

# *Imprint of Exotic Matter on Quasi-normal modes of Neutron Stars*

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# *Imprint of Exotic Matter on Quasi-normal modes of Neutron Stars*

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2. Numerical Method
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  - a) Axial and polar fundamental wI mode
  - b) Axial wII mode
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# 1. Introduction

## 1. Introduction

- *We consider an static spherical space-time*

$$ds^2 = - e^{2\psi(r)} dt^2 + e^{2\lambda(r)} dr^2 + r^2 [d\theta^2 + \sin^2\theta d\varphi^2]$$

- *Two regions*

*Interior of the star perfect fluid with eos  $p = p(\rho)$   
Empty outside*

- *Following Chandrasekhar and Ferrari we make perturbations over this metric, expanding in tensor harmonics, and taking into account both axial and polar perturbations:*

*Under parity transformation  $\theta \rightarrow \pi - \theta$ ,  $\varphi \rightarrow \varphi + \pi$*

*Axial transform like  $(-1)^{l+1}$ .*

*Polar transform like  $(-1)^{l+2}$*

$$g_{\mu\nu} = g_{\mu\nu}^{(0)} + g_{\mu\nu}^{(1)}$$

$$T_{\mu\nu} = T_{\mu\nu}^{(0)} + T_{\mu\nu}^{(1)}$$

- We solve Einstein Equations up to first order in perturbation theory.

For static functions:

$$\frac{dM}{dr} = 4\pi r^2 \rho \quad \frac{dp}{dr} = \frac{(4\pi r^3 p + M)(\rho + p)}{(r - 2M)r} \quad \frac{dv}{dr} = \frac{-1}{(\rho + p)} \frac{dp}{dr}$$

First order perturbation: AXIAL

Regge-Wheeler equation

$$\frac{d^2 Z^{lm}}{dr_*^2} + [\omega - V^l(r)] Z^{lm} = 0$$

$$V^l(r) = \frac{e^{2v}}{r^3} [l(l+1)r + 4\pi r^3(\rho + p) - 6M]$$

Axial modes do not couple to matter perturbations

First order perturbation: POLAR

Inside the star:

Lindblom equations  $\{V^{lm}, H_{(1)}^{lm}, K^{lm}, W^{lm}\}$

Outside the star:

Zerilli equation  $\frac{d^2 Z^{lm}}{dr_*^2} + [\omega - V^l(r)] Z^{lm} = 0$

$$V^l(r) = \frac{2(r - 2M)}{r^4(nr + 3M)^2} \left[ n^2(n + 1)r^3 + 3Mn^2r^2 + 9M^2nr + 9M^3 \right]$$

$$n = \frac{(l - 1)(l - 2)}{2}$$

Polar modes couple to pressure and density perturbations

- A general solution will be a *superposition of an outgoing wave and an ingoing wave*:

$$\lim_{r \rightarrow \infty} Z_{out} \sim e^{-i\omega r_*}$$

$$\lim_{r \rightarrow \infty} Z_{in} \sim e^{i\omega r_*}$$

- The outgoing contribution diverges at radial infinity
- The ingoing contribution tends exponentially to zero at radial infinity



*A small contamination of the outgoing wave behaviour faraway of the star gives rise to an important ingoing wave contribution nearby the star.*

*Families of modes for static (realistic) stars.*

*Space-time modes: Axial and Polar*

*wl modes: curvature modes. Infinite in number. ~10 khz ~0.1 ms*

*wll modes: interface modes. ~1 khz ~0.1 ms*

*Fluid modes: Polar*

*f mode: fundamental mode. Only one. ~1 khz ~0.1 s*

*p modes: pressure modes. Infinite in number. ~1 khz ~1 s*

*Other modes we do not consider (gravity modes, rotational modes...)*

## 2. Numerical Method

*Exterior part of the solution:*

$$\lim_{r \rightarrow \infty} Z_{out} \sim e^{-i\omega r_*} = e^{-i\omega R^{r_*}} e^{r_*/\tau}$$

- *Two main problems for outgoing quasi-normal modes:*

*Infinite oscillations*

*Divergence towards infinity*

- *We can combine two methods in order to deal with these issues:  
(Andersson et al 1995)*

*Integration of the phase of the Regge-Wheeler/Zerilli function*

*Exterior Complex Scaling*

## 2. Numerical Method

Studying the phase function we deal with the oscillations

$$g = \frac{Z'}{Z}$$

$$\partial_r g + g^2 \frac{\omega^2 r^2}{(2 - 2M)^2} + \frac{2M}{r(r - 2M)} g - \frac{l(l + 1)r - 6M}{r^4(r - 2M)} = 0 \quad (\text{axial})$$

Without oscillations we can compactify the radial coordinate.

Far from the star:

$$g \sim i\omega \frac{Ae^{i\omega r} - Be^{-i\omega r}}{Ae^{i\omega r} + Be^{-i\omega r}}$$

Incoming wave:  $g \sim i\omega$

Mixture:  $g \sim -i\omega$

We can not distinguish between mixture and outgoing wave behaviour.

## 2. Numerical Method

*Exterior complex scaling:*

*(J. Aguilar and J. Combes, Commun. Math. Phys. 22 269 (1971))*

*(G. Alvarez et al, PRA 44, 3060 (1991))*

$$r = a + ye^{-i\alpha}$$

*Parametrize the radius along a line in the complex plane.*

*The boundary condition is dependent of  $\alpha$ .*

*We use  $\alpha$  as a free parameter, allowing us to make:*

$$\omega_R \sin \alpha + \omega_I \cos \alpha > 0$$

*Outgoing wave:  $g \sim -i\omega$*

*Mixture:  $g \sim i\omega$*

*Using  $\alpha$  as a free parameter for the integration path, the boundary condition for the phase fixes completely the **outgoing wave behaviour***

### *Realistic Equations of State:*

*Important constraint to realistic equations of state after the discovery of the 2  $M_{\text{sun}}$  pulsar PSR J164-2230*

- 1. Piece-wise polytrope interpolation (Read et al 2008)*
- 2. Hermite interpolation of tabulated equations of state*

- Hyperon matter.*

*GNH3 - N. Glendenning, APJ 293, 470 (1985)*

*H4 - B. D. Lackey, M. Nayyar, and B. J. Owen, PRD 73 (2006)*

*WCS1-2 - S. Weissenborn, D. Chatterjee, and J. Schaner-Bielich, PRC 85 (2012)*

*BHZBM - Bednarek, I., Haensel, P., Zdunik, J. L., Bejger, M., and Manka, R., A&A 543 (2012)*

- *Quark matter*

### *Hybrid stars:*

*ALF2-4* - M. Alford, M. Braby, M. Paris, and S. Reddy, *APJ* 629 (2005)

*WSPHS3* - S. Weissenborn, I. Sagert, G. Pagliara,  
M. Hempel. and J. Schaner-Bielich, *APJL* 740 (2011)

### *Hybrid stars with hyperons:*

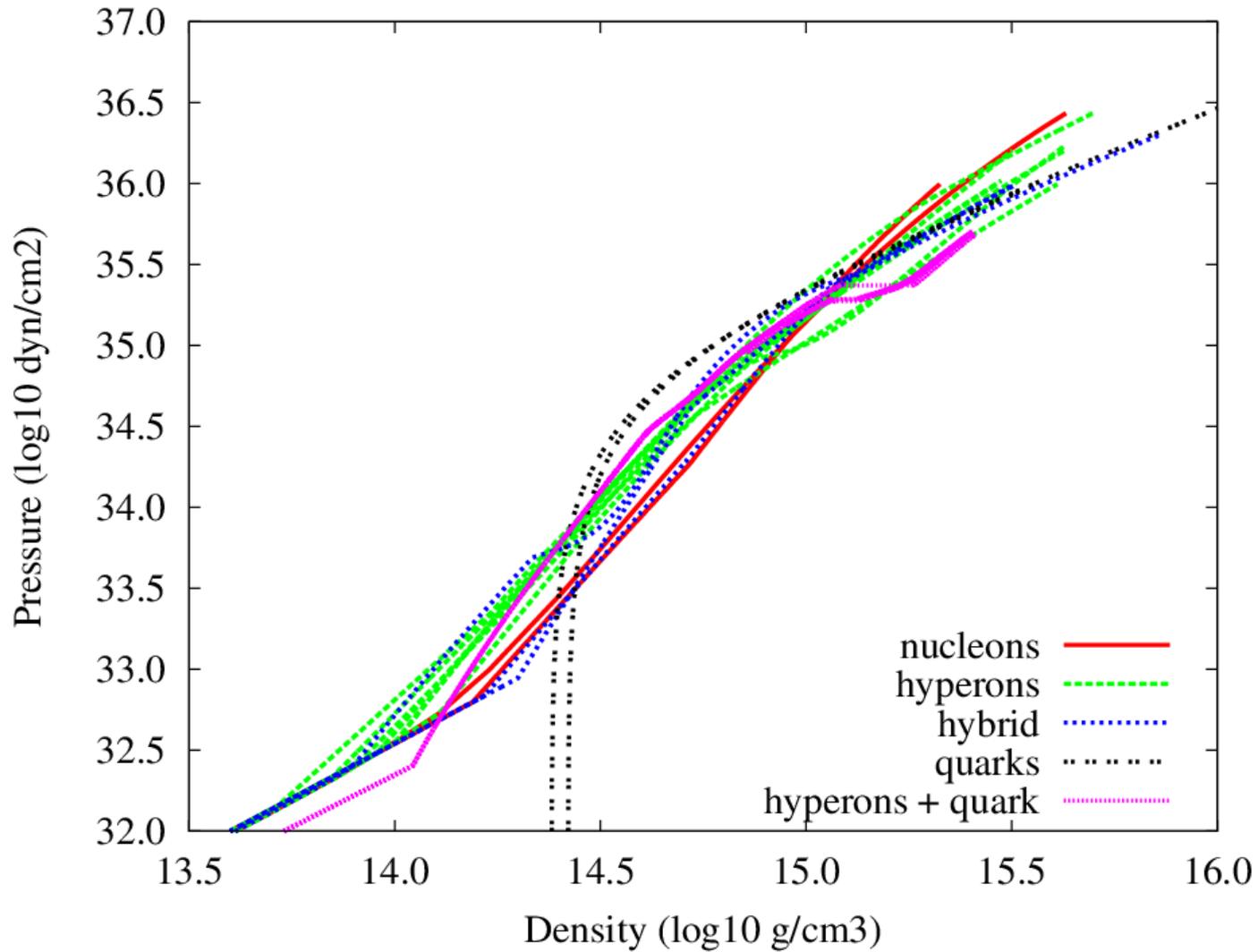
*BS1-4* - Bonanno, Luca and Sedrakian, Armen, *A&A* 539 (2012)

