

Introduction

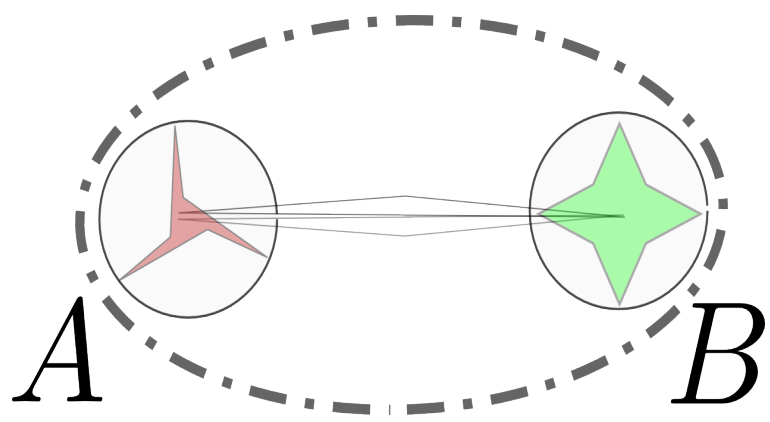
Entanglement: main concepts

Definition: a quantum state ρ in a composite Hilbert space $\mathcal{H} = \mathcal{H}_A \otimes \mathcal{H}_B$ is separable if $\rho = \sum_i p_i \rho_i^A \otimes \rho_i^B$ ($\sum_i p_i = 1$)

Otherwise the state is **entangled**; generalization to N subsystems gives definition of **multipartite** entangled state

Peres-Horodecki criterion: If a state is separable: **partial transpose** over one subsystem *has positive eigenvalues*.

Entanglement witness: observable that signals the presence of many-body entanglement when its value lies above some threshold



Formalism: Gaussian states

Provides tools for studying ground state properties of Hamiltonians which are **second order polynomials of bosonic operators**: all properties of the system can be deduced from expectation values of a **finite** set of operators O or **moments**:

$$O \in \{x_n, p_n, x_n x_m, p_n p_m, x_n p_m\}$$

In particular, for correlation properties, such as entanglement, we only need one ingredient, **the covariance matrix** σ

$$\sigma_{ij} = \frac{1}{2} \langle R_i R_j + R_j R_i \rangle - \langle R_i \rangle \langle R_j \rangle$$

$$\text{where } \mathbf{R}^T := (x_1, x_2, \dots, x_N, p_1, p_2, \dots, p_N)$$

$$\text{and } x_i = \frac{1}{\sqrt{2}}(a_i + a_i^\dagger), \quad p_i = \frac{1}{\sqrt{2}}i(a_i^\dagger - a_i)$$

$$a_n := \text{bosonic operators} \quad \text{C. Weedbrook et al.} \\ \text{Reviews of Modern Physics 84, 621 (2012)}$$

• Alternative:

Use an **entanglement witness** generalizing the idea of *Duan et al.* (2 modes),

$$W = \langle (\Delta(x_1 - x_2))^2 \rangle_\rho + \langle (\Delta(p_1 + p_2))^2 \rangle_\rho$$

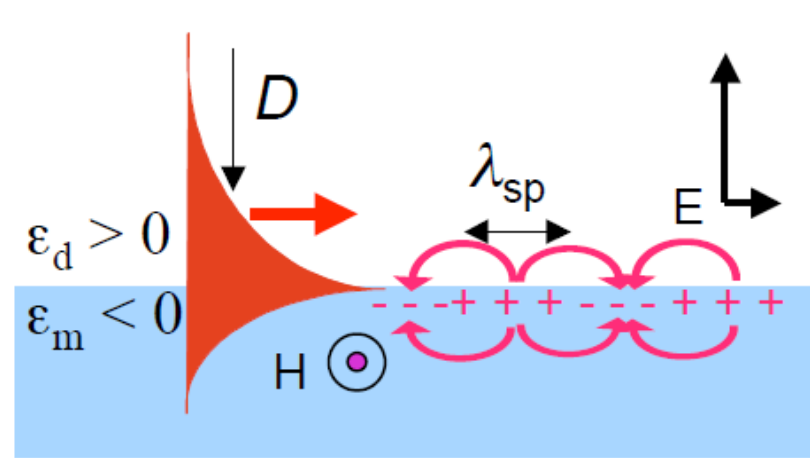
($W < 1 = \text{entanglement}$)

