

# Quantum Darwinism in practice

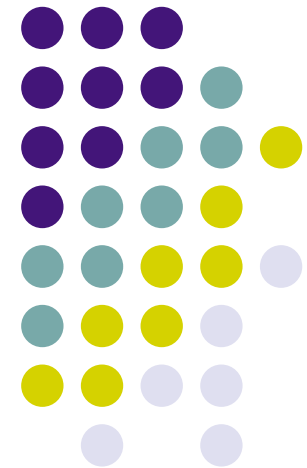
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w/ H. T. Quan, W. H. Zurek

Theoretical Division

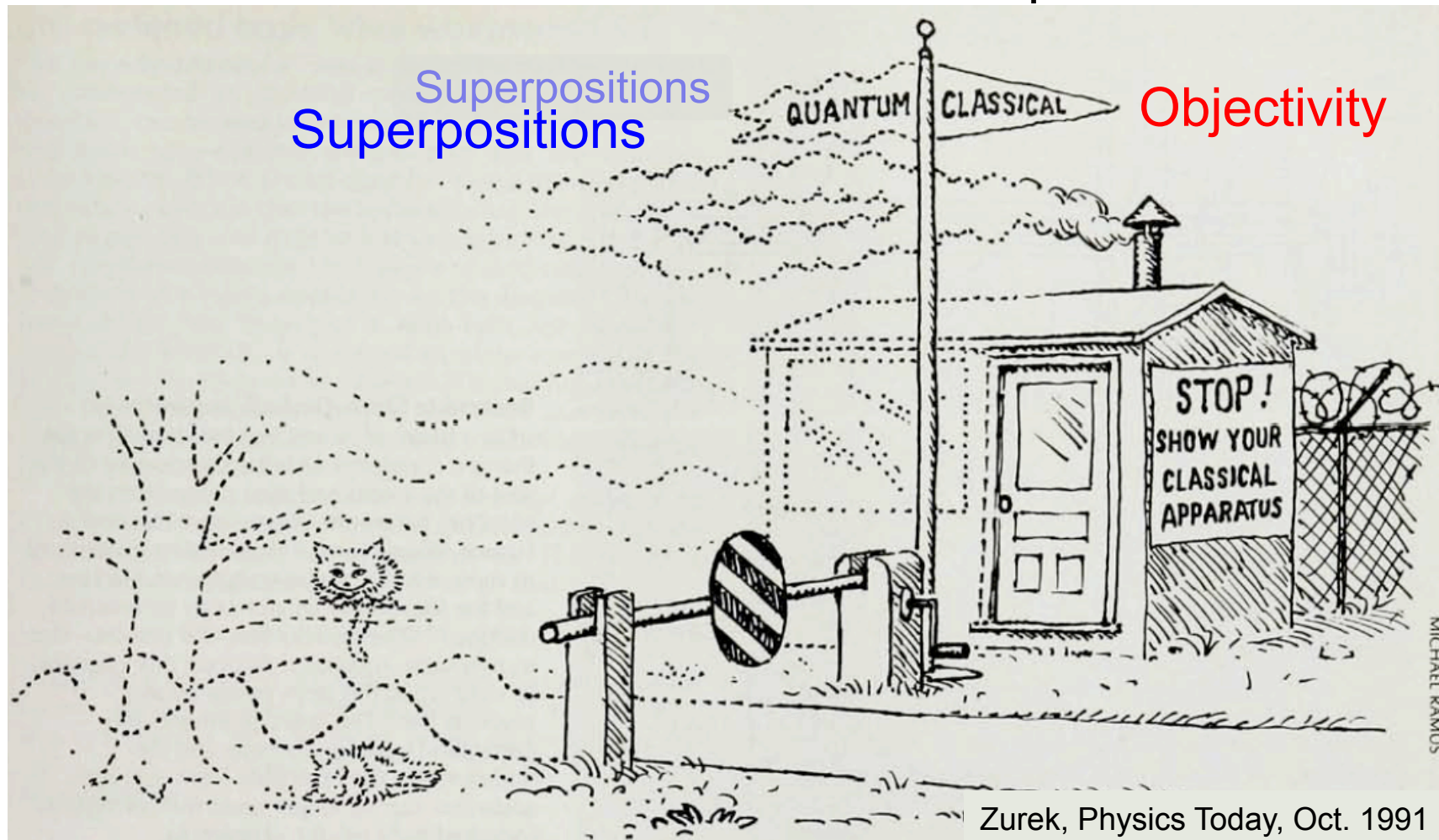
Los Alamos National Laboratory



# Quantum vs. Classical



How does the classical world arise from the quantum substrate?

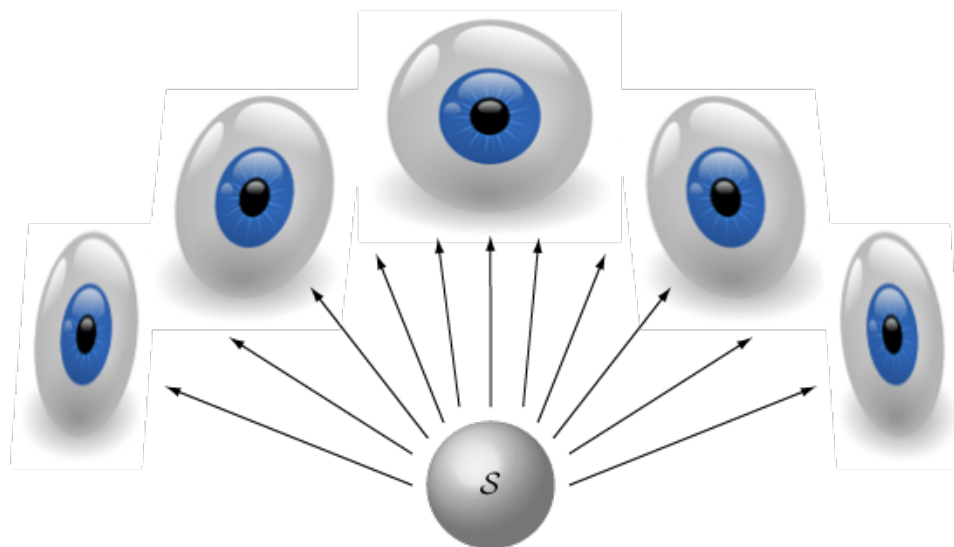




# Quantum Darwinism

Q-Darwinism recognizes the environment's role as a communication channel

- Many observers independently gather information about a system indirectly by intercepting some small fragment of the environment
- Redundant information in the environment gives rise to objectivity





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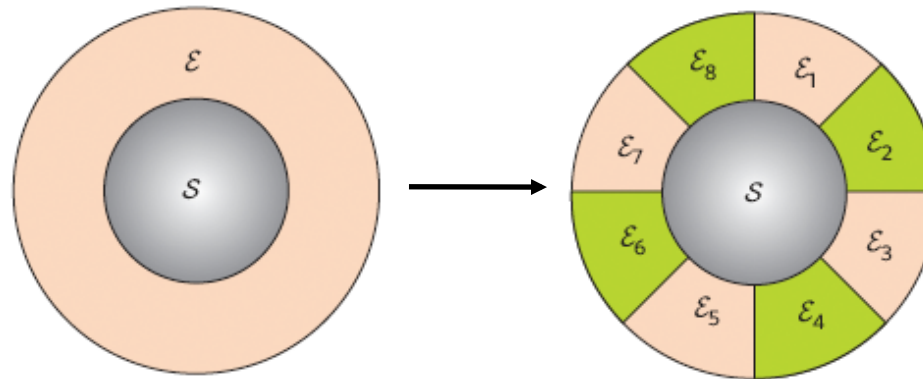
## Why Darwinism?

- Classicality is the result of a selection of preferred states – the pointer states - and the redundant proliferation of information
- Fitness = ability of a state to survive and “procreate” (i.e., spawn the most information theoretic progeny)

# Decoherence vs. Q-Darwinism

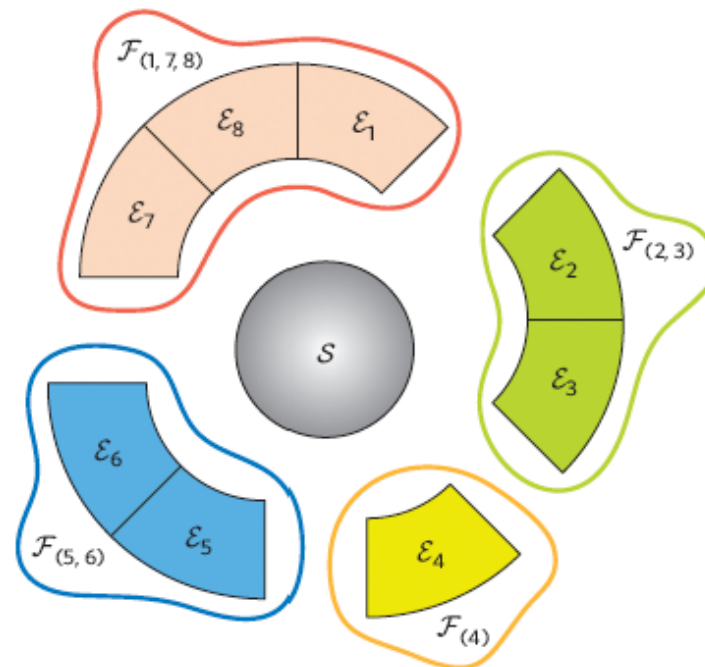


Decoherence  
Paradigm



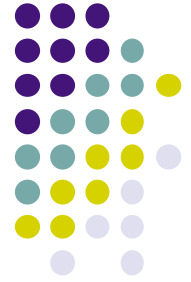
Redundancy  
Paradigm

What fragment  
size gives  
“complete”  
information about  
the system?



Under  
decoherence  
many fragments  
give complete  
information

Ollivier et al.,  
PRL 93, 220401 (2004)  
Blume-Kohout, Zurek,  
PRA 73, 062310 (2006)



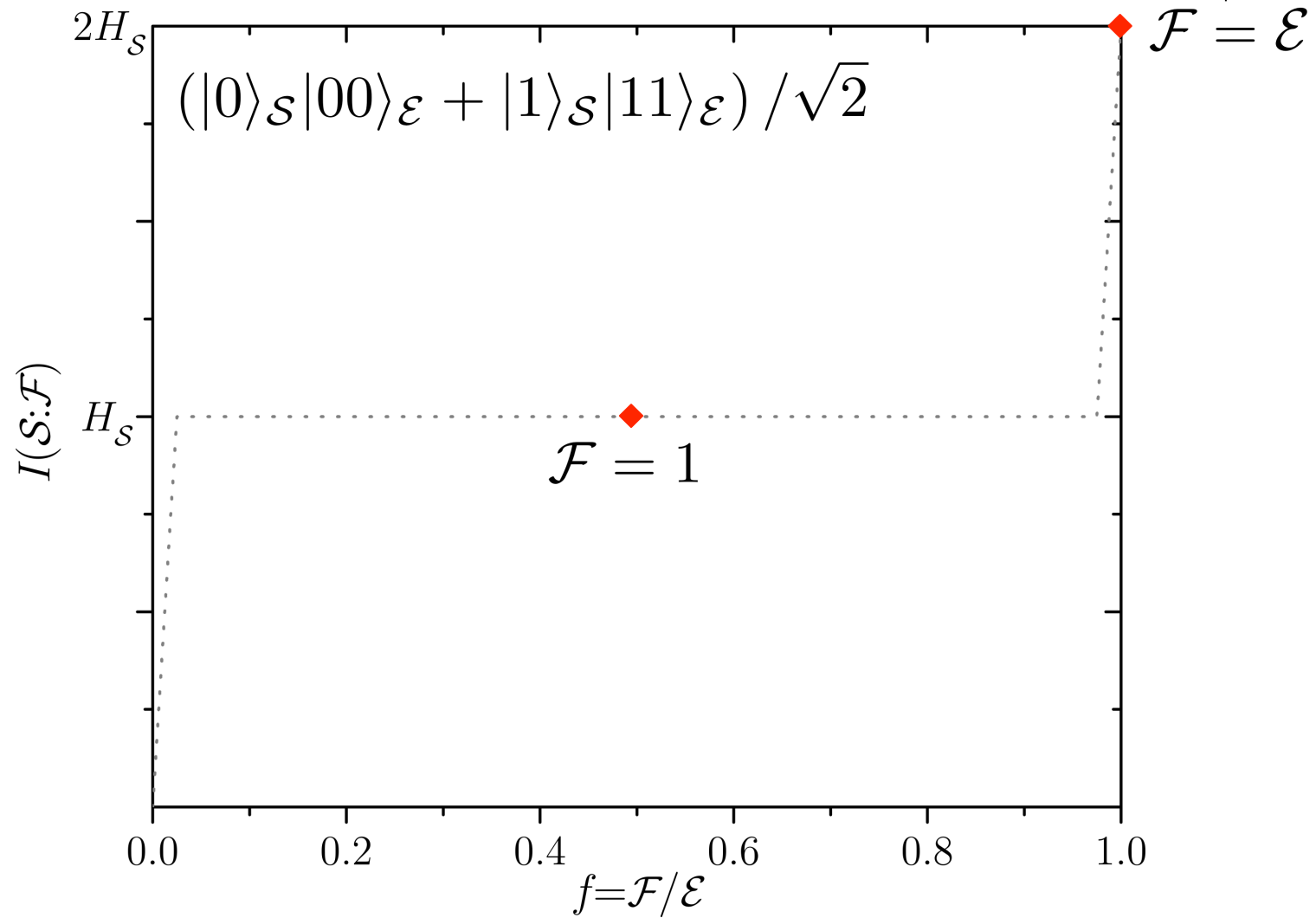
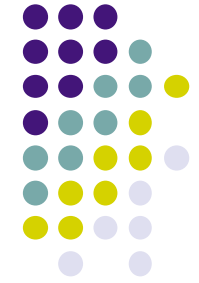
# Correlated qubits

$$\begin{aligned} & (|0\rangle_S + |1\rangle_S) \overset{\mathcal{E}_1}{\downarrow} \overset{\mathcal{E}_2}{\downarrow} |00\rangle_{\mathcal{E}} / \sqrt{2} \\ \xrightarrow{\text{CNOT}} & (|0\rangle_S |00\rangle_{\mathcal{E}} + |1\rangle_S |11\rangle_{\mathcal{E}}) / \sqrt{2} \\ \xrightarrow[\text{tr}_{\mathcal{E}/\mathcal{F}}]{\mathcal{F} = \mathcal{E}_1 \text{ or } \mathcal{E}_2} & \rho_{S\mathcal{F}} = \frac{1}{2} |00\rangle\langle 00| + \frac{1}{2} |11\rangle\langle 11| \end{aligned}$$

