

# Current cosmic string constraints based on the Sussex-Imperial-Geneva defect simulations

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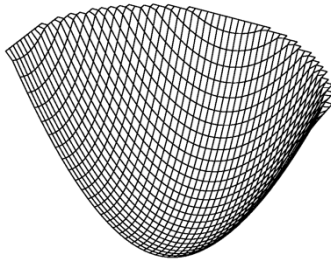
# overview

- what are topological defects?
- why are they interesting?
- defect simulations
- from simulations to the CMB
- current power spectra results
- outlook: beyond the CMB
- conclusions

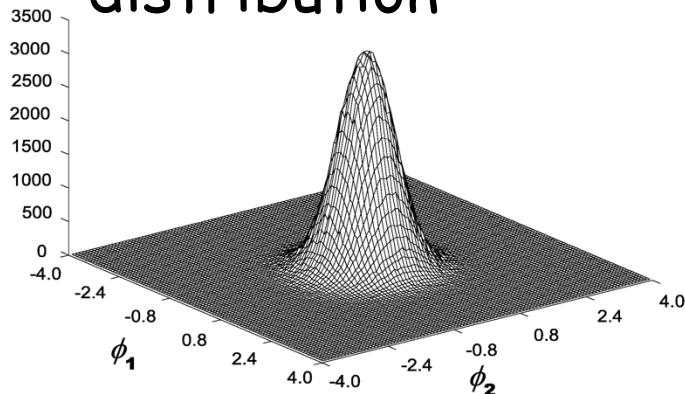
# what are topological defects?

## *High temperature*

- potential has single minimum

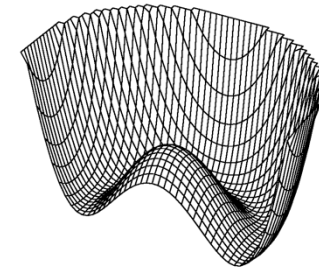


- symmetric field distribution

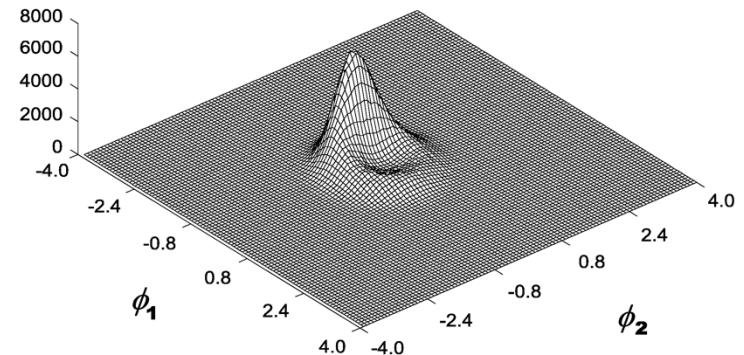


## *low temperature*

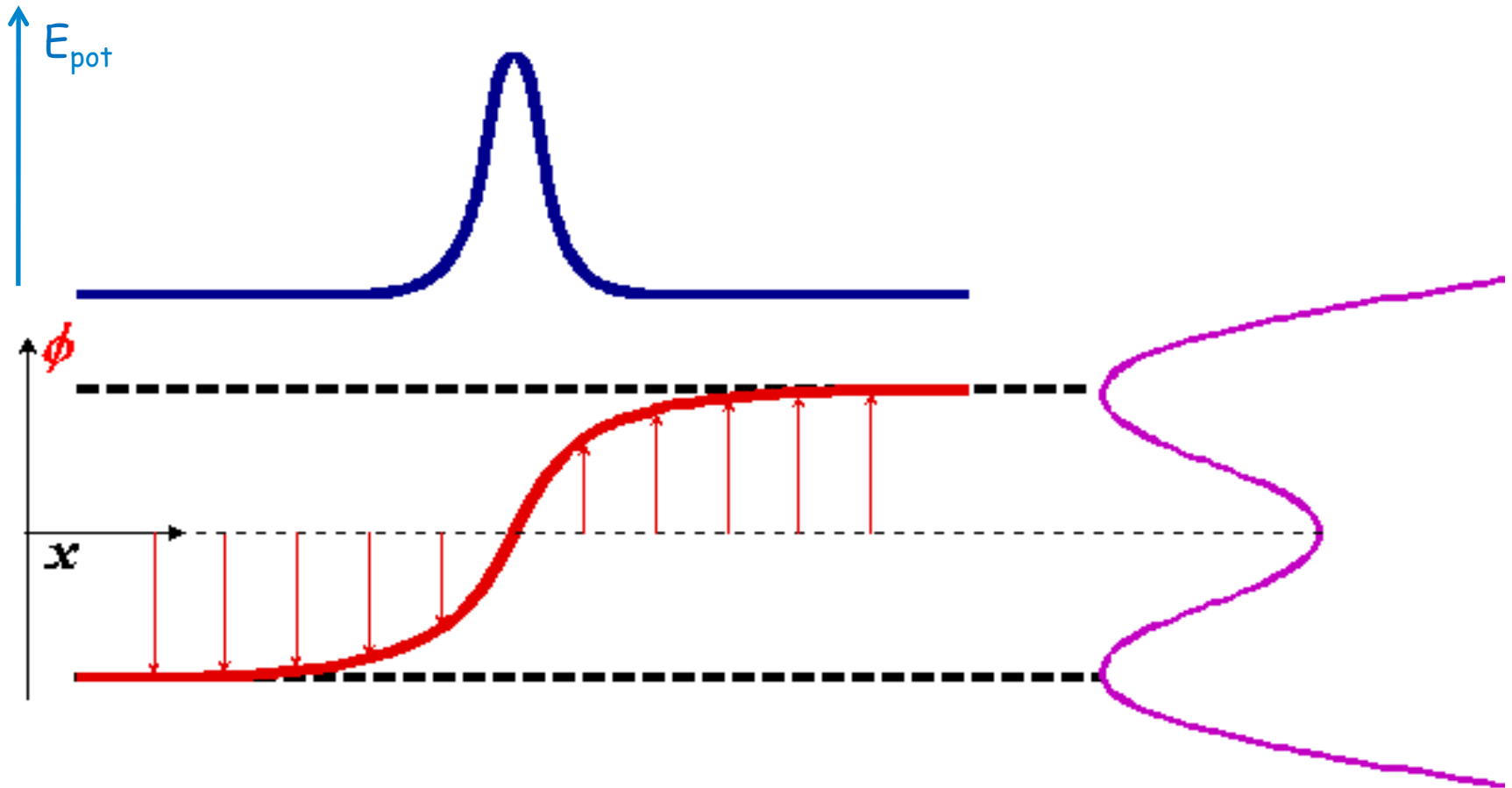
- potential has several minima



- asymmetric field distribution



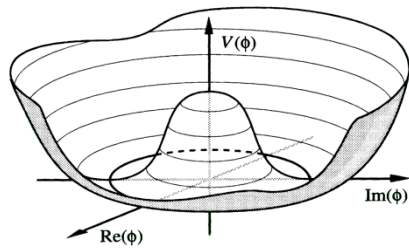
# what are topological defects?



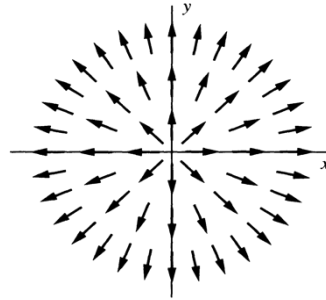
- crossing of  $\phi=0$  protected by topology of  $V_{\text{min}}$
- field has to balance potential and gradient energy

# zoology I

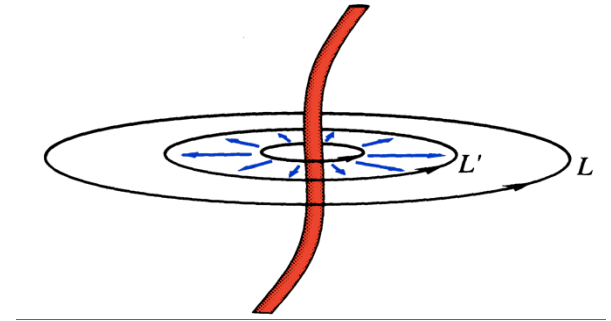
*strings:*



potential

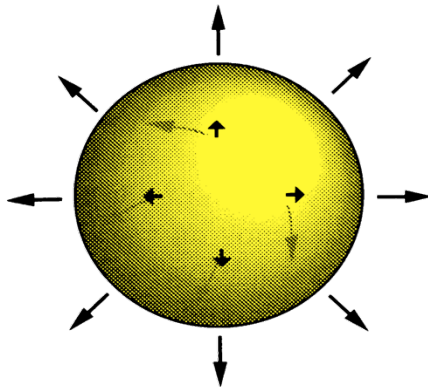


field configuration

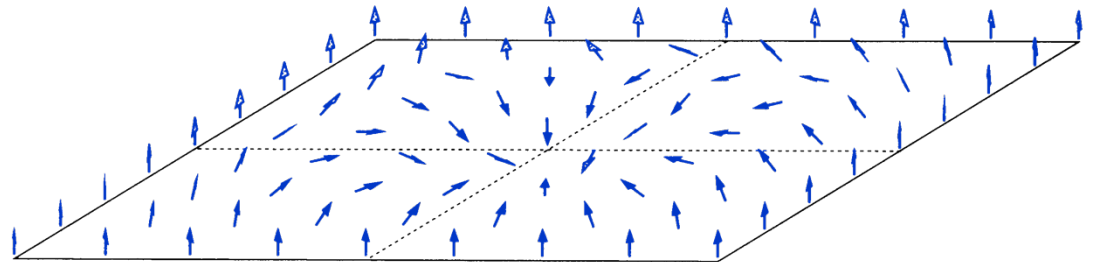


string in 3D

*monopoles:*



*texture:*



(illustrations from Vilenkin&Shellard)

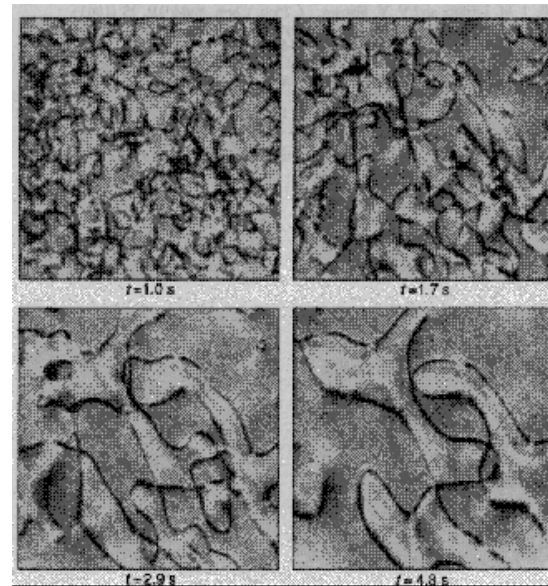
# zoology II

- **global defects**
  - big, fat, fluffy things
  - just one (N-component) scalar field  $\rightarrow$  long range forces
  - gradient and potential energy balance
- **local defects**
  - scalar and gauge fields
  - gauge field can remove gradients, no long range force
  - small
  - only strings relevant (others disappear or dominate)
- **other strings**
  - semilocal strings: not topologically stable, can disappear
  - fundamental strings
  - $(p,q)$  strings
  - many more possibilities! (eg strings ending on monopoles, etc)

# do topological defects exist?

**Yes!** At least in “normal” physics:

- vortices in superfluid helium
- flux lines in superconductors
- polymers
- liquid crystals
- ...



defects in liquid crystals  
(Chuang, Durrer, et al)

# why think about defects?

- *topological defects exist in nature*
- defects are created naturally in phase transitions
- phase transitions are natural in the universe
- cosmic defects would perturb the universe
- the universe shows perturbations at the level of  $10^{-5}$  (e.g. in the CMB)
- the scale would be just right for a "GUT" (Grand Unified Theory)
- did the defects cause the initial perturbations?