$$\text{L. Duan et al.} \\ \text{PRL 84, 2722 (2000)}$$

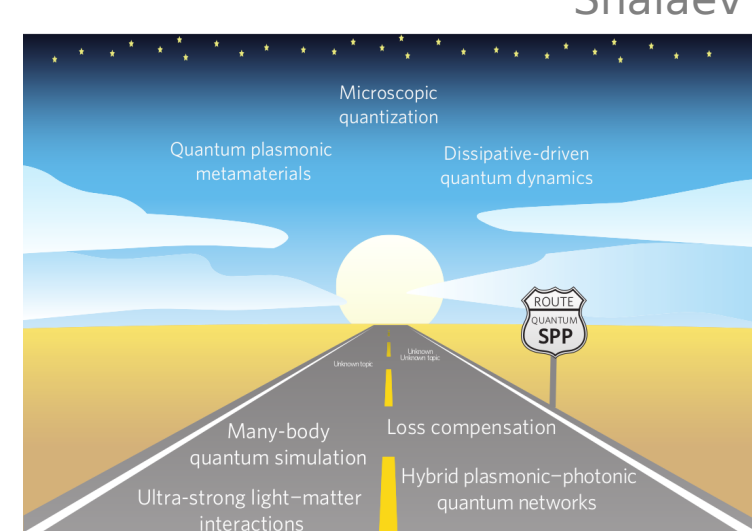
Why use plasmons

• Achievements of Quantum Plasmonics:

Many quantum properties of photons are preserved when these photons are coupled in and out of **plasmonic modes** (plasmons used as carriers)



from lecture by V. Shalaev



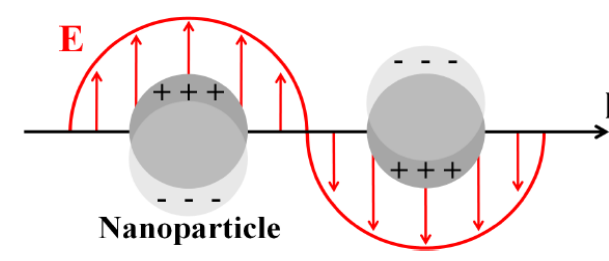
from M. S. Tame et al. Nature Physics 9, 329 (2013)

D. E. Chang, et al., Phys. Rev. Lett. 97, 053002 (2006).
D. Dzotjjan, et al., Phys. Rev. B 82, 075427 (2010).
A. Gonzalez-Tudela et al. (PRL 106, 020501 (2011))
E. Altewischer et al. (Nature 418,304 (2002))

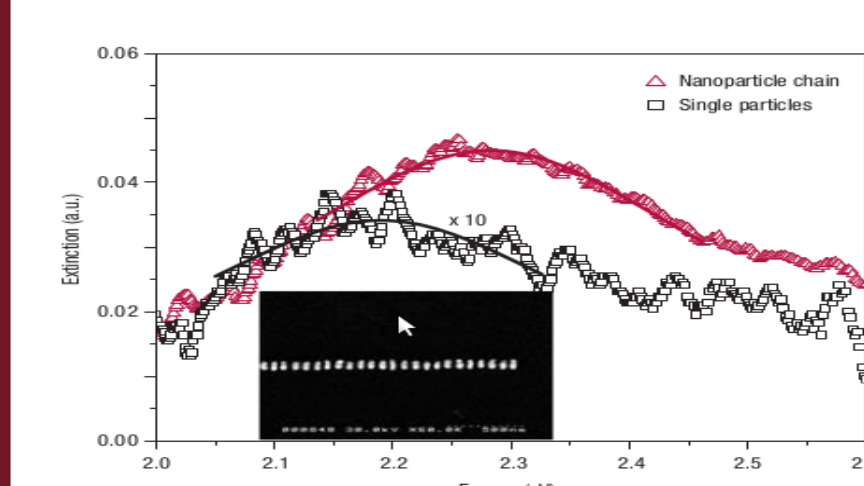
• **Our proposal:** use of quantum properties of plasmons *themselves*; localized plasmonic resonances (**SPPs**) as entangled parties.