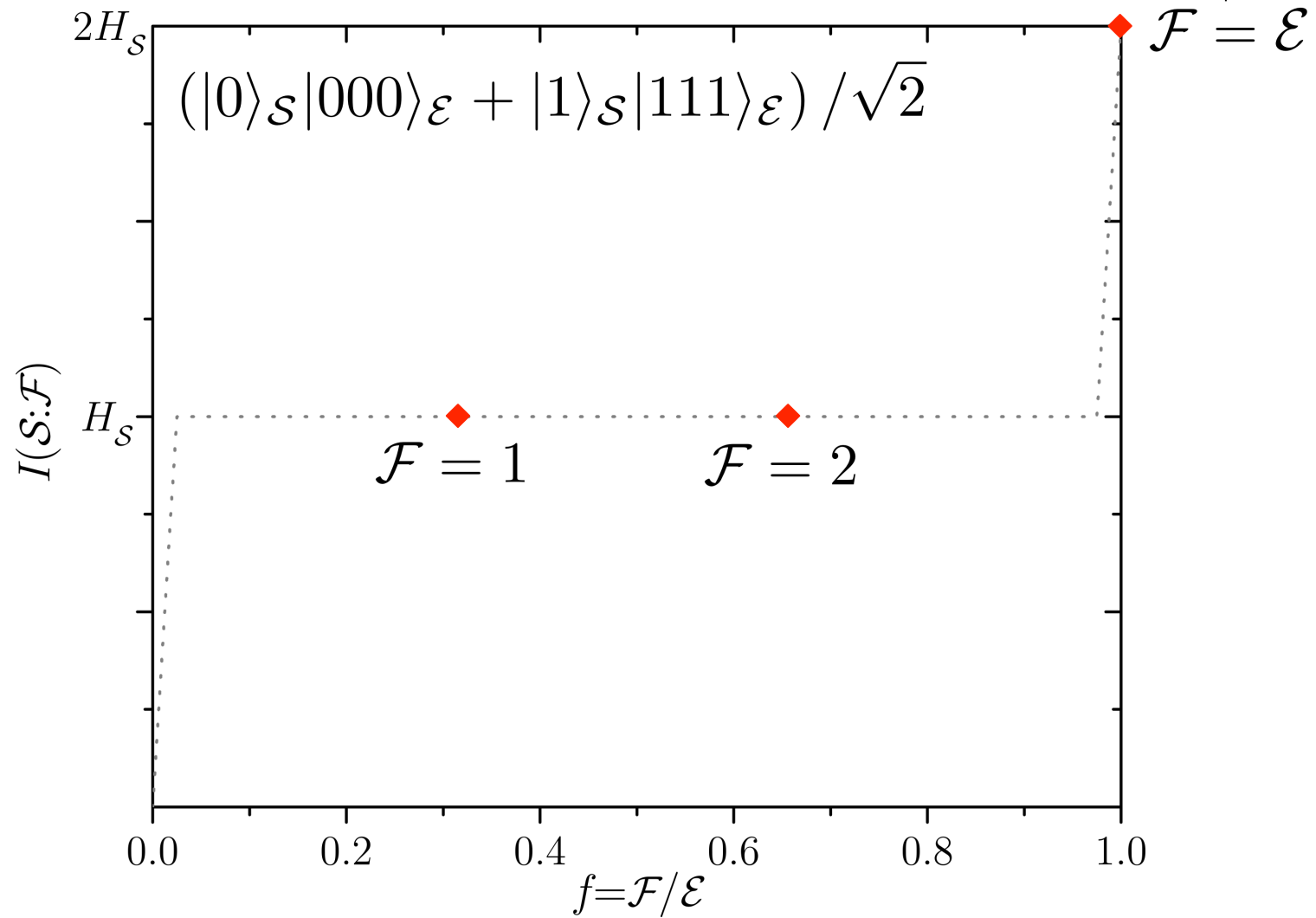
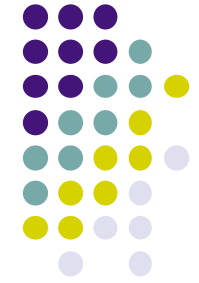
$S, \mathcal{F}$  are classically correlated:

$$p_{0,0} = \frac{1}{2}; p_{1,1} = \frac{1}{2}$$

# Information plot

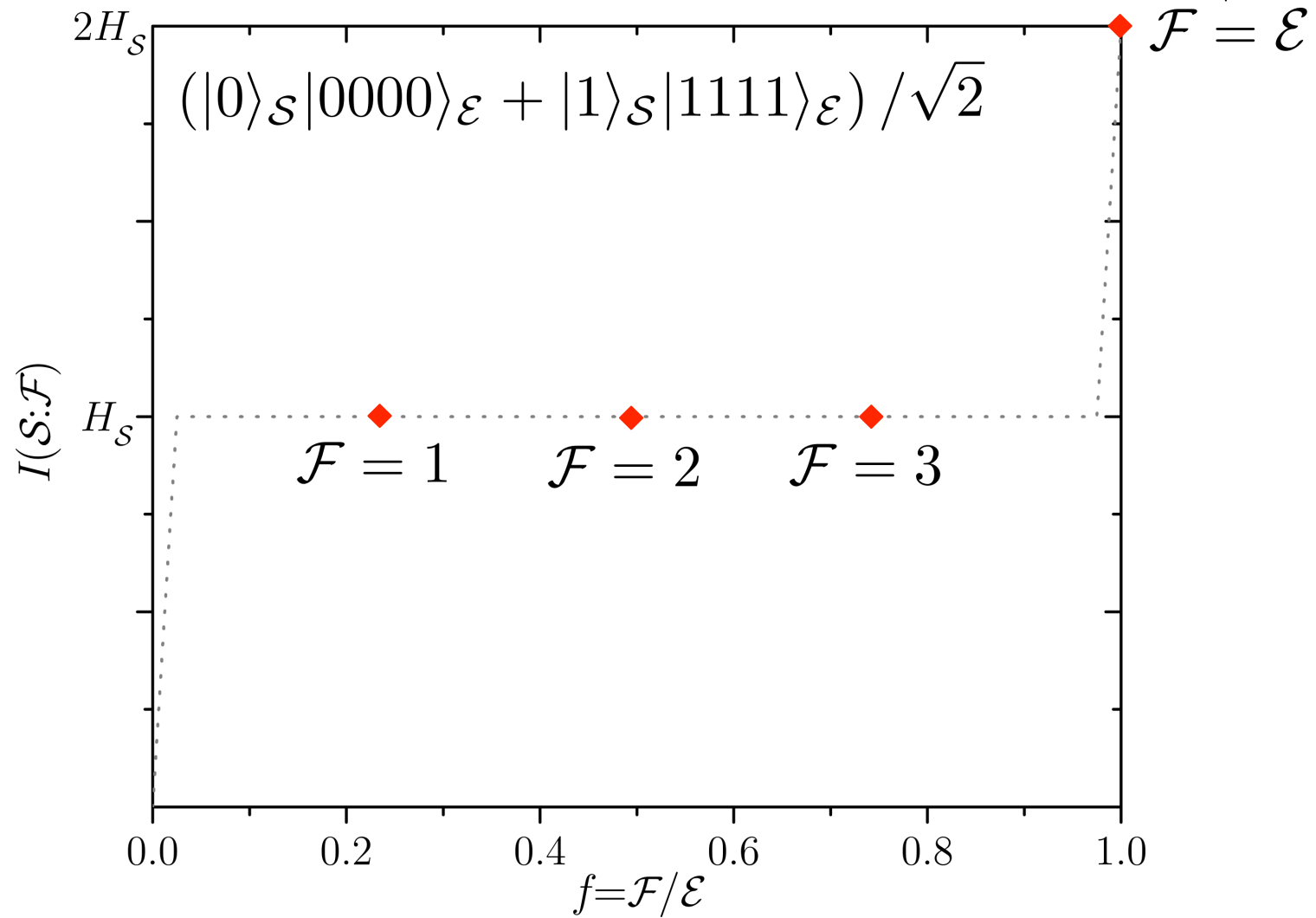
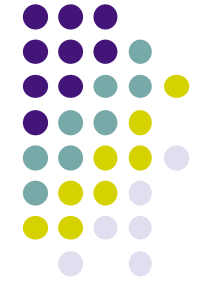


