
String theory

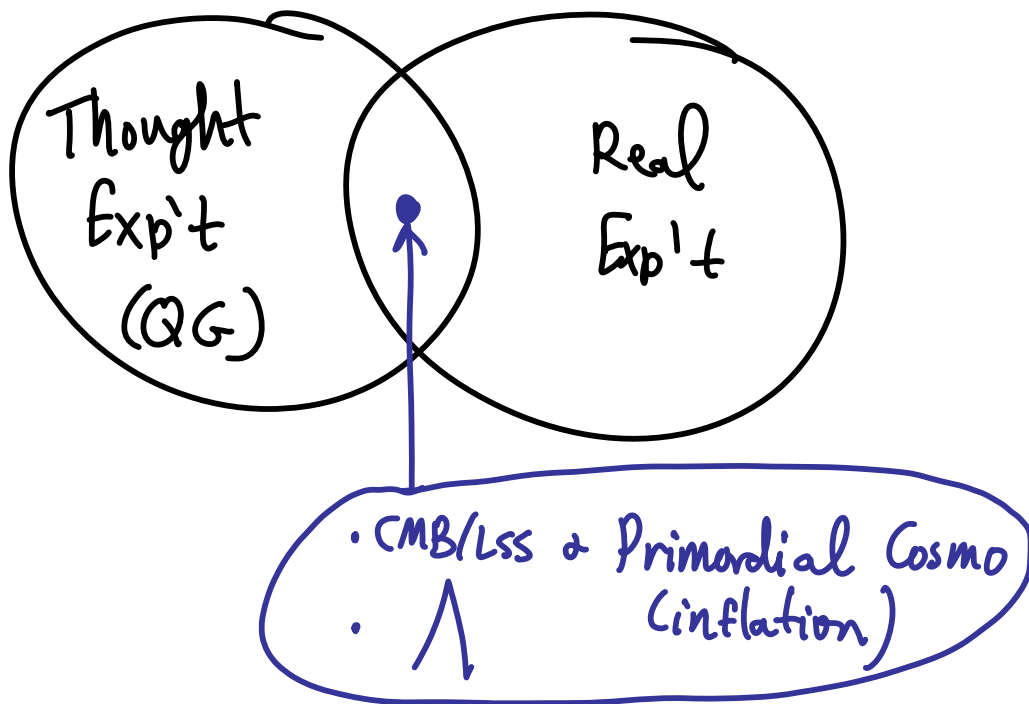
&

astrophysics

Unsolved Problems

String theory: UV completion of GR

- Reproduces GR at long distance, no UV divergences
- "Holographic": Number of degrees of freedom $\sim \frac{\text{Area}}{4G_N}$
- QFT results (formal & model-building mechanisms)
- Math



Plan

I. Thought Experiments

II. Real Experiments

"The scientific theorist is not to be envied. For Nature, or more precisely experiment, is an exorable and not very friendly judge of his work. It never says "yes" to a theory. In the most favorable cases it says "Maybe," and in the great majority of cases simply "No." If an experiment agrees with a theory it means for the latter "Maybe," and if it does not agree it means "No." Probably every theory will some day experience its "No" - most theories, soon after conception."

--Entry into memory book for Professor Kammerling-Onnes, November 11, 1922; quoted in Dukas and Hoffmann, *Albert Einstein, the Human Side*, p.18.

I. Black Hole Physics:

$$\Delta A \geq 0 \text{ in classical GR}$$

↑
area of
event horizon

$$\Delta M = K \Delta A \text{ in classical GR}$$

↑ surface gravity

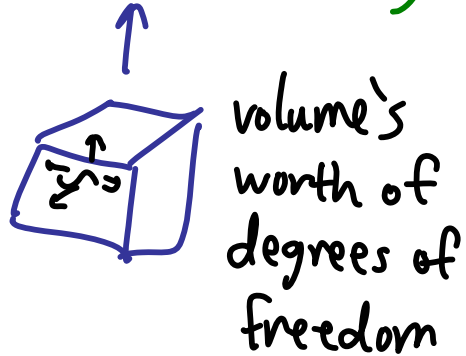
$$T \sim K \text{ Quantum mechanically}$$

Dynamics of BHs \Leftrightarrow Thermodynamics

$$\frac{\text{Area}}{G_N} \sim (\text{Area } M_p^2) \Leftrightarrow \text{Entropy } S$$

Gravity \neq Local QFT
(d dimensions) (d dimensions)

↑
IF you pack
too much entropy
in local region
 \Rightarrow form a BH, $S \sim$ Area



String theory: BH Stat. Mech.

