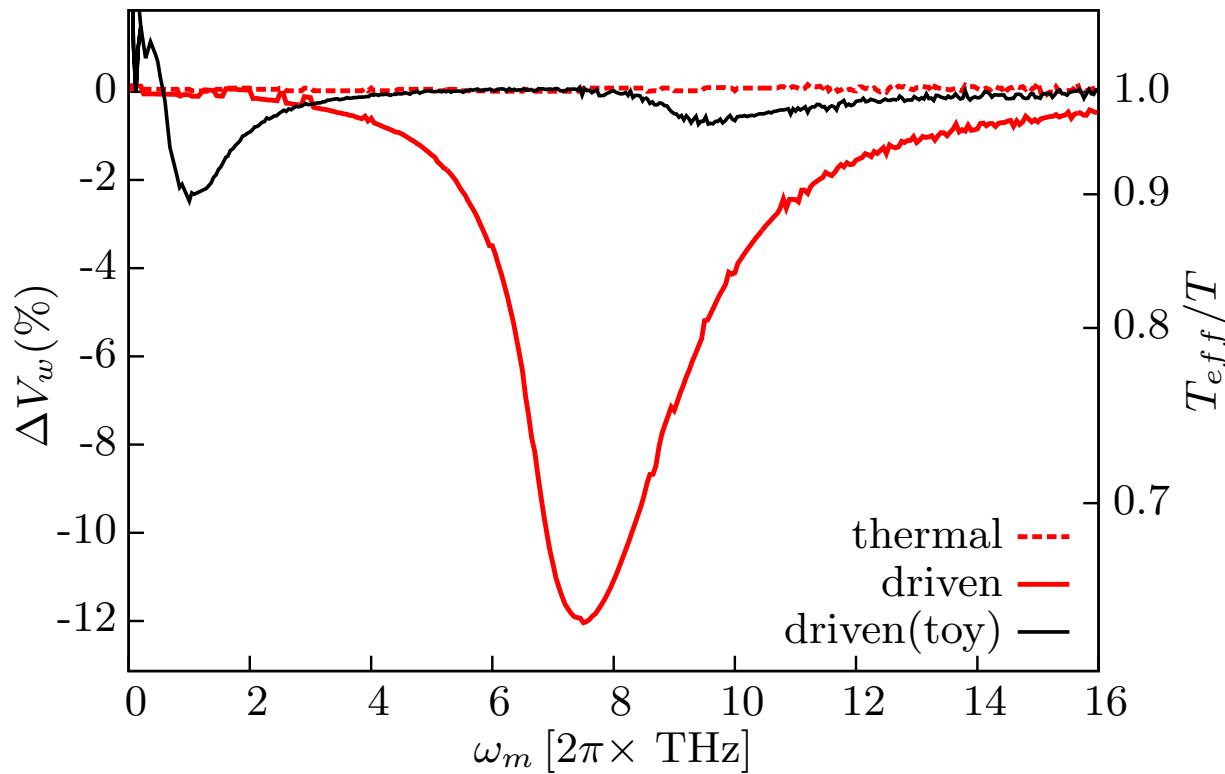
Lattice: 128 \* 128 \* 4

# Time evolution of the bulk system



- ✓ Qualitatively similar to the toy model
- ✓ Larger reduction of phase fluctuations

## Frequency dependence



- ✓ Strong reduction near high frequency plasmon
- ✓ Feature due to direct driving washed out