Nanoparticle arrays

- **Nanoparticles: localized SP resonance** strong absorption/scattering strong local field enhancement high losses



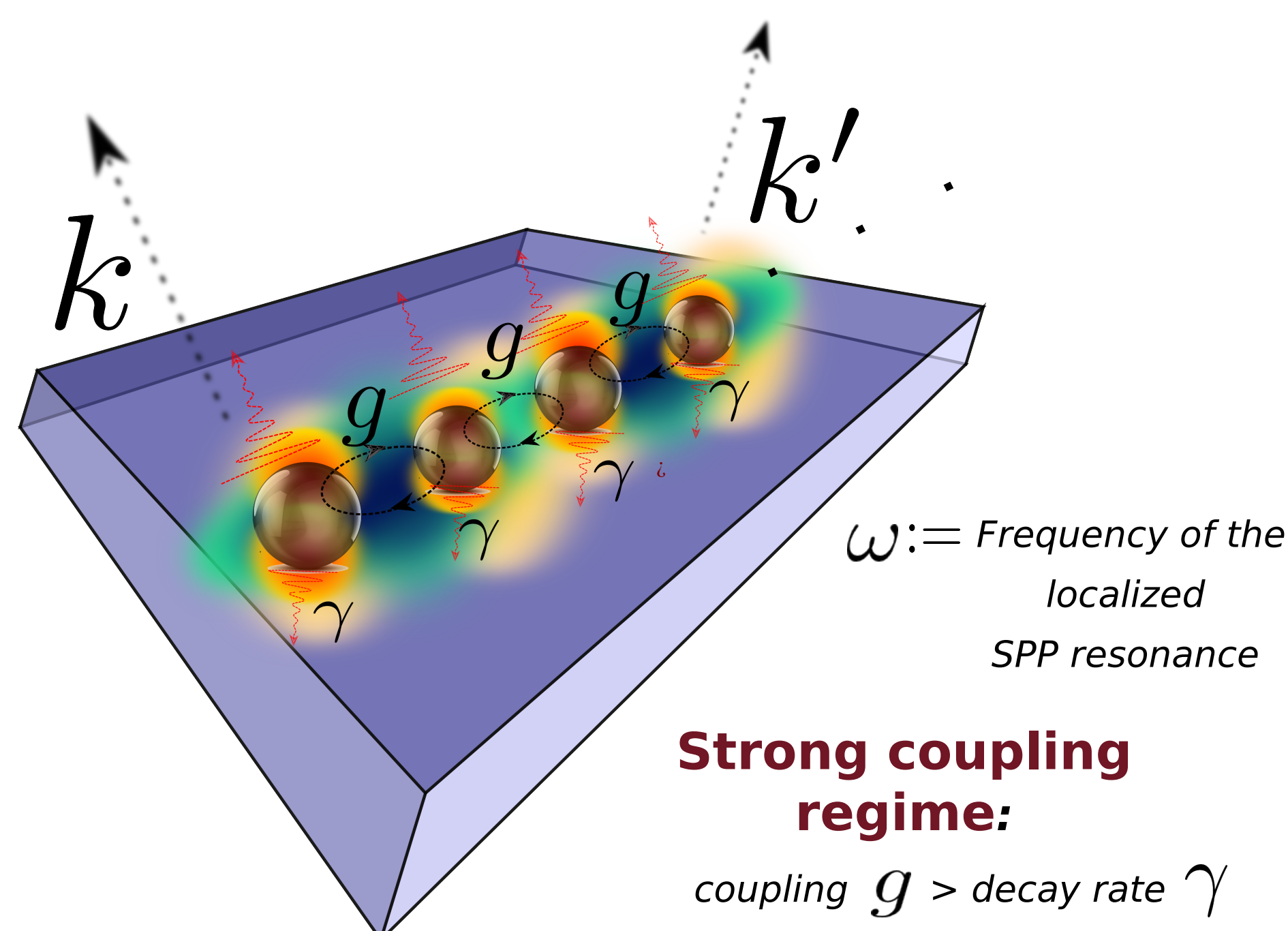
- **Nanoparticles arrays:** Earlier experiments revealed **short propagation lengths**, discouraging the use of such arrays for transport of quantum information