# Information plot

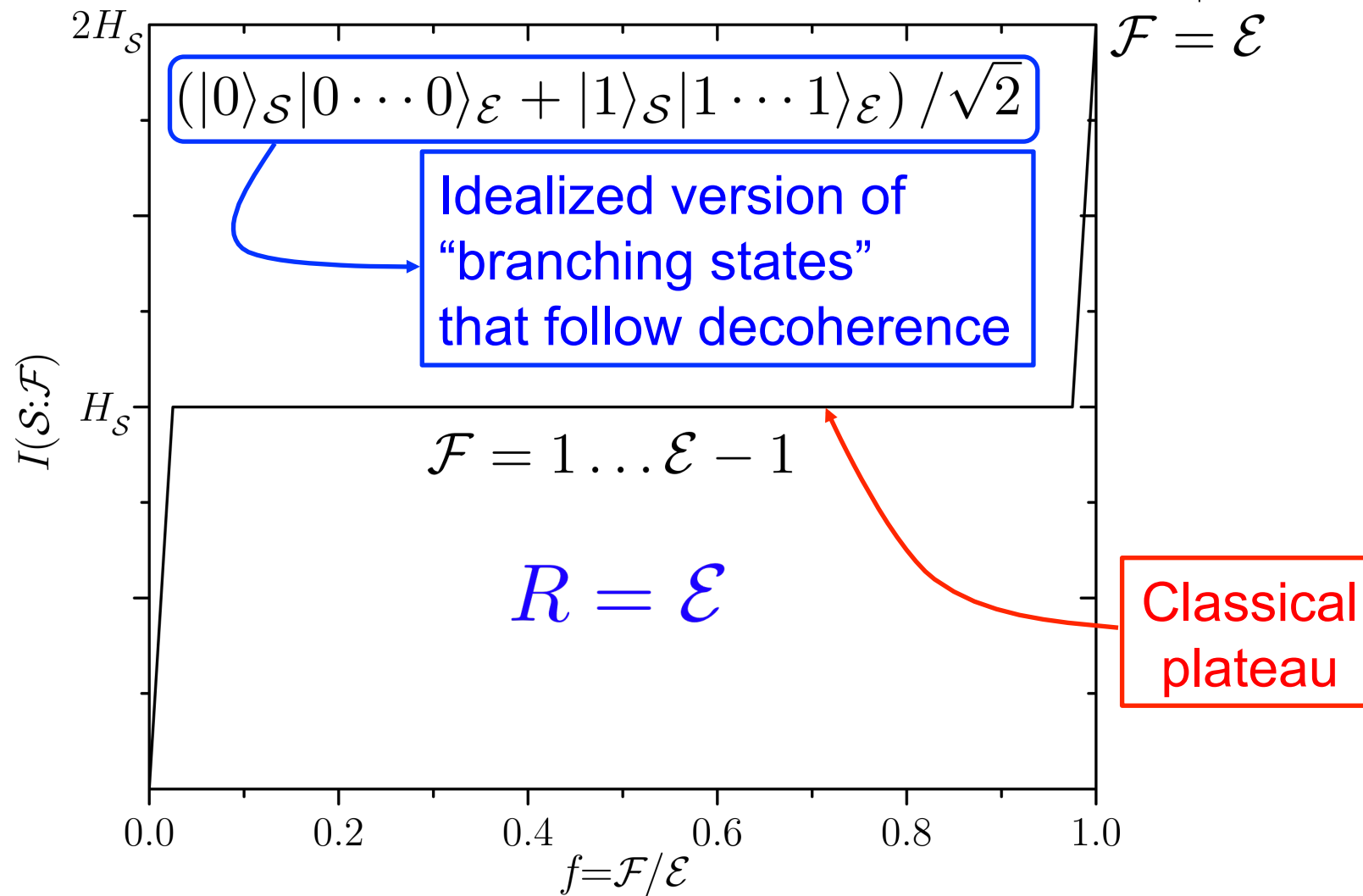




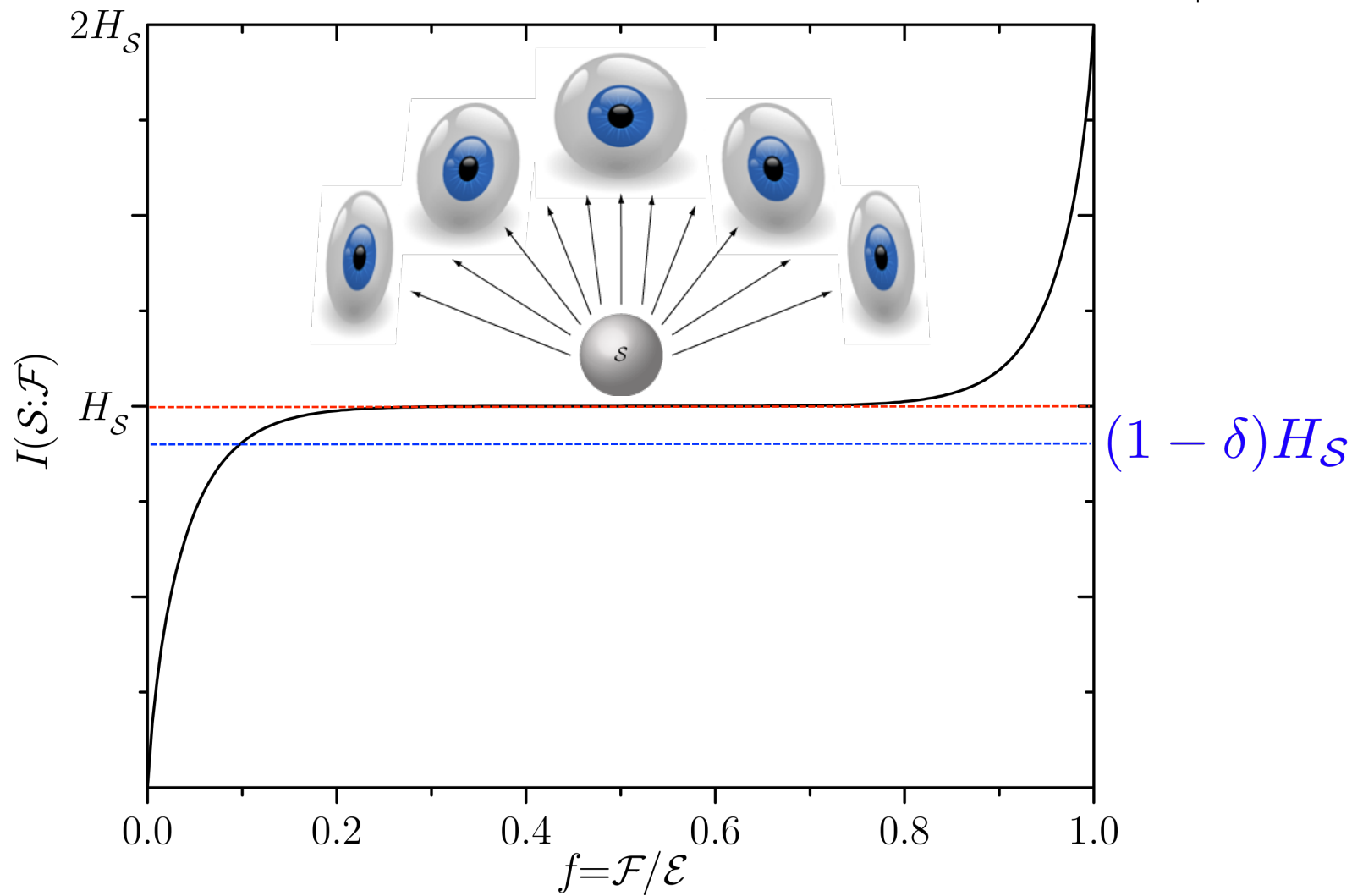
# Information plot



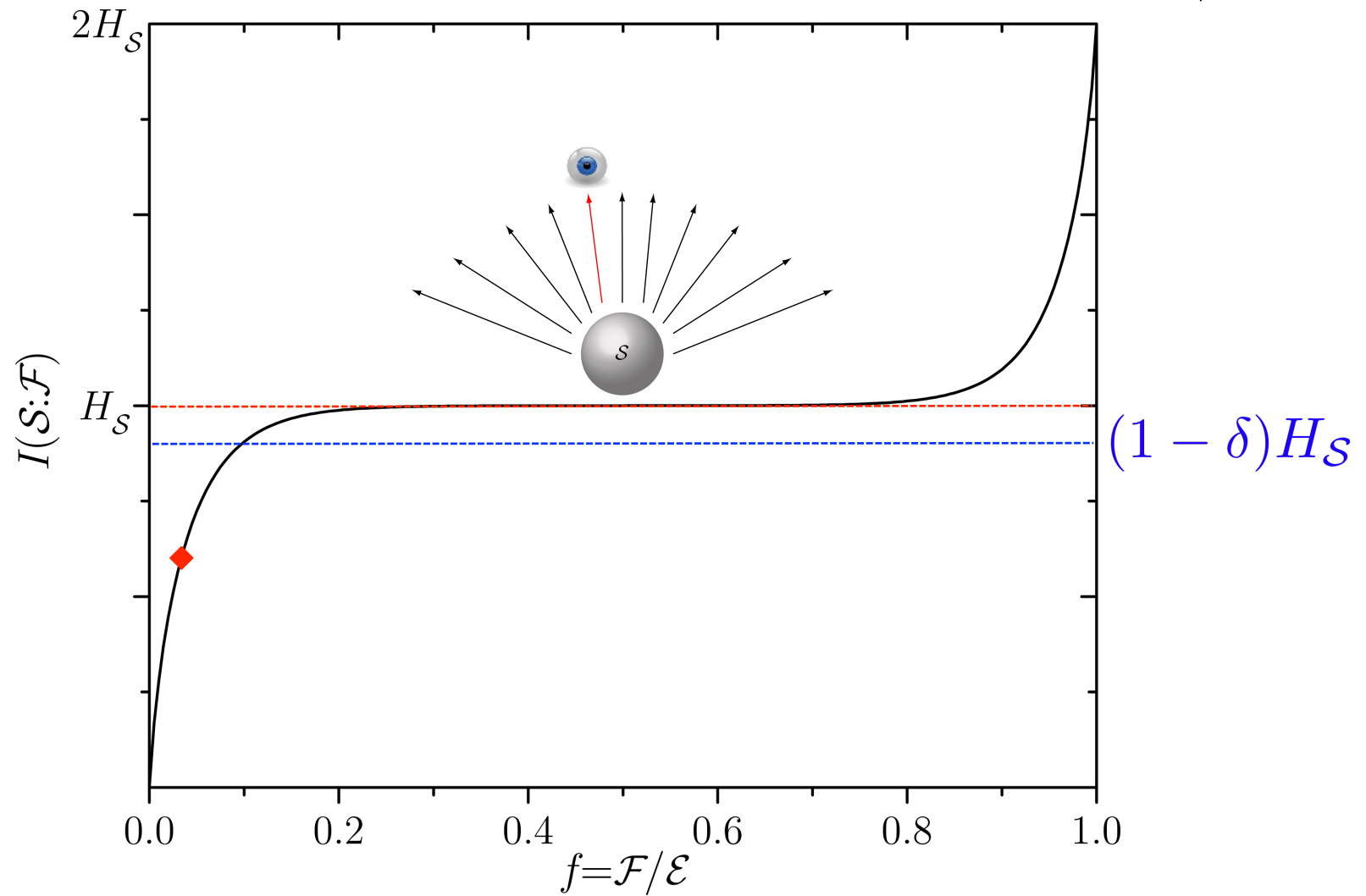
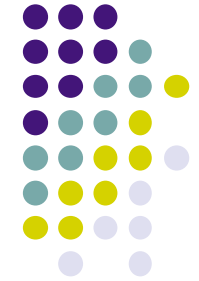
# Information plot



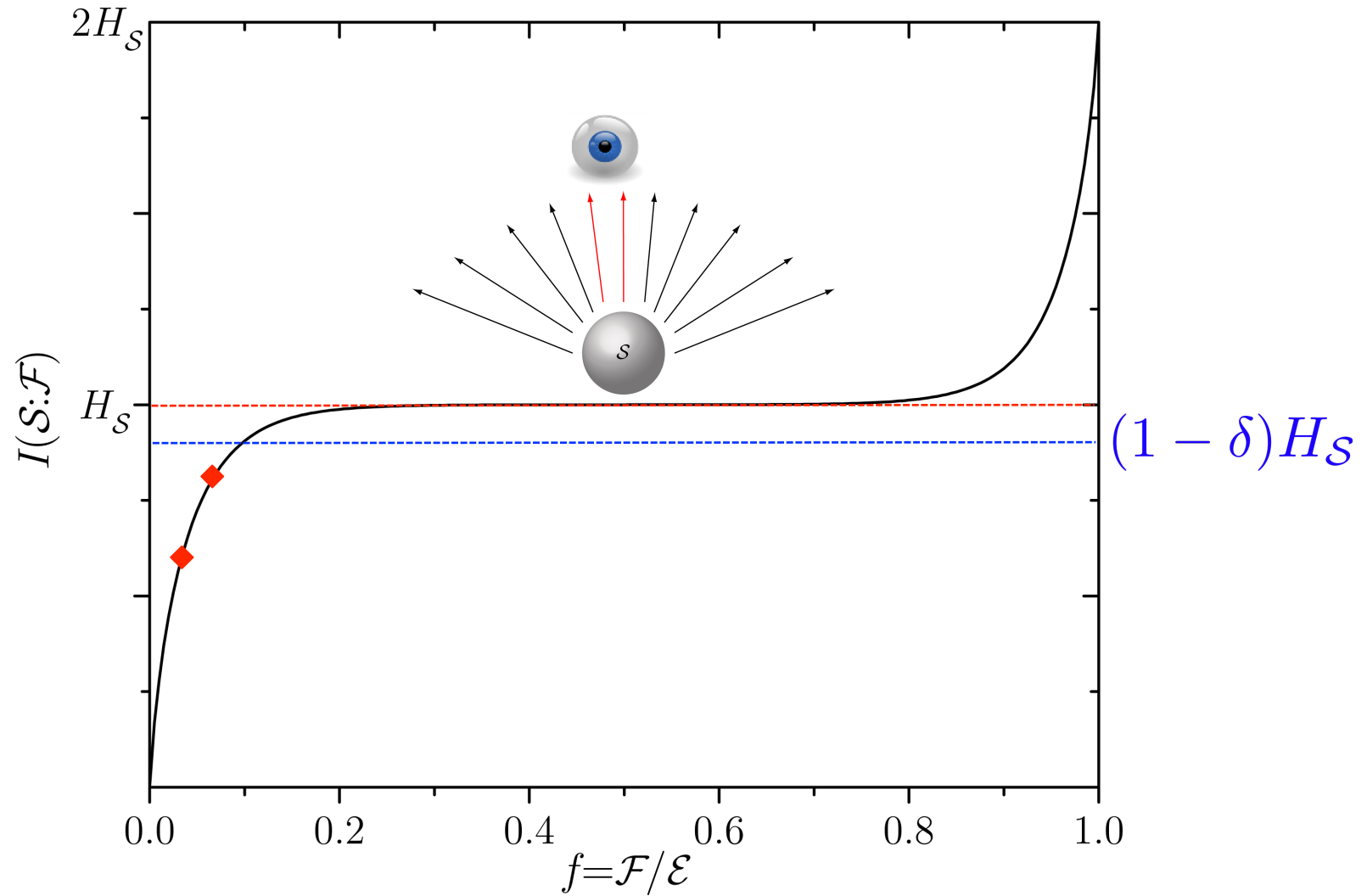
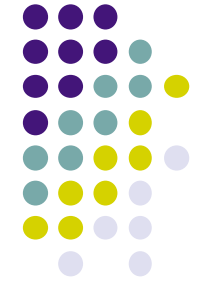
# Information plot