### *Quark stars:*

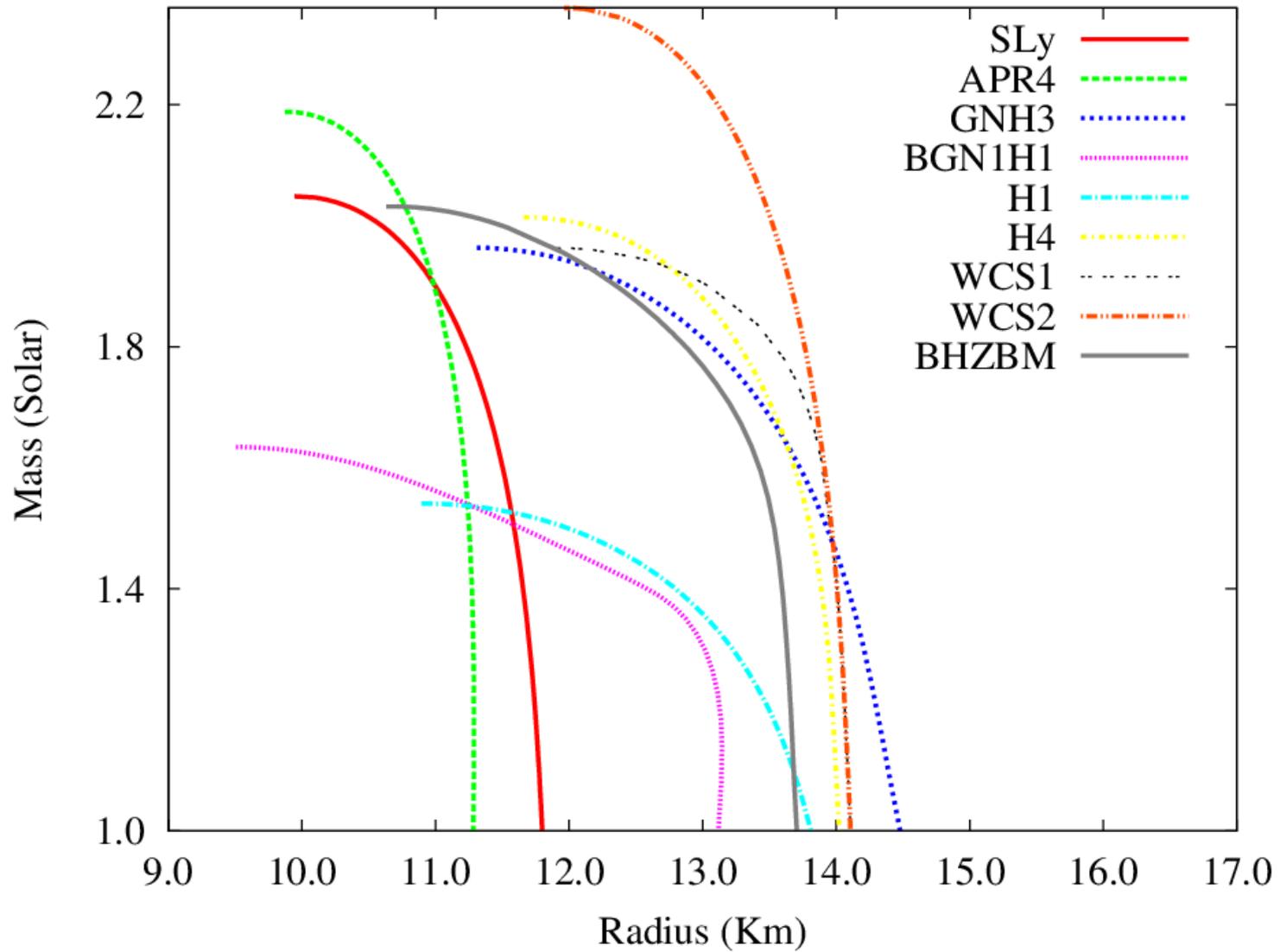
*WSPHS1-2* - S. Weissenborn, I. Sagert, G. Pagliara,  
M. Hempel. and J. Schaner-Bielich, *APJL* 740 (2011)

