

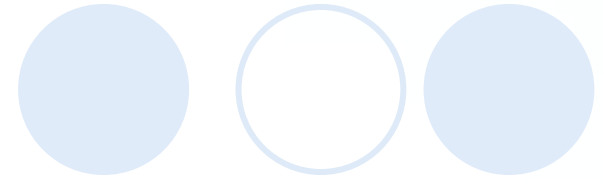
Non-Gaussianity & Inflation



Detection of NG would prove wrong at least one of these assumptions:

- single scalar field inflation
- canonical kinetic term
- slow roll
- Bunch-Davies vacuum

Non-Gaussianity & Inflation



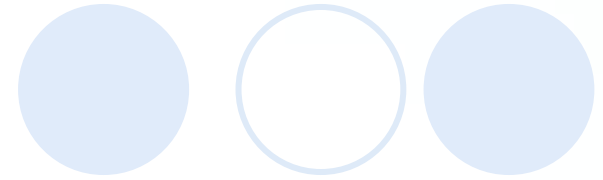
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Measurements so far leave room for some local non-Gaussianity

WMAP7	$-10 < f_{NL} < 74$	at 95% level	Komatsu et al., AJS (2011)
SDSS	$-29 < f_{NL} < 70$	at 95% level	Slosar et al., JCAP (2008)
PLANCK	$\Delta f_{NL} \simeq 5$		Tauber et al., A&A (2010)

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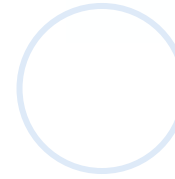
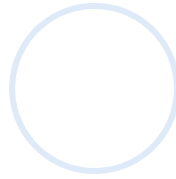
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Are we ready for that?

Contaminants



Any non-linearity naturally produces non-Gaussian features

- Galactic foregrounds
- SZ effect
- Cosmic reionization
- Detector-induced noise
- Lensing – ISW correlation
- etc.

e.g. Komatsu, CQG, 2010
Liguori et al., AiA 2010

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We shall focus on particularly tricky one

Non-linear dynamics at recombination

Second-order perturbation theory

At first order the initial conditions are linearly propagated:

$$a_{lm} = \Phi(\mathbf{k}) \mathcal{T}_l(k) Y_{lm}(\mathbf{k})$$

Non-Gaussianity is in the potential

Second-order perturbation theory

At first order the initial conditions are linearly propagated:

$$a_{lm} = \Phi(\mathbf{k}) \mathcal{T}_l(k) Y_{lm}(\mathbf{k})$$

Non-Gaussianity is in the potential

... non vanishing bispectrum linked to primordial non-Gaussianity .

$$\langle \Phi(k_1) \Phi(k_2) \Phi(k_2) \rangle = 0 \quad \Rightarrow \quad \langle a_{l_1 m_1} a_{l_2 m_2} a_{l_3 m_3} \rangle = 0$$

Second-order perturbation theory

At second order:

*Quadratic in the
primordial fluctuations*

$$a_{lm}^{(2)}(\mathbf{k}) = \mathcal{K} \left[\mathcal{T}_{lm}(\mathbf{k}, \mathbf{k}_1, \mathbf{k}_2) \Phi^{(1)}(\mathbf{k}_1) \Phi^{(1)}(\mathbf{k}_2) \right]$$

Second-order transfer function

Second-order perturbation theory

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Second-order transfer function

... Gaussian initial conditions are propagated nonlinearly into the observed CMB anisotropies .

$$\langle \Phi(k_1) \Phi(k_2) \Phi(k_2) \rangle = 0$$



$$\langle a_{l_1 m_1} a_{l_2 m_2} a_{l_3 m_3} \rangle = 0$$

Second-order perturbation theory



It is crucial to predict the shape and amplitude of second-order effects in order to subtract them from the data