# defects *and* inflation

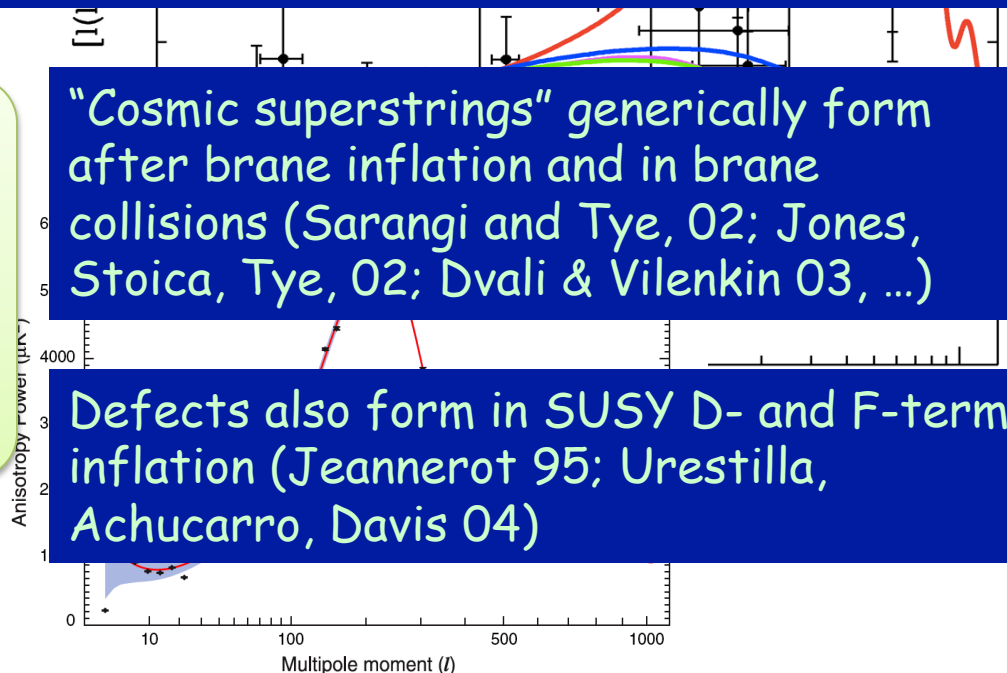
C  
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† "We consider gauge groups having a rank between 4 and 8. We examine all possible spontaneous symmetry breaking patterns from the GUT down to the standard model gauge group. Assuming standard hybrid inflation, we select all the models which can solve the GUT monopole problem, lead to baryogenesis after inflation and are consistent with proton lifetime measurements. We conclude that in all acceptable spontaneous symmetry breaking schemes, cosmic string formation is unavoidable."  
(Jeannerot, Rocher & Sakellariadou, hep-ph/0308134)

**BUT:**

many "realistic"  
inflationary models  
create generically  
defects after inflation!

"Cosmic superstrings" generically form after brane inflation and in brane collisions (Sarangi and Tye, 02; Jones, Stoica, Tye, 02; Dvali & Vilenkin 03, ...)

Defects also form in SUSY D- and F-term inflation (Jeannerot 95; Urestilla, Achucarro, Davis 04)



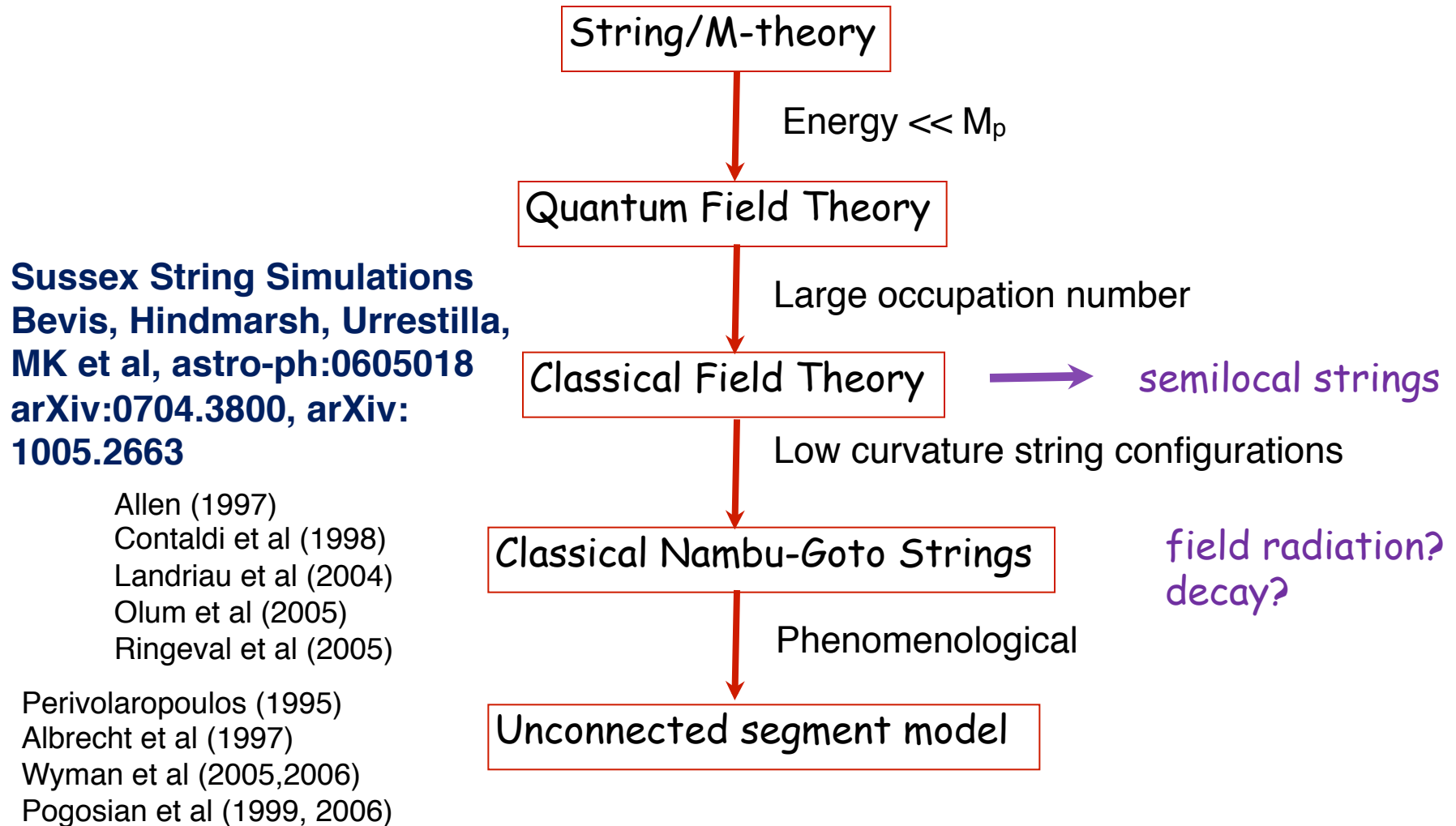
# how do we study defects?

- run defect simulation (field theory on a grid)
- record energy-momentum tensor
- defects are highly non-linear objects, so assume that they are stable against small perturbations
  - > they perturb the universe, but are not perturbed by it
- add their perturbation of the metric to the perturbations from everything else
- run CMBEasy (or CAMB, CLASS or cmbfast)

# progress in computing

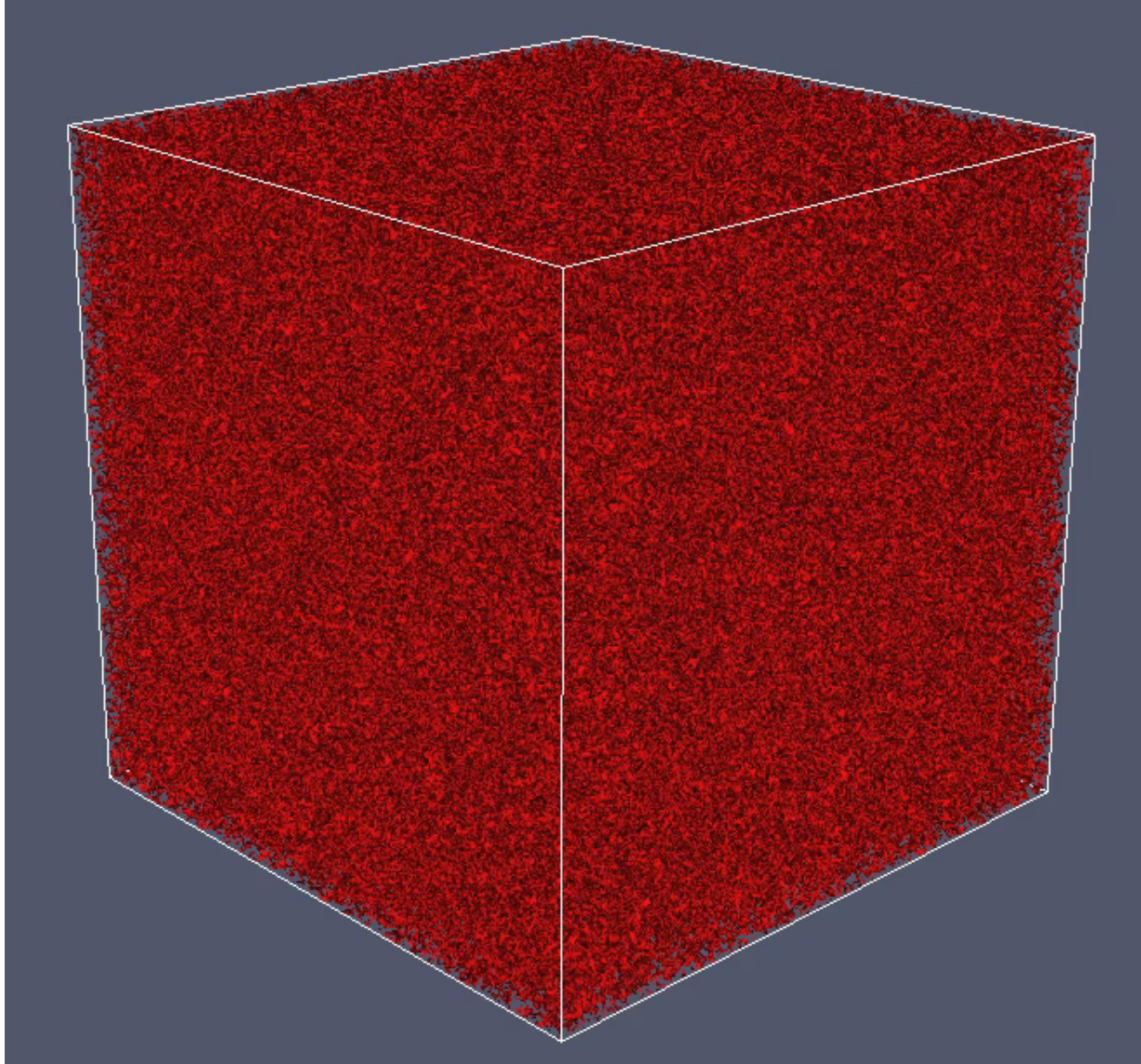
- **ca 1998: global defects,  $400^3$  grid**
  - single vector processor
- **ca 2005: cosmic strings,  $512^3$  grid**
  - MPI code w/ '1D' parallelisation (FFT issue)
- **ca 2009: cosmic strings,  $1024^3$  grid**
  - bigger computer (some other improvements)
- **2012: cosmic strings,  $4096^3$  grid**
  - '2D' parallelisation (MPI), scales to  $>10^5$  cores
  - 11 TB of field-data in memory...
  - improved parallel I/O -> **David Daverio's talk**