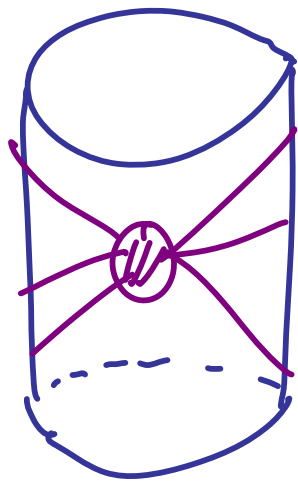


Explicit microstate count
for tractable BH's ... \rightarrow

Gravity = Local QFT !
(d dimensions, (d-1 dimensions)
e.g. AdS) \Rightarrow Unitary BH evolution

$\Lambda < 0$ (AdS_4) has a precise
non-perturbative formulation.

4d GR + strings \downarrow QFT₃, no gravity

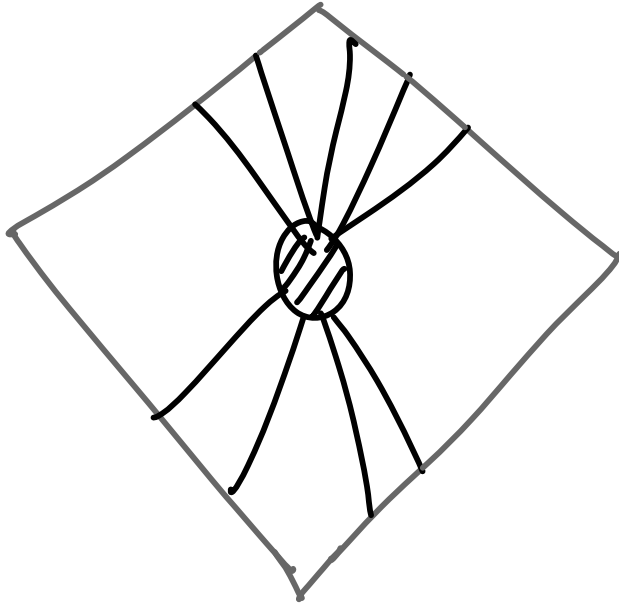


\cong CFT₃ |
I

\uparrow many distinct
examples, don't
mix, no horizon

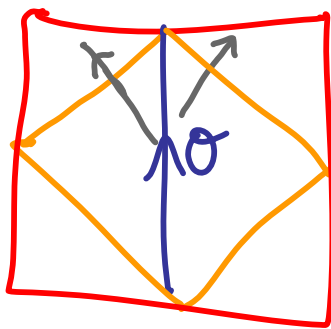
Observables = QFT Correlation
Functions.

For $\Lambda = 0$, the general framework
(an S-matrix is also known



An observer can causally collect
the results.

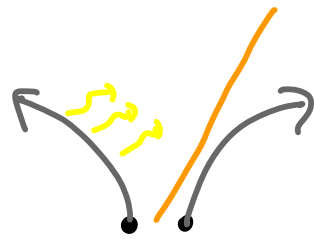
In contrast,
in dS :



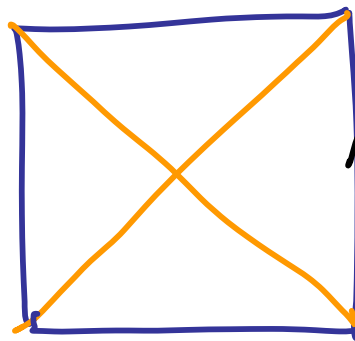
Classically
appears as
though info
lost (cf Black
Holes)

Mysteries of Dark Energy

$w < -\frac{1}{3} \Rightarrow$ horizon

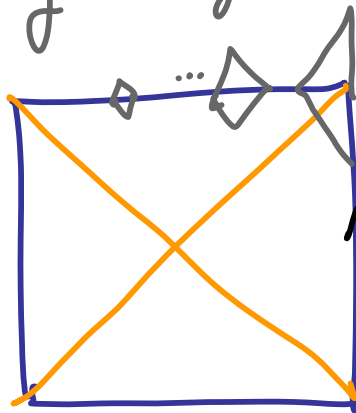


e.g. dS $w = -1$



observer \mathcal{O}

Including decays:



cf decay to $\Lambda = 0$

observer \mathcal{O}

still causally disconnected in future

Dynamical connections between different Λ , & different D , topology, fluxes, ...