- *Nuclear matter: SLy and APR4*

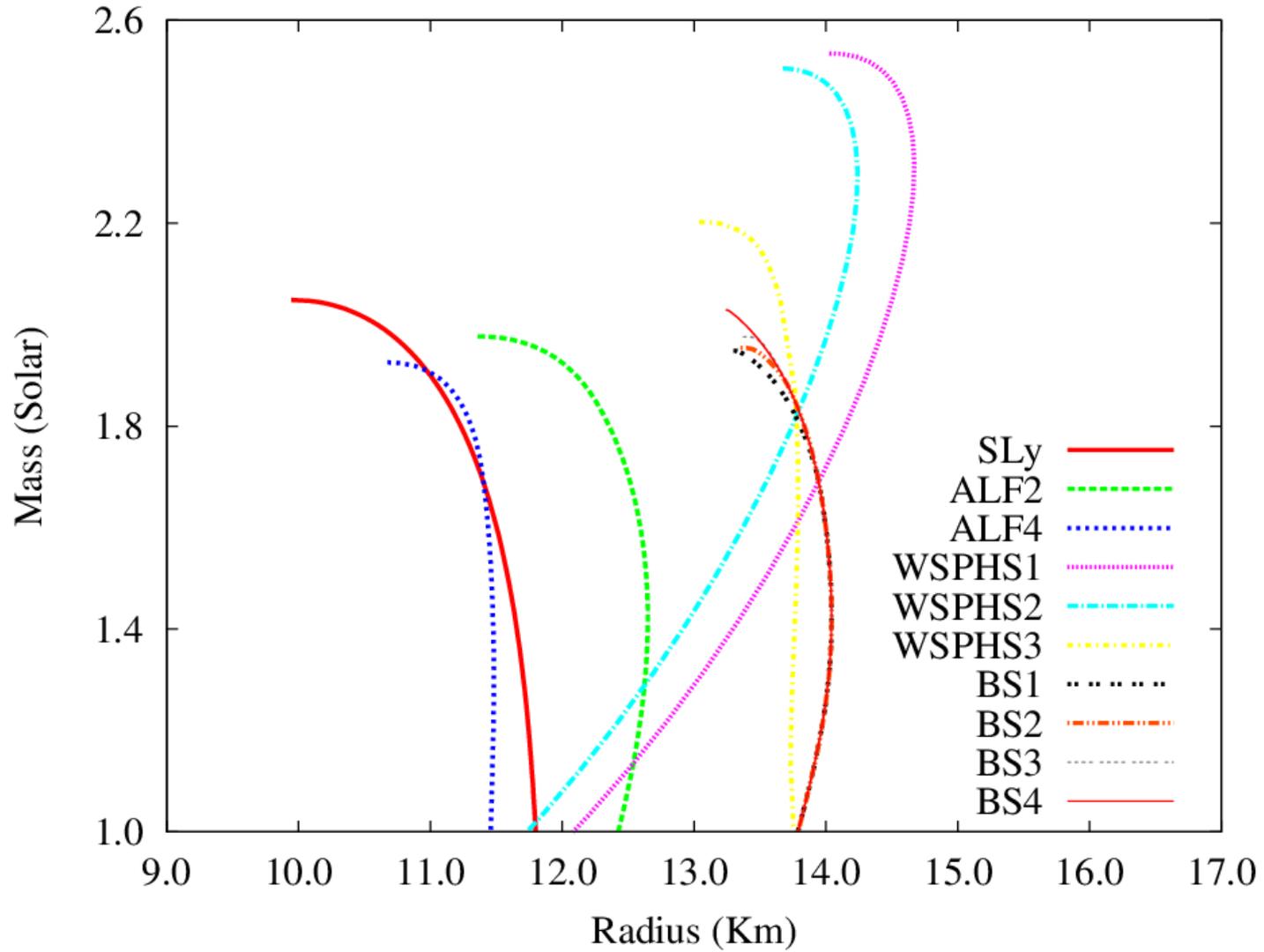
## 2. Numerical Method



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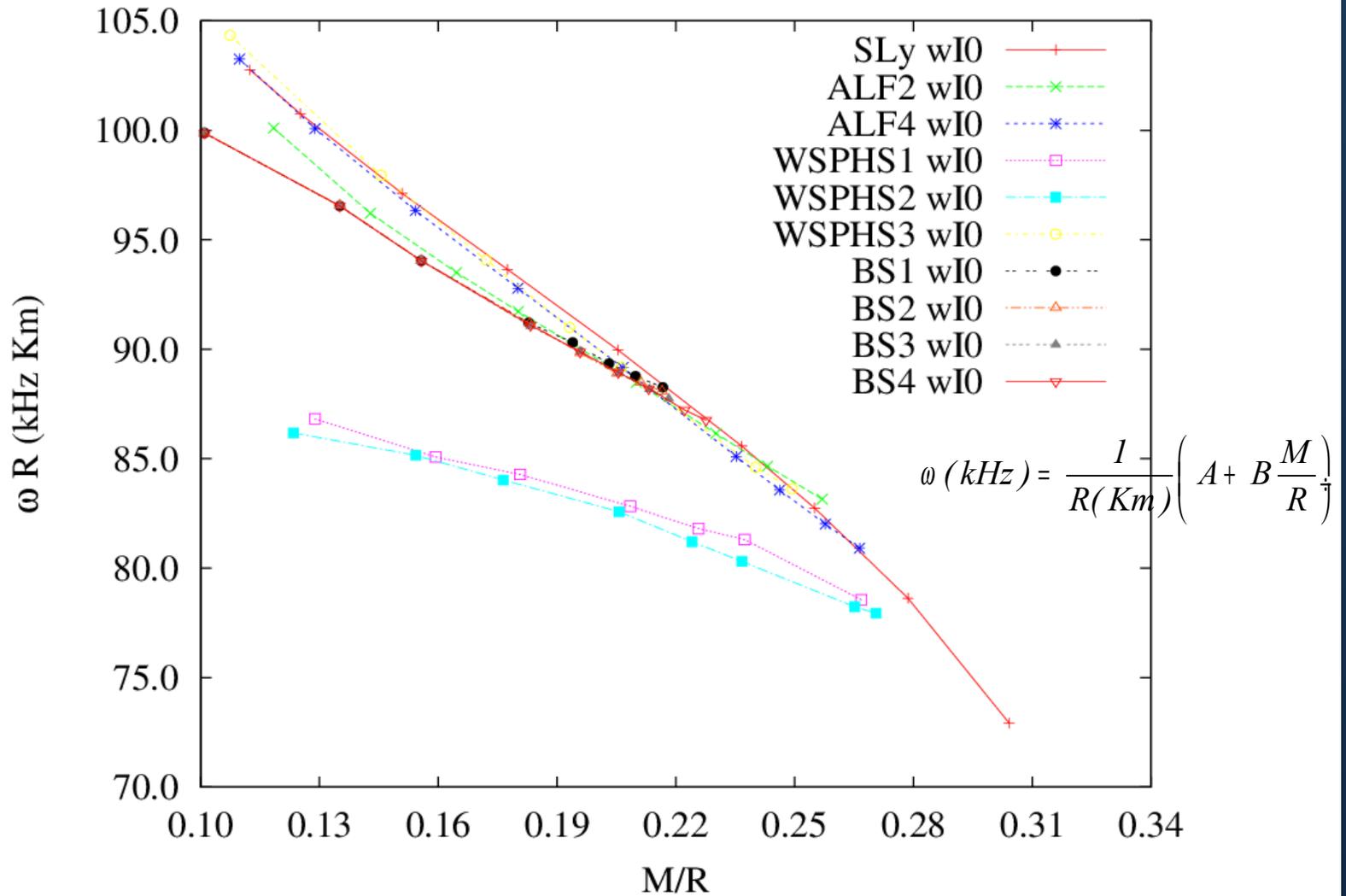
## 3. Results