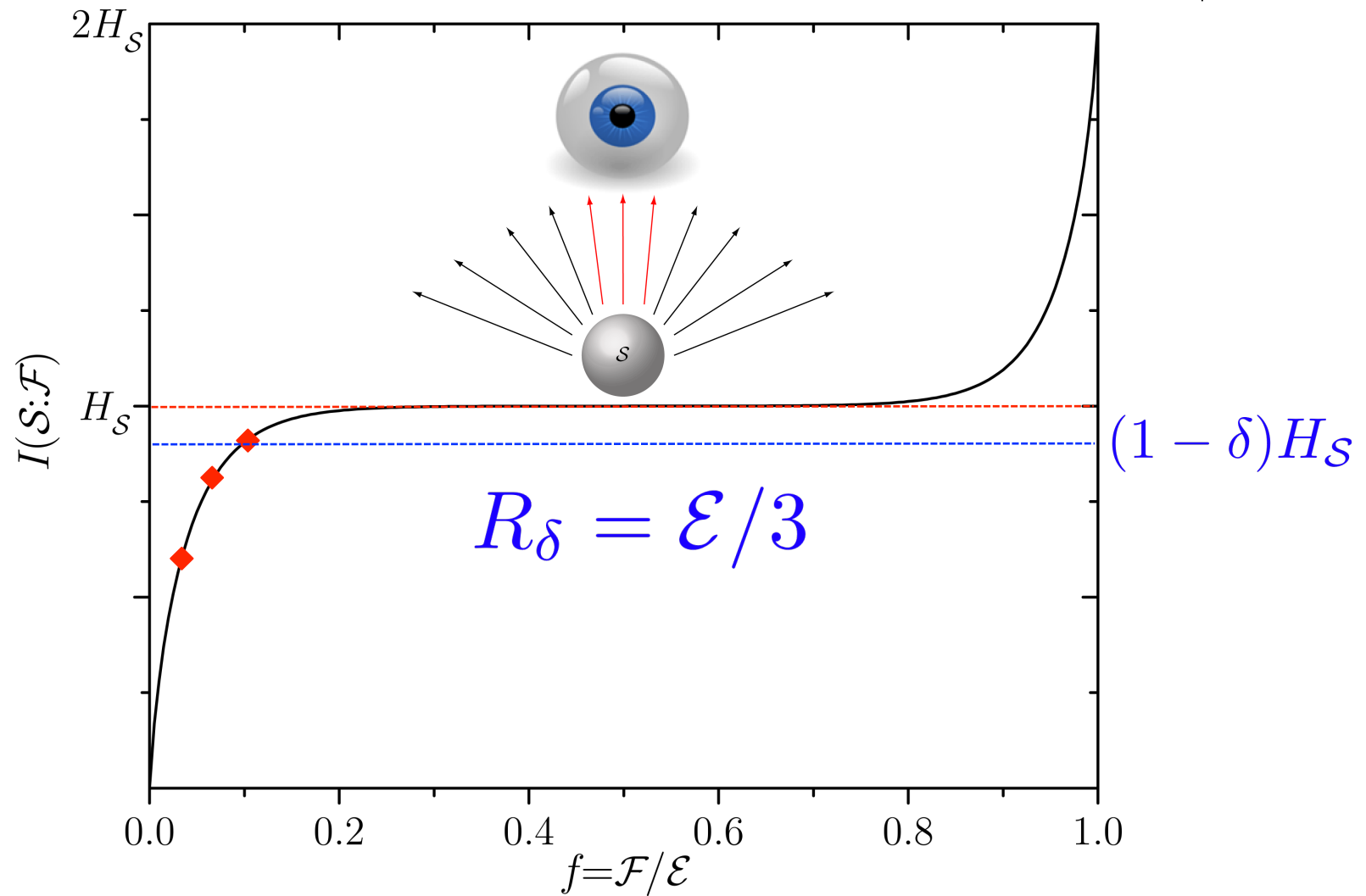
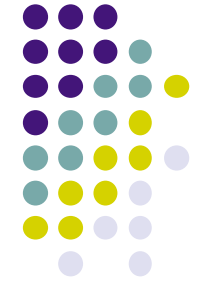
# Redundancy



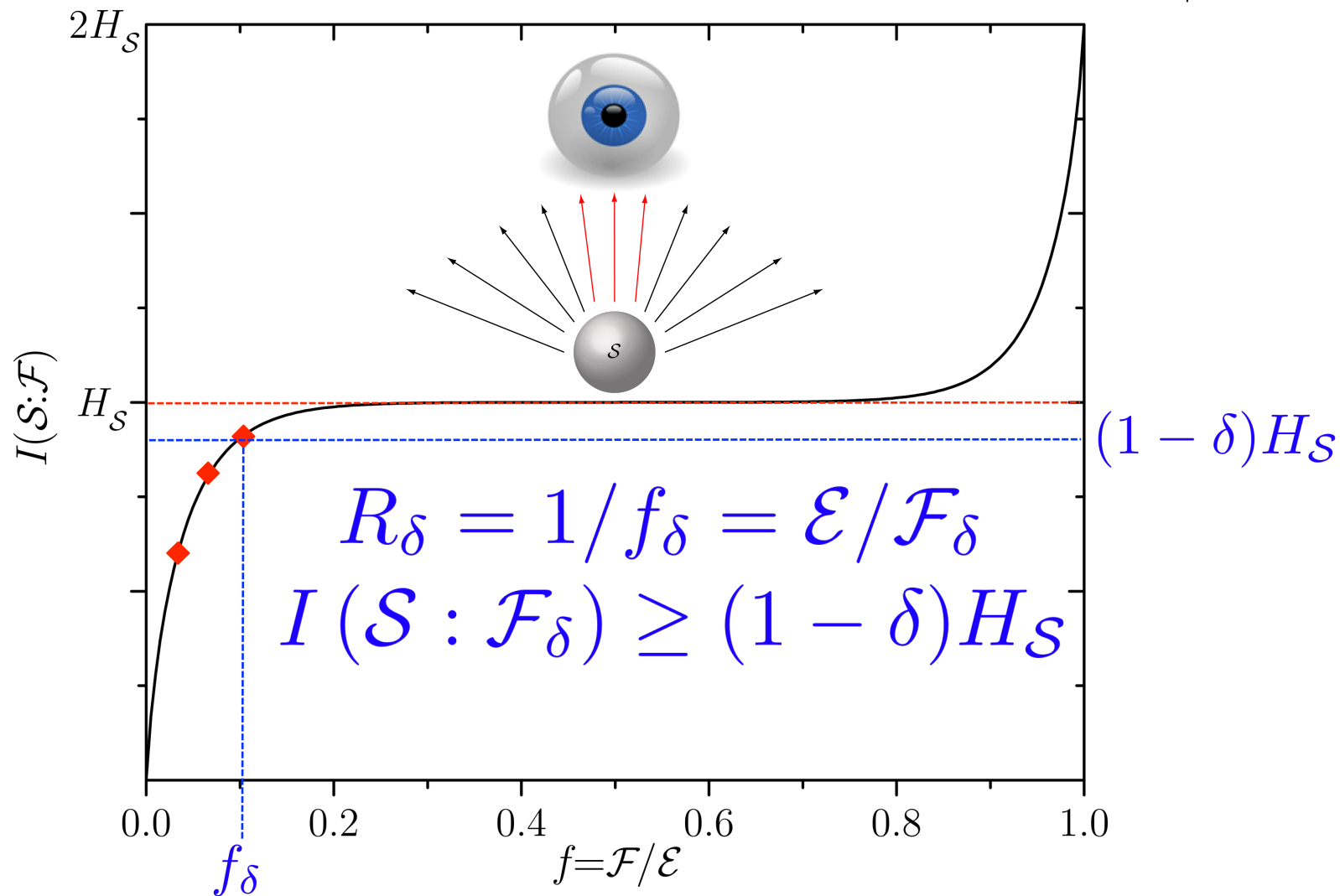
# Redundancy



# Redundancy



# Redundancy





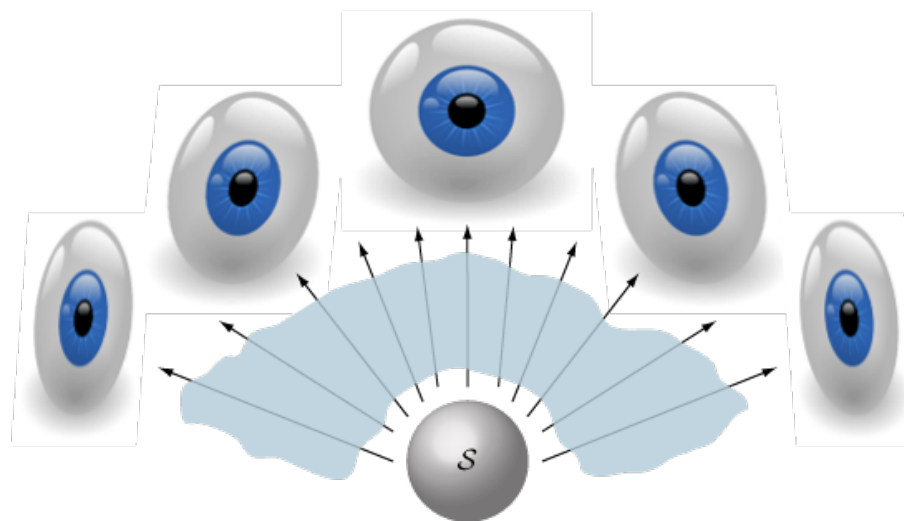
# Quantum Darwinism

## Environment as a communication channel

- Preferred states that survive and redundantly proliferate information into the environment
- **Observers can independently determine the state of the system and reach consensus, aka, classicality**

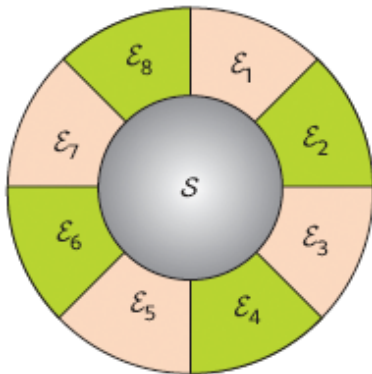
## Q-Darwinism in practice

1. Example system in non-ideal environments
2. State distinguishability and Decohering interactions





# 1. Example: symmetric spin $\mathcal{E}$



$$\mathbf{H}_{S\mathcal{E}} = \frac{1}{2} \sigma_S^z \sum_{k=1}^{\mathcal{E}} \sigma_k^z$$

$$\rho(0) = \rho_S(0) \otimes \rho_r^{\otimes \mathcal{E}}$$

Purely decohering  
Hamiltonian with  
independent  
environment  
components

Decoheres in pointer basis

$$\rho_S(t) = \begin{pmatrix} s_{00} & s_{01} \Lambda_{\mathcal{E}}(t) \\ s_{10} \Lambda_{\mathcal{E}}^*(t) & s_{11} \end{pmatrix} \longrightarrow \rho_S = \begin{pmatrix} s_{00} & 0 \\ 0 & s_{11} \end{pmatrix}$$

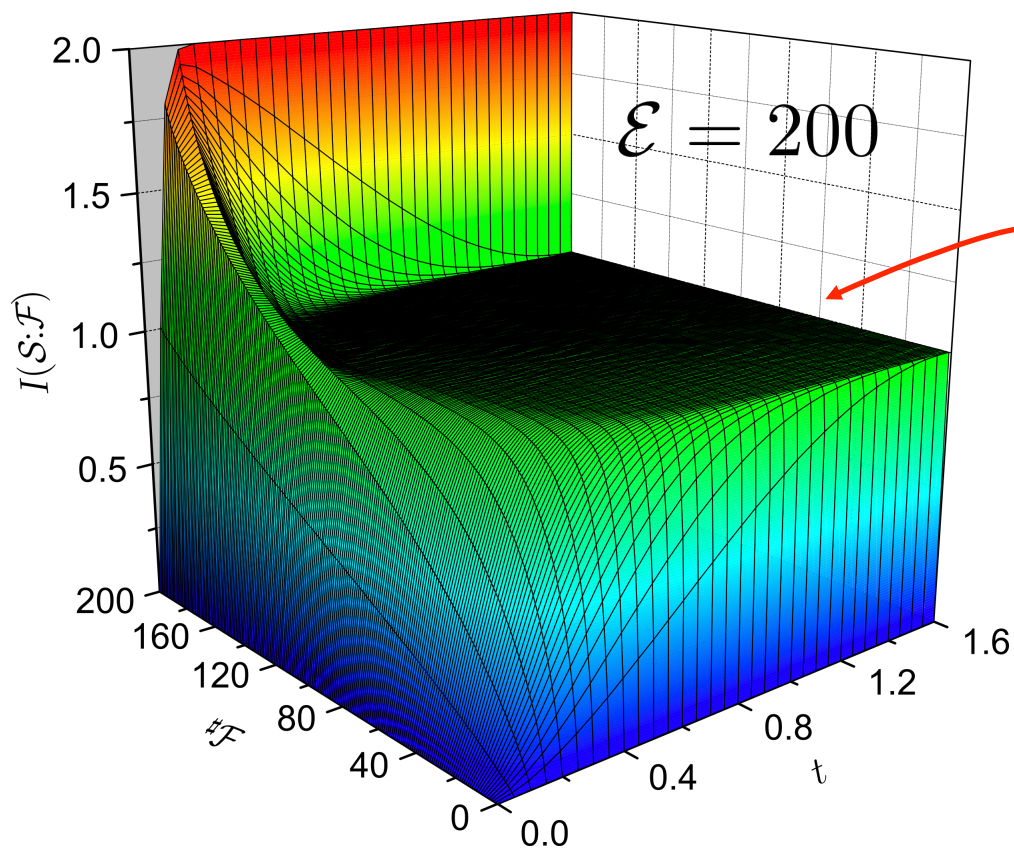
$$\Lambda_{\mathcal{E}}(t) = [\Lambda_k(t)]^{\mathcal{E}} = [\cos(t) - i\sigma \sin(t)]^{\mathcal{E}}$$

$$\sigma = r_{00} - r_{11} \quad \rho_r = \begin{pmatrix} r_{00} & r_{01} \\ r_{10} & r_{11} \end{pmatrix}$$