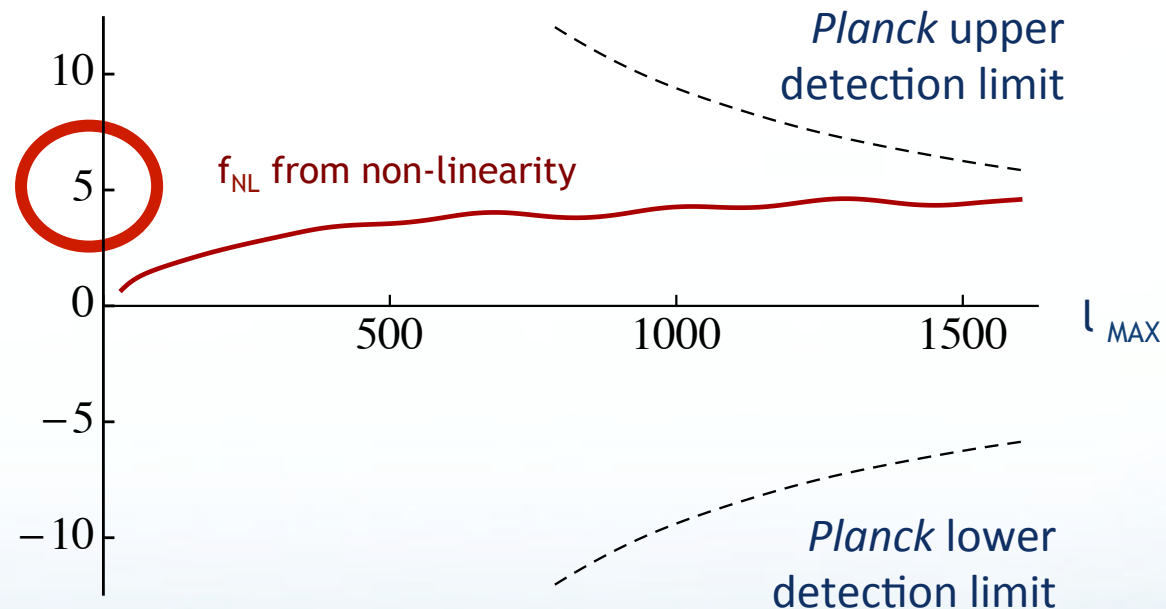
Previous attempts

Existing approximations either neglect part of the source terms or focus on a particular configuration

Approximation	Reference	f_{NL} contamination
Perturbed recombination	Khatri et al. 2009	- 1
Perturbed recombination	Senatore et al. 2009	- 3.5
Only quadratic terms	Nitta et al. 2009	+ 1
Squeezed limit	Bartolo et al. 2011	negligible for $L_{LONG} < 100$
Squeezed limit	Creminelli et al. 2011	
Squeezed limit	Lewis 2012	