a) Axial and polar fundamental wI mode



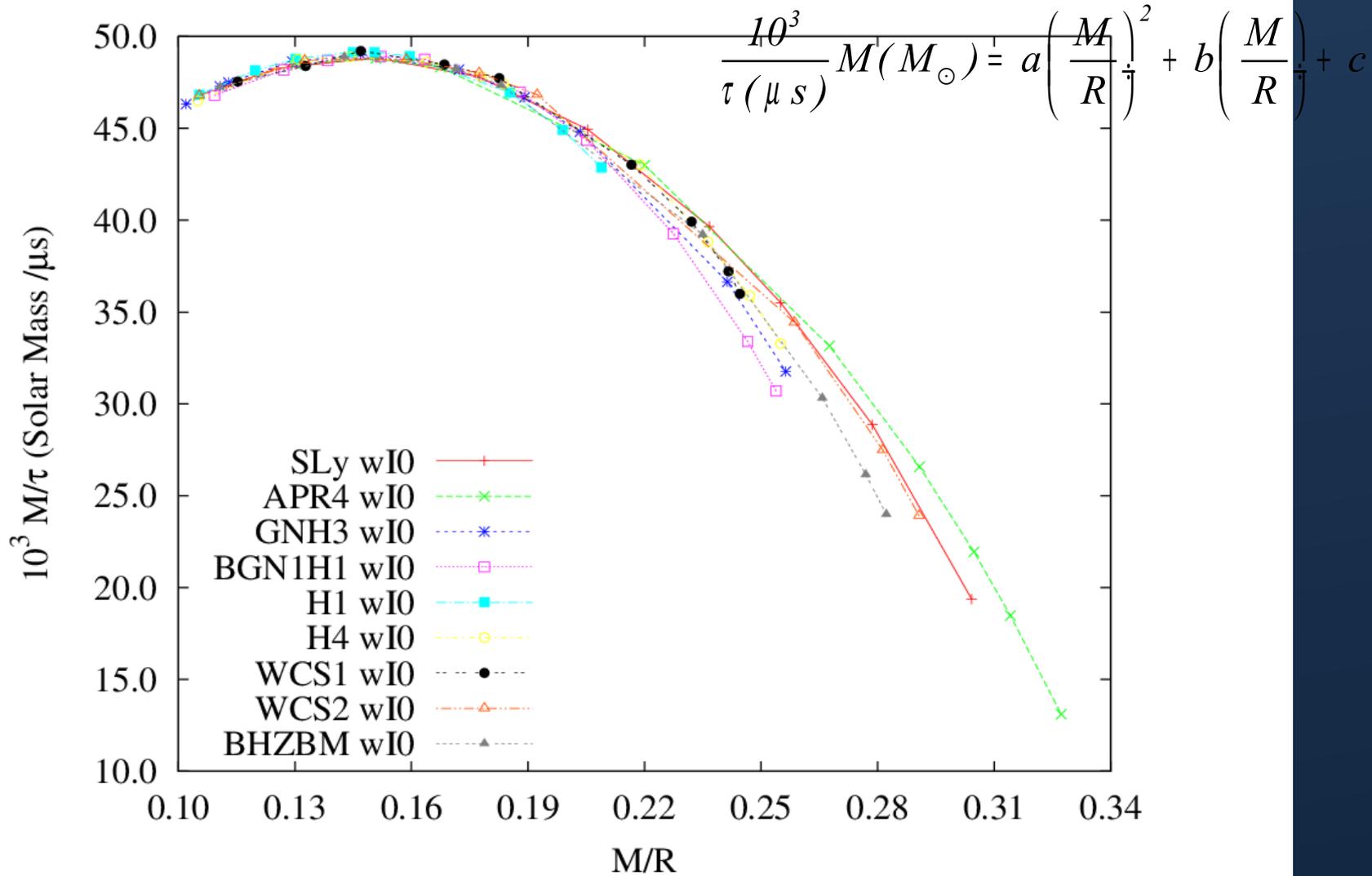
### 3. Results: a) wI0 modes

Axial fundamental wI mode: *Hybrid – quark matter*



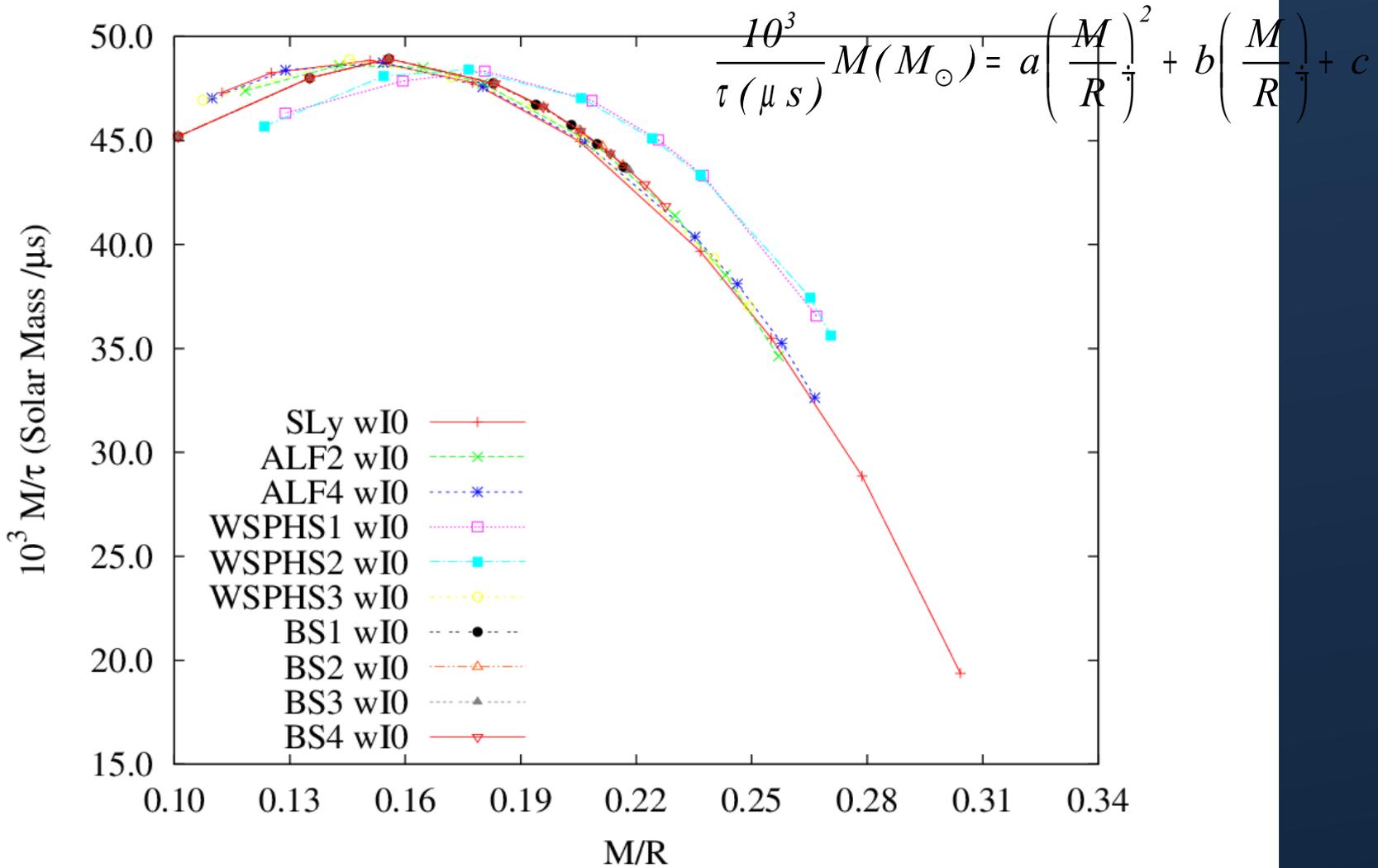
3. Results: a) wI0 modes

Axial fundamental wI mode: *Hyperon matter*



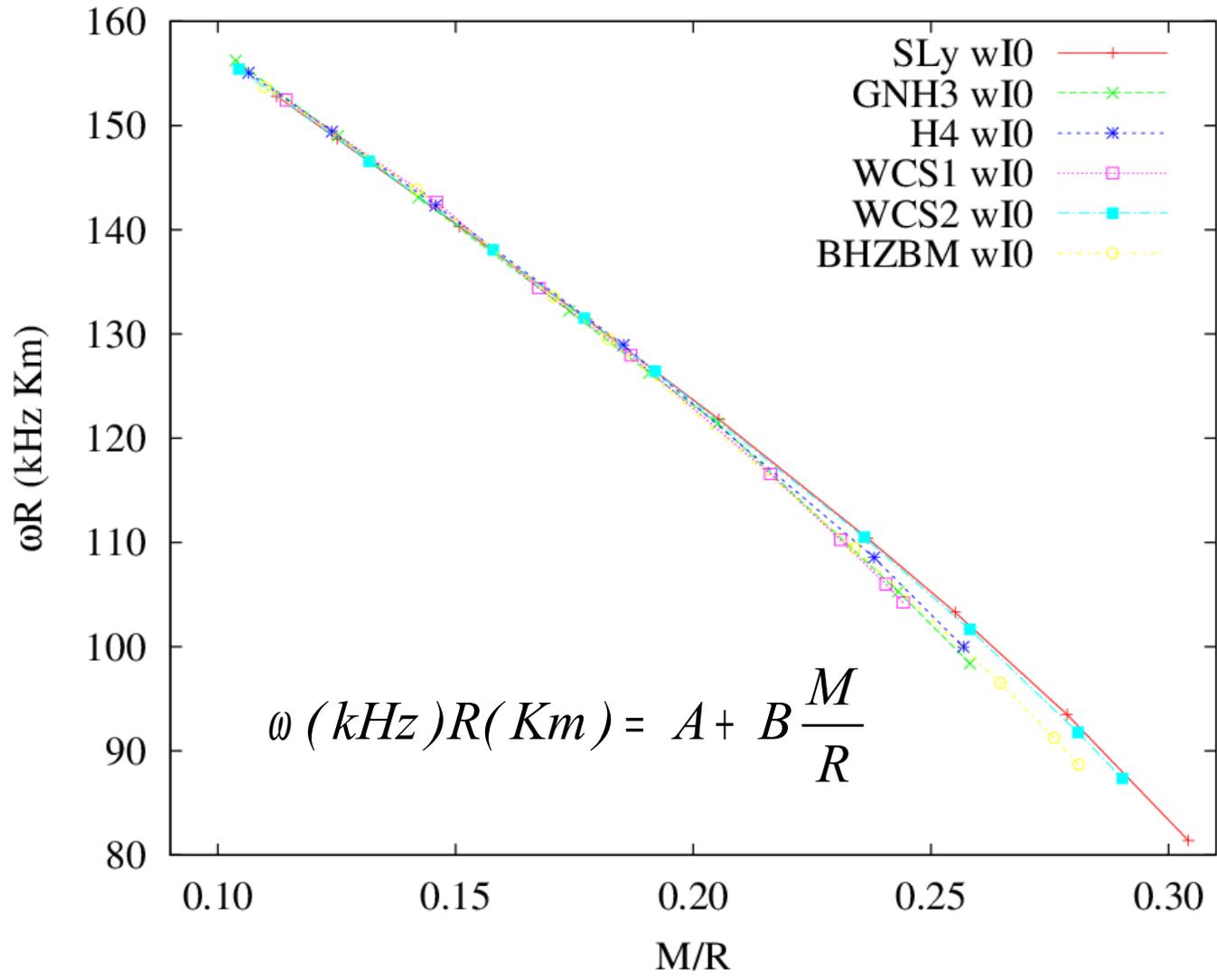
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3. Results: a) wI0 modes

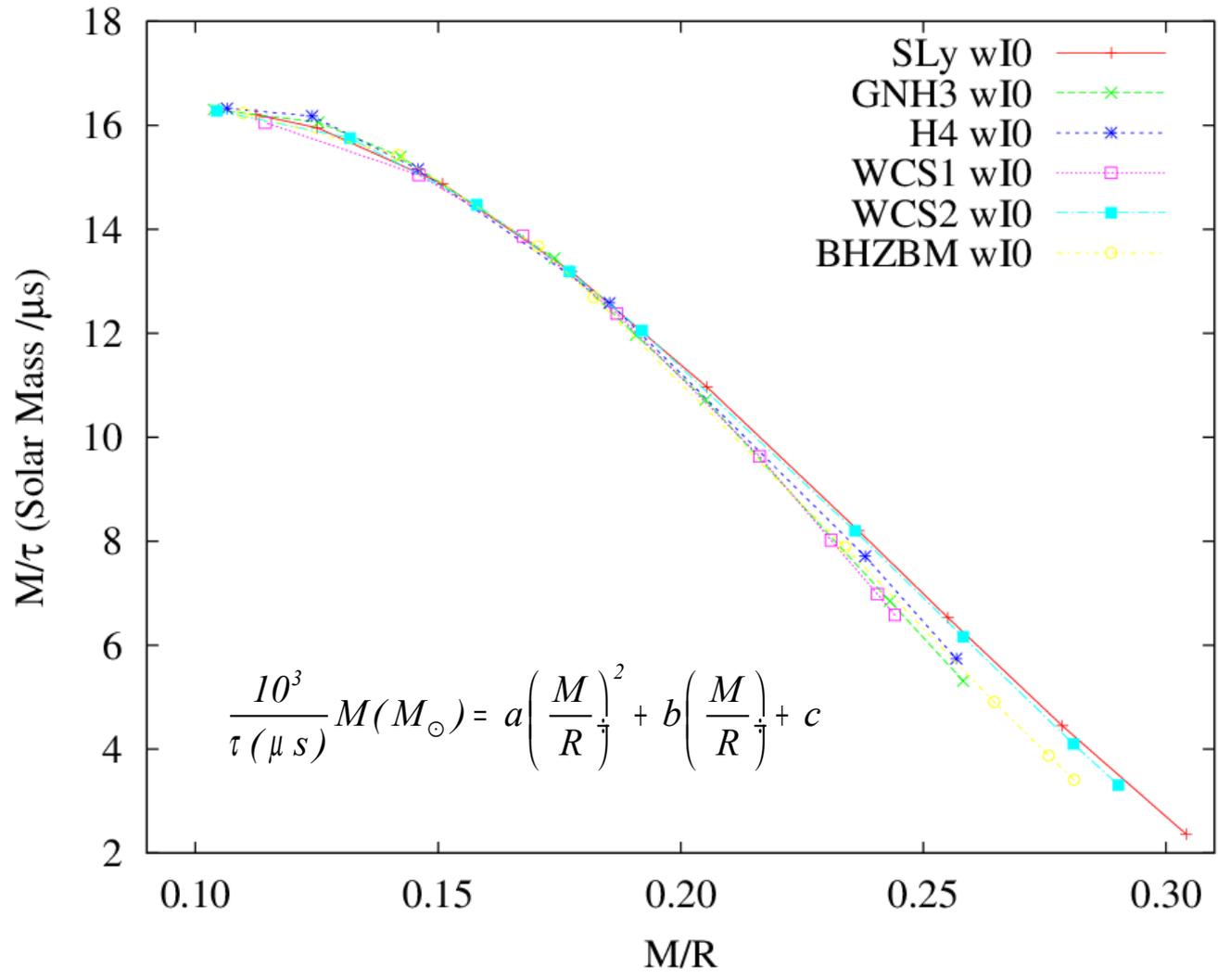
Polar fundamental wI mode: *Hyperon matter*