# Mutual Information



Information flow into the environment



$$|\psi(0)\rangle = |+\rangle \otimes |+\rangle^{\otimes \mathcal{E}}$$

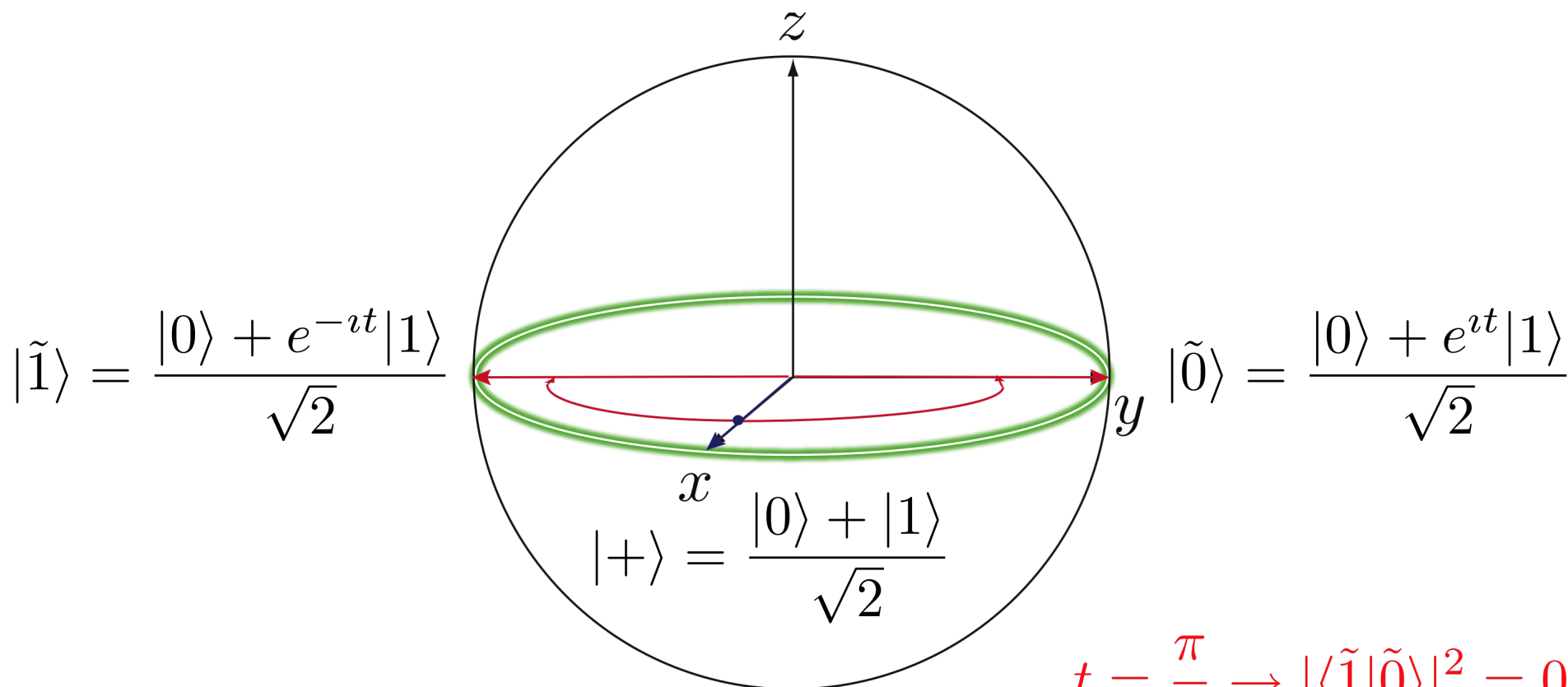
Classical plateau

$$t \rightarrow \frac{\pi}{2}, R_\delta \rightarrow \mathcal{E}$$

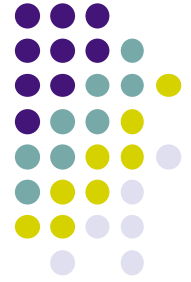
# Single environment qubit



$$|\psi(0)\rangle = |+\rangle|+\rangle^{\otimes \mathcal{E}} \rightarrow |\psi(t)\rangle = \frac{1}{\sqrt{2}}|0\rangle|\tilde{0}(t)\rangle^{\otimes \mathcal{E}} + \frac{1}{\sqrt{2}}|1\rangle|\tilde{1}(t)\rangle^{\otimes \mathcal{E}}$$



$$t = \frac{\pi}{2} \rightarrow |\langle \tilde{1} | \tilde{0} \rangle|^2 = 0$$



## Non-ideal $\mathcal{E}$

- Another extreme:

$$\rho(0) = \rho_S(0) \otimes \bar{I}^{\otimes \mathcal{E}}$$

$$\longrightarrow I(S : \mathcal{F}) = 0$$

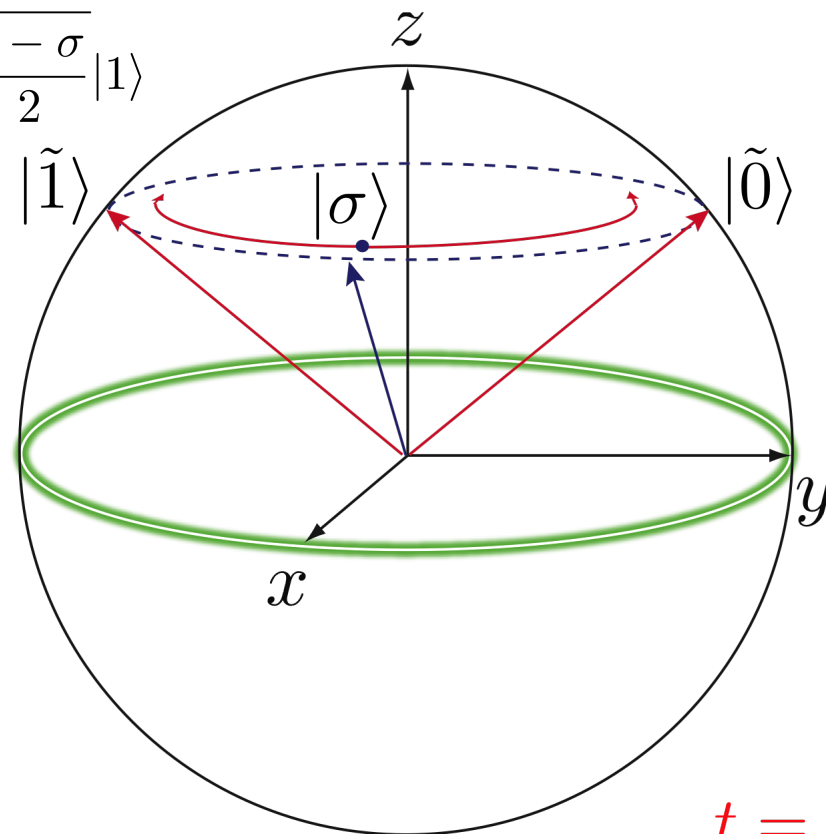
- How does this picture change when starting with different environment states?

# Single environment qubit



$$|\psi(0)\rangle = |+\rangle |\sigma\rangle^{\otimes \mathcal{E}} \rightarrow |\psi(t)\rangle = \frac{1}{\sqrt{2}} |0\rangle |\tilde{0}(t)\rangle^{\otimes \mathcal{E}} + \frac{1}{\sqrt{2}} |1\rangle |\tilde{1}(t)\rangle^{\otimes \mathcal{E}}$$

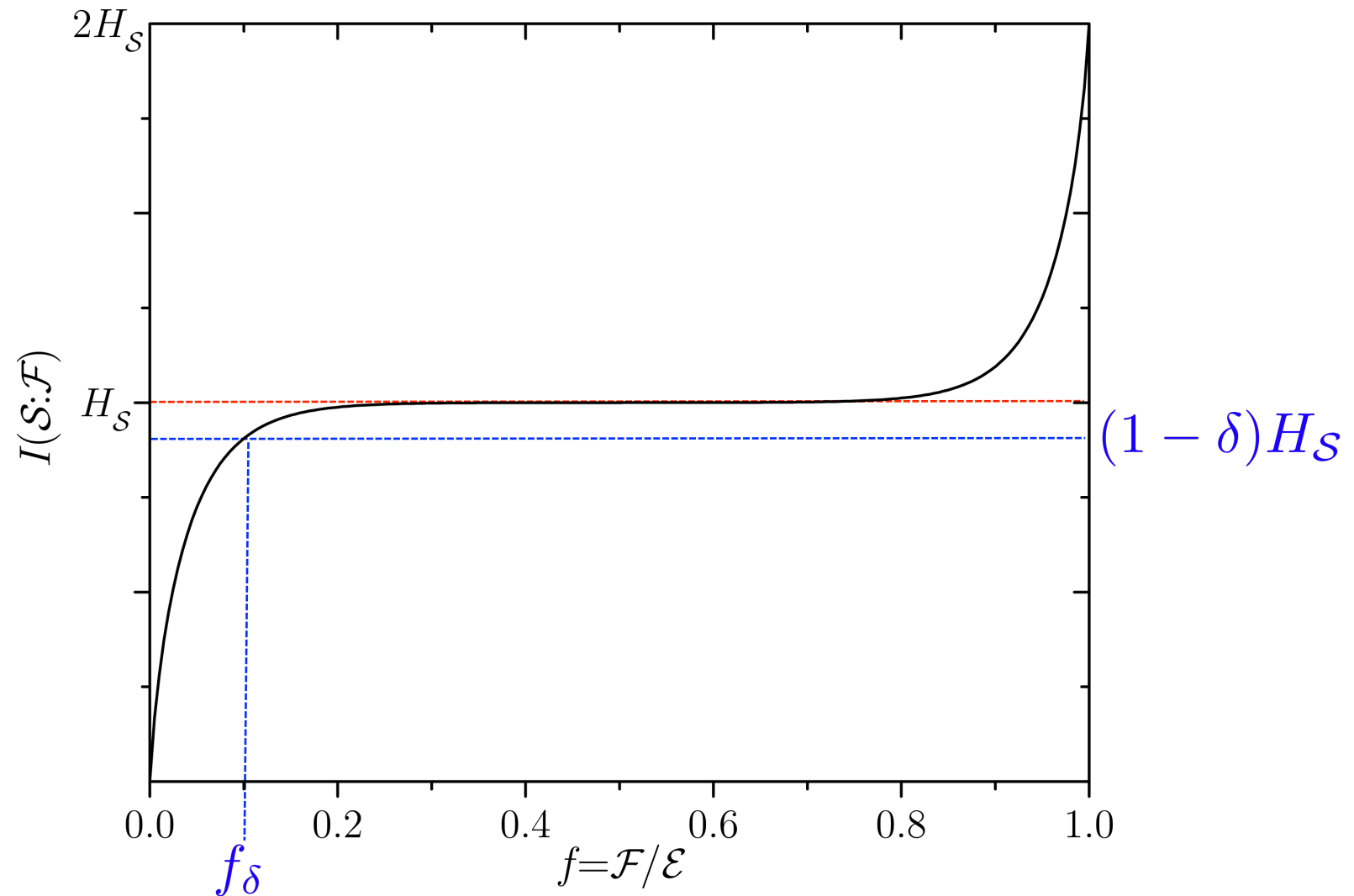
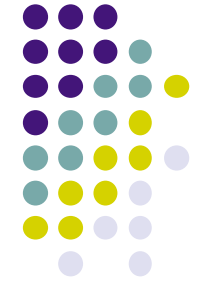
$$|\sigma\rangle = \sqrt{\frac{1+\sigma}{2}} |0\rangle + \sqrt{\frac{1-\sigma}{2}} |1\rangle$$