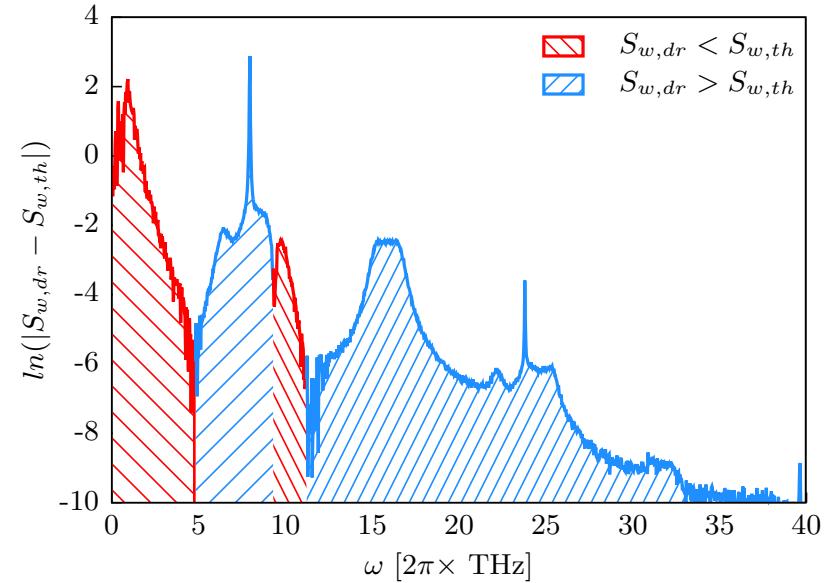
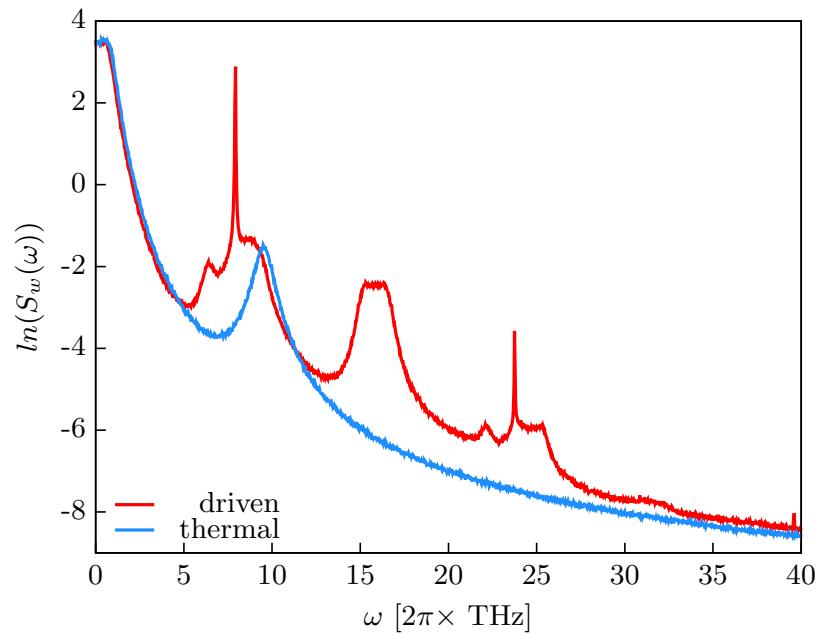
# Power spectrum