First non-approximated result

... found by Pitrou et al. 2010, who wrote the Mathematica code CMBQuick

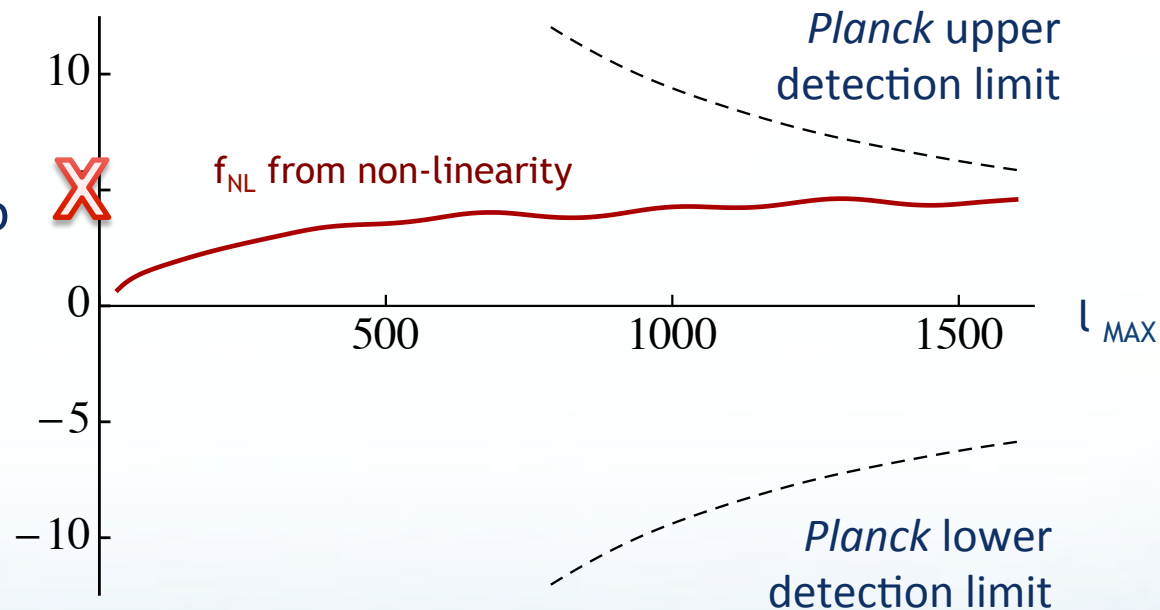


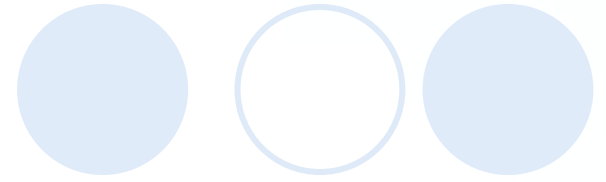
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Recently
corrected to

3





Why is it so complicated?

Recipe for a 2nd-order Boltzmann code

- ✓ Derive Boltzmann equation at 2nd order

Deriving Boltzmann equation

Beneke & Fidler, 2010

Pitrou et al. 2009

Senatore, Tassev & Zaldarriaga, 2009

Bartolo, Matarrese & Riotto, 2006-7

Theoretical complications (often forgotten at first order):

- Tetrad formalism for Boltzmann equation

$$G^{\mu}_{\nu} = \kappa e_{\underline{A}}^{\mu} e^{\underline{B}}_{\nu} T^{\underline{A}}_{\underline{B}}$$

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- Fluid limit & multipoles

$$\frac{1}{\mathcal{P}} T_{(n)}^{\underline{0}} = - {}_0 \Delta_{00}^{(n)}$$

$$\frac{1}{\mathcal{P}} T_{(n)}^{\underline{i}\underline{0}} = \frac{i}{3} \sum_m (-1)^m \binom{i}{-m} {}_1 \Delta_{1m}^{(n)}$$

$$\frac{1}{\mathcal{P}} T_{(n)}^{\underline{i}\underline{j}} = \frac{\delta_{ij}}{3} {}_2 \Delta_{00}^{(n)} - \frac{1}{5} \sum_m (-1)^m \chi_{2-m}^{ij} {}_2 \Delta_{2m}^{(n)}$$

Deriving Boltzmann equation

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$$\begin{aligned} \frac{1}{\mathcal{P}} T_{(n)}^{\underline{0}}_{\underline{0}} &= -\Delta_{00}^{(n)} \\ \frac{1}{\mathcal{P}} T_{(n)}^{\underline{i}}_{\underline{0}} &= \frac{i}{3} \sum_m (-1)^m \zeta_{-m}^i \Delta_{1m}^{(n)} \\ \frac{1}{\mathcal{P}} T_{(n)}^{\underline{i}}_{\underline{j}} &= \frac{\delta_{ij}}{3} \Delta_{00}^{(n)} - \frac{1}{5} \sum_m (-1)^m \chi_{2-m}^{ij} \Delta_{2m}^{(n)} \end{aligned}$$

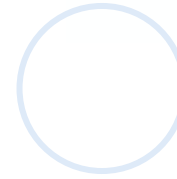
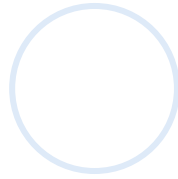
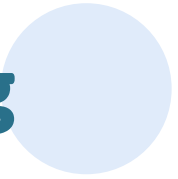
- Horrible angular projections

$$\sum_{l_1=|l-1|}^{l+1} \sum_{m_1=-l_1}^{l_1} i^{l-l_1} \begin{pmatrix} l_1 & 1 & l \\ -s & 0 & -s \end{pmatrix} \begin{pmatrix} l_1 & 1 & l \\ m_1 & m_2 & m \end{pmatrix} f_{ab, l_1 m_1}^{(2)}(\mathbf{k}).$$

Recipe for a 2nd-order Boltzmann code

- ✓ Derive Boltzmann equation at 2nd order
- ✓ Solve the Einstein-Boltzmann differential system at 2nd order

Mode coupling



At second order, convolution integrals appear:

$$A(\mathbf{x}) B(\mathbf{x}) \rightarrow \mathcal{K} \left[\tilde{A}(\mathbf{k}_1) \tilde{B}(\mathbf{k}_2) \right]$$