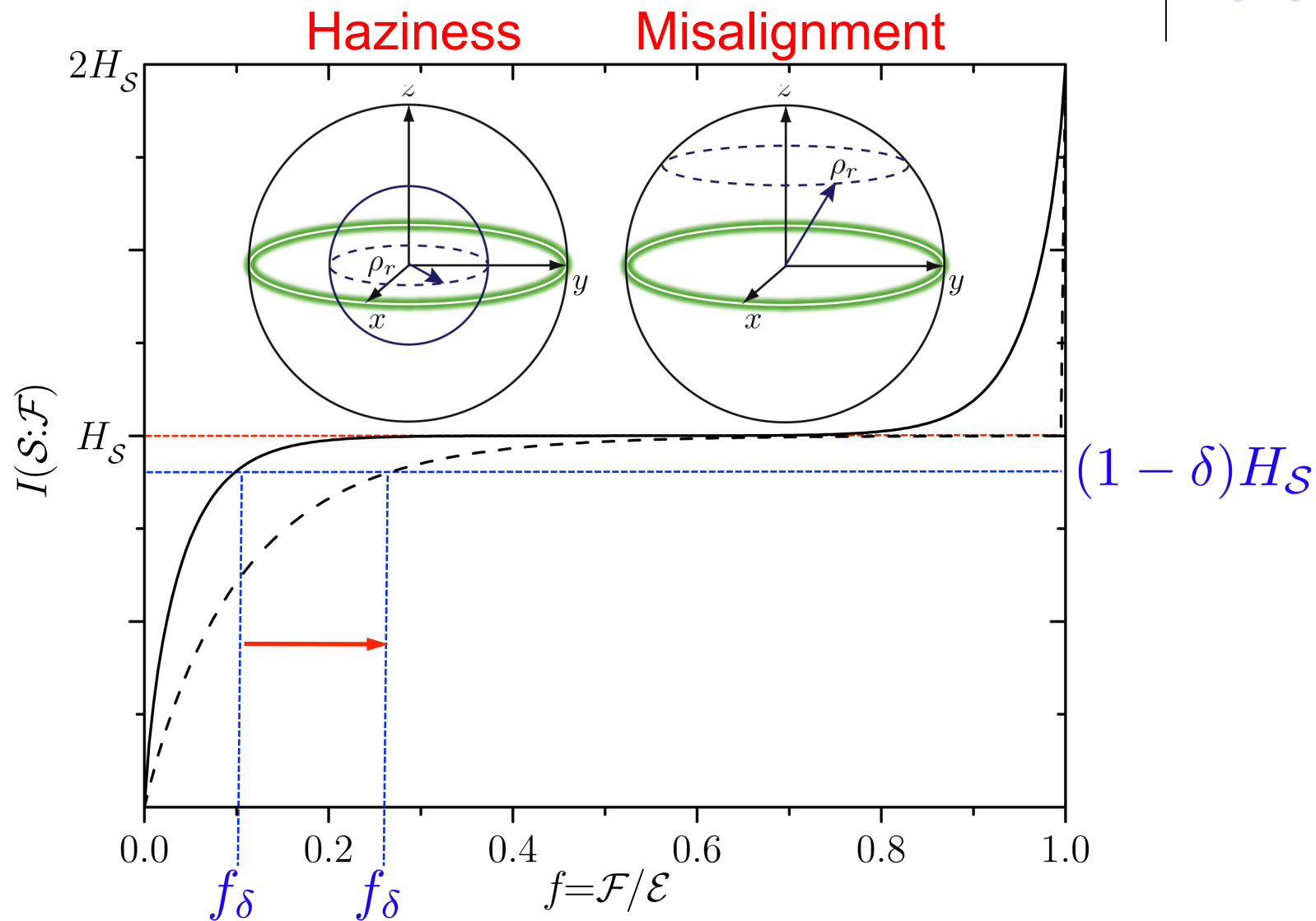
“misalignment”

$$t = \frac{\pi}{2} \rightarrow |\langle \tilde{1} | \tilde{0} \rangle|^2 = \sigma^2$$

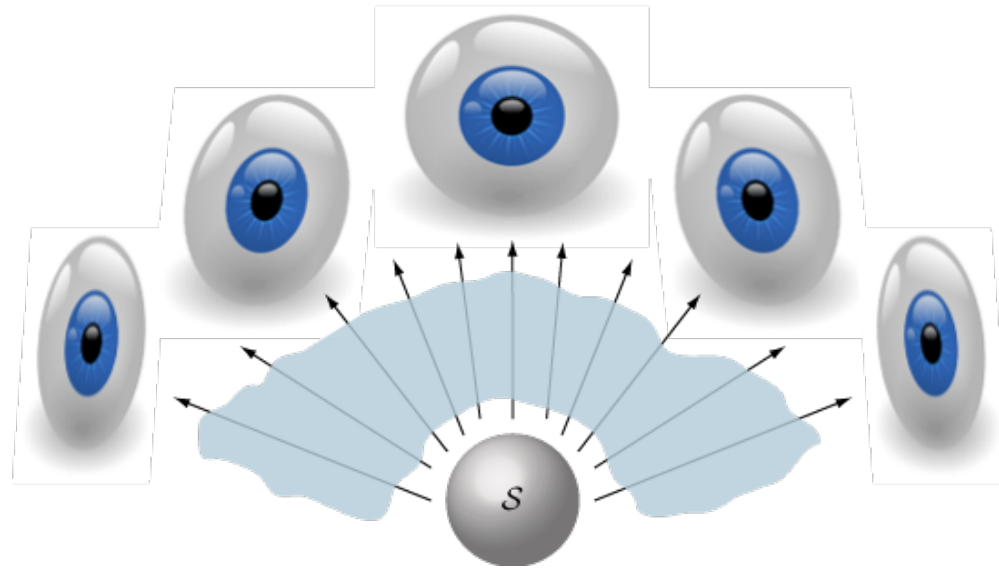
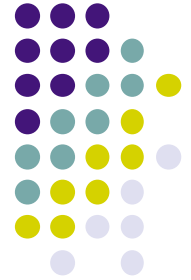
# Redundancy



# Redundancy



# Hazy $\mathcal{E}$



Haziness  $h \equiv H(\rho_r)$

The initial entropy of an environment component

Extreme:  $\rho(0) = \rho_S(0) \otimes \bar{I}^{\otimes \mathcal{E}} \longrightarrow I(S : \mathcal{F}) = 0$





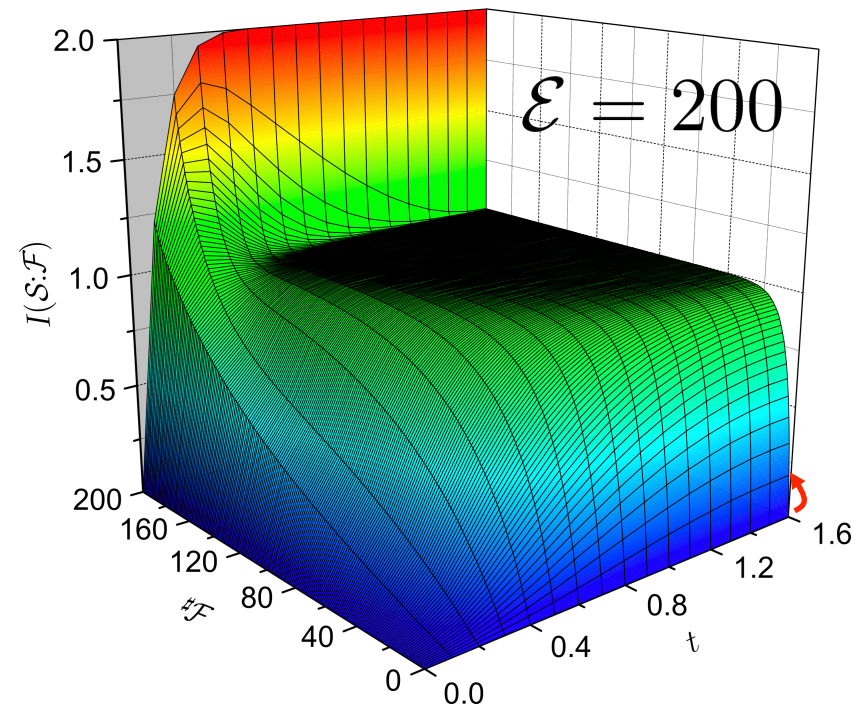
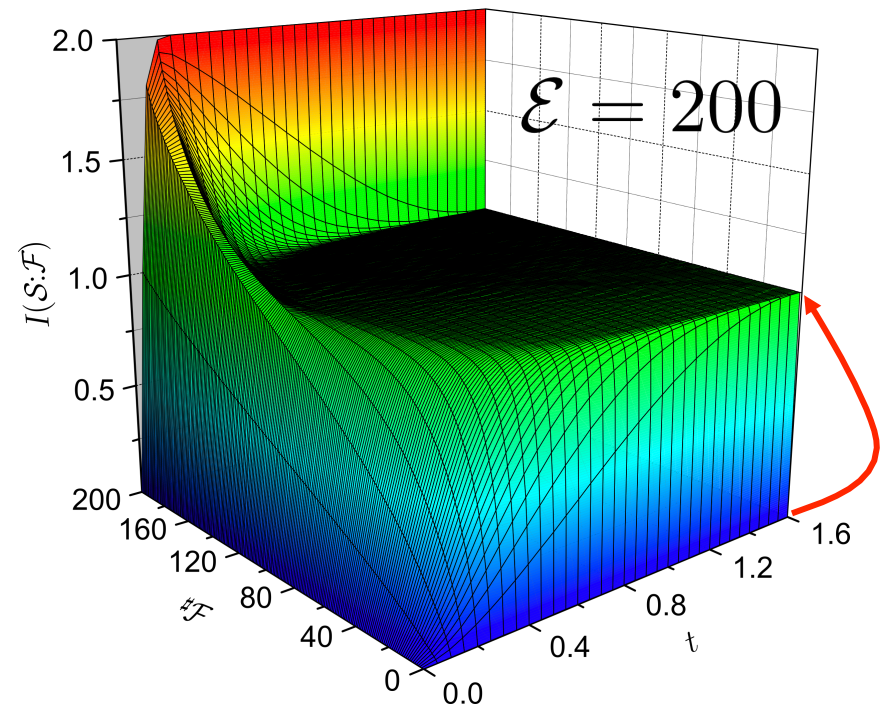
# Hazy $\mathcal{E}$

$$\rho(0) = \rho_S(0) \otimes \rho_r^{\otimes \mathcal{E}}$$

Pure environment,  $h=0$

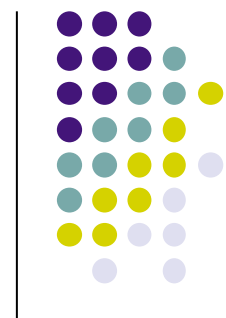
$$h \equiv H(\rho_r)$$

Hazy environment,  $h=0.8$



Classical plateau still forms even for quite hazy environments  
Initial rate of information gain is  $1-h$

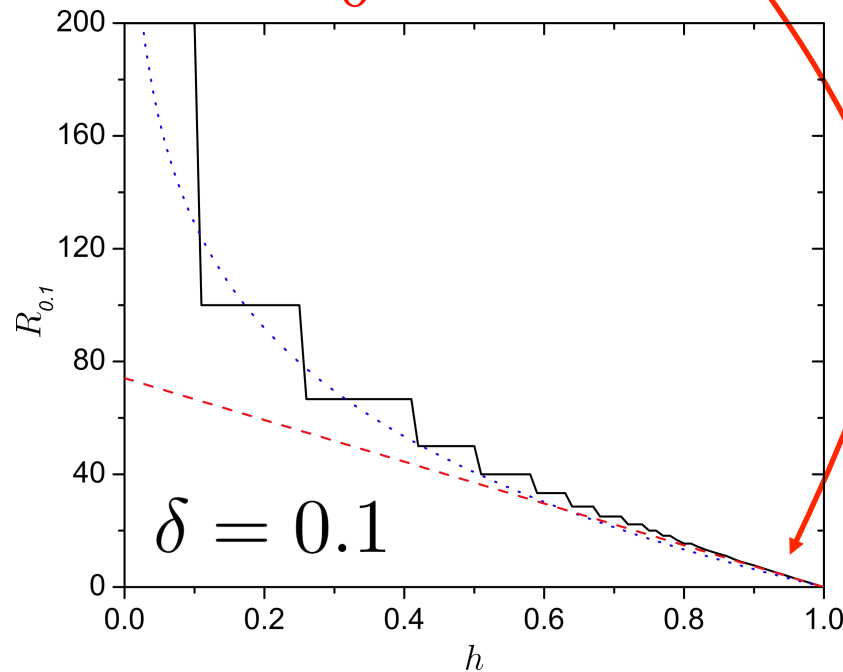
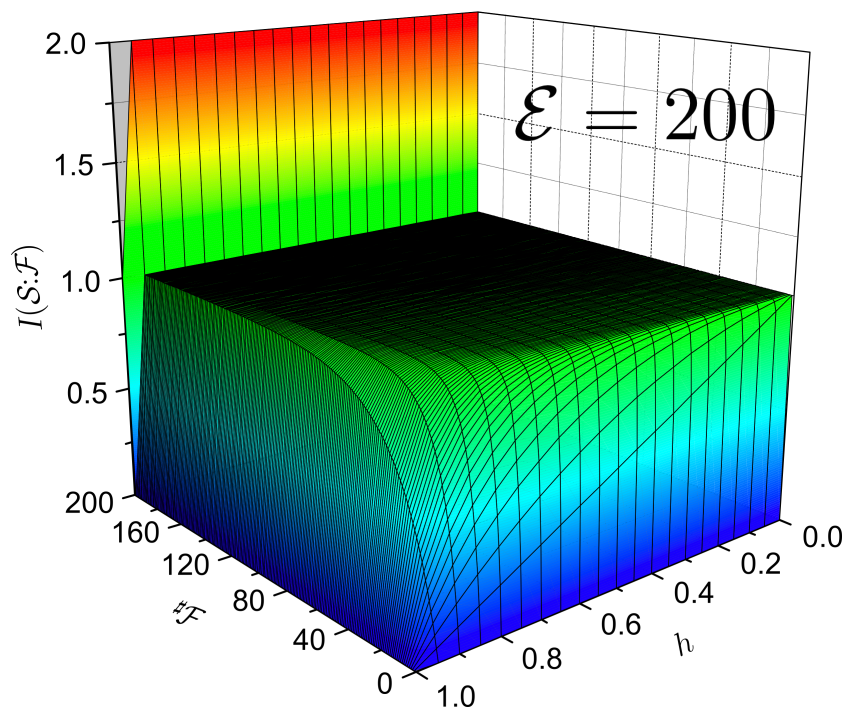
# Hazy $\mathcal{E}$



$$t = \frac{\pi}{2}$$

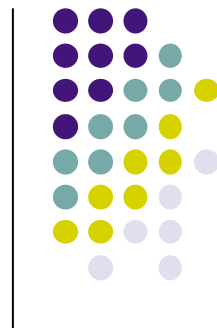
Noisy channel

$$R_\delta \sim 1 - h$$



$$R_\delta = \frac{\mathcal{E}}{\mathcal{F}_\delta} \quad \mathcal{F}_\delta \rightarrow I(\mathcal{S} : \mathcal{F}) \geq (1 - \delta)H_{\mathcal{S}}$$