$$S(\omega) = \langle j_{w,tot}(-\omega) j_{w,tot}(\omega) \rangle$$

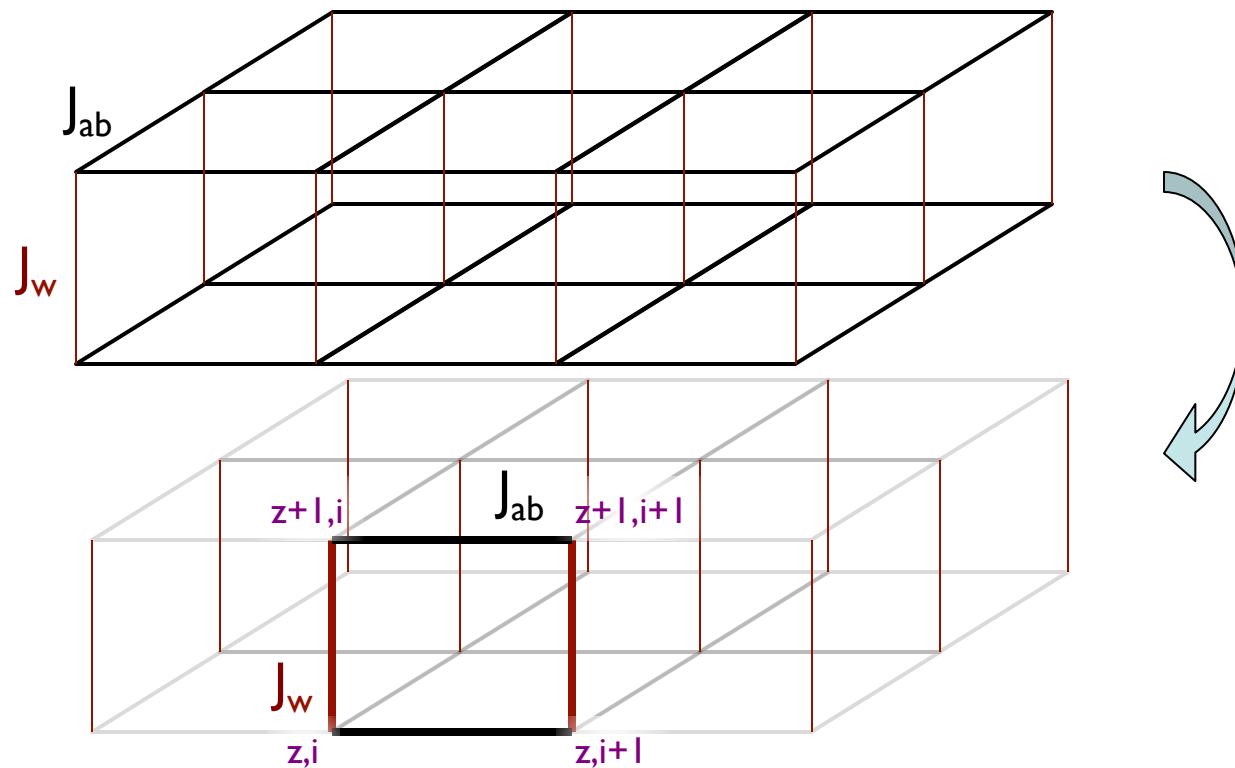
Sum over a layer

$$j_{w,tot} = \sum_{ij} 2J_w \sin(\theta_{ij})$$

- ✓ Suppression of low-frequency modes
- ✓ Enhancement of high frequencies



# Effective single junction model



$$\begin{aligned}
 H &\approx -J_{ab} (\cos(\theta_{z+1,i+1} - \theta_{z+1,i}) + \cos(\theta_{z,i+1} - \theta_{z,i})) \\
 &\approx -J_{eff}(T) \cos((\theta_{w,i+1} - \theta_{w,i})/2) \\
 &\approx -J_{eff}(T) (\cos(\theta_{w,i}/2) \langle \cos(\theta_{w,i+1}/2) \rangle + \sin(\theta_{w,i}/2) \langle \sin(\theta_{w,i+1}/2) \rangle)
 \end{aligned}$$

