Numerical Investigation Aeroelasticity Ice accretion

PDE's, optimal design and numerics Benasque 2013

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Knowledge for Tomorrow



Unsteady harmonic motions





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Linearization of Unsteady Reynolds-Averaged Navier-Stokes (URANS) equations

time-dependent coordinate: $\mathbf{x}(t) = [x(t), y(t), z(t)]^T$ velocity field: $\mathbf{U}(\mathbf{x}(t), t) = [u(\mathbf{x}, t), v(\mathbf{x}, t), w(\mathbf{x}, t)]^T$

nondimensionalized integral strong conservation form:

$$\frac{\mathrm{d}}{\mathrm{d}t}\int_{\Omega(t)} \mathbf{W}\mathrm{d}|\Omega| + \int_{\partial\Omega(t)} (\mathbf{f}_c \cdot \mathbf{n} - \mathbf{f}_v \cdot \mathbf{n} - \mathbf{W}\dot{\mathbf{x}} \cdot \mathbf{n}) \,\mathrm{d}|\partial\Omega| = \int_{\Omega(t)} \mathbf{Q}\mathrm{d}|\Omega|, \quad \Omega(t) \subset \mathbb{R}^3$$

with
$$\mathbf{W} := [\rho, \rho u, \rho v, \rho w, \rho E, \rho \check{\mathbf{v}}]^T$$

isolating the time derivative

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$$\mathbf{R}(\mathbf{W}, \mathbf{x}, \dot{\mathbf{x}}) := \int_{\partial \Omega(t)} (\mathbf{f}_c \cdot \mathbf{n} - \mathbf{f}_v \cdot \mathbf{n} - \mathbf{W} \dot{\mathbf{x}} \cdot \mathbf{n}) \, \mathrm{d} |\partial \Omega| - \int_{\Omega(t)} \mathbf{Q} \, \mathrm{d} |\Omega|$$



semi-discrete form (a):

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$$\frac{\mathrm{d}(\mathbf{M}\mathbf{W})}{\mathrm{d}t} + \mathbf{R}(\mathbf{W}, \mathbf{x}, \dot{\mathbf{x}}) = 0$$

M being the mass matrix $\mathbf{M} := \mathbf{M}(\mathbf{x}) := diag(|\Omega_i(t)|), i = 1, ..., N$

linearization around steady state (Taylor-series) - (b)

$$\begin{split} \mathbf{W}(\mathbf{x}, \dot{\mathbf{x}}) &\approx \overline{\mathbf{W}}(\bar{\mathbf{x}}) + \widetilde{\mathbf{W}}(\mathbf{x}, \tilde{\mathbf{x}}, , \dot{\tilde{\mathbf{x}}}) + ..., \quad \|\widetilde{\mathbf{W}}\| \ll \|\overline{\mathbf{W}}\| \\ \mathbf{x}(t) &\approx \bar{\mathbf{x}} + \tilde{\mathbf{x}}(t) + ..., \quad \dot{\bar{\mathbf{x}}} = 0 \\ \mathbf{M}(\mathbf{x}) &\approx \overline{\mathbf{M}}(\bar{\mathbf{x}}) + \frac{\partial \mathbf{M}}{\partial \mathbf{x}} \bigg|_{\bar{\mathbf{x}}} = \overline{\mathbf{M}}(\bar{\mathbf{x}}) + \widetilde{\mathbf{M}}(\bar{\mathbf{x}}, \tilde{\mathbf{x}}) + ..., \\ \mathbf{R}(\mathbf{W}, \mathbf{x}, \dot{\mathbf{x}}) &\approx \mathbf{R}(\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}) + \frac{\partial \mathbf{R}}{\partial \mathbf{W}} \bigg|_{\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}} \widetilde{\mathbf{W}} + \frac{\partial \mathbf{R}}{\partial \mathbf{x}} \bigg|_{\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}} \tilde{\mathbf{x}} + \frac{\partial \mathbf{R}}{\partial \dot{\mathbf{x}}} \bigg|_{\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}} \dot{\bar{\mathbf{x}}} + ... = \\ &= \overline{\mathbf{R}}(\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}) + \widetilde{\mathbf{R}}(\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}, \widetilde{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}) + ... \end{split}$$

after (b) into (a) we obtain:

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$$\overline{\mathbf{M}}(\overline{\mathbf{x}}) + \widetilde{\mathbf{M}}(\overline{\mathbf{x}}, \widetilde{\mathbf{x}}) \frac{\mathrm{d}}{\mathrm{d}t} \left(\overline{\mathbf{W}}(\overline{\mathbf{x}}) + \widetilde{\mathbf{W}}(\overline{\mathbf{x}}, \widetilde{\mathbf{x}}) \right) + \left(\overline{\mathbf{W}}(\overline{\mathbf{x}}) + \widetilde{\mathbf{W}}(\overline{\mathbf{x}}, \widetilde{\mathbf{x}}) \right) \frac{\partial \mathbf{M}(\mathbf{x})}{\partial \mathbf{x}} \Big|_{\overline{\mathbf{x}}} \dot{\overline{\mathbf{x}}} + \overline{\mathbf{R}} \left(\overline{\mathbf{W}}, \overline{\mathbf{x}}, \dot{\overline{\mathbf{x}}} \right) + \widetilde{\mathbf{R}} \left(\overline{\mathbf{W}}, \overline{\mathbf{x}}, \dot{\overline{\mathbf{x}}}, \widetilde{\mathbf{W}}, \widetilde{\mathbf{x}}, \dot{\overline{\mathbf{x}}} \right) = 0$$

using discrete Fourier-series (c) for harmonic excitations - transformation into frequency domain

$$\widetilde{\mathbf{W}} := \sum_{n=1}^{\infty} \widehat{\mathbf{W}}_n e^{in_h \omega t}, \ \widehat{\mathbf{W}}_n \in \mathbb{C}$$
$$\widetilde{\mathbf{x}} = \sum_{n=1}^{\infty} \widehat{\mathbf{x}}_n e^{in_h \omega t}, \ \widehat{\mathbf{x}}_n \in \mathbb{C}$$

x is known which is a mode (bending, torsion, ...)