Gibbons/Hawking : de Sitter

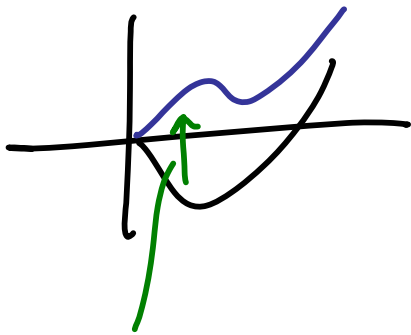
$$\text{Horizon} \Rightarrow S \sim \frac{M_p^2}{H^2}$$

Involves a microstate count.

↳ Get parametrically by building up from AdS/CFT

add ingredients to "uplift"

interpret new ingredients



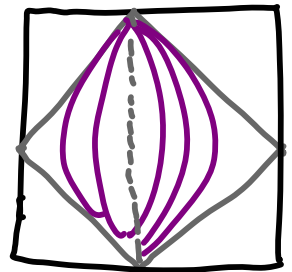
Magnetic Matter

heavy branes overcoming AdSxS⁵ curvature

- Attempts to build up from AdS/CFT make some progress, but leave propagating Gravity in the dual

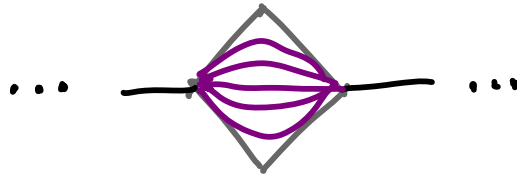
dS/dS $ds_{dS_4}^2 = \sin^2 \frac{w}{L} ds_{dS_3}^2 + dw^2$

$\Rightarrow \cdot 2 \text{QFT}_3 + \text{GR}_3$



- UV-complete construction, concrete dual \rightarrow dS₄ entropy (order of magnitude)

FRW/CFT



Euclidean AdS₃/CFT₂
+ 2d Gravity

- Attempts at probability "measures" interesting but still ad hoc.

II. Thought Exp't \wedge Real Exp't

General Relativity describes gravity accurately
at long distances

$$S = \int d^4x \sqrt{g} \frac{R}{16\pi G_N} + S_{\text{matter}} \rightarrow R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G_N T_{\mu\nu}$$

GR breaks down for $\lambda_G \rightarrow 1$ (or before)

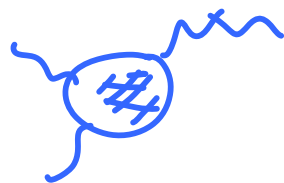
Quantum fluctuations / classical UV physics \rightarrow

$$S = \int \left(\frac{R}{16\pi G_N} - V(\phi) \right) \left(1 + R \left(\frac{c_1}{M_*^2} + \tilde{c}_1 G_N \right) + \dots \right)$$

$$+ \int (\partial\phi)^2 + k_1 \frac{(\partial\phi)^4}{M_*^2} + \dots$$

M_*^2 \leftarrow scale of "new physics"

with corrections sensitive to
short-distance physics

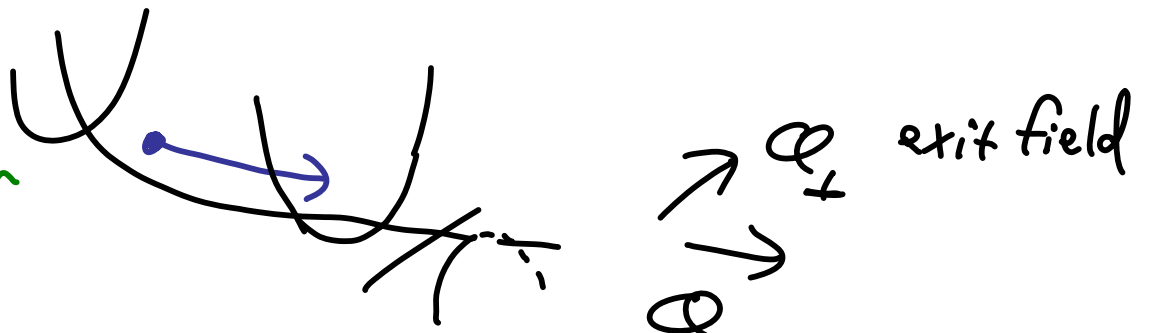


2) These corrections matter
for inflation

e.g. A seemingly simple way to obtain
inflation is to postulate a very flat
potential for the inflaton $\phi(x)$.