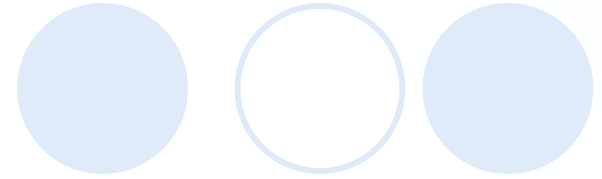
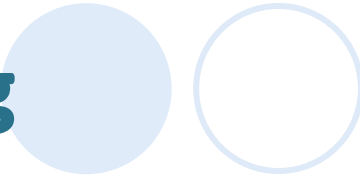
Typical equation:

Purely second-order terms

=

Convolution of quadratic
source terms

Mode coupling



Example Anisotropic stresses in Newtonian gauge

$$\Psi(\mathbf{k}) - \Phi(\mathbf{k}) = 0$$

Mode coupling



Example Anisotropic stresses in Newtonian gauge

$$\Psi(\mathbf{k}) - \Phi(\mathbf{k}) = \mathcal{K} \left[\begin{aligned} & - \frac{2}{k^4} \Psi(k_1) \Phi(k_2) (2\mu^2 k_1^2 k_2^2 + 3\mu k_1^3 k_2 + \mu k_1 k_2^3 + k_1^2 k_2^2 + k_1^4) + \text{symm} \\ & + \frac{1}{k^4} \Phi(k_1) \Phi(k_2) (15\mu^2 k_1^2 k_2^2 + 14\mu k_1^3 k_2 + 14\mu k_1 k_2^3 + 5k_1^2 k_2^2 + 4k_1^4 + 4k_2^4) \\ & + \frac{1}{k^4} \Psi(k_1) \Psi(k_2) (7\mu^2 k_1^2 k_2^2 + 6\mu k_1^3 k_2 + 6\mu k_1 k_2^3 + k_1^2 k_2^2 + 2k_1^4 + 2k_2^4) \end{aligned} \right]$$

Solve the system for a (k_1, k_2, k) grid

Differential system



- Four relativistic hierarchies (I, E, B, N) with custom L_{MAX}
+ 2 non-relativistic ones (baryons and CDM)
- Full Boltzmann equation, including collision term & lensing
- Approx. 100 equations on a grid in (k_1, k_2, k)
- We use the first-order sources & integrator from CLASS
Blas, Lesgourgues & Tram, JCAP 2011
- Same results using transfer functions or Green functions

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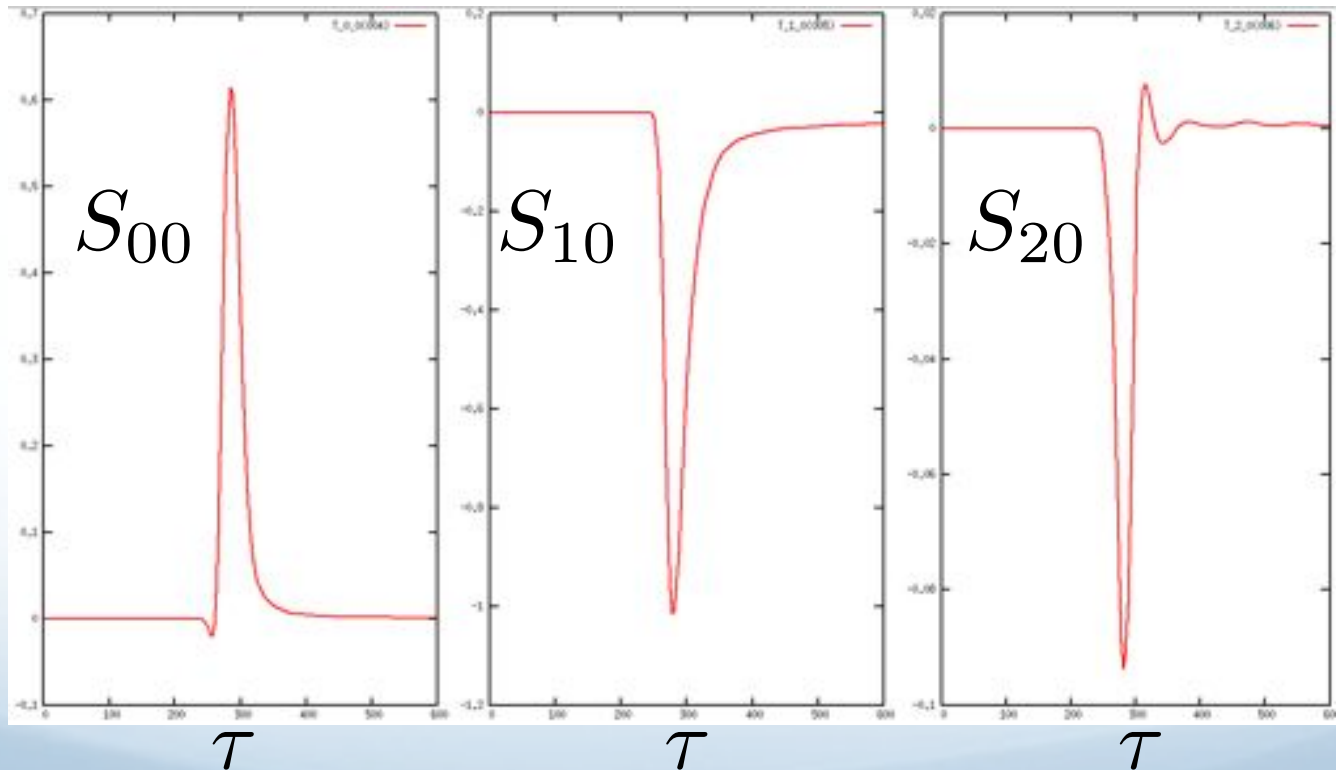
Computation time: 40 minutes