N. Harris et al. J. Phys. Chem. C 113, 2784 (2009);
S. A. Maier et al, PRB 65, 193408 (2002), Phys. Lett. 81,1714 (2002).
M. L. Brongersma et al., PRB 62, R16356 (2000).

Effective model

1D Array of metallic nanoparticles (Ag, Au): coupling of localized plasmons



- **Coupled particle plasmons system:** oscillating dipoles forming a 1D array (open boundary conditions), which interact through a nearest-neighbor dipole coupling ($\hbar = 1$)

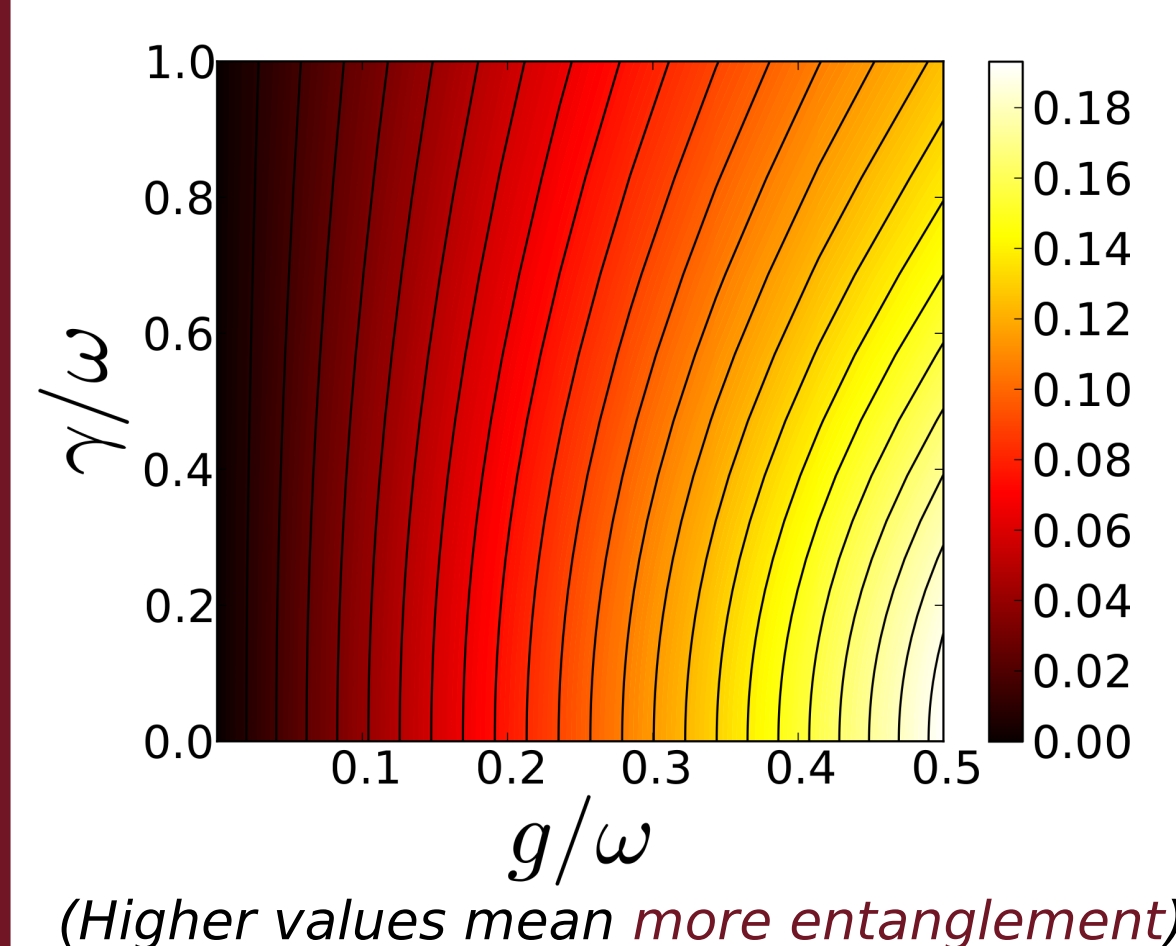
$$H = \sum_{n=1}^N \frac{\omega}{2} (x_n^2 + p_n^2) + \sum_{\langle n,m \rangle} g x_n x_m + \sum_n f_n(t) x_n$$