# levels of approximations



**"less modelling, more physics" : fully field-theoretic cosmic strings!**

# how do they look like?

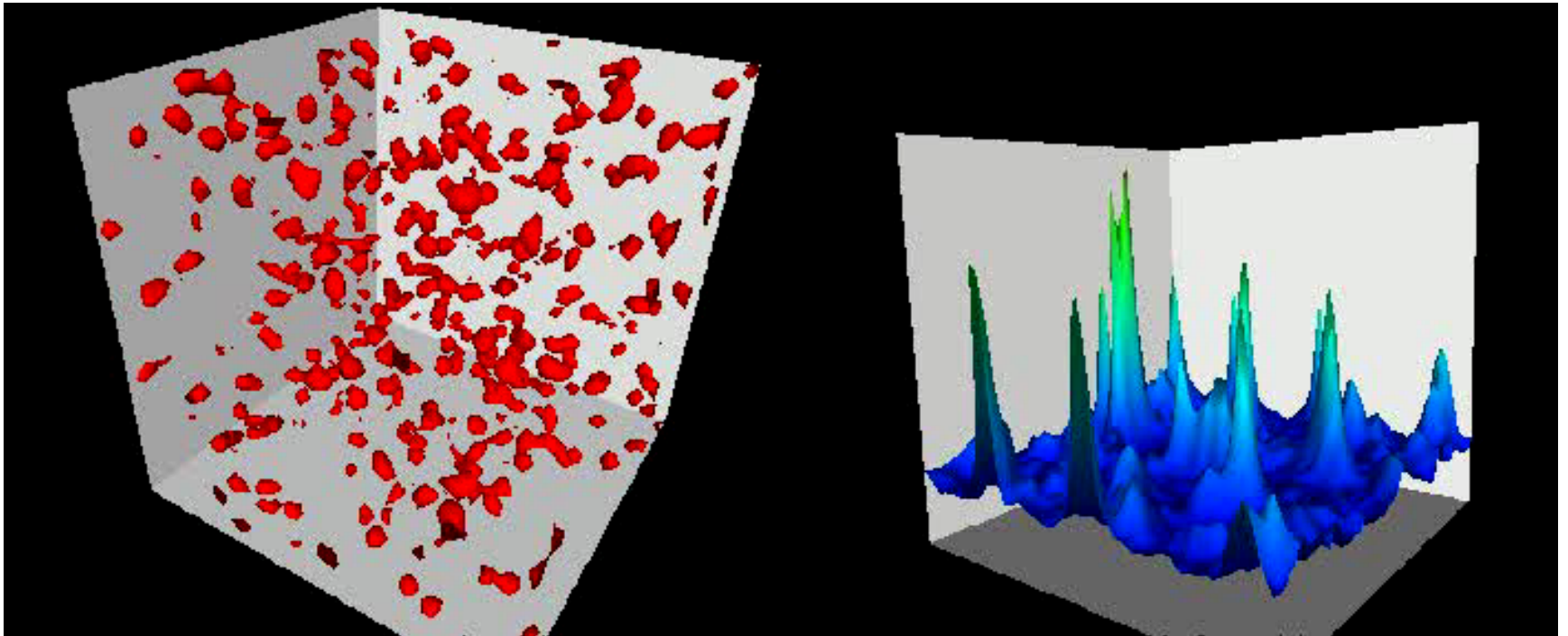


animation by David Daverio



# how do they look like?

global texture



$$\mathcal{L} = -\frac{1}{4e^2} F_{\mu\nu} F^{\mu\nu} + (D_\mu \phi)^* (D^\mu \phi) - \frac{\lambda}{4} (|\phi|^2 - \phi_0^2)^2$$

### ***Abelian Higgs strings***

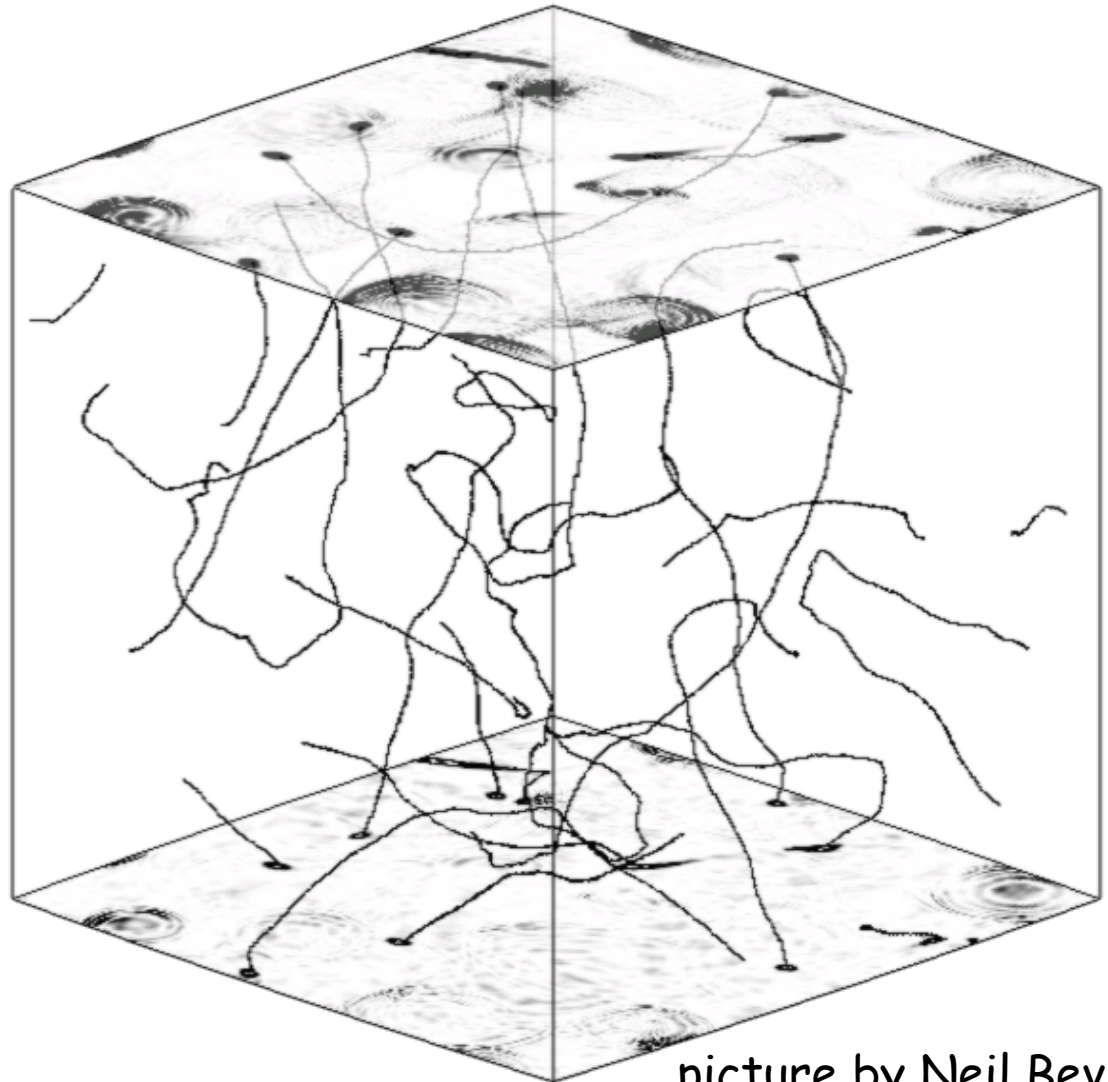
(scalar + electromagnetic fields)

Horizon roughly fills box

**lines:** string centres  
(from winding of field)

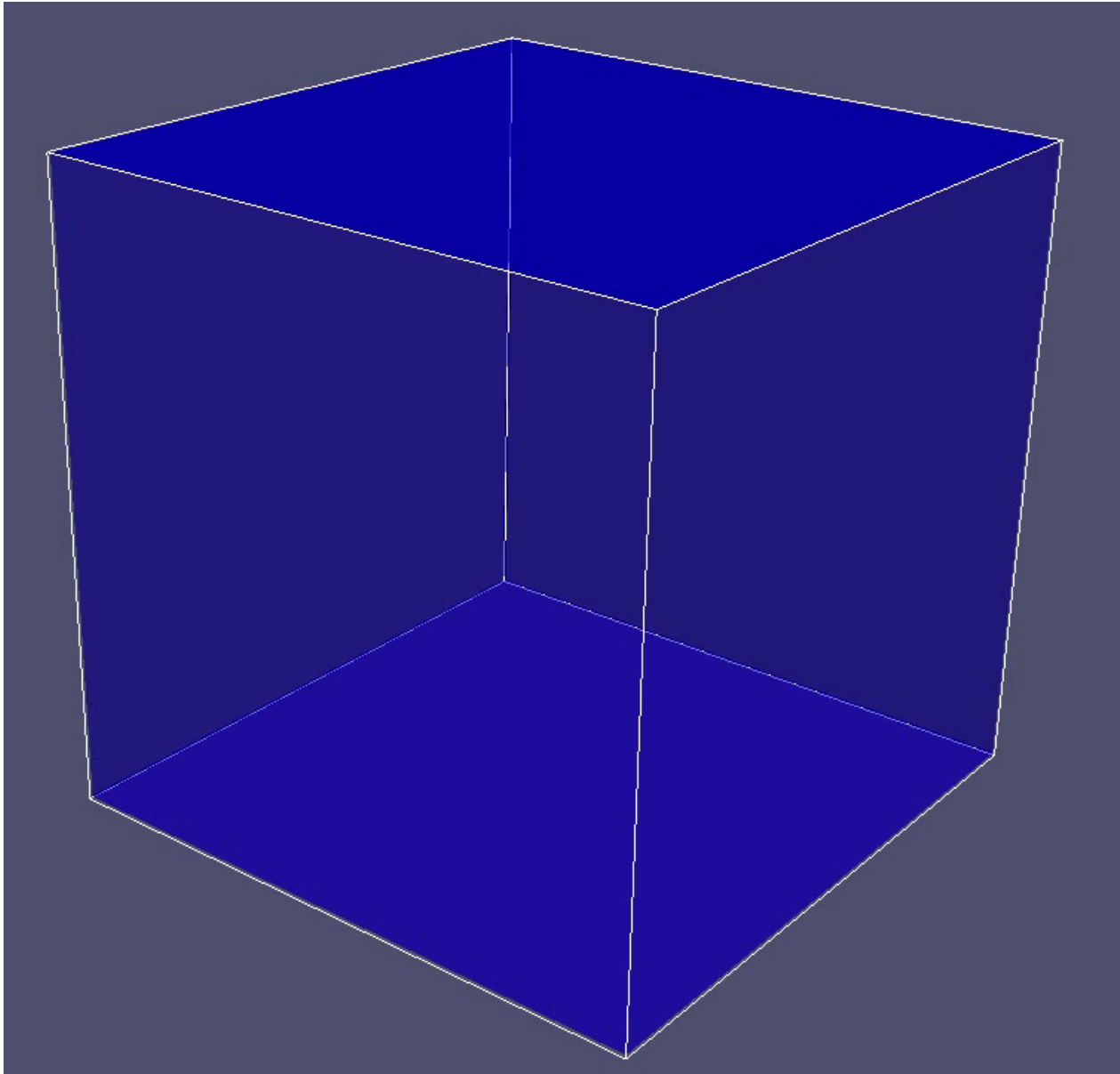
**bottom:** scalar field  
energy density