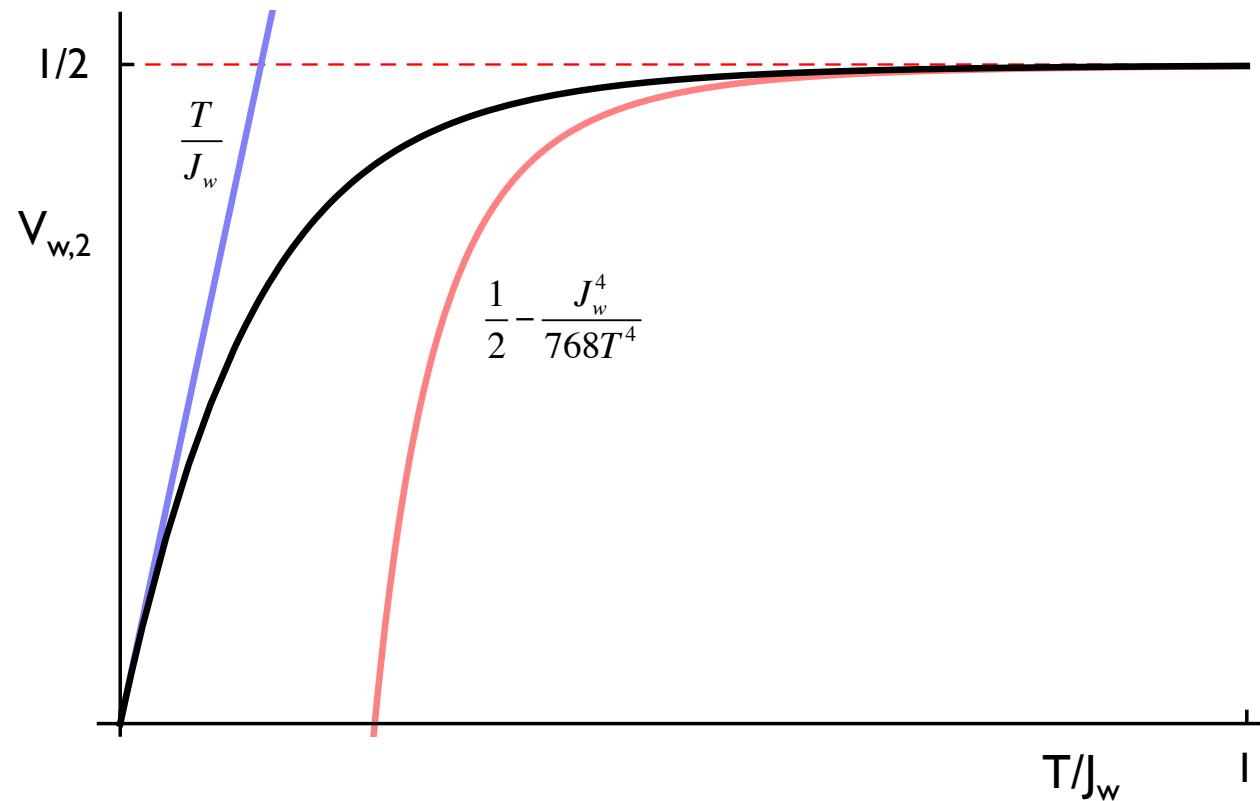
Effective single junction      
 Effective driving

→  $\cos(\theta/2)$  instead