$$x_n := \text{dipole operator} \quad p_n := \text{conjugate momentum} \\ f_n(t) := \text{external driving}$$

- **Local dissipative dynamics** (spontaneous emission, absorption, etc) modeled through a **master equation** in Lindblad form

$$\partial_t \rho = -i[\rho, H] + \sum_{n=1}^N \frac{\gamma}{2} (2a_n \rho a_n^\dagger - a_n^\dagger a_n \rho - \rho a_n^\dagger a_n) \\ \text{where } a_n = \frac{1}{\sqrt{2}}(x_n + ip_n)$$

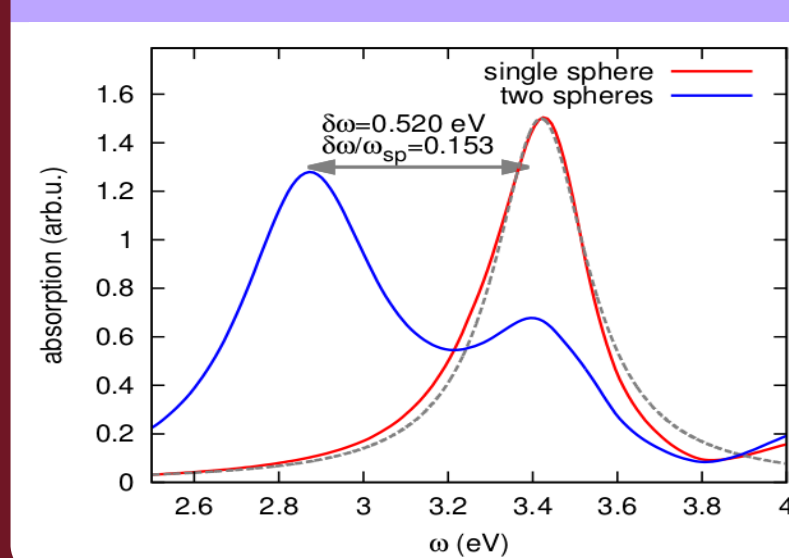
From **steady state (SS)** σ under **no driving**, (**no effect on quantum correlations for Gaussian states**) we calculate amount of entanglement: **logarithmic negativity (LN)**, between 2 partitions of 10 consecutive particles.