**top:** EM field energy  
density (curl around  
string cores from  
counteracting scalar field  
gradients)



picture by Neil Bevis

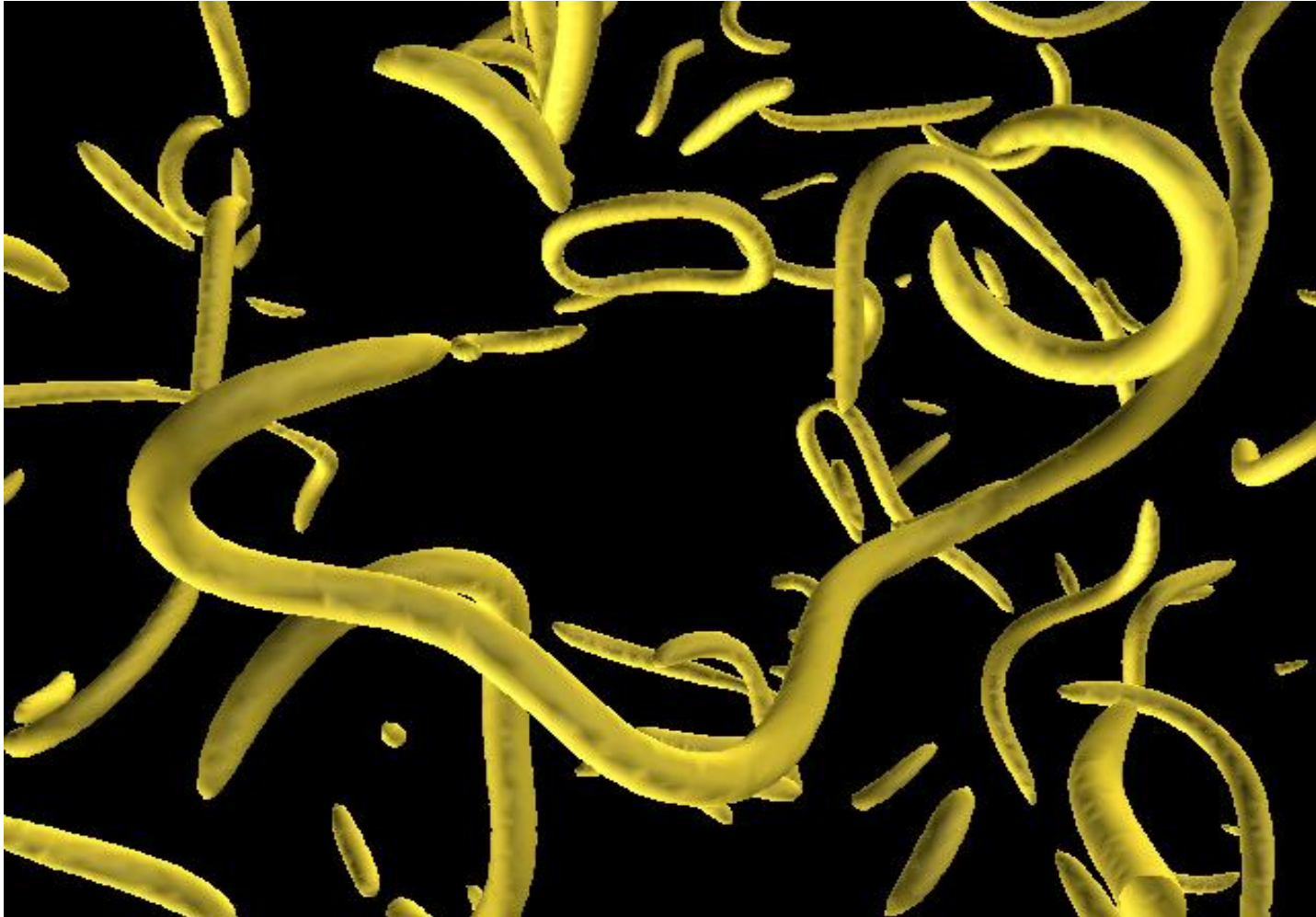
# how do they look like?



animation by David Daverio



# how do they look like?



semilocal strings  
**need** field theory simulations

Achucarro, Borrill, Liddle

# slim strings, fat strings

- the universe is expanding!
- comoving coordinates expand with the universe
- no problem for big, fluffy global defects
- but **cosmic strings shrink!** width  $\sim 1/a$
- very quickly they are no longer resolved
- **limits dynamical range (already quite small)**
- need to **blow them up** artificially
- check with different blow-up parameters (including no fattening), find little systematic effect

# UETC formalism

How can we get the defects into cmbfast/CAMB/CLASS/CMBEasy?

Typical structure of CMB code:

$$C_\ell \sim \left\langle \int dk \left| \int d\eta X(k, \eta) j_\ell(k(\eta_0 - \eta)) \right|^2 \right\rangle$$
$$\sim \int dk d\eta d\eta' \langle X(k, \eta) X(k, \eta') \rangle j_\ell(k(\eta_0 - \eta)) j_\ell(k(\eta_0 - \eta'))$$

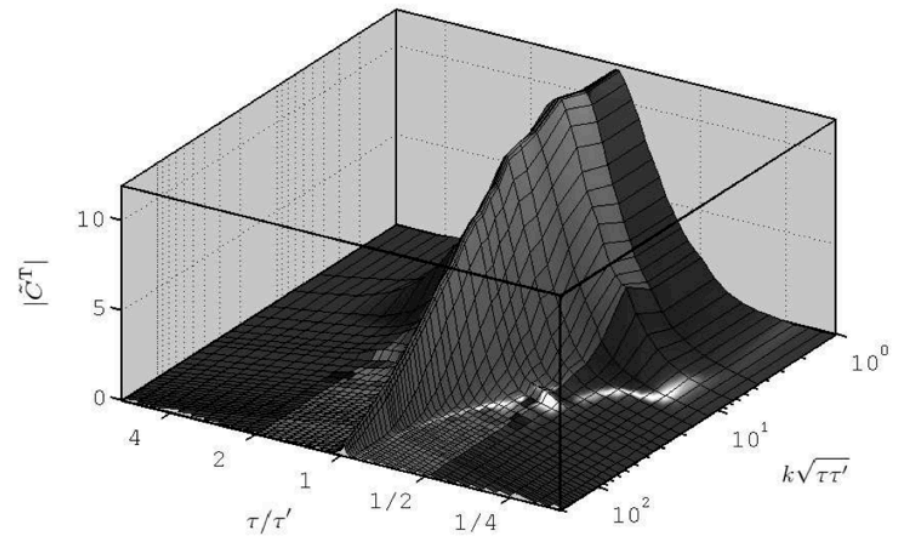
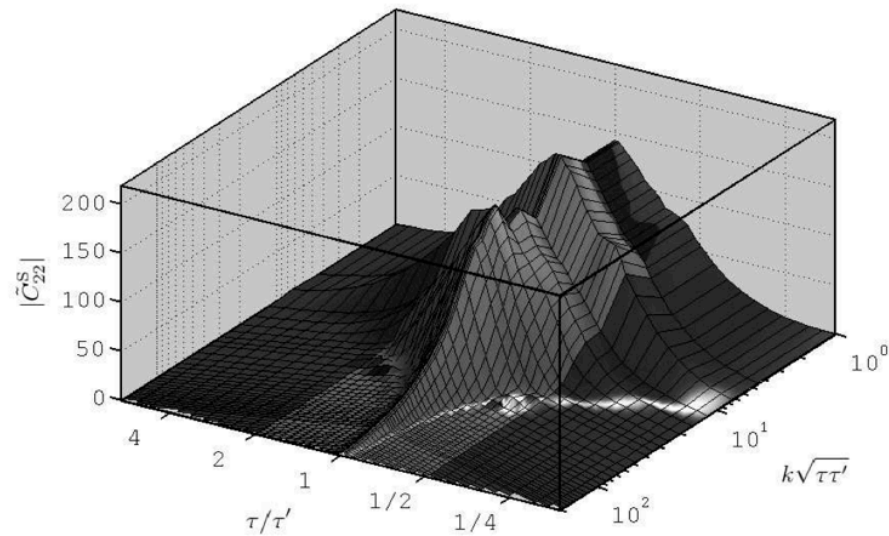
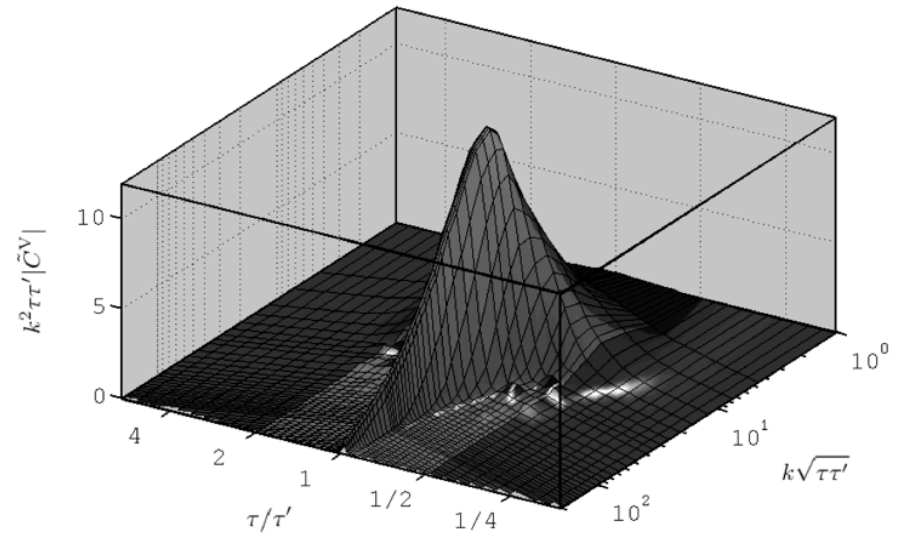
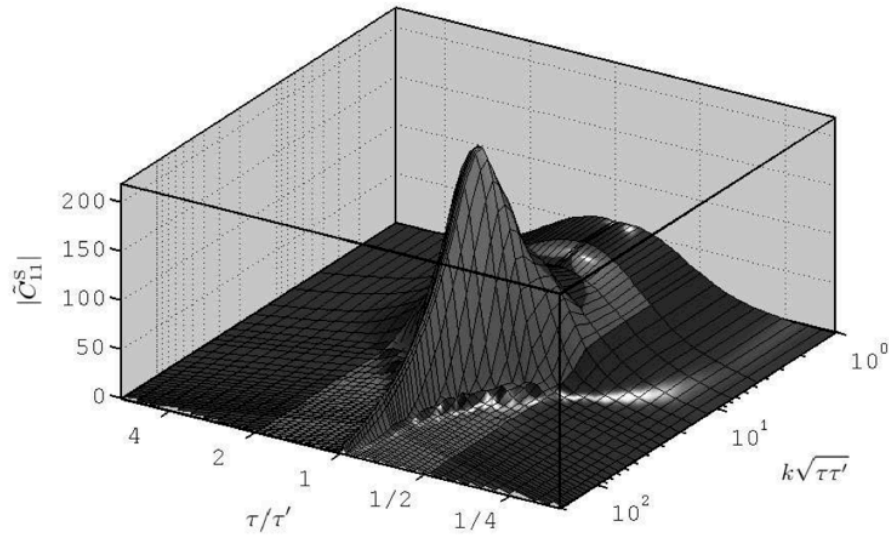
$X(k, \eta)$  are quantities like  $\delta_\gamma$ ,  $V_b$ ,  $\Phi$ , etc, which are solutions of a system of linear differential equations, sourced by the defects.