Linde '93

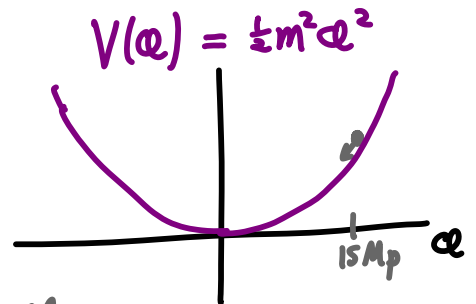
hybrid inflation



$$\epsilon \equiv \frac{M_p^2}{2} \left(\frac{V'}{V} \right)^2 \ll 1 \quad \eta \equiv M_p^2 \left| \frac{V''}{V} \right| \ll 1$$

However, corrections from the UV
physics can generate substructure in
 $\mathcal{L}(\phi, \partial\phi)$: $\frac{V(\phi - \phi_0)^2}{M_p^2} \rightarrow \Delta\eta \sim 1$

This UV sensitivity is greatest in the case of "chaotic inflation" A. Linde '83 where the inflaton ϕ ranges over more than a distance M_p e.g.



$$\left\{ \begin{array}{l} \epsilon = \frac{1}{2} \left(\frac{V'}{V} M_p \right)^2 \\ \eta = M_p^2 \frac{V''}{V} \end{array} \right\} \sim \left(\frac{M_p}{\phi} \right)^2 \Rightarrow \phi \sim 15 M_p$$

In General:

★ "Lyth Bound": $\frac{\Delta \phi}{M_p} \sim \left(\frac{r}{0.01} \right)^{\frac{1}{2}}$

UV sensitive if ≥ 1 \ observable

→ Control with approximate shift symmetry (Wilsonian 'natural')

UV Sensitivity of Inflation

① Terms of order $V \cdot \frac{(\mathcal{Q} - \mathcal{Q}_0)^2}{M_p^2}$
(dimension 6)

in the effective action can ruin inflation

② $\frac{\Delta \mathcal{Q}}{M_p} \simeq r^{\frac{1}{2}} \frac{N_e}{\sqrt{24}}$ (Lyth)

GUT-scale inflation (with observable

tensor modes) $\Leftrightarrow \Delta \mathcal{Q} > M_p$

③ General Single-field inflation involves higher derivative terms which affect solution & perturbations

④ $g^2 \mathcal{Q}^2 \chi^2$ couplings \Rightarrow temporarily Non-Gaussianity
light fields affect evolution. . . .

5

(mass $> H$)

Heavy fields affect results in a different way: they adjust in response to inflationary potential energy.

QFT toy model

$$V(\phi_L, \phi_H) = g^2 \phi_L^2 \phi_H^2 + m^2 (\phi_H - \phi_0)^2$$

$$\frac{\partial V}{\partial \phi_H} \equiv 0 \Rightarrow V = \frac{g^2 \phi_L^2}{g^2 \phi_L^2 + m^2} m^2 \phi_0^2$$

(ϕ_H^2 term subdominant)

flatter: energetically favorable.