(Higher values mean more entanglement)

- **Conclusion:** both parties are **entangled**
- We check that any other bipartition is **entangled as well:** **Many-body multipartite entanglement**
- **Problem if N is large** Needs sufficiently accurate reconstruction of σ (Experimentally difficult)

System parameters: Coupling, dissipation



Lorentzian fit of absorption spectrum for Ag nanopart. (radius $R = 25$ nm, separation $\Lambda = 2$ nm)

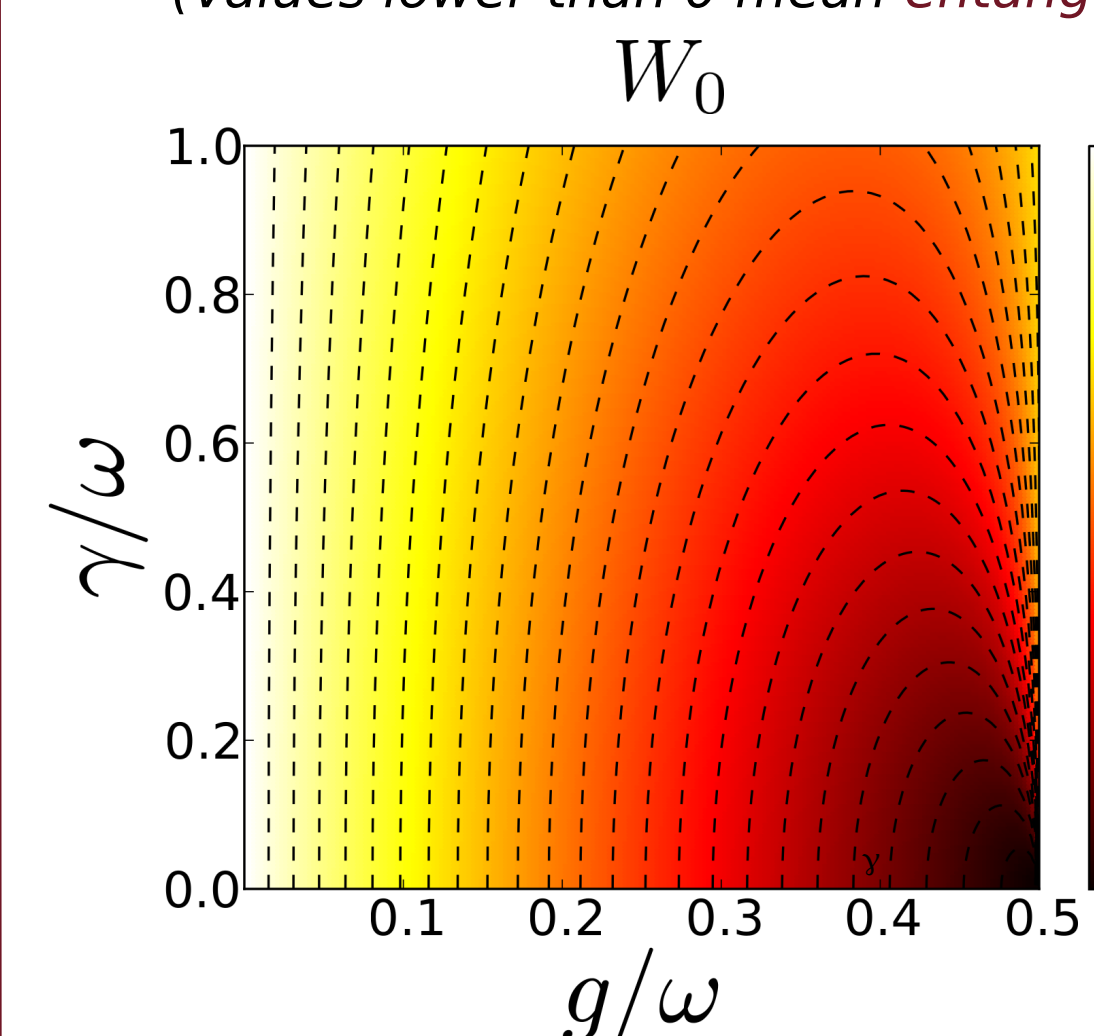
- Decay rate $\frac{\gamma}{\omega} \simeq 0.078$
- Coupling $\frac{g}{\omega} \simeq 0.153$