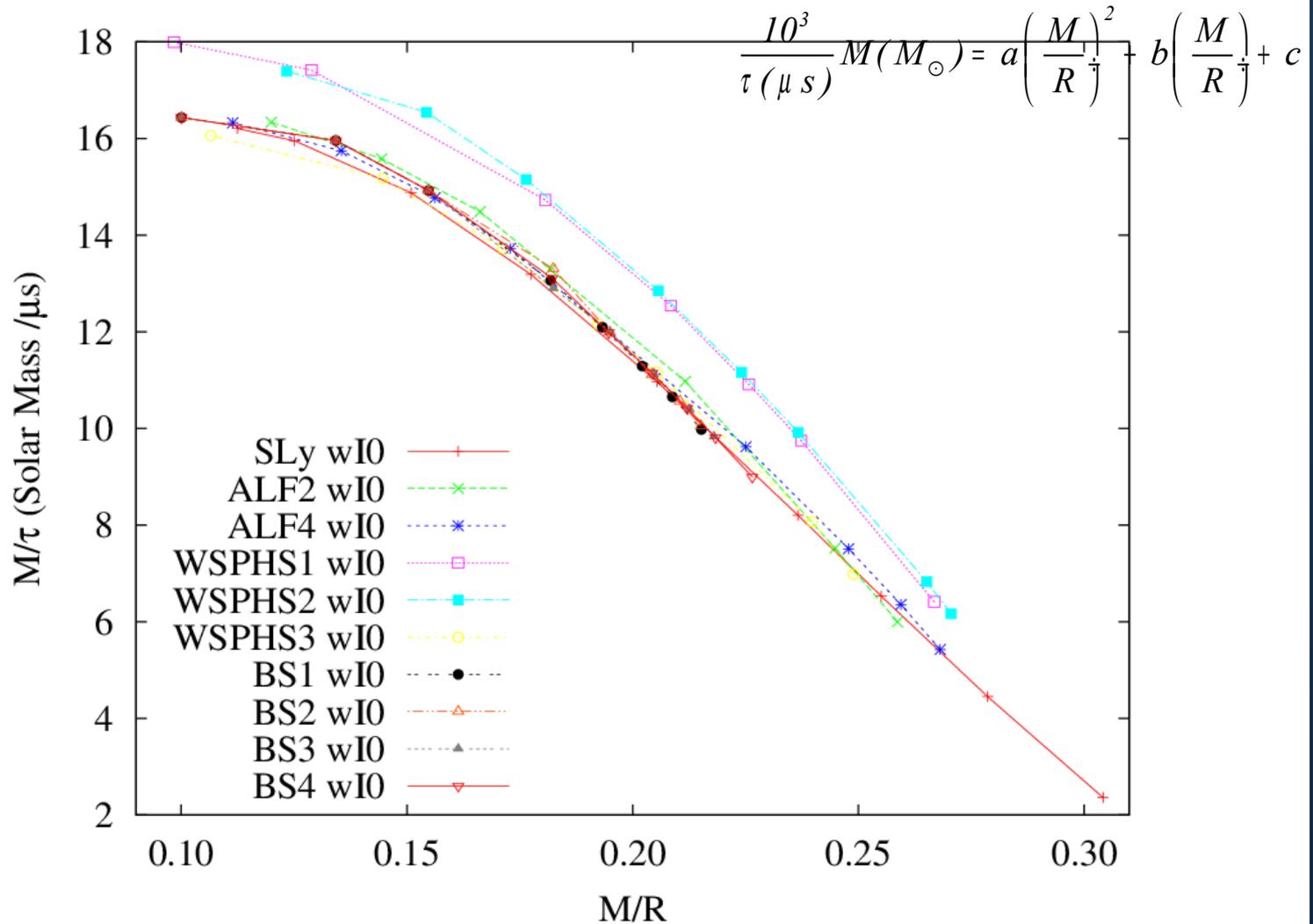
### 3. Results: a) wI0 modes

Polar fundamental wI mode: *Hyperon matter*



### 3. Results: a) wI0 modes

Polar fundamental wI mode: *Hybrid – quark matter*

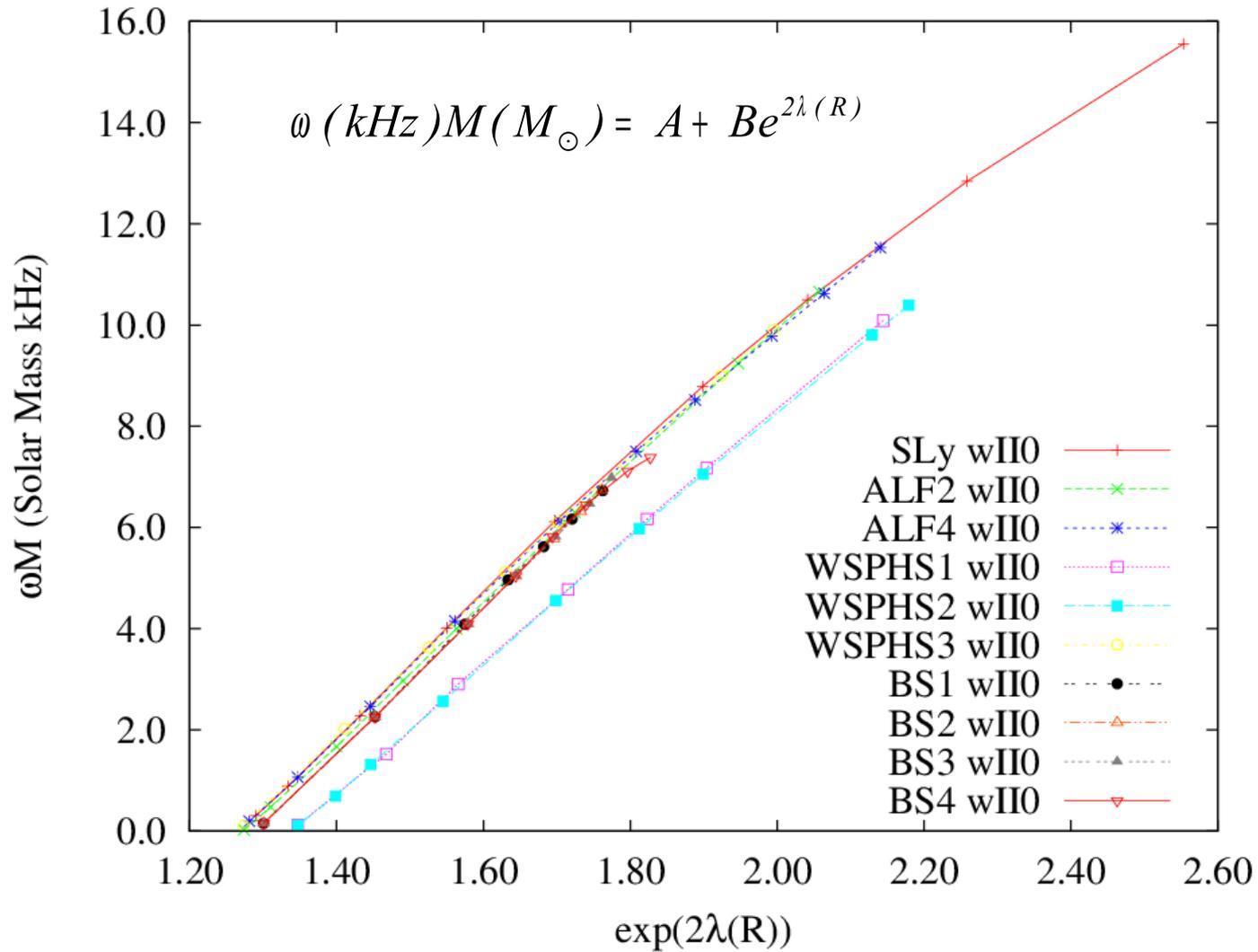


b) Axial fundamental wII mode



### 3. Results: b) wII0 modes

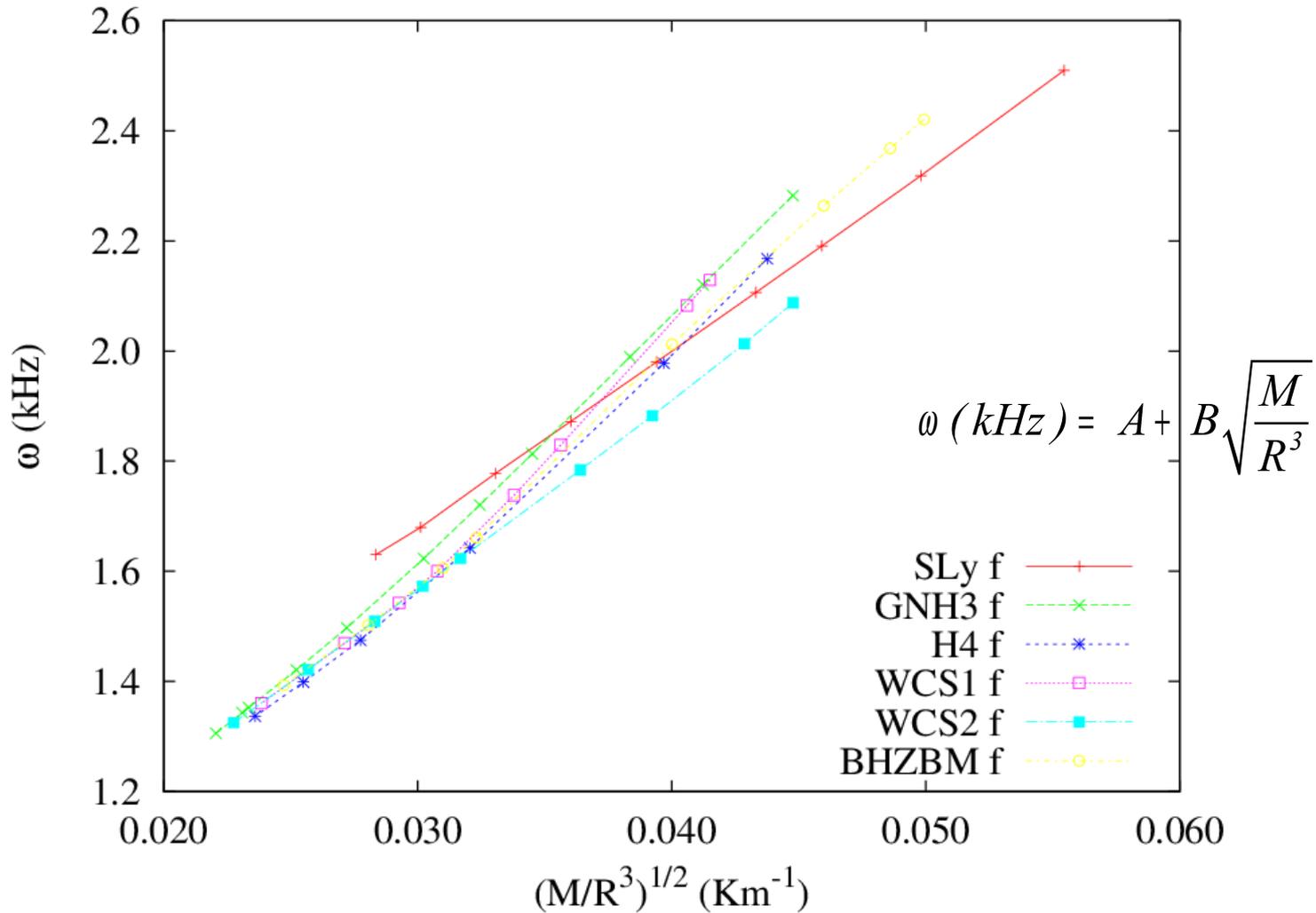
Axial fundamental wII mode: *Hybrid – quark matter*