$$X(k, \eta) = \int d\eta' S(k, \eta') \mathcal{G}(k, \eta, \eta')$$

$$\langle X(k, \eta) X(k, \eta') \rangle = \int d\eta_1 d\eta_2 \underbrace{\langle S(k, \eta_1) S(k, \eta_2) \rangle}_{\text{UETC}} \mathcal{G}(\dots) \mathcal{G}(\dots)$$

$$\text{UETC} : F(k, \eta_1, \eta_2) = F(k\eta_1, k\eta_2)$$

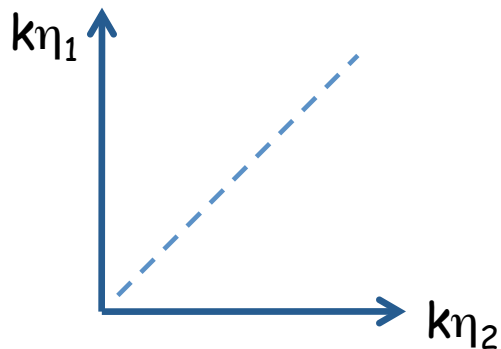
# big universe, small computer



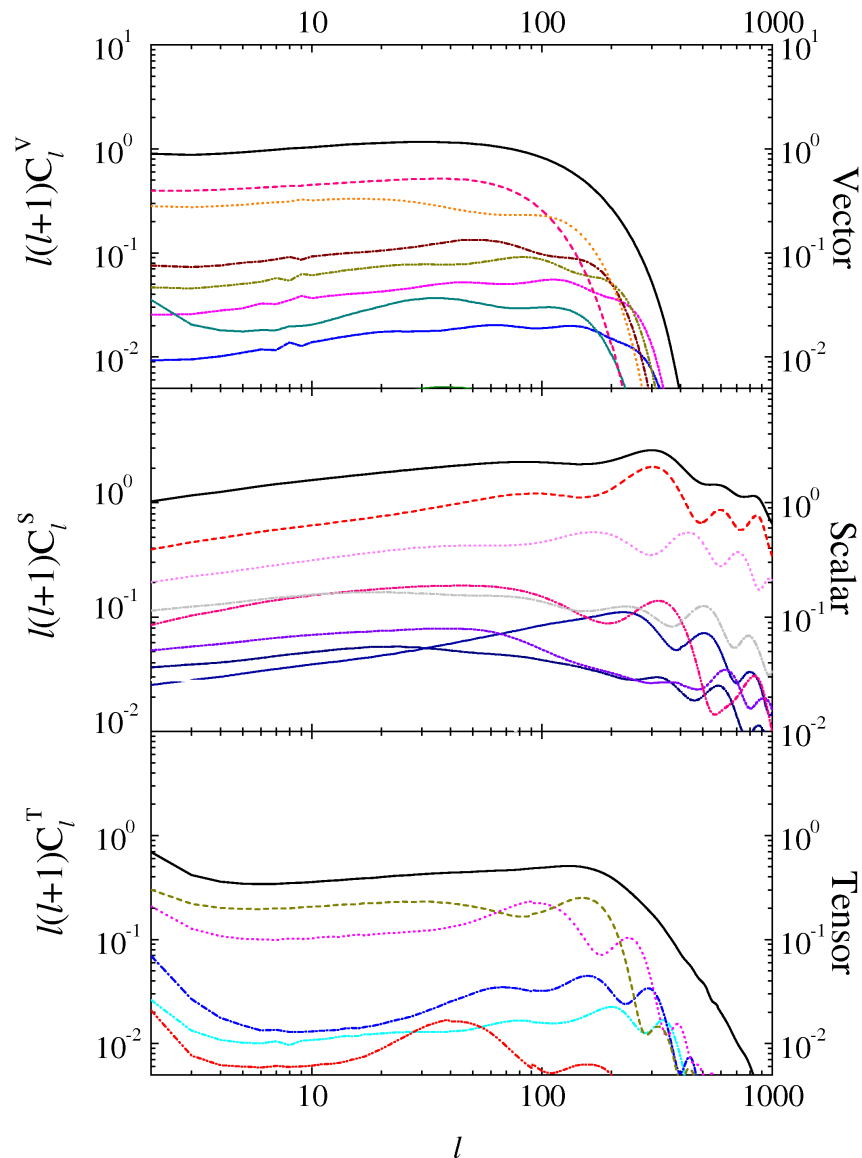
# why there are no peaks

We can think of the two-point function  $F(k\eta_1, k\eta_2)$  as a matrix that can be diagonalised. Each eigenvector  $v_i(k, \eta_1)$  becomes a source.

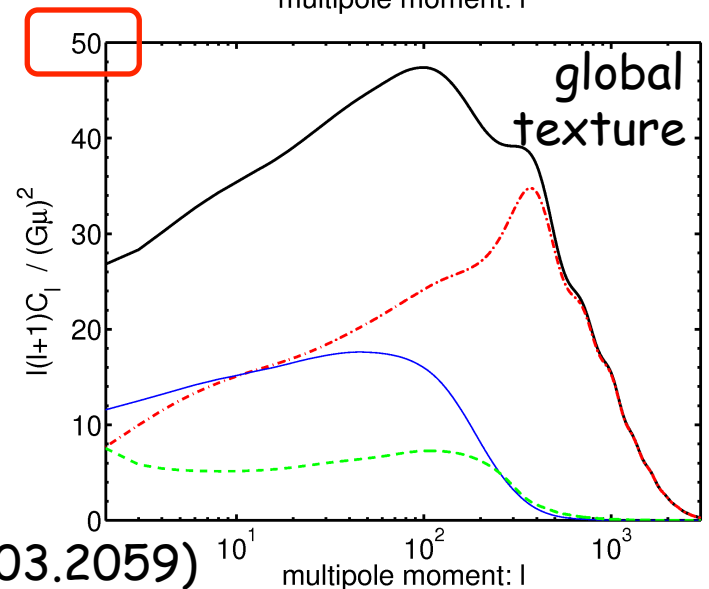
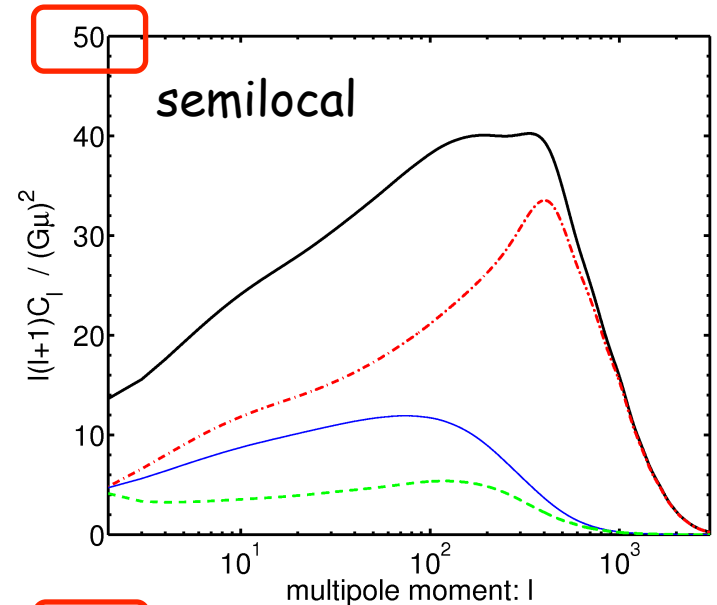
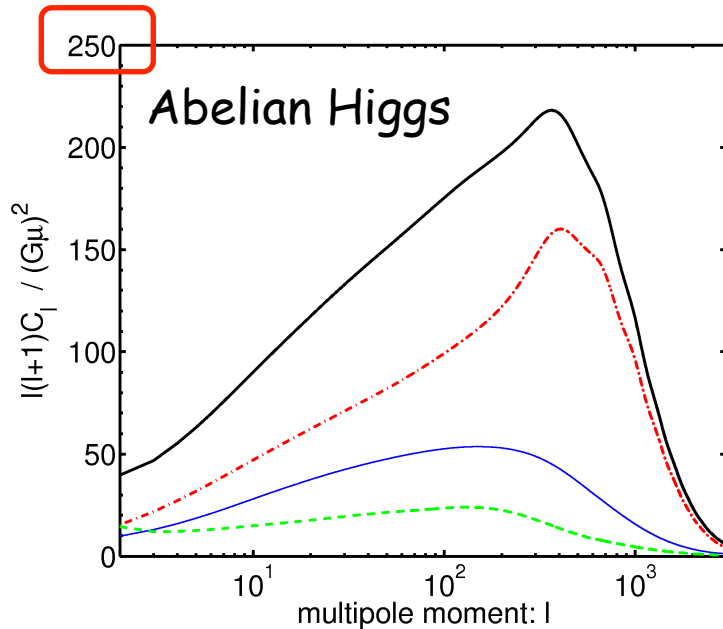
$$C_\ell \sim \sum_i \lambda_i \underbrace{\int dk \left| \int d\eta v_i(k, \eta) \mathcal{G}(\dots) \right|^2}_{C_\ell^{(i)}}$$



Each source has (small) oscillations, but they do not add coherently and the peaks get smeared out.



# CMB 'TT' power spectra



total (sum)

scalar (~ density perturbations)

vector (~ vorticity)

tensor (gravitational waves)

(defects generically have all three contributions!)