6 Exit physics: defects, strings, oscillons, ...?

7 Entry Physics ?? bubble nucleation ..

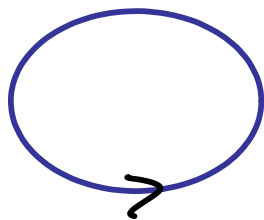
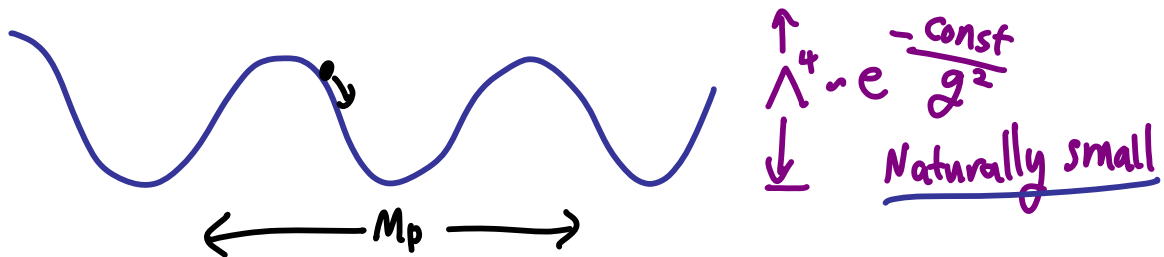
Tensor Modes

Large $\frac{\Delta Q}{M_p}$.

Axions naturally respect an (approximate)

shift symmetry $Q \rightarrow Q + \alpha$
(couple via their derivatives)

→ "Natural Inflation"



$a \cong a + (2\pi)^2$
 $Q_a = f_a a$ — canonical scalar field

→ Does $\frac{\Delta Q}{M_p} \gtrsim 1$, protected by shift symmetry, arise in string theory?

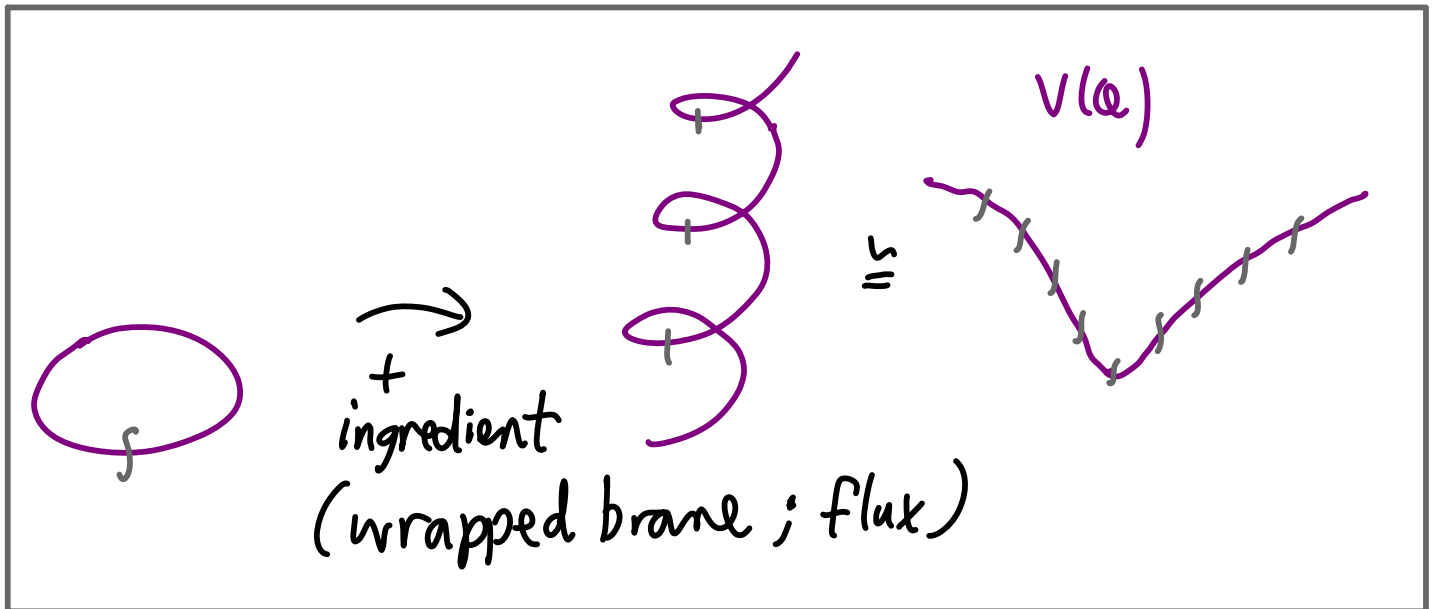
* Basic period small compared to M_p

For axions, $f \ll M_p$
 in currently controlled regions of
 the landscape. (size $L \gg M_p^{-1}$)

$$\int d^4x \sqrt{-g} \underbrace{|dC_p|^2}_{g^{i_1 j_1} \dots g^{i_{p+1} j_{p+1}}} = \int d^4x \sqrt{-g} \frac{M_p^2}{(LM_p)^{2p}} (\partial\theta)^2$$