Line-of-sight source functions

$$\dot{\Delta}_{lm} + i \mathbf{k} \cdot \mathbf{n} \Delta_{lm} + \kappa \Delta_{lm} = S_{lm}$$

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Recipe for a 2nd-order Boltzmann code

- ✓ Derive Boltzmann equation at 2nd order
- ✓ Solve the Einstein-Boltzmann differential system at 2nd order
- ✓ Derive transfer functions using the line-of-sight formalism

Line-of-sight integral

At first order is a convolution over spherical Bessels:

Seljak & Zaldarriaga, 1996

$$\mathcal{T}_l(\mathbf{k}) = \int_0^{\tau_0} d\tau e^{-\kappa(\tau, \tau_0)} j_l(k(\tau - \tau_0)) S(k)$$

*First-order
line-of-sight sources
One scalar function!*

Line-of-sight integral

Oscillatory integral over complicated projection functions

Beneke, Fidler & Klingmuller, 2011

*Combination of spherical
Bessels and 3j's*

$$\mathcal{T}_{lm}(\mathbf{k}, \mathbf{k}_1, \mathbf{k}_2) = \int_0^{\tau_0} d\tau e^{-\kappa(\tau, \tau_0)} \sum_{L=0}^{l_{\max}} J_{Llm}(k(\tau - \tau_0)) S_{Lm}(\mathbf{k}, \mathbf{k}_1, \mathbf{k}_2)$$