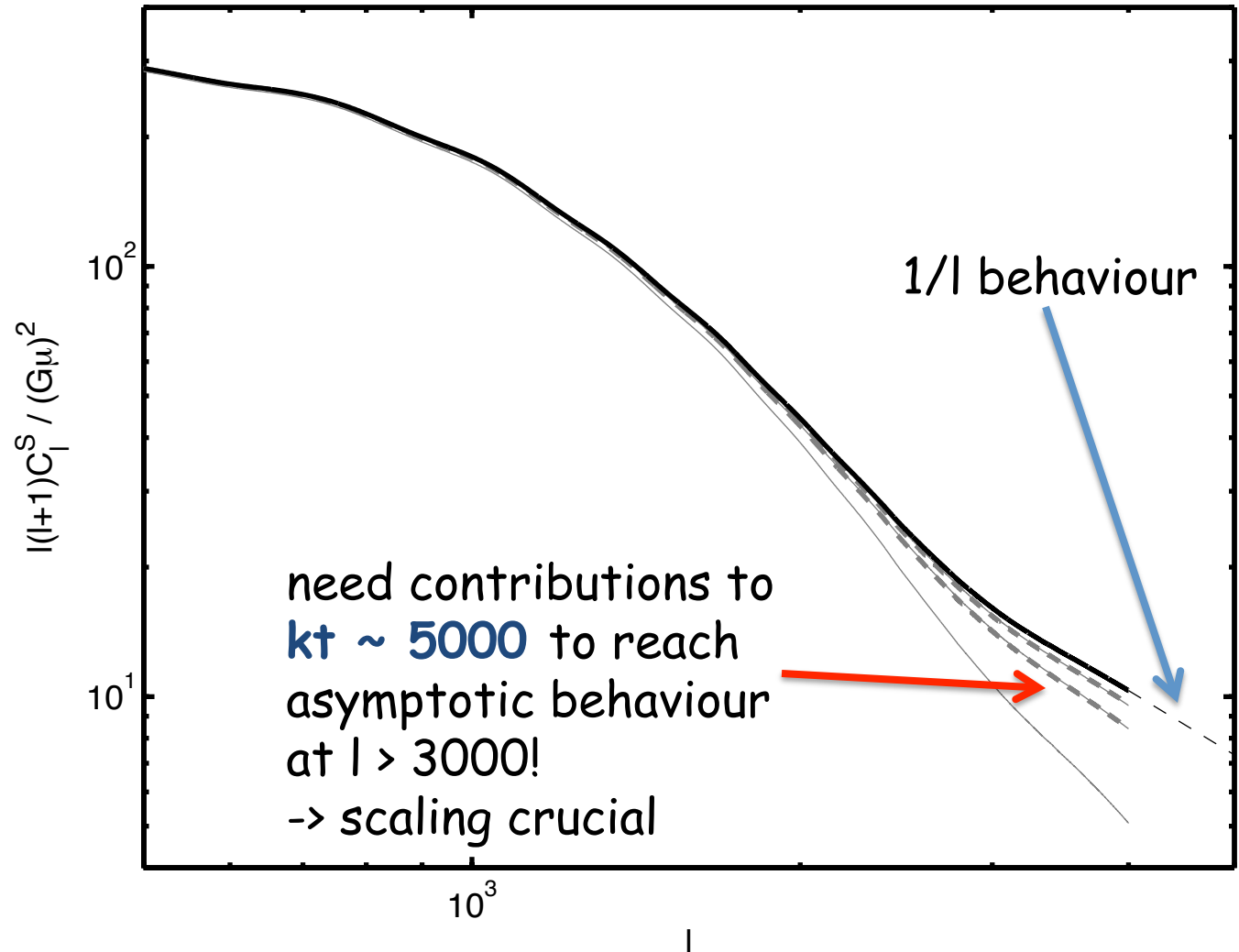
(Urrestilla et al, arXiv:0803.2059)

# high- $l$ behaviour

cosmic strings expected to scale like  $l(l+1)C_l \sim 1/l$

primary CMB  
drops very  
fast at high  $l$

so maybe  
strings will  
dominated  
eventually?



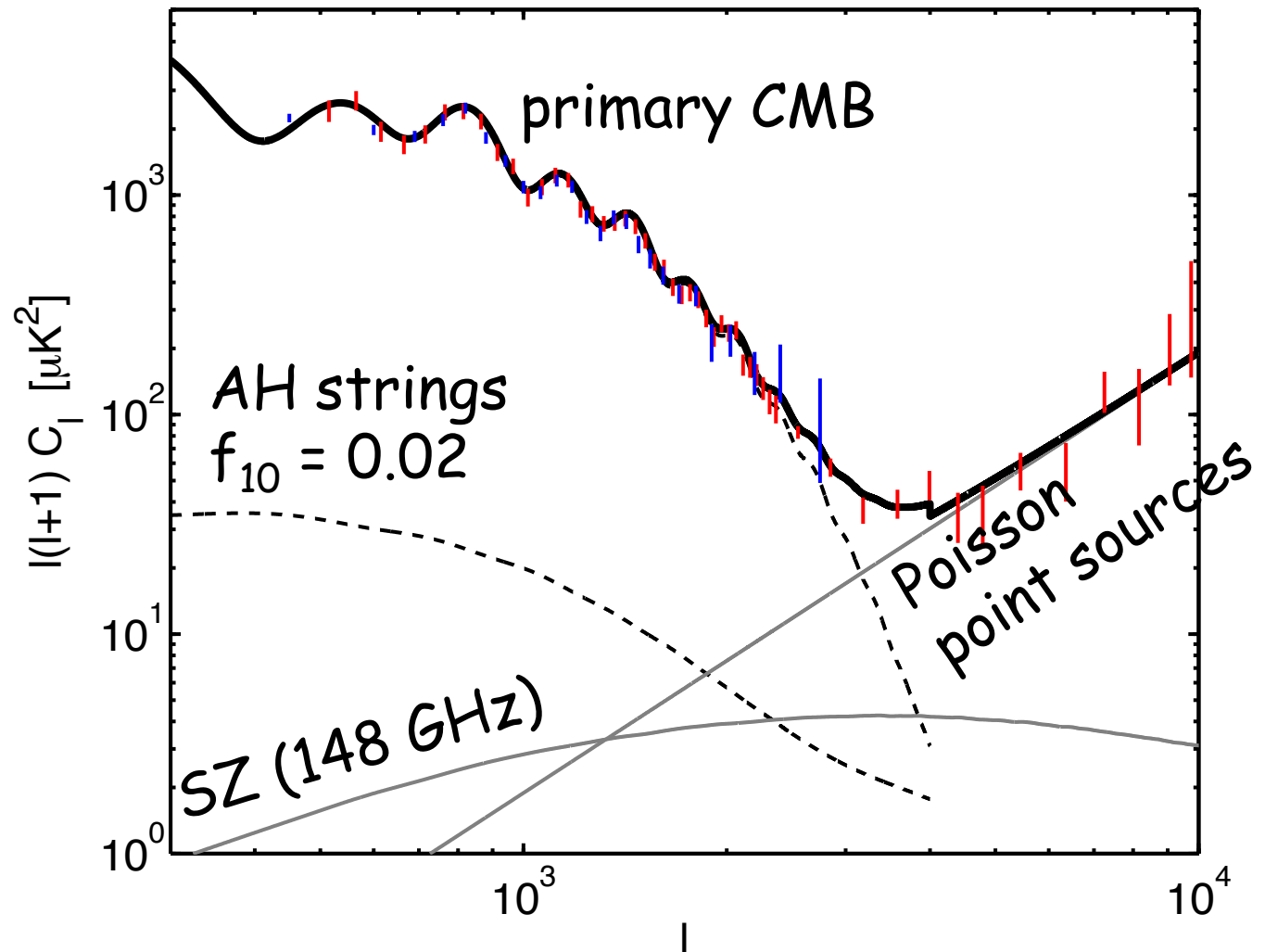
# high- $l$ behaviour

high- $l$  region strongly dominated by point sources  
completely buries string (and SZ) signal

data:

ACT

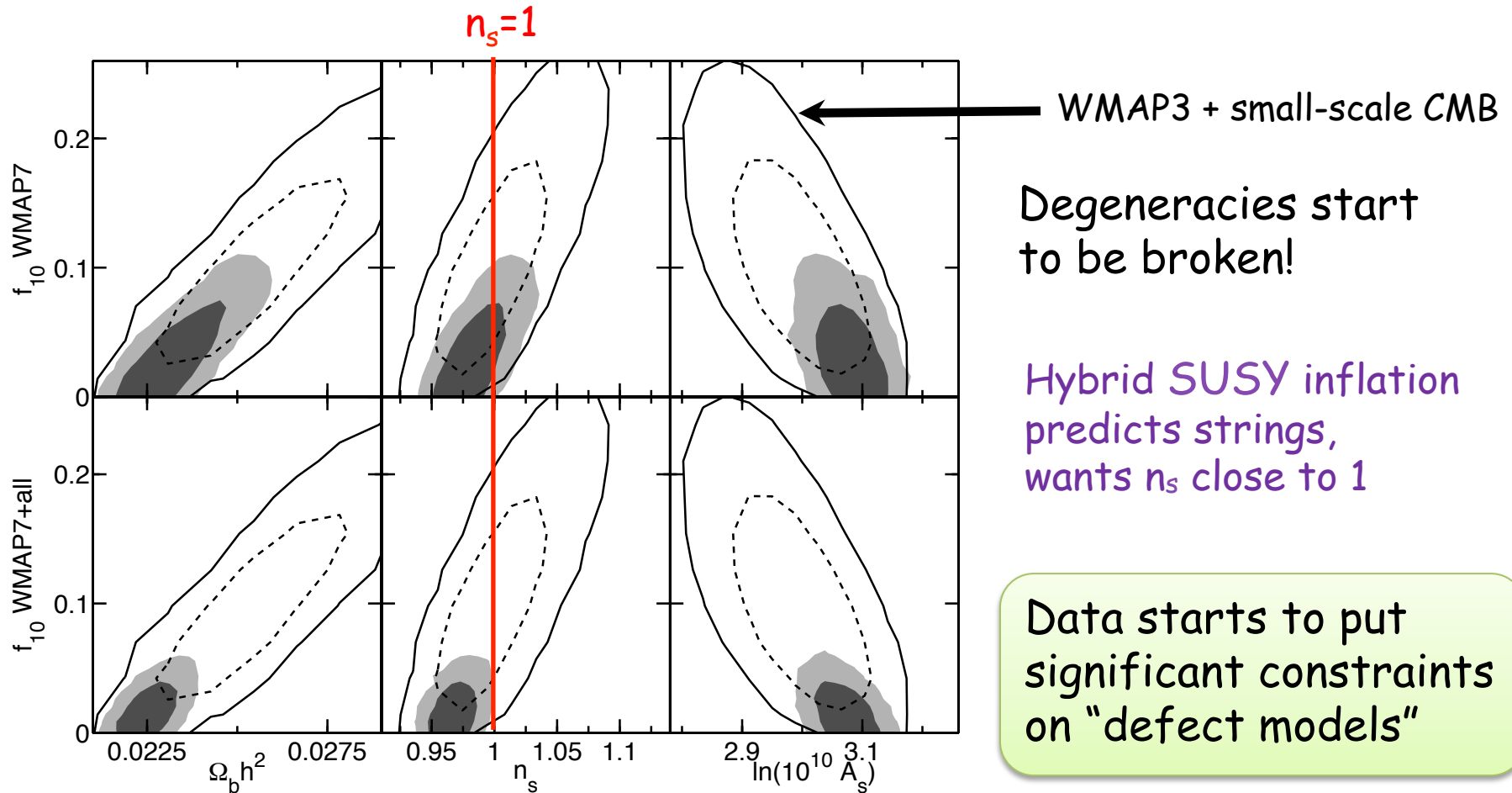
ACBAR





# did WMAP measure a tilt?

MCMC with CMB (WMAP7, ACBAR+QUAD+ACT) for AH Urrestilla et al, arXiv:1108.2730



( $n_s=1$  & strings when varying e.g.  $N_{\text{eff}}$ , see Lizarraga, Sendra, Urrestilla 2012)

# cosmic string constraints

$G\mu$ : string scale  
(1 = Planck scale)

$f_{10}$ : ratio of  $C_l$  from inflation  
and defects at  $l = 10$

| Model          | Data set                      | $10^6 G\mu$ (95%) | $f_{10}$ (95%) |
|----------------|-------------------------------|-------------------|----------------|
| AH [25]        | WMAP3+BOOMERANG+CBI+ACBAR+VSA | 0.7               | 0.11           |
| AH (this work) | WMAP7                         | 0.57              | 0.095          |
| AH (this work) | WMAP7 + ACBAR + QUAD + ACT    | 0.42              | 0.048          |
| USM-AH [35]    | WMAP5                         | 0.68              | 0.11           |
| USM-NG [35]    | WMAP5                         | 0.28              | 0.054          |
| USM-NG [5]     | WMAP7+ACT                     | 0.16              |                |

- slight preference for strings seen in WMAP3 has gone away
- notice model dependence: AH vs NG  
→ mostly due to higher string density in NG simulations (cannot model decay into massive radiation)