after (c) into (b) we obtain:

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$$\begin{split} \left[i\omega \overline{\mathbf{M}}(\bar{\mathbf{x}}) + \frac{\partial \mathbf{R}}{\partial \mathbf{W}} \Big|_{\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}} \right] \widehat{\mathbf{W}} e^{i\omega t} + \\ + \left[\frac{\partial \mathbf{R}}{\partial \mathbf{x}} \Big|_{\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}} + i\omega \left(\frac{\partial \mathbf{R}}{\partial \dot{\mathbf{x}}} \Big|_{\overline{\mathbf{W}}, \bar{\mathbf{x}}, \dot{\bar{\mathbf{x}}}} + \overline{\mathbf{W}}(\bar{\mathbf{x}}) \frac{\partial \mathbf{M}}{\partial \mathbf{x}} \Big|_{\bar{\mathbf{x}}} \right) \right] \hat{\mathbf{x}} e^{i\omega t} + \\ + \frac{\partial \mathbf{M}(\mathbf{x})}{\partial \mathbf{x}} \Big|_{\bar{\mathbf{x}}} \left(-\omega^2 + i\omega \right) \hat{\mathbf{x}} \widehat{\mathbf{W}} e^{i2\omega t} = 0 \end{split}$$

Implementation! - neglect higher order term - solve for hat W

$$\left[i\omega\overline{\mathbf{M}}(\bar{\mathbf{x}}) + \frac{\partial \mathbf{R}}{\partial \mathbf{W}}\Big|_{\overline{\mathbf{W}},\bar{\mathbf{x}},\dot{\bar{\mathbf{x}}}}\right]\widehat{\mathbf{W}} = -\left[\frac{\partial \mathbf{R}}{\partial \mathbf{x}}\Big|_{\overline{\mathbf{W}},\bar{\mathbf{x}},\dot{\bar{\mathbf{x}}}} + i\omega\left(\frac{\partial \mathbf{R}}{\partial \dot{\mathbf{x}}}\Big|_{\overline{\mathbf{W}},\bar{\mathbf{x}},\dot{\bar{\mathbf{x}}}} + \overline{\mathbf{W}}(\bar{\mathbf{x}})\left.\frac{\partial \mathbf{M}}{\partial \mathbf{x}}\Big|_{\bar{\mathbf{x}}}\right)\right]\widehat{\mathbf{x}}$$



DLR F12 - RANS Ma = 0.85, Re = 1.28 mill., α = 0.0°, amp = 0.5°, red. f = 0.068

Dynamic response



DLR

Simulation	Magnitude [rad]	Phase [rad/s]
URANS	8.9966	-4.20E-02
LFD	8.9519	-4.24E-02
URANS	4.5619	-2.580
LFD	4.5436	-2.585



for higher harmonics nh we obtain:

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$$\begin{split} \left[\left(in_{1}\omega\overline{\mathbf{M}} + \frac{\partial \mathbf{R}}{\partial \mathbf{W}} \right) \widehat{\mathbf{W}}_{1} + \left[\frac{\partial \mathbf{R}}{\partial \mathbf{x}} + in_{1}\omega\left(\frac{\partial \mathbf{R}}{\partial \dot{\mathbf{x}}} + \overline{\mathbf{W}}\frac{\partial \mathbf{M}}{\partial \mathbf{x}}\right) \right] \widehat{\mathbf{x}}_{1} \right] e^{in_{1}\omega t} + \\ + \dots + \\ + \left[\left(in_{h}\omega\overline{\mathbf{M}} + \frac{\partial \mathbf{R}}{\partial \mathbf{W}} \right) \widehat{\mathbf{W}}_{n_{h}} + \left[\frac{\partial \mathbf{R}}{\partial \mathbf{x}} + in_{h}\omega\left(\frac{\partial \mathbf{R}}{\partial \dot{\mathbf{x}}} + \overline{\mathbf{W}}\frac{\partial \mathbf{M}}{\partial \mathbf{x}}\right) \right] \widehat{\mathbf{x}}_{n_{h}} \right] e^{in_{h}\omega t} + \\ + \frac{\partial \mathbf{M}}{\partial \mathbf{x}} \left((-\omega^{2} + i\omega)n_{1}\widehat{\mathbf{x}}_{1}\widehat{\mathbf{W}}_{1} \right) e^{in_{2}\omega t} + \\ + \dots + \\ + \frac{\partial \mathbf{M}}{\partial \mathbf{x}} \left(-n_{h}^{2}\omega^{2} + n_{h}i\omega \right) \widehat{\mathbf{x}}_{n_{h}} \widehat{\mathbf{W}}_{n_{h}} e^{i(2n_{h})\omega t} = 0 \end{split}$$