Detecting entanglement (SS)

- To measure LN is experimentally difficult
- We want to detect the **presence** of Entanglement with the **least number of measurements**, while being **robust to noise and imperfections**: We propose a witness, in the normal-mode (momentum) basis

$$W_k := \min \left\{ 0, \langle \Delta X_k^2 \rangle + \langle \Delta P_{(\pi/\Lambda) - k}^2 \rangle - 1 \right\}$$

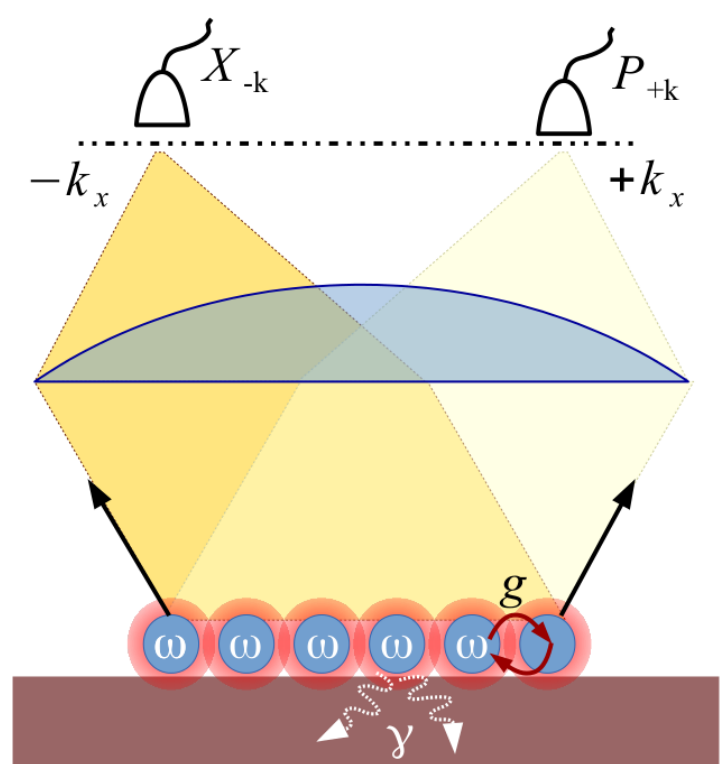
(Values lower than 0 mean *entanglement*)



- Also valid for **non-gaussian states**: irrespective of our model
- We also proof: **weakly dependence** on Temperature and also on k

- **Squeezing in the light with opposite momenta** is a signature of entanglement. Using fitted g and γ , we obtain $\simeq 12\%$ of **squeezing** which is **measurable**

- **Experimental setup:** measurements of fluctuations (homodyne) from **emitted light** collected by a lens at different points in the focal plane



Transport: propagation length

- We describe the **transport** of plasmonic excitations that are introduced by driving a single plasmon close to resonance