of  $\cos(\theta)$

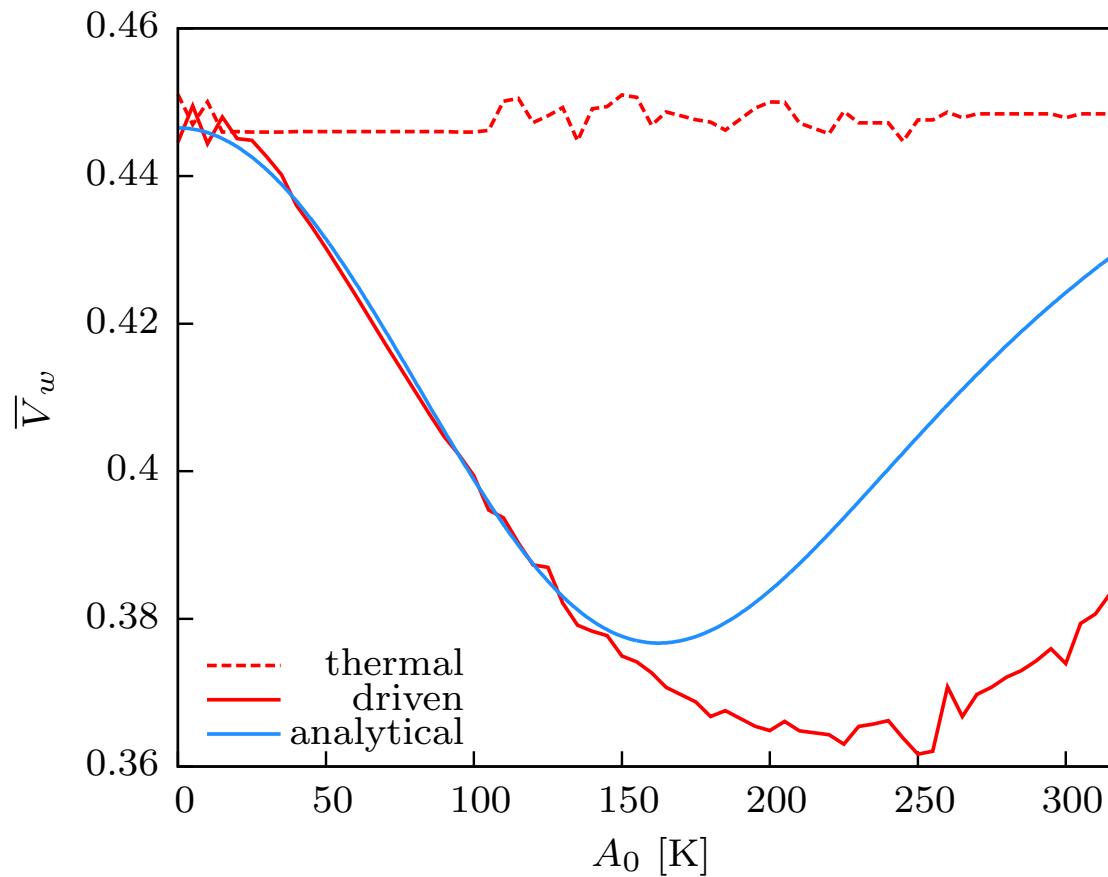
# Effective single junction model for the bulk system