*Second-order
line-of-sight sources*

Line-of-sight integral

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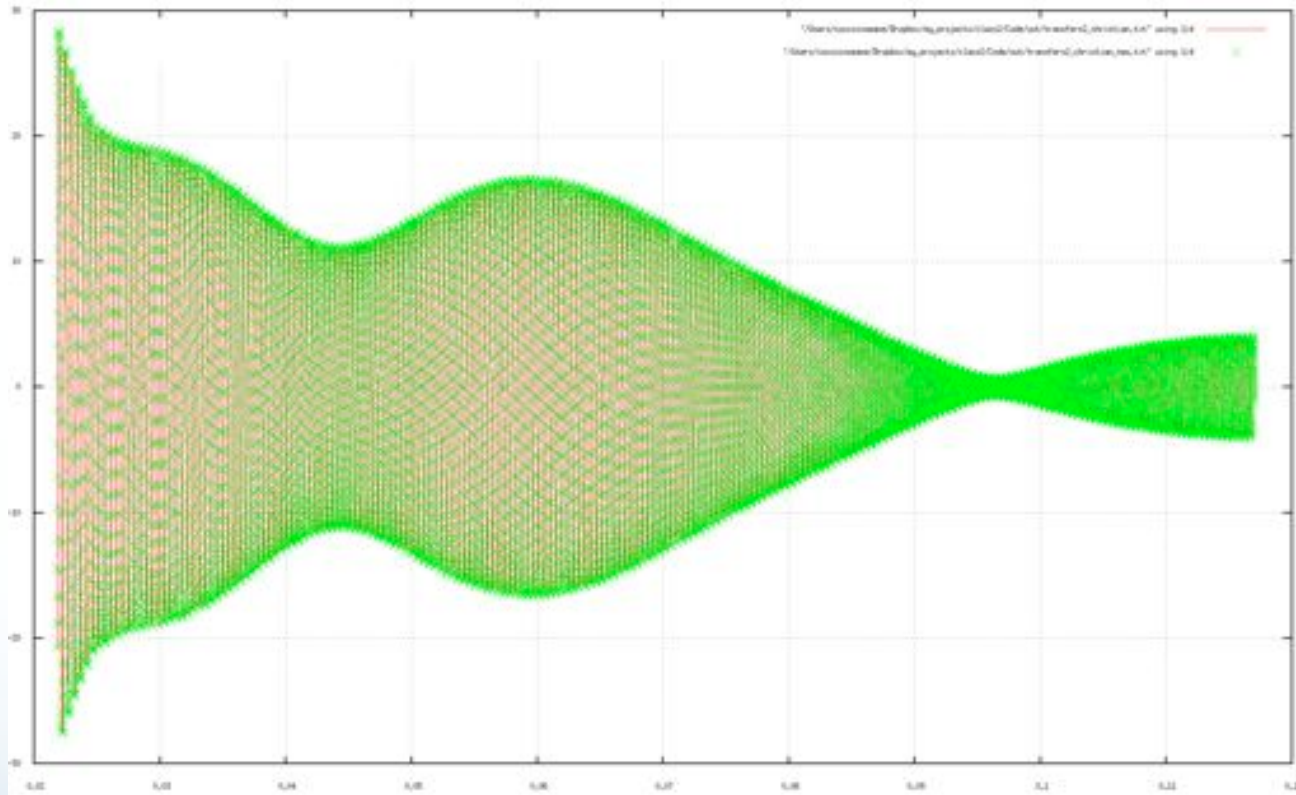
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*Second-order
line-of-sight sources*

Computation time: 15 minutes

Line-of-sight integral

\mathcal{T}_{100}

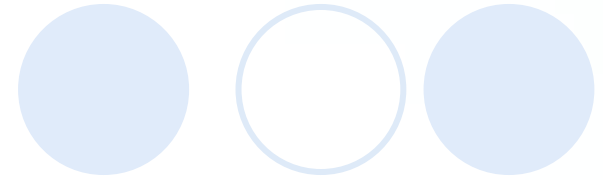


k [Mpc/h]

Recipe for a 2nd-order Boltzmann code

- ✓ Derive Boltzmann equation at 2nd order
- ✓ Solve the Einstein-Boltzmann differential system at 2nd order
- ✓ Derive transfer functions using the line-of-sight formalism
- ✓ Compute the temperature full-sky bispectrum

Angular power spectrum



Convolution of an oscillatory function with the smooth power spectrum

$$C_l \propto \int_0^{\infty} dk k^2 P(k) |T_l(k)|^2$$

Full-sky bispectrum

Involved integral on 6 highly oscillatory functions

Fergusson & Shellard, 2007, 2009
Fergusson, Regan & Shellard 2011

$$b_{l_1 l_2 l_3} = \int dr r^2 \int dk_1 k_1^2 T_{l_1}(k_1) P(k_1) j_{l_1}(r k_1) \int dk_2 k_2^2 T_{l_2}(k_2) P(k_2) j_{l_2}(r k_2) \int dk_3 k_3^2 T_{l_3}(k_1, k_2, k_3) j_{l_3}(r k_3)$$