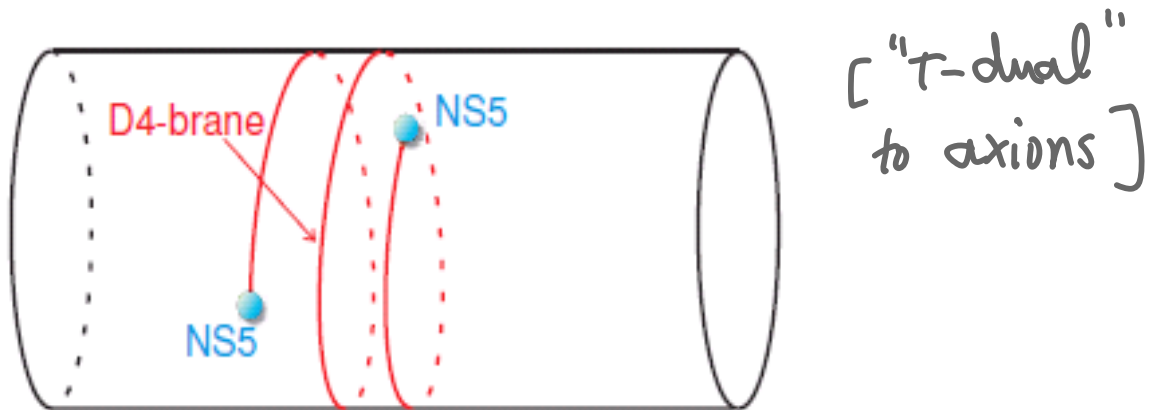
* Not "anything goes" in the
 landscape!

... But must take into account Monodromy
in string compactifications



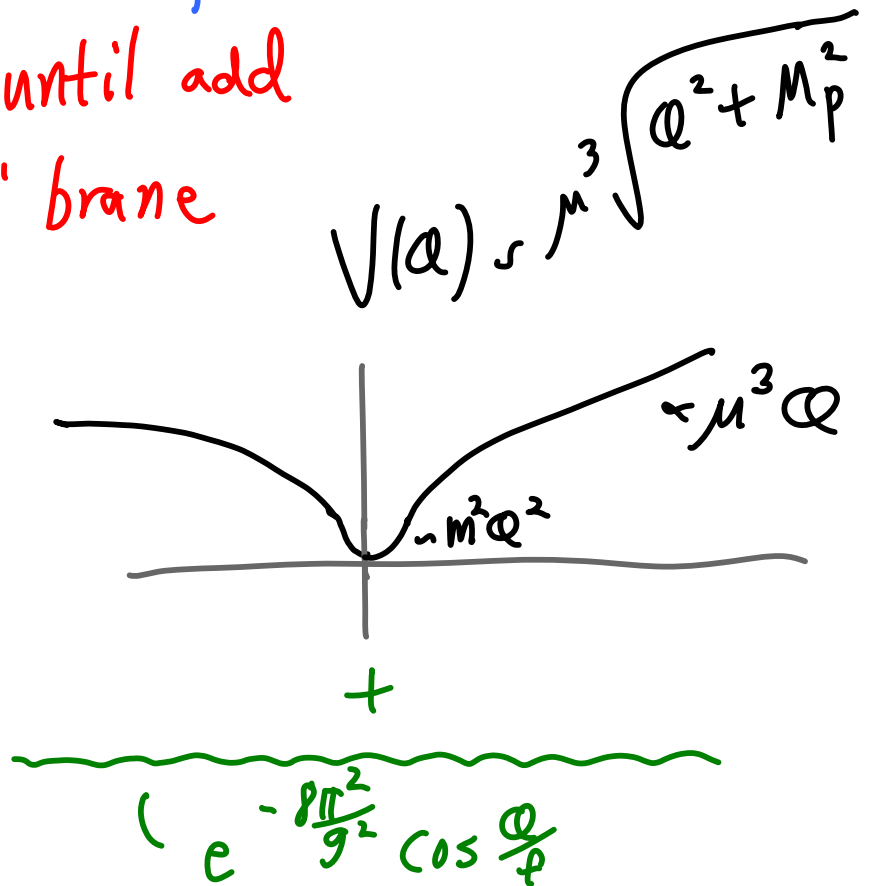
unwraps the would-be periodic
direction. \rightarrow Large field range
with distinctive potential, with
 $V(\phi > M_p) \sim \begin{cases} \phi^{2/3} & \text{twisted torus} \\ \phi & \text{axions} \end{cases}$
 the so far worked out examples.