Use  $\dot{\theta} = -\frac{\omega_w^2}{\gamma} \sin \theta + \frac{F_0 \sin(\omega_m t)}{\gamma} + \frac{\xi}{\gamma}$  with  $J_{\text{eff}}$  and effective  $F_0$

with  $V_{w,2} = \langle \sin^2 2\theta_w \rangle - \langle \sin 2\theta_w \rangle^2$



# Amplitude scan

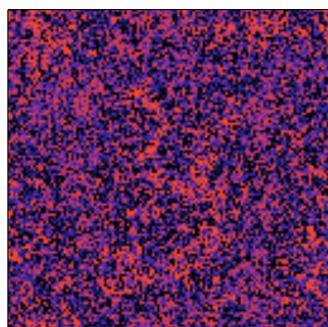


Large reduction captured by effective single junction model

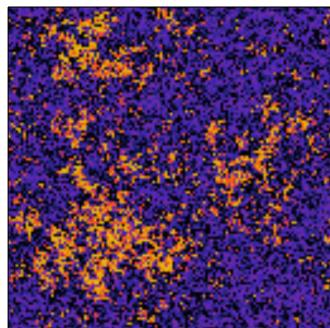
# In-plane dynamics

Current fluctuations  $\left(j_w(\vec{r}, t) - \bar{j}_w(t)\right)^2$

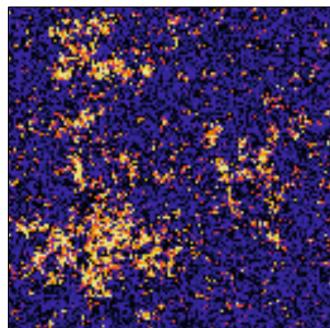
thermal



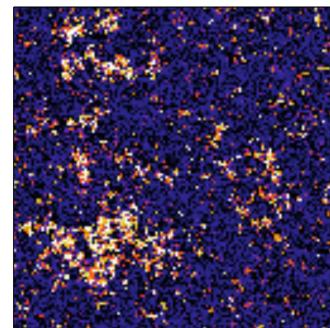
1/8 cycle



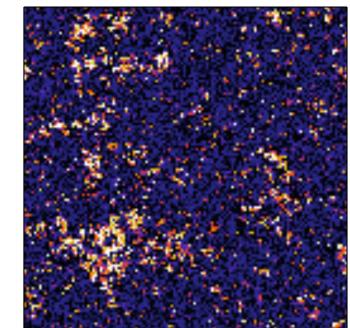
1/4 cycle



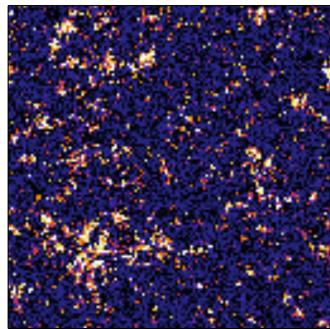
3/8 cycle



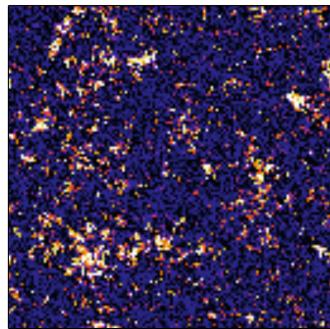
1/2 cycle



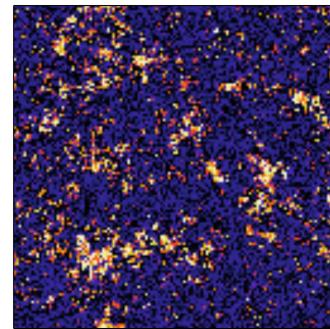
5/8 cycle



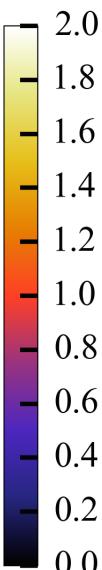
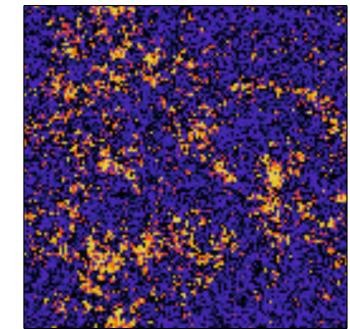
3/4 cycle



7/8 cycle



1 cycle



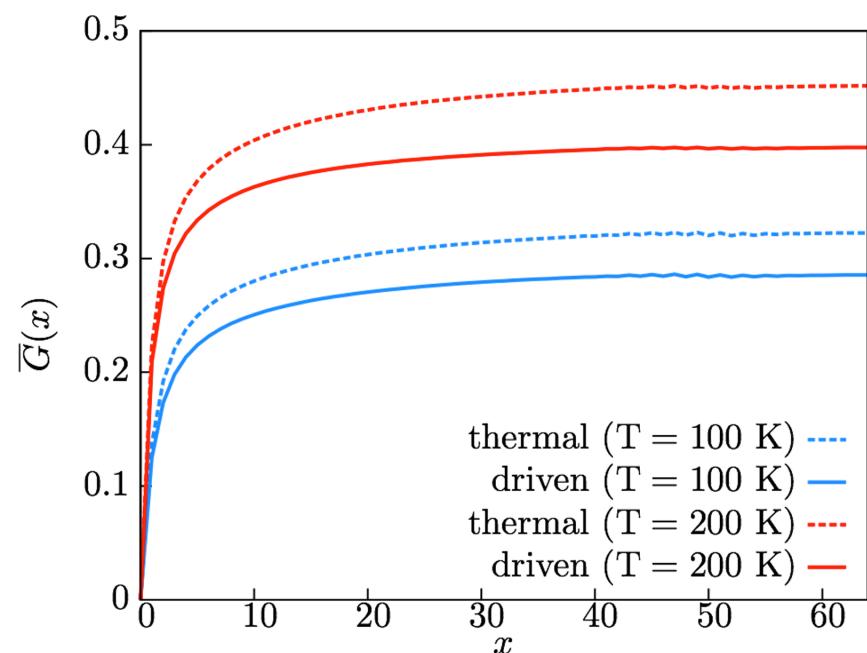
- ✓ Large regions of suppressed fluctuations, 'hot spots' of enhanced fluctuations
- ✓ Overall periodic breathing

# In-plane current correlations

Current fluctuations

$$G(\vec{r}, t) \sim \left\langle \left( j_w(0, t) - j_w(\vec{r}, t) \right)^2 \right\rangle$$