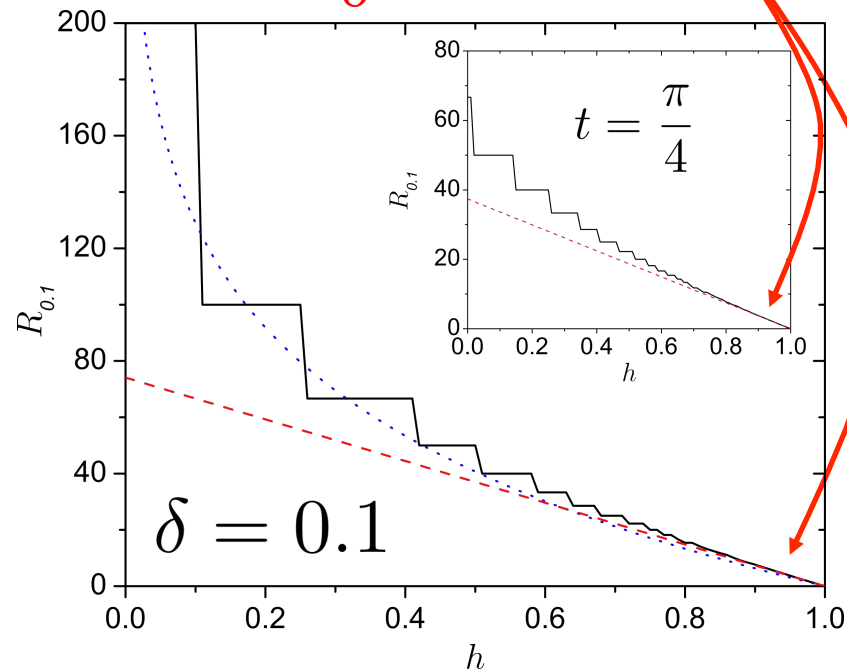
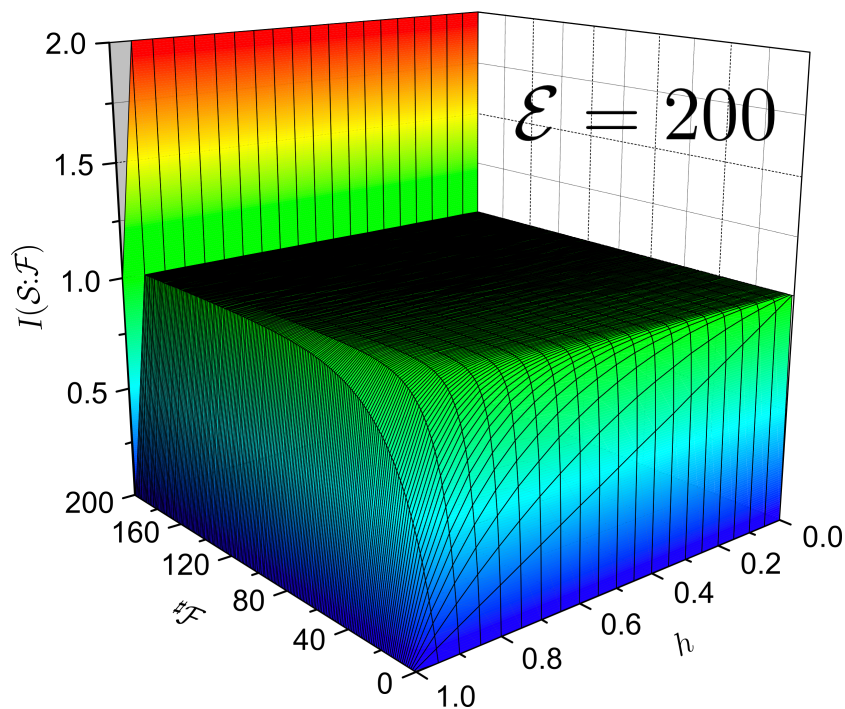
# Hazy $\mathcal{E}$



$$t = \frac{\pi}{2}$$

Noisy channel

$$R_\delta \sim 1 - h$$



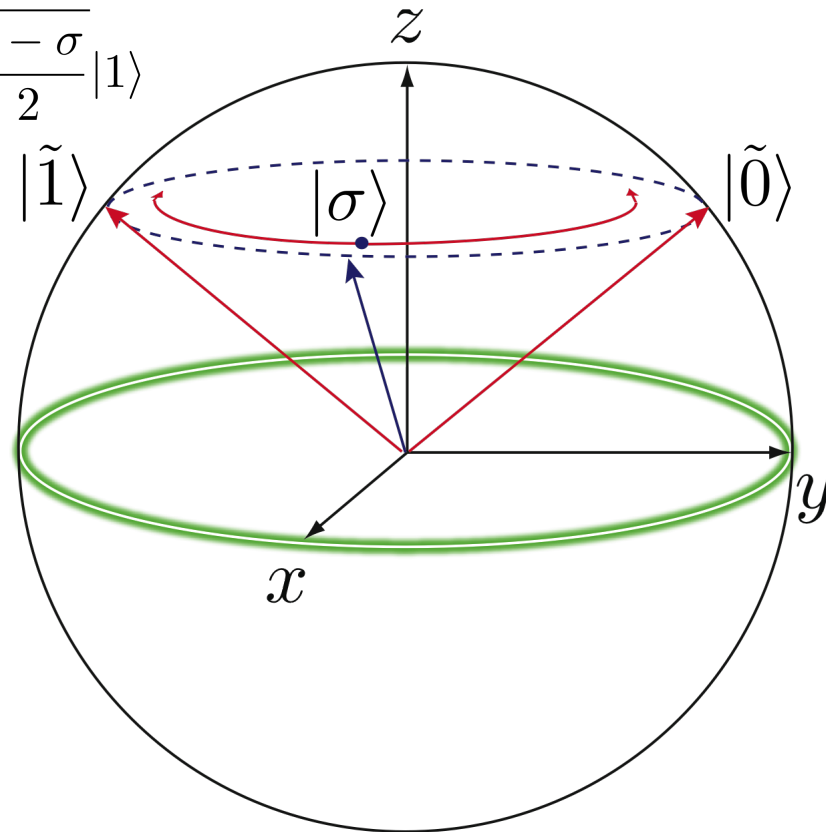
$$R_\delta = \frac{\mathcal{E}}{\mathcal{F}_\delta} \quad \mathcal{F}_\delta \rightarrow I(\mathcal{S} : \mathcal{F}) \geq (1 - \delta)H_{\mathcal{S}}$$

# Misaligned, pure $\mathcal{E}$



$$|\sigma\rangle = \sqrt{\frac{1+\sigma}{2}}|0\rangle + \sqrt{\frac{1-\sigma}{2}}|1\rangle$$

$$|\psi(0)\rangle = |+\rangle|\sigma\rangle^{\mathcal{E}}$$

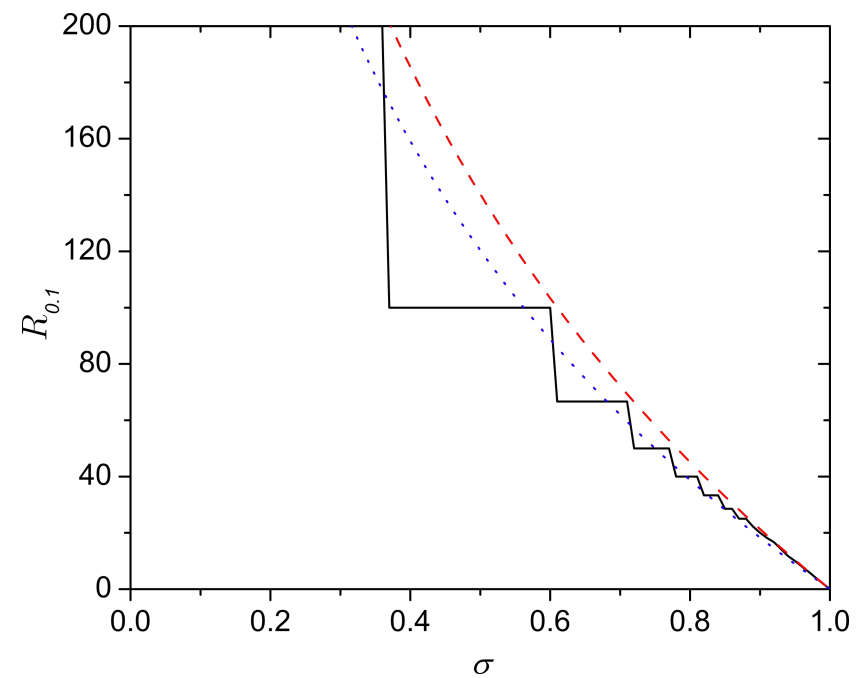
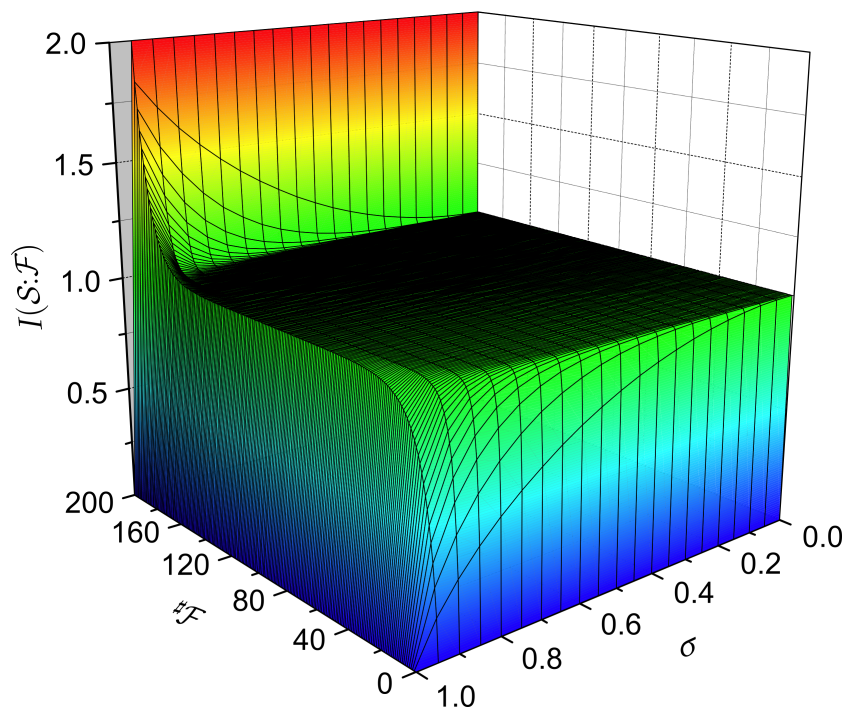
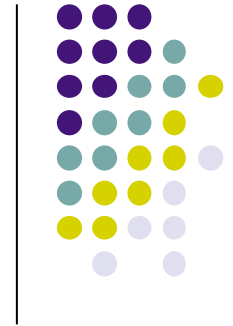