$$f_1(t) \sim \sin(\nu t)$$

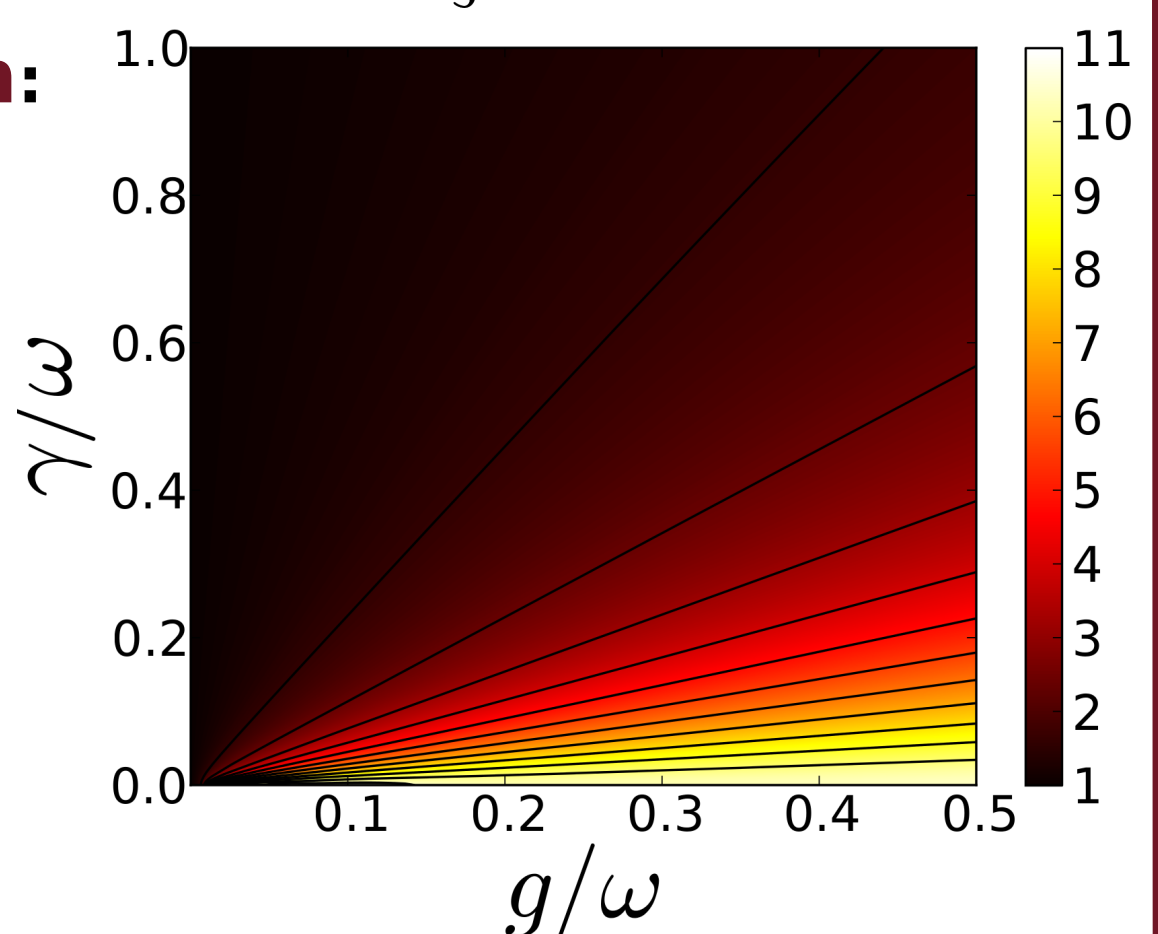
$$\xi(\nu = 0.99\omega)$$

Propagation length:

$$\xi = \frac{\sum_{n=1}^N n \Lambda |\langle x_n \rangle|}{\sum_{n=1}^N |\langle x_n \rangle|}$$

(Provided that, for big N)

$$\langle x_n \rangle \sim e^{-n\Lambda/\xi}$$



- Even when energy transport is not very efficient, (as observed in experiment) the array of plasmons **becomes entangled!**

Summary & Outlook

- We model a setup in which **multipartite entanglement** arises, estimating the strength of the measurement outcomes for realistic setups.
- We develop a quite general measure which detects it, **independent** of our modelling.
- Our predictions could be tested using **present-day state of the art technology**, and some of the ideas can be exported to nanophotonics, matter waves and coupled resonators in superconducting circuits.

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