The basic mechanism is very simple :

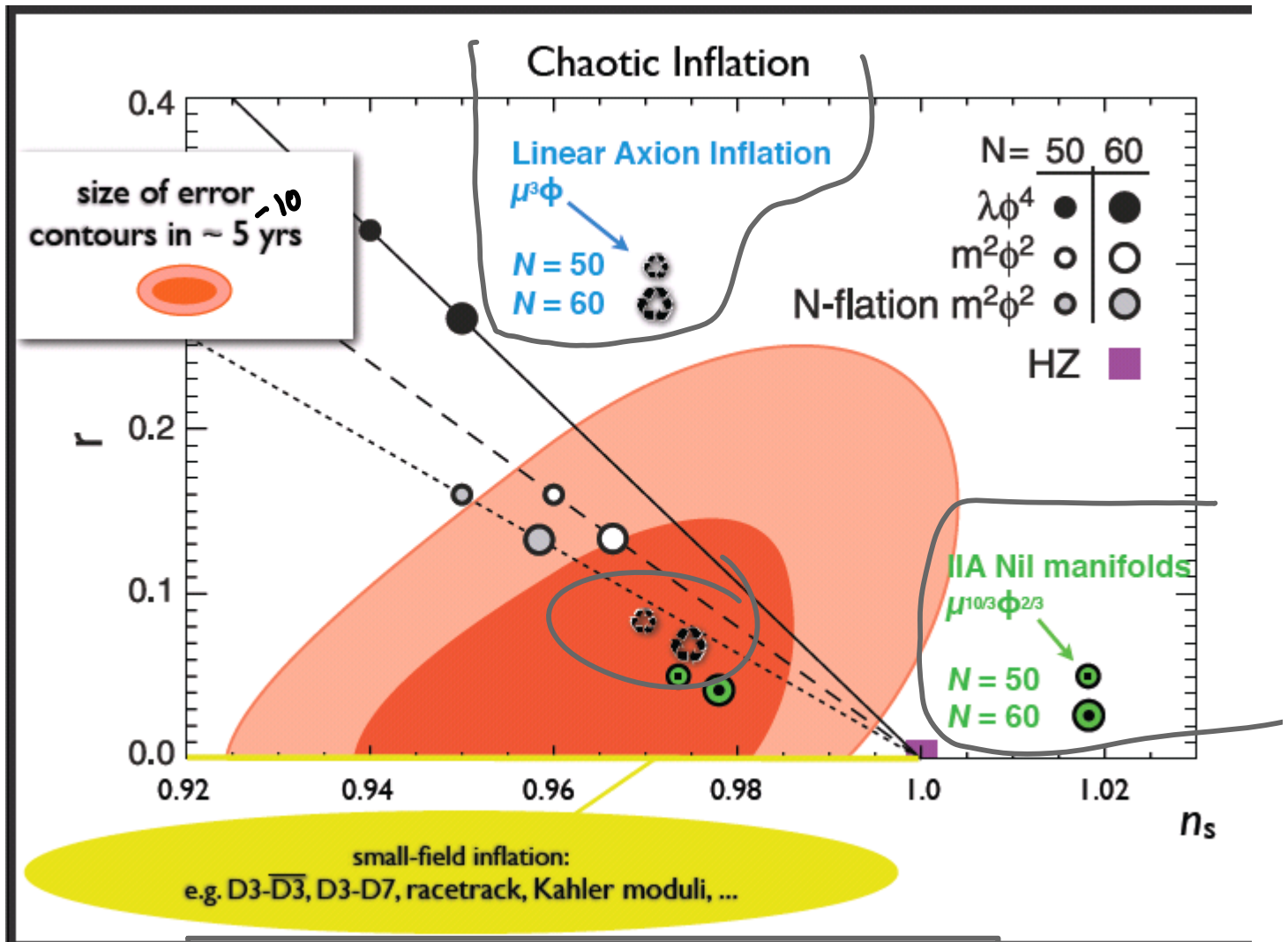


- "NS5" branes position periodic on this circle, until add stretched "D4" brane

→ Novel prediction for inflaton potential



Result: WMAP

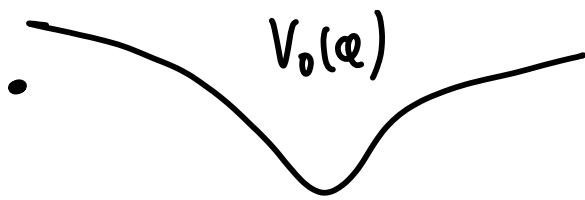


$$r = 0.07$$

$$n_s \approx 0.98$$

$$V(\phi) \approx \mu^3\phi + \Lambda^4 \cos\left(\frac{\phi}{2\pi f}\right)$$

Because of the symmetry, and oscillating nature of the (instanton-suppressed) corrections, these predictions are robust \Rightarrow falsifiable

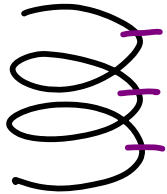


Flattened potential
is the result of
the adjustment
of heavy fields

- Model-dependent oscillations

$$V(\phi) = V_0(\phi) + \Lambda^4 \cos \frac{\phi}{f}$$

can lead to visible signatures
in $\langle s \dots s \rangle$

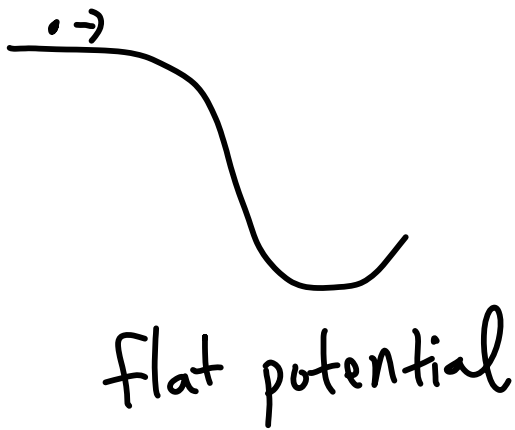
-  repeated production
of particles/strings
→ novel signatures

Open direction: new discovery windows,
Systematics of detailed signatures.

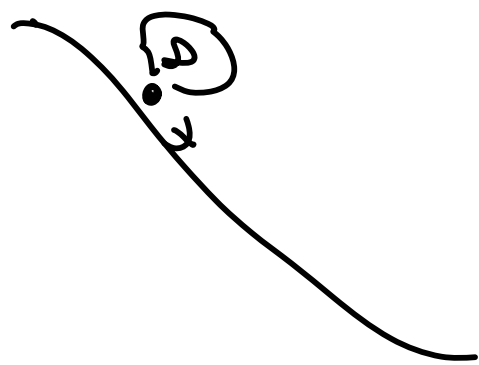
- Small-field inflation also possible,
in other directions in field space.

Non-Gaussianity Roughly Speaking,
2 classes of Inflation Mechanisms:

→ Slowly diluting potential energy



NG only from
substructure, e.g.
oscillations, in $V(\phi)$



Steep potential,
but interactions slow
the field. \Rightarrow NG.

↳ e.g. brane motion
limited by XD
Speed of light (DBI)

Now Systematic (EFT) understanding for
single-field; new effects for multiple
fields

Summary (part II)

- Inflation UV sensitive
- String theory \rightarrow simple mechanisms S
some (focused on here) with
distinctive signatures. e.g.

$$V \sim m^2 Q^2 \quad \leftarrow \text{distinguishable}$$

$$V \sim \mu^3 \sqrt{Q^2 + M_p^2} \quad \text{Pythagorean}$$

\leftarrow UV degrees of freedom

- Motivated systematic analysis in QFT & data analysis.
- Further opportunities – UV sources
of different NG shapes
 - resonant NG data analysis
 - strings

Summary (Part I)

- spacetime defined by $S \sim \frac{A}{4G_N}$
independent degrees of freedom

- In certain systems (symmetric BHs)
String theory has provided a
precise microstate count

- Unsolved Problem: extend to
a precise formulation of cosmology
(and of realistic BH's)

Recent progress (guided by AdS/CFT &
structure of string theory), not complete framework