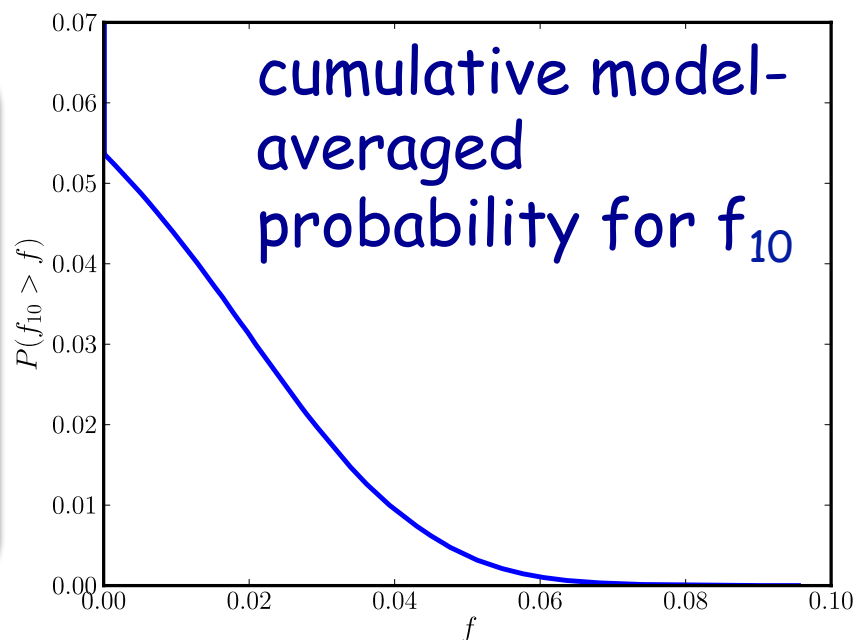
# how about “defect models”?

## Are “cosmic strings” in trouble?

- we can always only get an upper limit on  $G\mu$
- but we can address the question with Bayesian model comparison
- depends on priors ... we use flat priors  $0.75 < n_s < 1.25$ ;  $0 < f_{10} < 1$
- we compute the Bayes ratio  $B$  with the Savage-Dickey method
- 4 models: 'PL', 'HZ', 'PL+AH', 'HZ+AH'

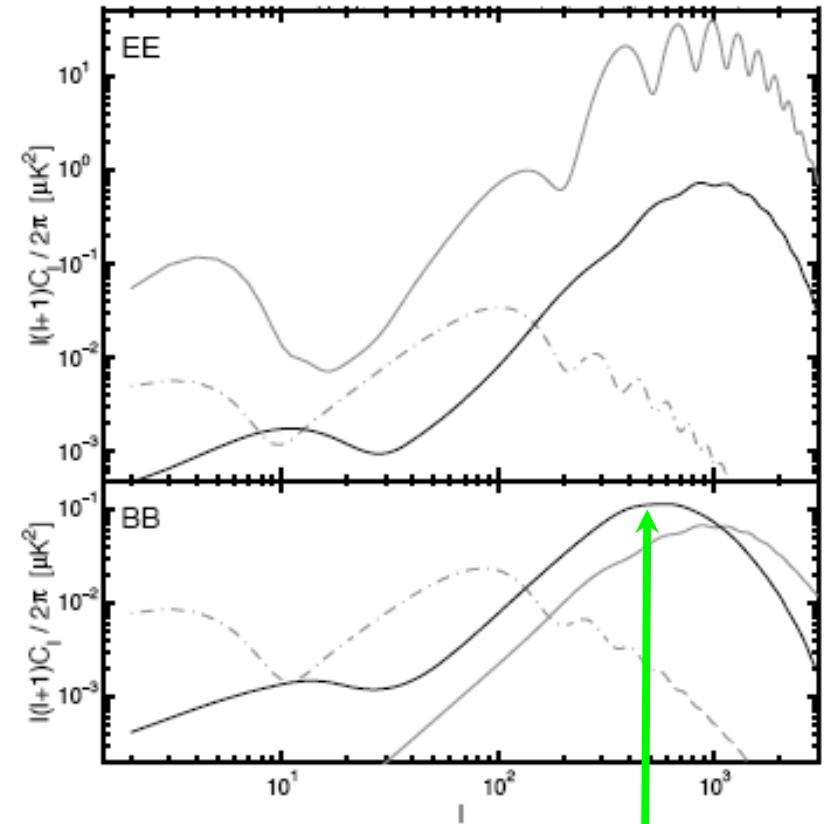
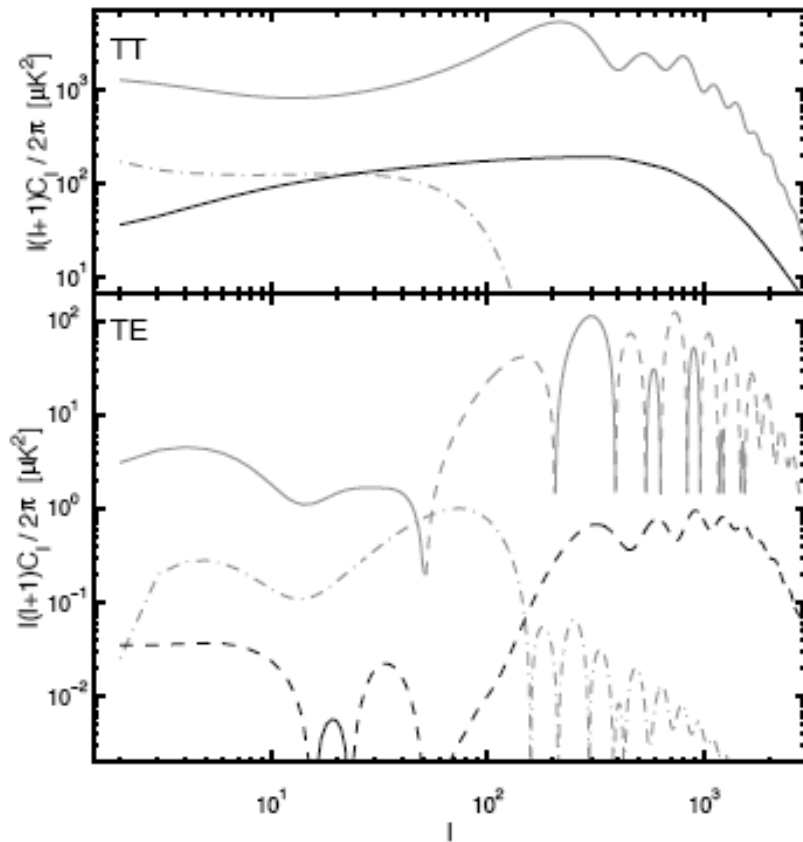
->  $\ln B [\text{AH}] \sim -3$   
-> presence of Abelian-Higgs strings moderately disfavoured!

-> nearly 95% of probability is in 'no-strings' models



Current CMB data starts to put pressure on strings!

# smoking gun: B-mode polarisation



Inflation  $r=0.4$  and strings  $f_{10}=0.1$

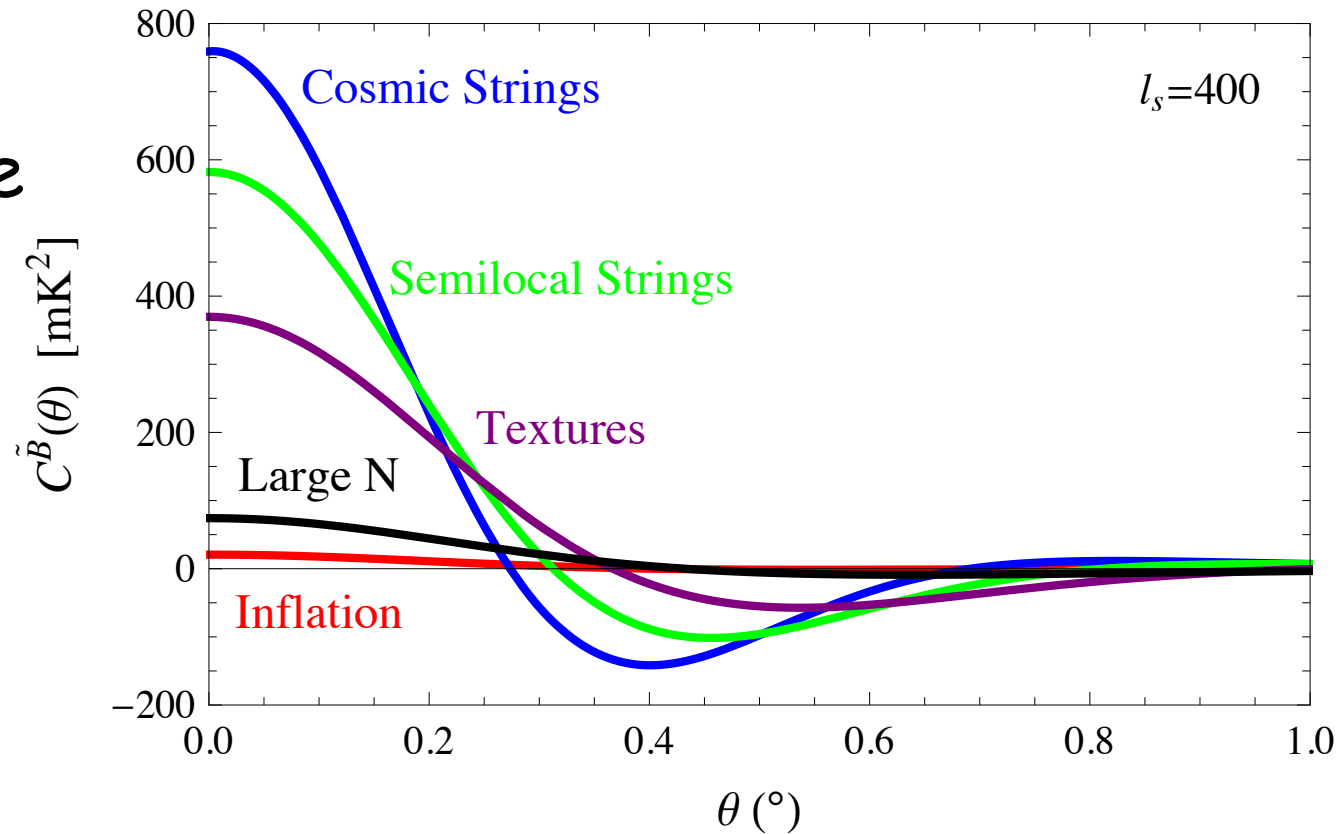
AH STRINGS!

topological defects can dominate in B-mode polarisation spectrum

# angular correlation function

other way to look at B-mode polarisation:  
angular correlation function of local B-modes!