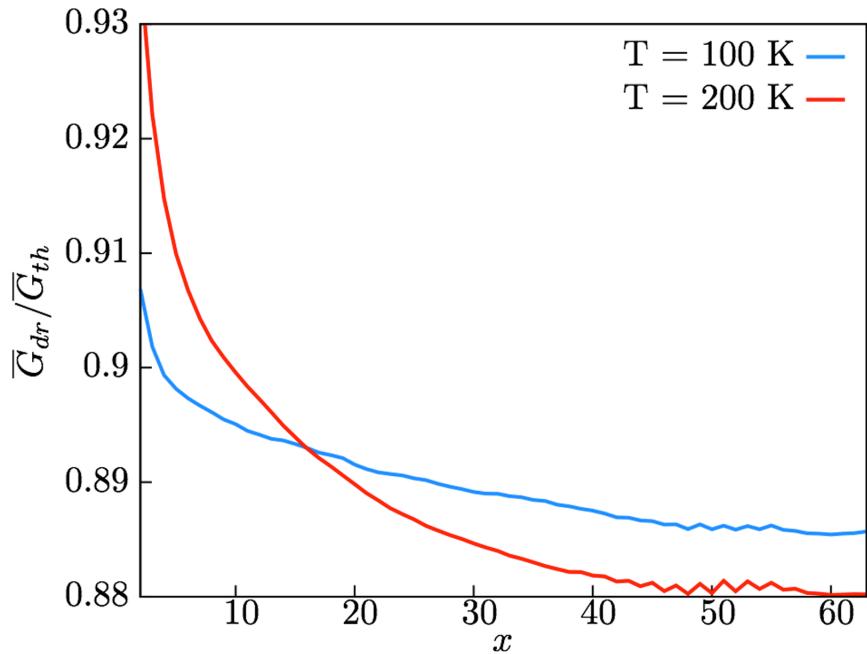
Overall reduction, esp. at long range



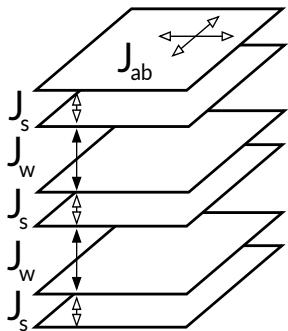
But: correlation length reduced

Consistent with:

- ✓ Suppression of low-frequencies
- ✓ Enhancement of high frequencies



# Conclusions

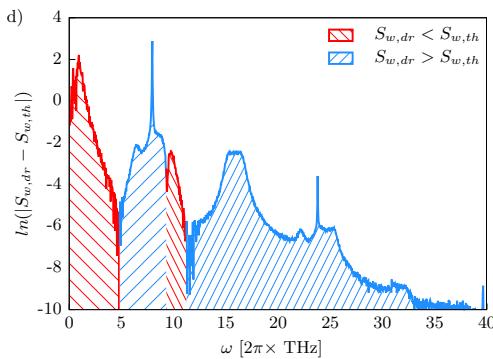
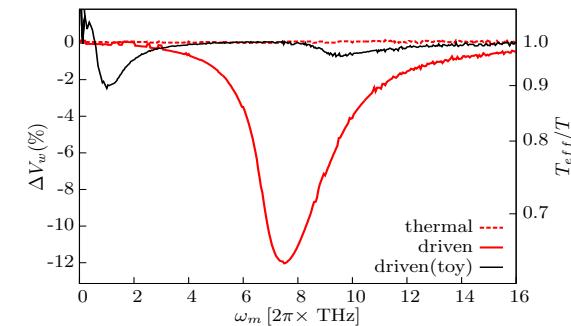


Driven superconductors modelled via layered system of Josephson junctions

Strong reduction of inter-layer phase fluctuations.

In the bulk system:  $T_{\text{eff}} / T \sim 0.6$

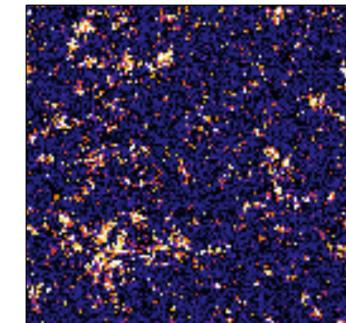
The strong junction serves as an amplifier