c) Polar f mode

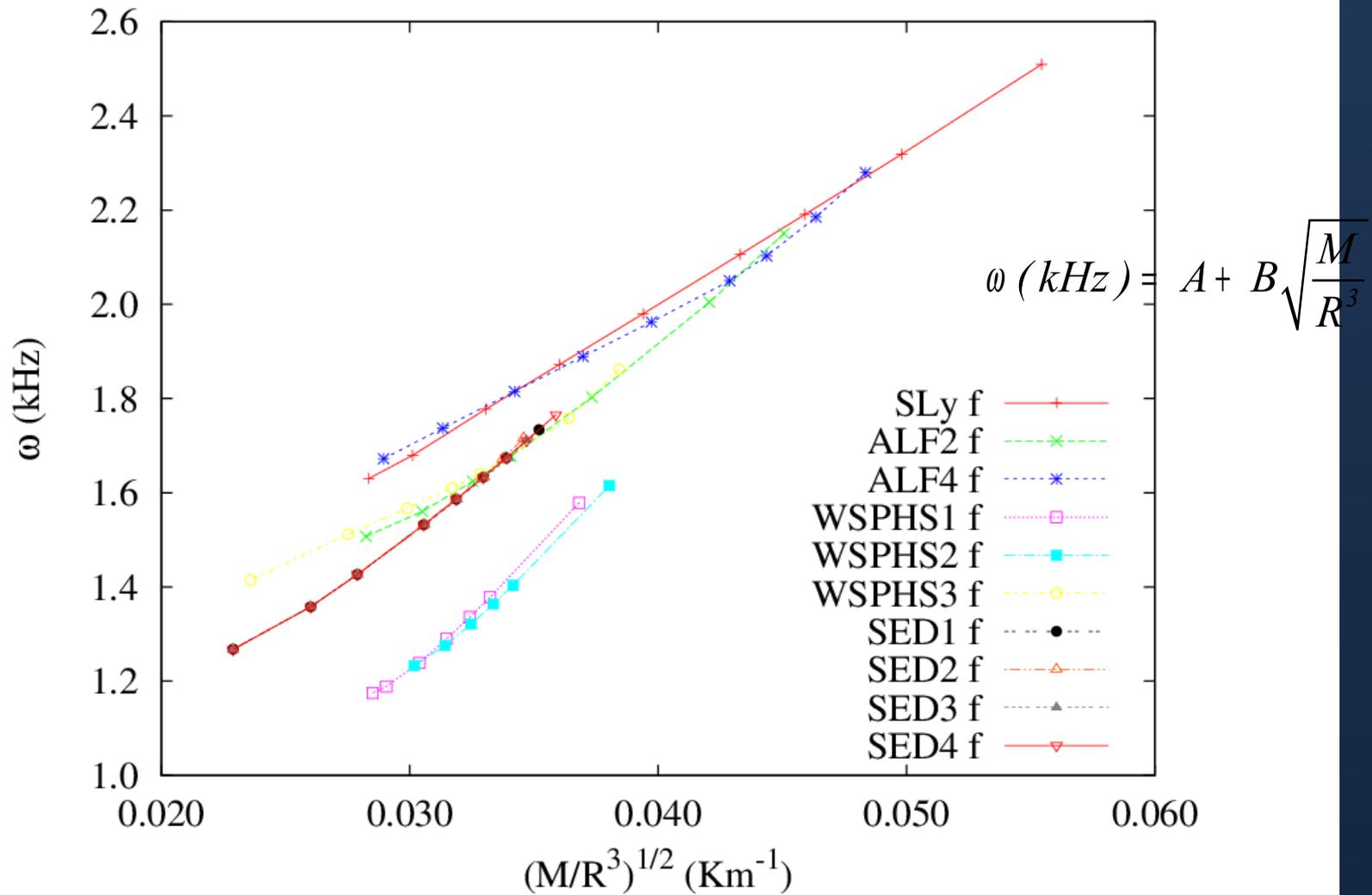
### 3. Results: c) f modes

Polar f mode: *Hyperon matter*



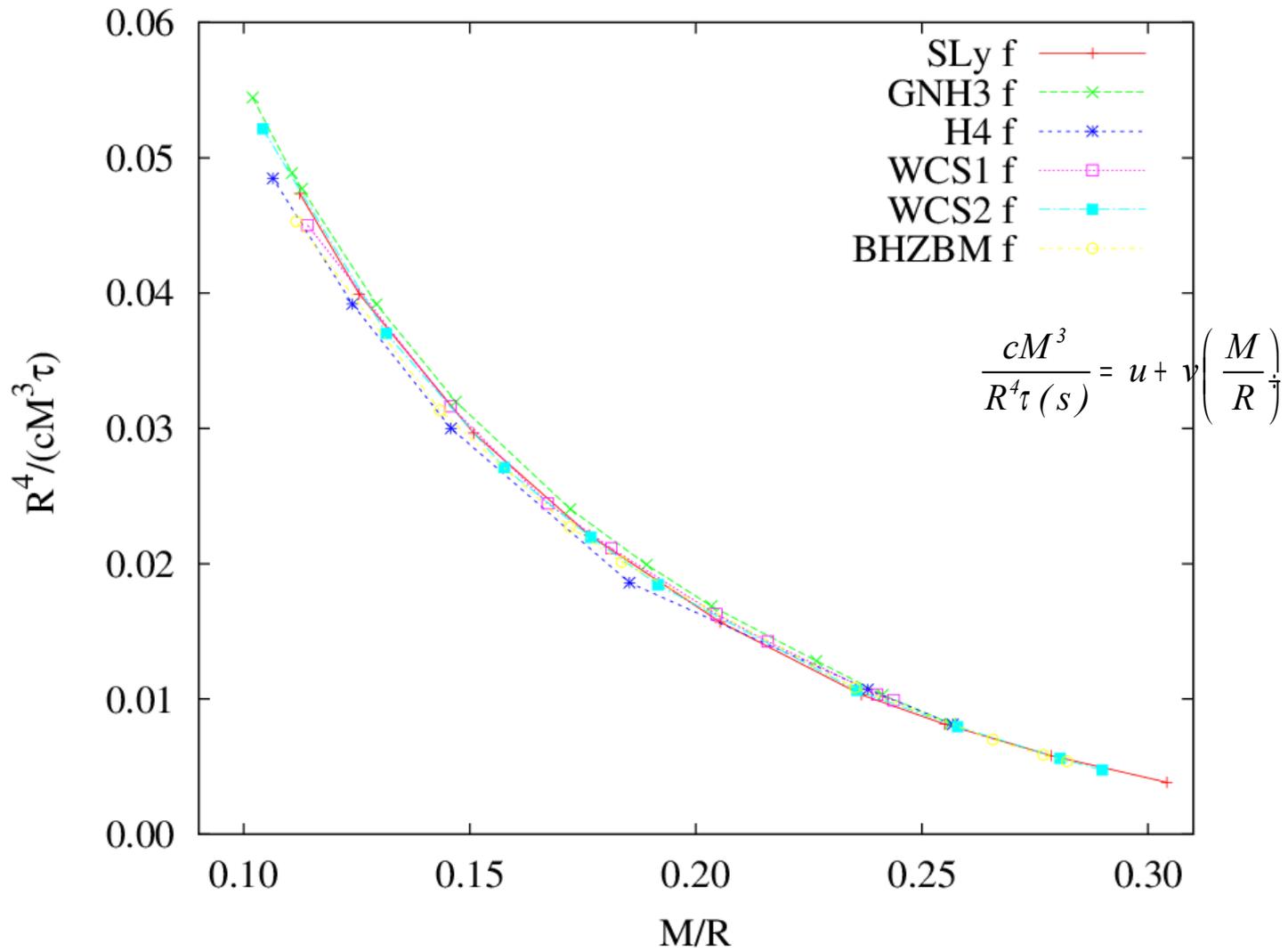
### 3. Results: c) f modes

Polar f mode: *Hybrid – quark matter*



### 3. Results: c) f modes

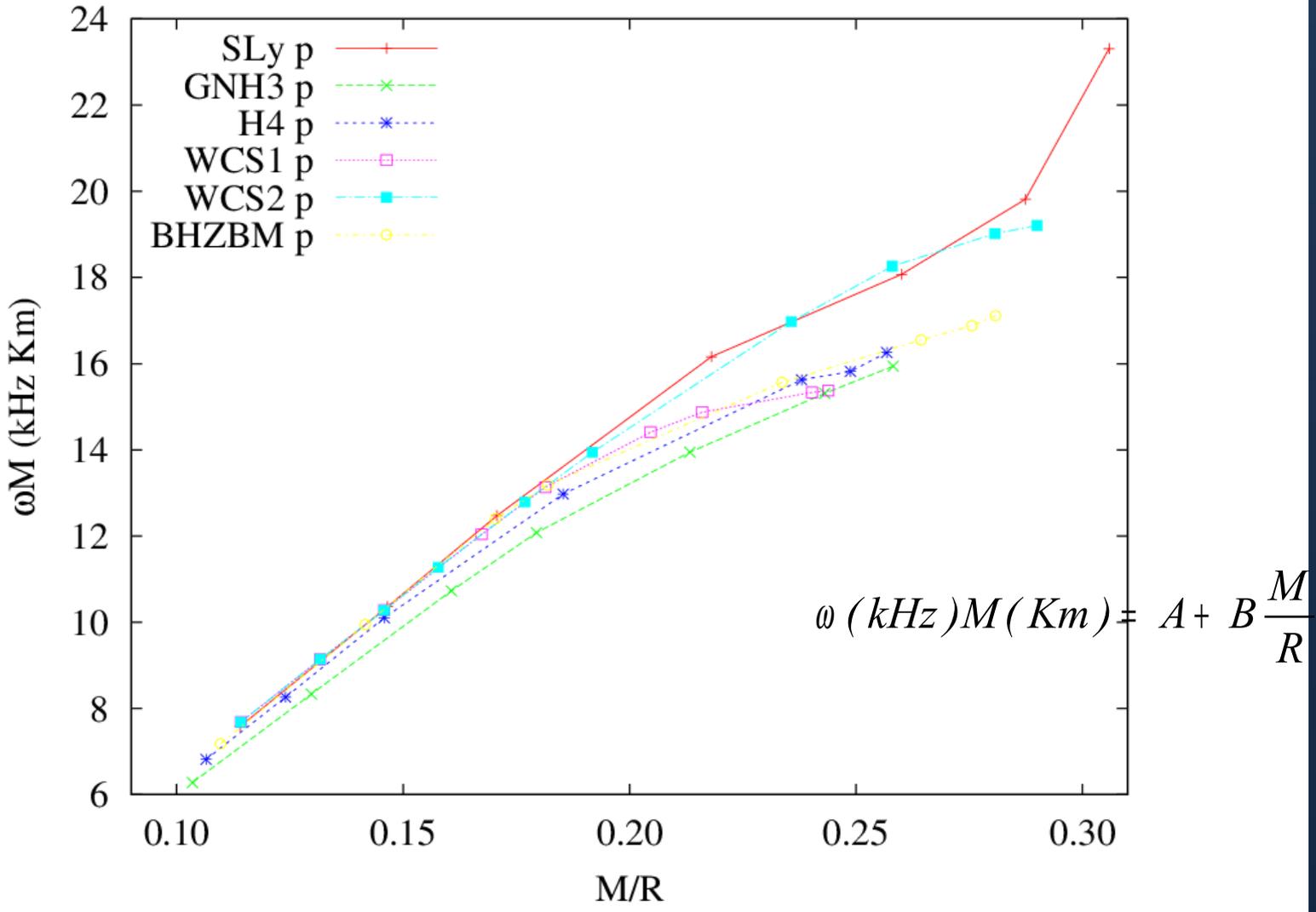
Polar f mode: *Hyperon matter*





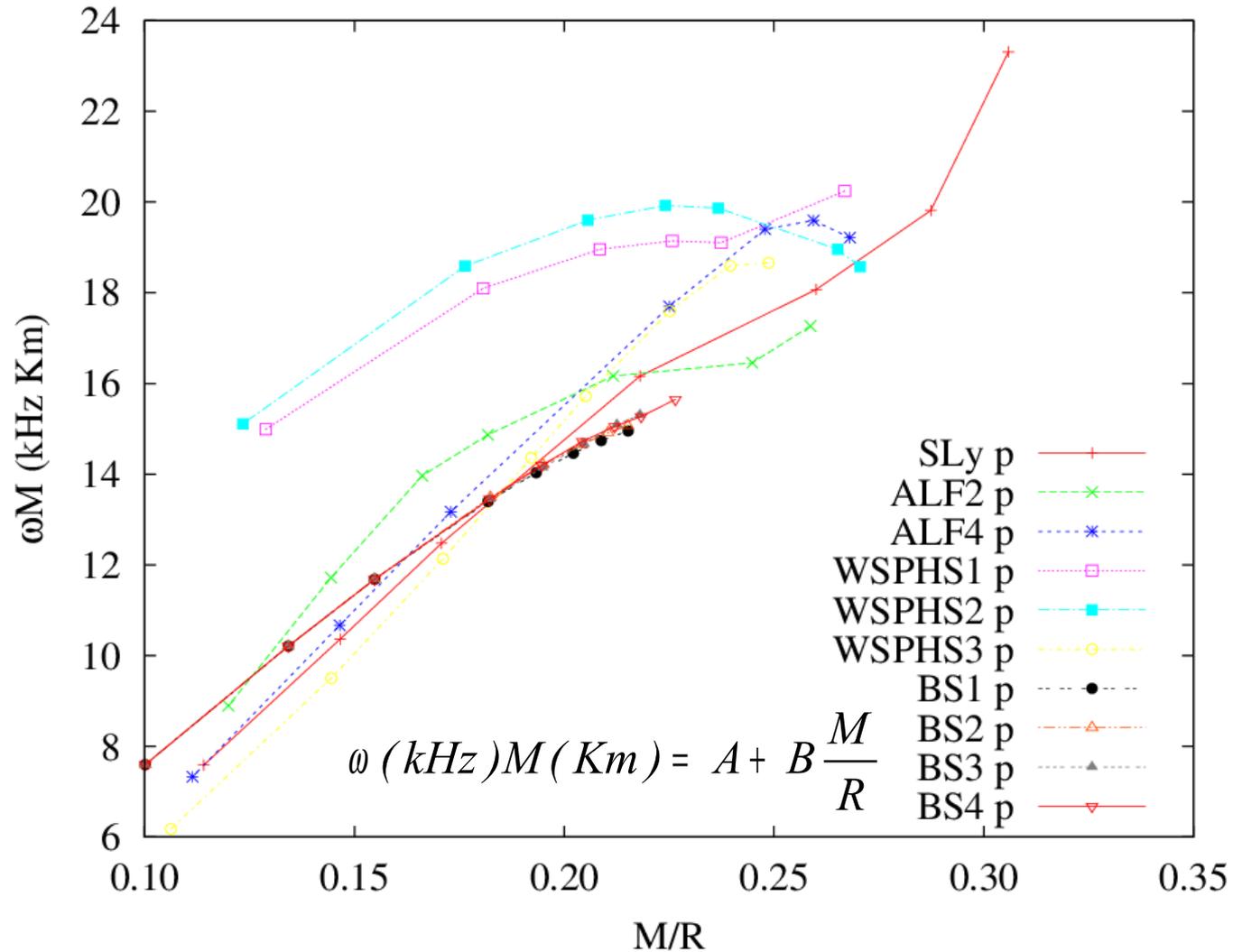
d) Polar p mode

Polar fundamental p mode: *Hyperon matter*



### 3. Results: d) p modes

Polar fundamental p mode: *Hybrid – quark matter*



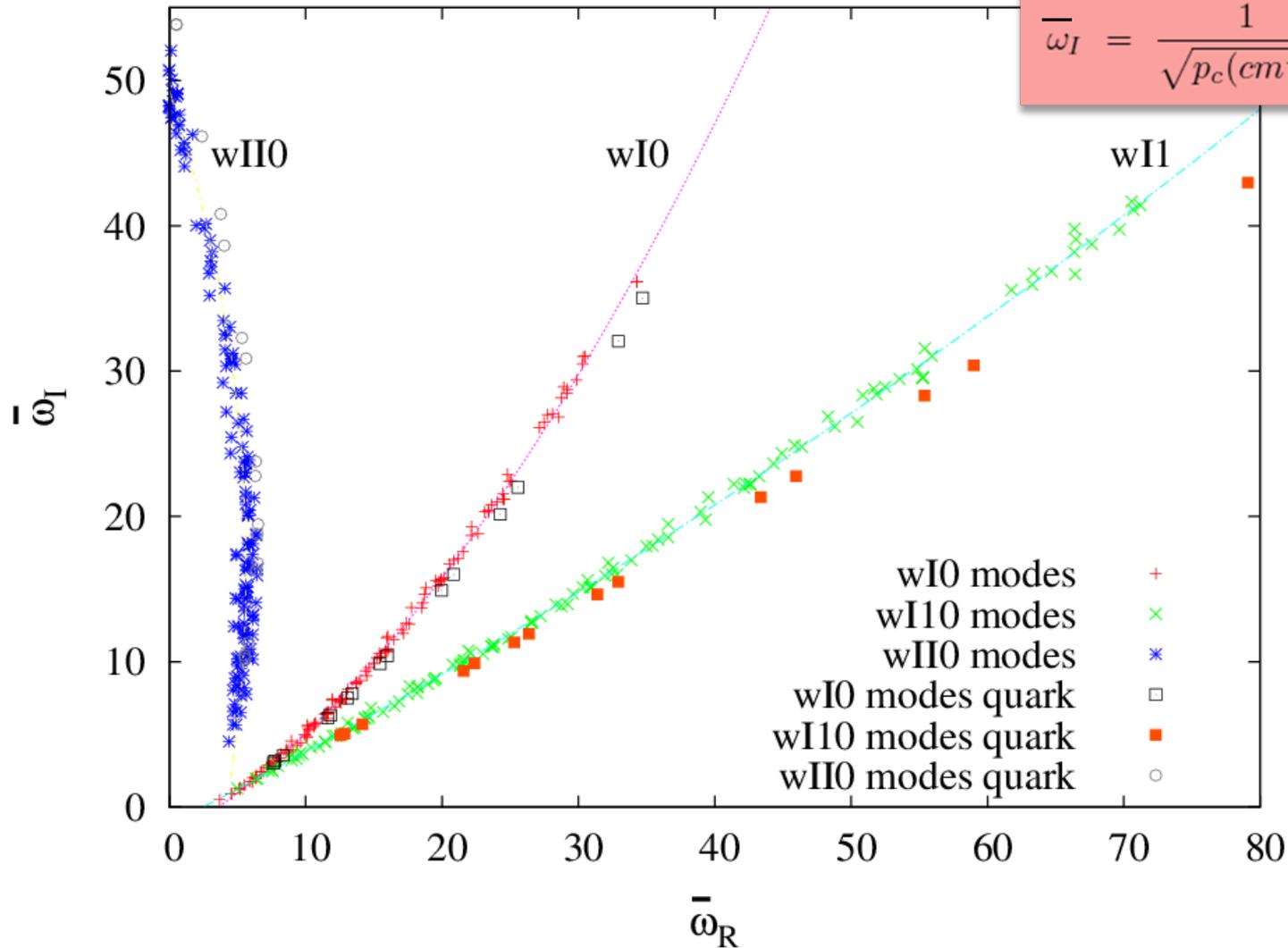
e) Empirical relations