Second-order transfer function

Full-sky bispectrum

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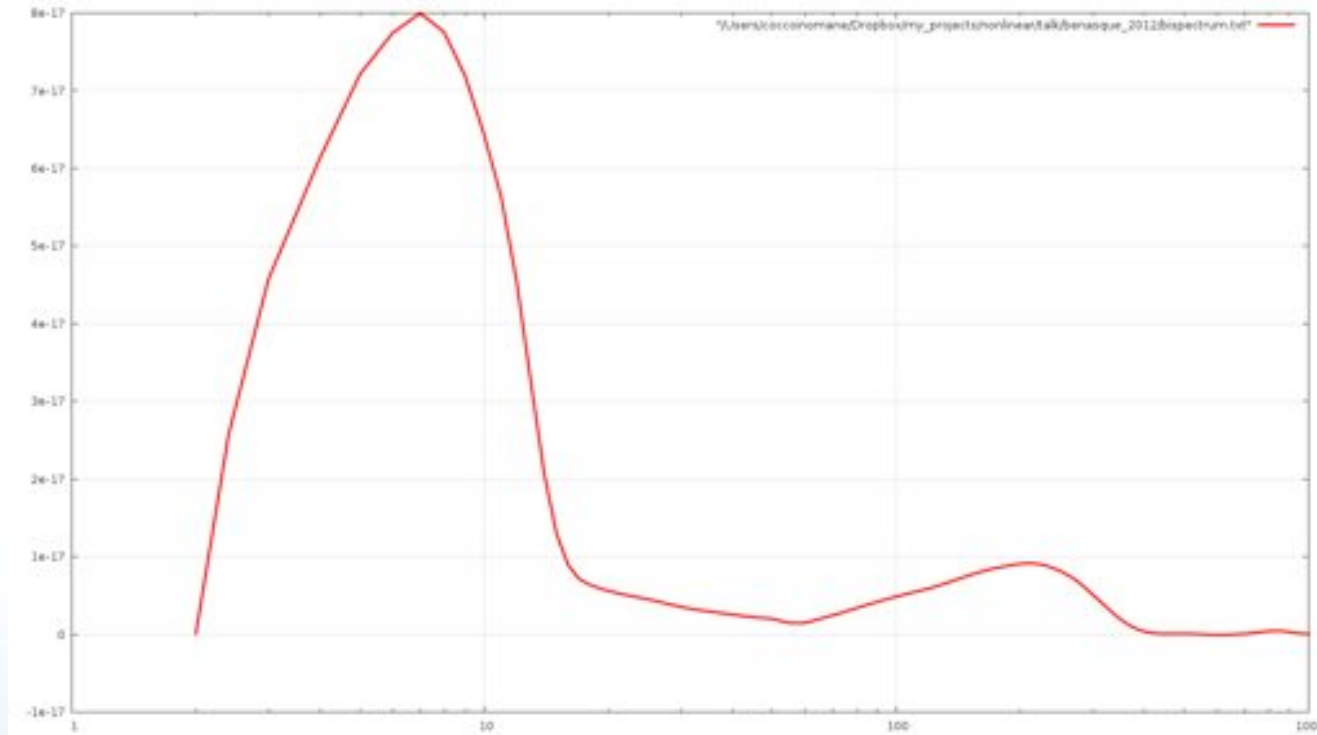
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Second-order transfer function

Computation time: 25 minutes

Full-sky bispectrum

b_{lll}



l

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- ✗ Compute f_{NL} contamination

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... wait until October



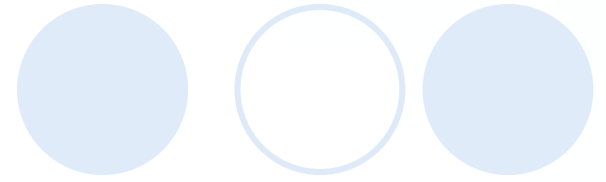
Conclusions



- Second-order corrections to the CMB bispectrum need to be computed
- The full second-order transfer function is needed to properly subtract the effect
- We have written a fast, parallel and flexible second-order Boltzmann code in Newtonian gauge
- From cosmology to bispectrum in less than 2 hours
- Expect the first science results on f_{NL} in October 2012

Current & future work

- Perform convergence tests on the bispectrum
- Check the bispectrum against squeezed analytical solutions
- Compute the bias on f_{NL}
- Include ISW in the line-of-sight sources
- Include lensing terms in the line-of-sight sources
- Implement a tight-coupling scheme
- Compute the bispectrum in a smarter way (e.g. make the integral separable using decomposition)
- Implement synchronous gauge



Thank you for your attention!