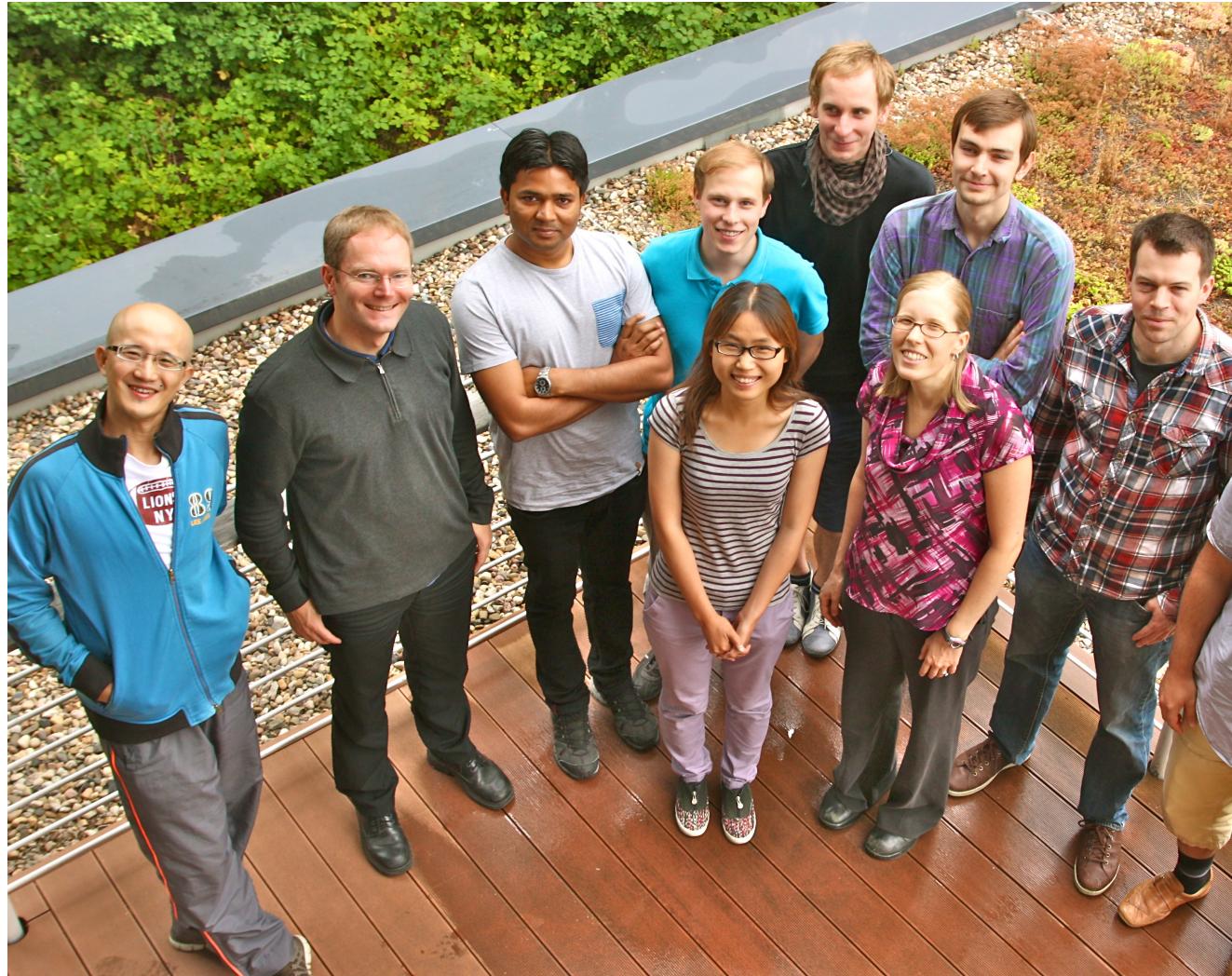
The reduction occurs at low frequencies.

Simultaneous increase of high frequency fluctuations:  
non-thermal state

In-plane correlations show increased coherence  
on long-scale



# Team



c-m/1406.3609

This project:

Robert Höppner  
Beilei Zhu  
Tobias Rexin

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W.-M. Huang, LM, V. Singh, M. Lahrz, B. Zhu, T. Rexin, A. Mathey, P. Janzen, R. Höppner