### 3. Results: e) Empirical relations

*Axial modes: scaling in units of the central pressure*

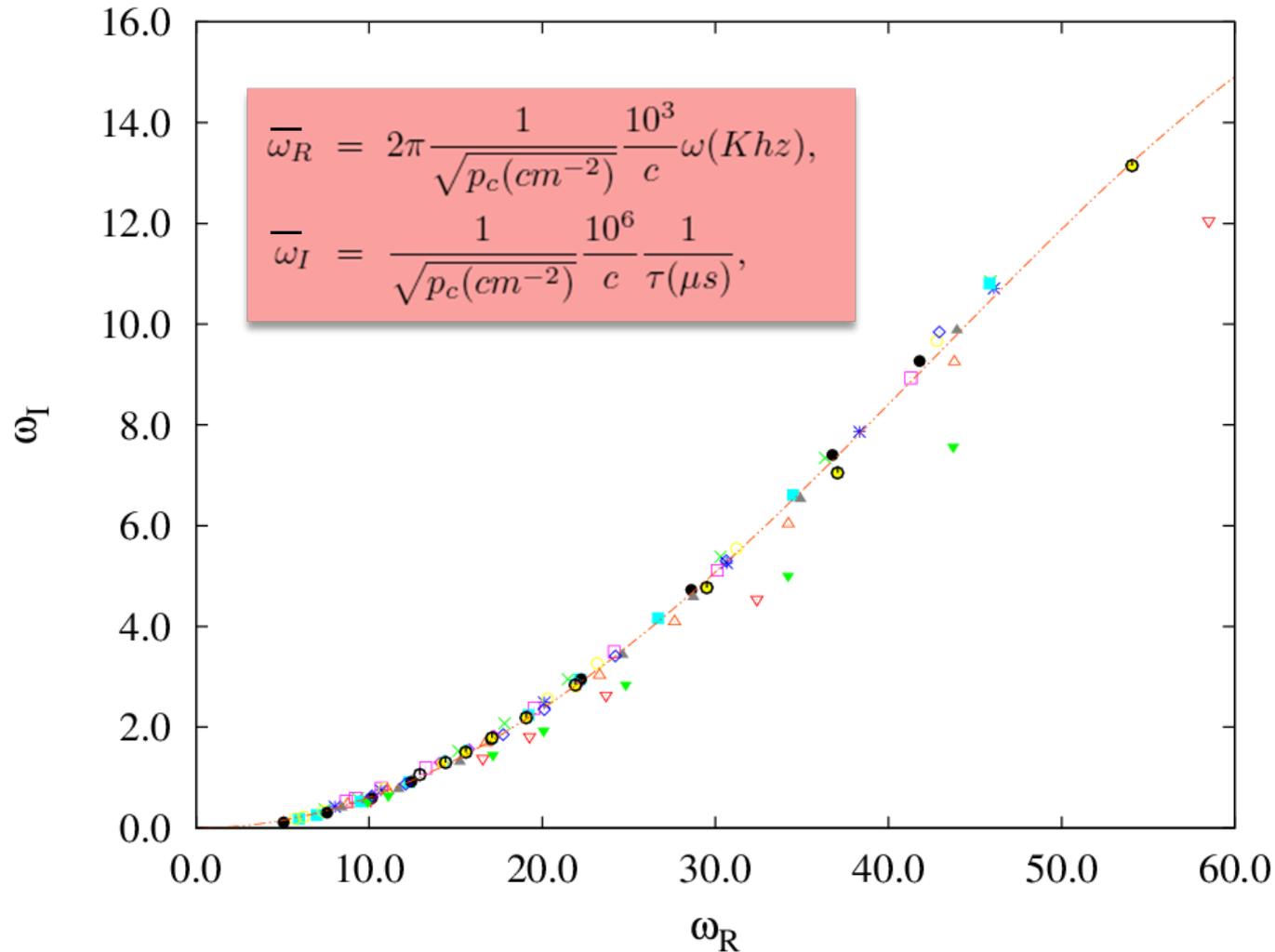
$$\bar{\omega}_R = 2\pi \frac{1}{\sqrt{p_c(\text{cm}^{-2})}} \frac{10^3}{c} \omega(\text{Khz}),$$

$$\bar{\omega}_I = \frac{1}{\sqrt{p_c(\text{cm}^{-2})}} \frac{10^6}{c} \frac{1}{\tau(\mu\text{s})},$$



### 3. Results: e) Empirical relations

*Similar por Polar wI0 modes:  
scaling in units of the central pressure*



*Axial fundamental wl modes*

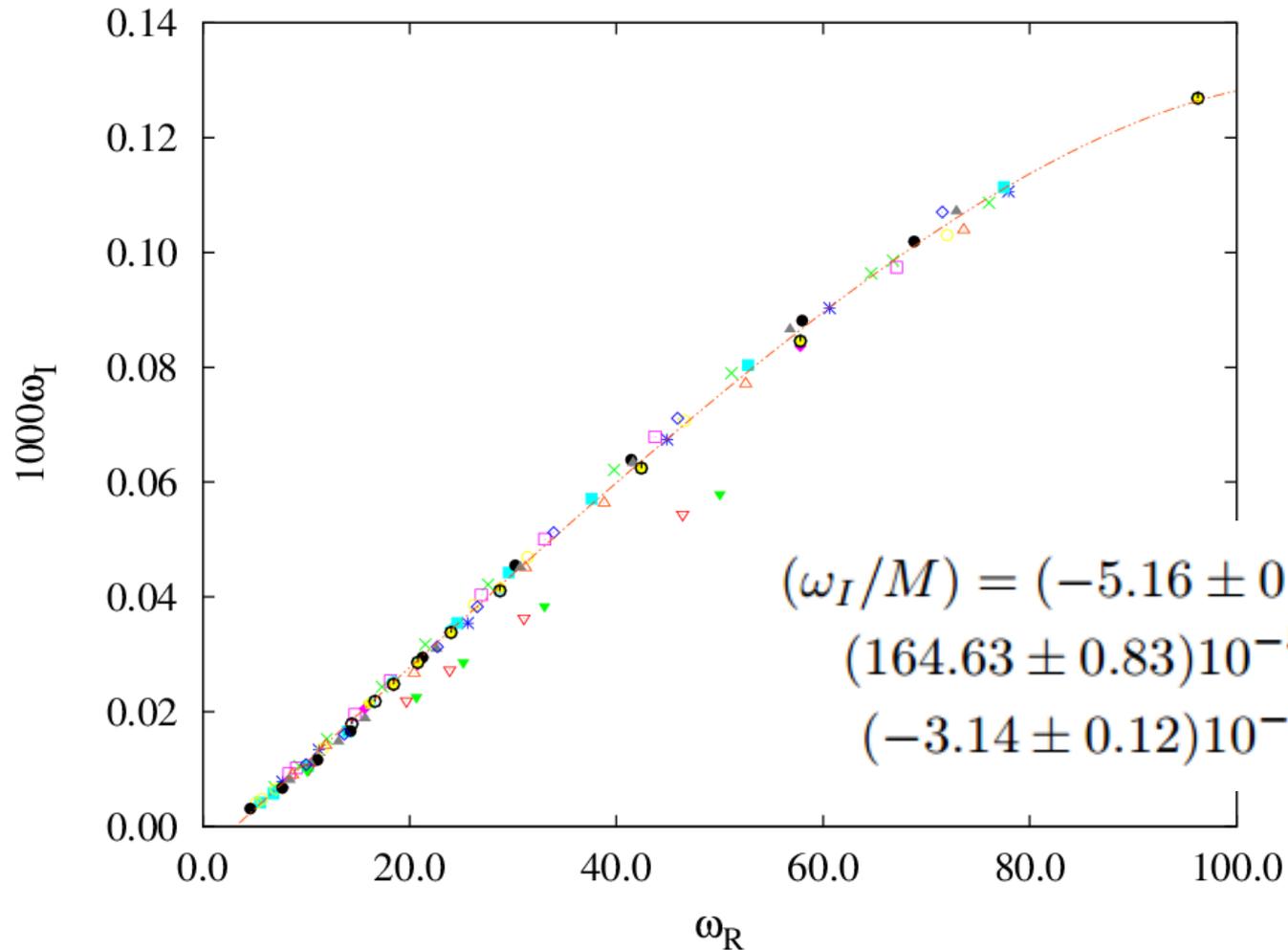
$$\bar{\omega}_I = (-2.30 \pm 0.13) + (0.58 \pm 0.16)\bar{\omega}_R + (0.0165 \pm 0.0004)\bar{\omega}_R^2$$

*Polar fundamental wl modes*

$$\bar{\omega}_I = (6.146 \pm 0.039)10^{-3}\bar{\omega}_R^2 + (-5.57 \pm 0.18)10^{-7}\bar{\omega}_R^4$$

### 3. Results: e) Empirical relations

*Polar f modes: scaling in units of Mass / Radius respectively*



## 4. Conclusions

- *We have studied axial and polar modes for 16 realistic equations of state*
- *These equations satisfy the 2 Msun condition*
- *Phenomenological relations useful for neutron star asteroseismology have been obtained*
- *New empirical relations have been obtained, valid for every equation of state, which allow to estimate the central pressure and radius*

**Thank you for your  
attention!**