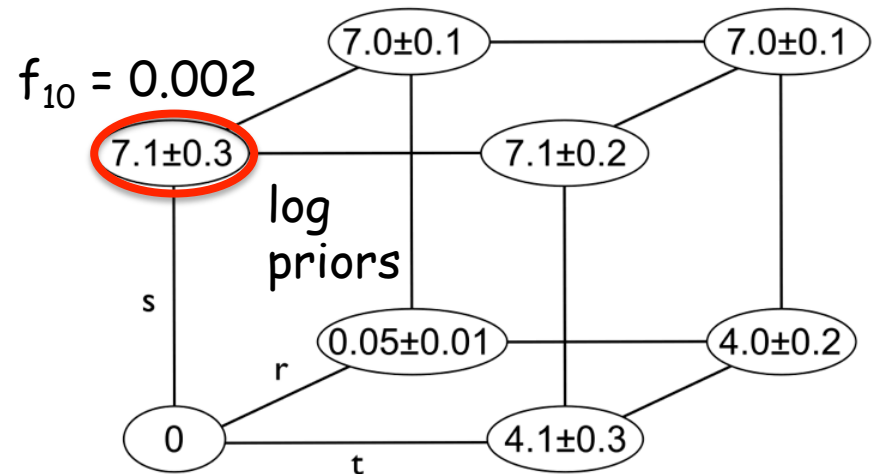
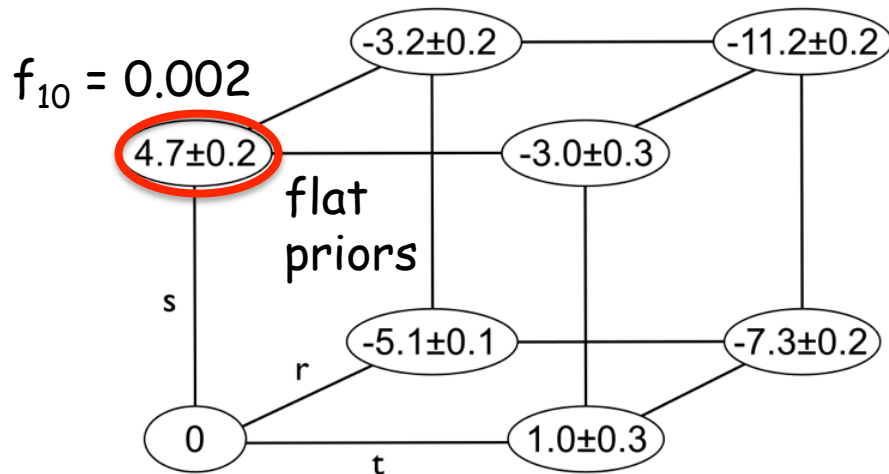
- $f_{10} = 0.1$
- defects have a huge peak in local B-mode correlation function!



(Garcia-Bellido, Durrer, Fenu, Figueroa, MK, arXiv:1003.0299)

# how about future CMB data?

- **Planck:** (Urrestilla et al, arXiv:0803.2059)
  - can distinguish between tensors and strings
  - should reach limit  $f_{10} \sim 0.01$  at 95%
- **CMBPol / COrE:** (Mukherjee et al, arXiv:1010.5662)
  - can distinguish between different defects
  - could reach  $f_{10} \sim 0.001$  at 95% ( $G\mu < 6 \times 10^{-8}$ )



# P(k) and N-body simulations

Do strings affect galaxy clustering?

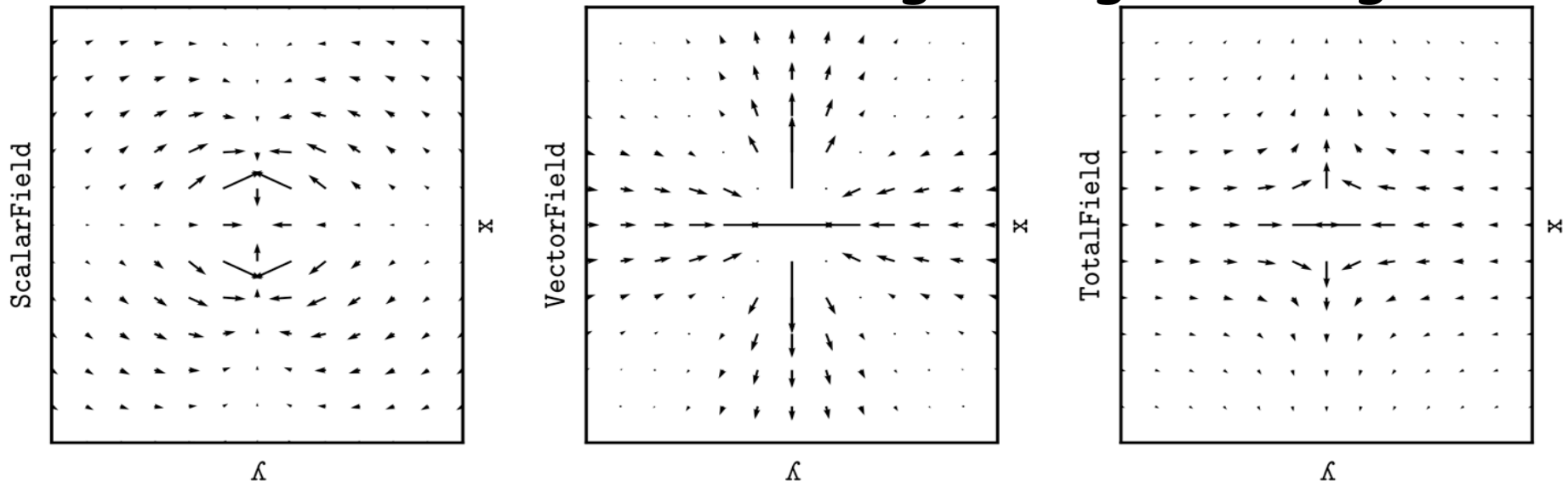
-> important question if we want to use P(k)

-> we include strings in N-body simulations, using

$$\ddot{x}_i + 2\frac{\dot{a}}{a}\dot{x}_i = -\frac{1}{a^2}\partial_i\psi + \frac{1}{a}\Sigma_i + \frac{\dot{a}}{a^2}\Sigma_i$$

(Obradovic et al,  
arXiv:1106.5866)

ex: acceleration field of long straight string



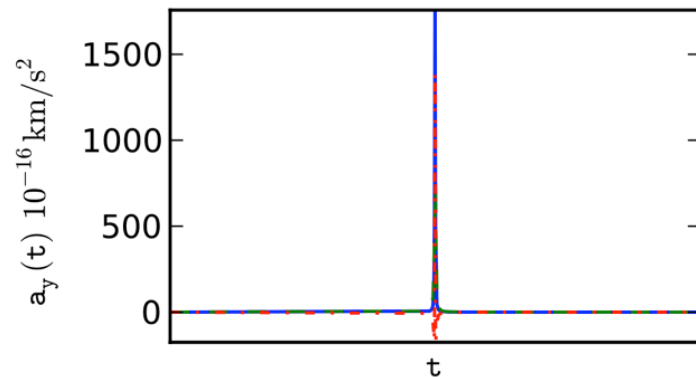
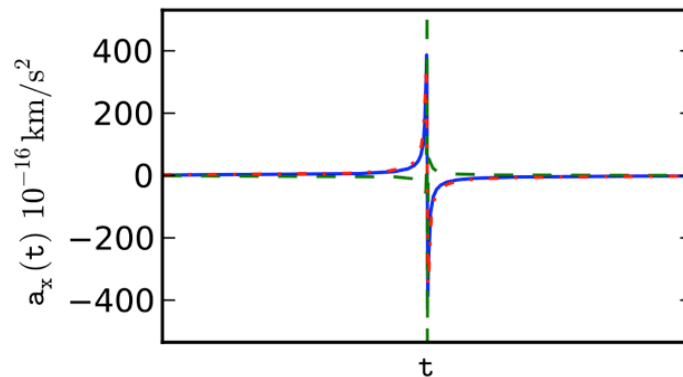
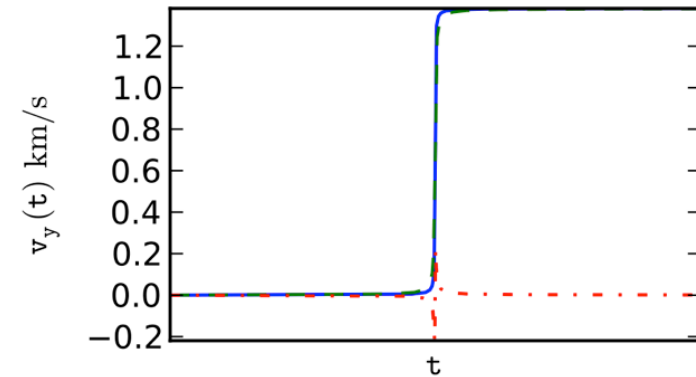
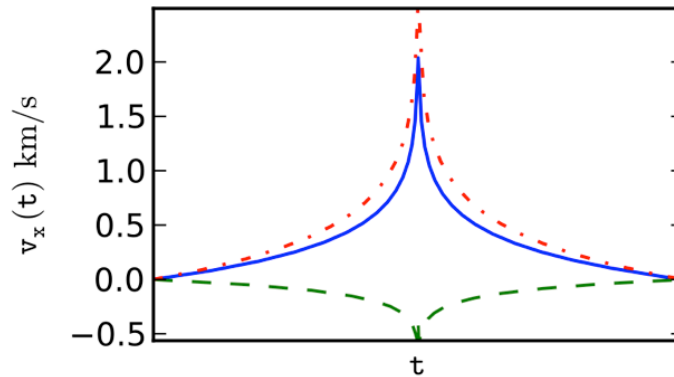
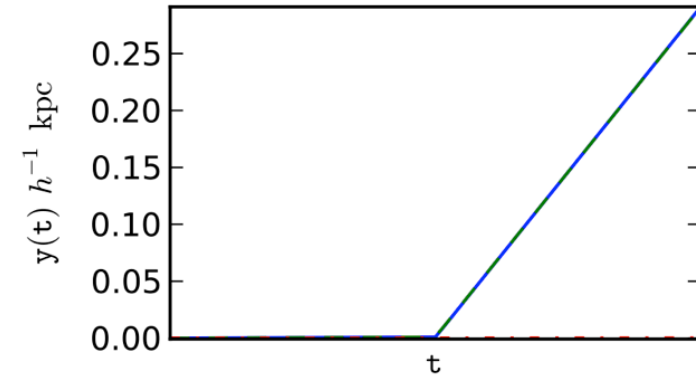
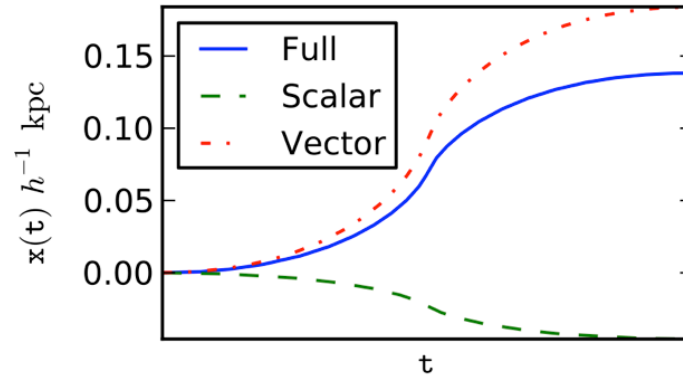
# velocity kick of straight string

analytical  
solution:  
velocity kick

reproduced  
correctly by  
N-body  
simulation

working on  
including  
actual string  
simulation

-> watch this  
space! ☺





# summary & outlook

- topological defects **exist!**
- at least in the lab...
- ... but also in many (most? all?) realistic inflation models
- **data starts to constrain defect models non-trivially**  
-> impact on inflation models?
- **(AH:  $f_{10} < 0.05$ ,  $G\mu < 0.4 \times 10^{-6}$  @ 95%)**
- defects generically have vector and tensor perturbations  
-> **B-mode polarisation** (and **non-Gaussianity**) in the CMB
- Planck should improve limits by a factor of 5-10 in  $f_{10}$   
-> waiting for spring 2013 ... ☺
- large-scale structure may provide an additional window  
-> work in progress