Misalignment: rotation out of x-y plane

Extreme:  $|\sigma| = 1 \longrightarrow I(\mathcal{S} : \mathcal{F}) = 0$

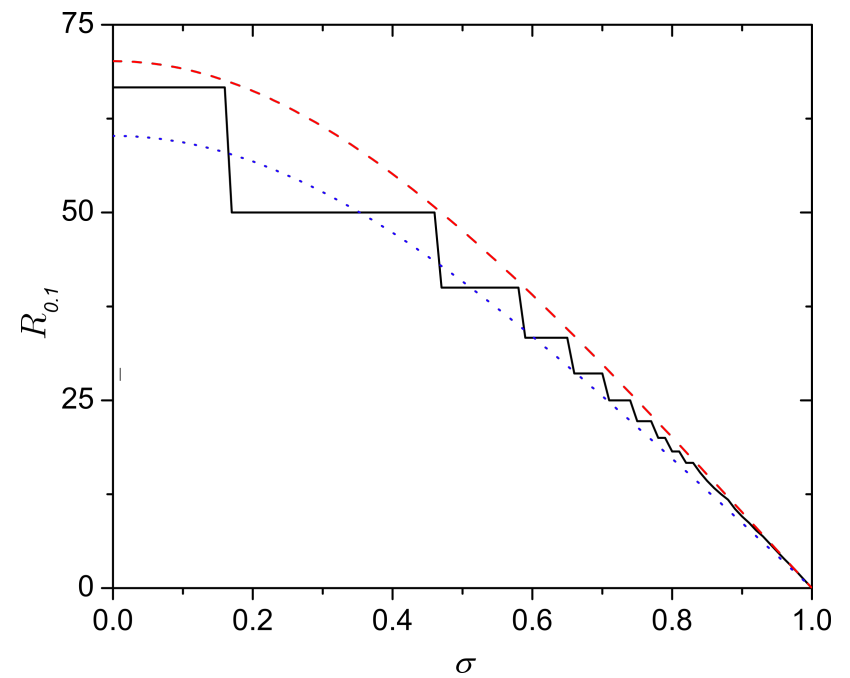
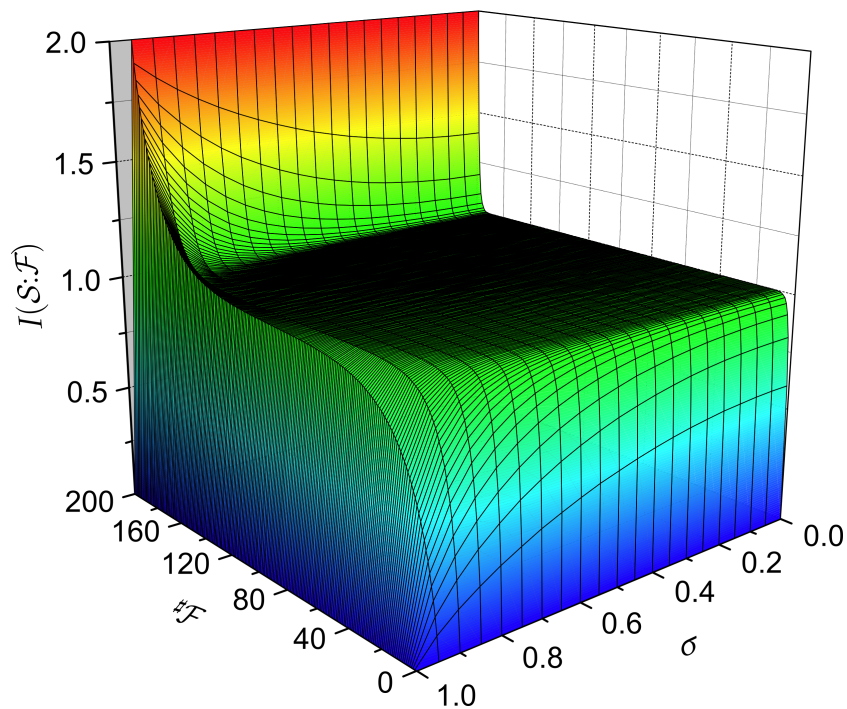
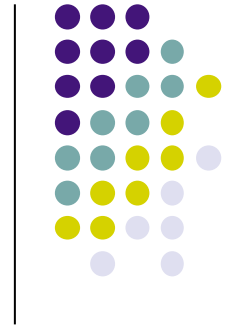
# Misaligned, pure $\mathcal{E}$

$$t = \frac{\pi}{2}$$

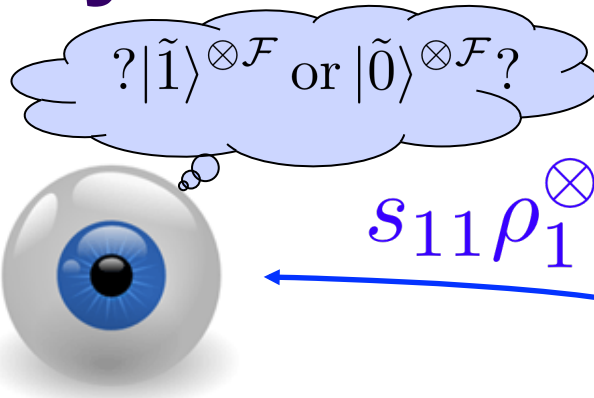


# Misaligned, pure $\mathcal{E}$

$$t = \frac{\pi}{4}$$



# 2. Inferring the system's state

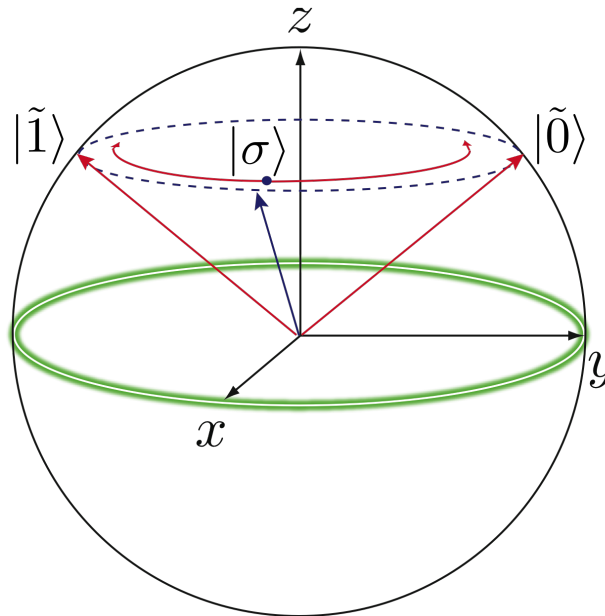


$$s_{00} \rho_0^{\otimes \mathcal{F}}$$

$$s_{11} \rho_1^{\otimes \mathcal{F}}$$

$$\frac{1}{\sqrt{2}} |1\rangle_S |\tilde{1}\rangle^{\otimes \mathcal{F}} |\tilde{1}\rangle^{\otimes \mathcal{E}-\mathcal{F}}$$

$$\frac{1}{\sqrt{2}} |0\rangle_S |\tilde{0}\rangle^{\otimes \mathcal{F}} |\tilde{0}\rangle^{\otimes \mathcal{E}-\mathcal{F}}$$



$$\frac{|1\rangle_S + |0\rangle_S}{\sqrt{2}} |\sigma\rangle^{\otimes \mathcal{F}} |\sigma\rangle^{\otimes \mathcal{E}-\mathcal{F}}$$

# Proliferation of information under decohering interactions



$$I(\mathcal{S} : \mathcal{F}) = H_{\mathcal{S}}(t) + H_{\mathcal{F}}(t) - H_{\mathcal{S}\mathcal{F}}(t)$$
$$\mathbf{H}_{\mathcal{S}\mathcal{E}} = \sum_{k=1}^{\mathcal{E}} \Pi_{\mathcal{S}} \Upsilon_k \quad \rho(0) = \rho_{\mathcal{S}}(0) \otimes \left[ \bigotimes_{k=1}^{\mathcal{E}} \rho_k(0) \right]$$

A purely decohering Hamiltonian with independent environment components



# Proliferation of information under decohering interactions



$$I(\mathcal{S} : \mathcal{F}) = H_{\mathcal{S}}(t) + H_{\mathcal{F}}(t) - H_{\mathcal{S}\mathcal{F}}(t)$$

$$I(\mathcal{S} : \mathcal{F}) = \chi(\mathcal{S} \rightarrow \mathcal{F}) + \delta(\mathcal{S} : \mathcal{F})_{\{\Pi_{\mathcal{S}}\}}$$

Quantum Discord

$$\chi(\mathcal{S} \rightarrow \mathcal{F}) = H\left(\sum_i s_{ii} \rho_i^{\otimes \mathcal{F}}\right) - \sum_i s_{ii} H(\rho_i^{\otimes \mathcal{F}})$$

Holevo Bound: Maximum classical information transmittable by a quantum channel



# Quantum Discord

$$\delta(\mathcal{S} : \mathcal{F})_{\{\Pi_{\mathcal{S}}\}} = H_{\mathcal{S}d\mathcal{E}}(t) - H_{\mathcal{S}d(\mathcal{E}/\mathcal{F})}(t)$$

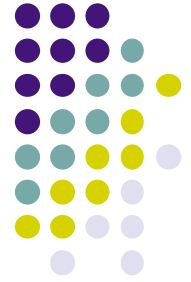
System decohered only by  $\mathcal{E}/\mathcal{F}$

“Good decoherence” – when  $\mathcal{E}$  and  $\mathcal{E}/\mathcal{F}$  suffice to decohere  $\mathcal{S}$

$$H_{\mathcal{S}d\mathcal{E}}(t) - H_{\mathcal{S}d(\mathcal{E}/\mathcal{F})}(t) \approx 0$$

Mutual information is only the classical information

$$I(\mathcal{S} : \mathcal{F}) \approx \chi(\mathcal{S} \rightarrow \mathcal{F})$$



# Classical Information

$$\chi(\mathcal{S} \rightarrow \mathcal{F}) = H \left( \sum_i s_{ii} \rho_i^{\otimes \mathcal{F}} \right) - \sum_i s_{ii} H(\rho_i^{\otimes \mathcal{F}})$$

$$\chi(\mathcal{S} \rightarrow \mathcal{F}) = H_{\mathcal{F}}(t) - H_{\mathcal{F}}(0)$$

Information transmitted is suppressed by initial entropy

$$H(\mathcal{S} \leftrightarrow \mathcal{F}) \geq \chi(\mathcal{S} \rightarrow \mathcal{F}) \geq I(\Phi_D[\mathcal{S}] : \Phi_M[\mathcal{F}])$$

Exchange entropy

Measure the fragment

**Information transmitted is tightly bounded**

Roga, Fannes, Życzkowski, PRL 105, 040505 (2010)

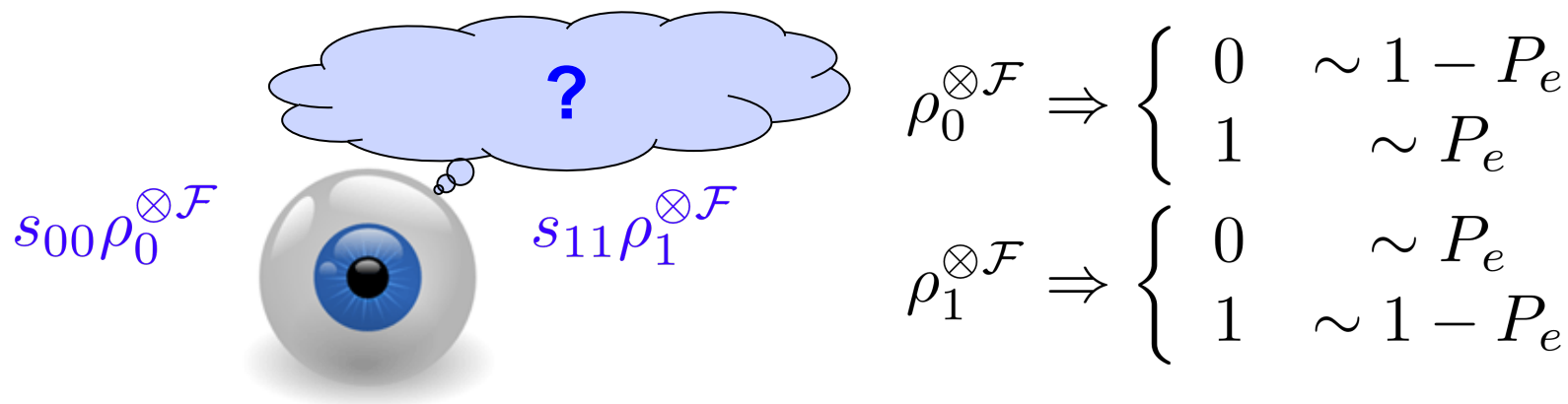


# Lower bound

$$\chi(\mathcal{S} \rightarrow \mathcal{F}) \geq I(\Phi_D[\mathcal{S}] : \Phi_M[\mathcal{F}])$$

$\Phi_D[\mathcal{S}]$  kills any remaining coherence in the system

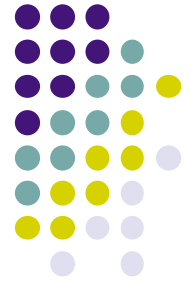
$\Phi_M[\mathcal{F}]$  distinguishes the states of the fragment:



$$P_e = \frac{1}{2} (1 - \text{tr}|s_{11}\rho_1^{\otimes \mathcal{F}} - s_{00}\rho_0^{\otimes \mathcal{F}}|) \quad \text{\`a la Helstrom}$$

$$\chi(\mathcal{S} \rightarrow \mathcal{F}) \geq I(\Phi_D[\mathcal{S}] : \Phi_M[\mathcal{F}]) \sim H_S - H(P_e)$$

# Classical Information, $D_S = 2$



$$H\left(\left(\frac{1 - \sqrt{(s_{00} - s_{11})^2 + 4P_e}}{2}\right)\right) \geq \chi(\mathcal{S} \rightarrow \mathcal{F}) \geq H_S - H(P_e)$$

$$\rightarrow H_S - \mathcal{O}(P_e^2) \geq \chi(\mathcal{S} \rightarrow \mathcal{F}) \geq H_S - \mathcal{O}(P_e \log P_e)$$

$$P_e = \frac{1}{2} (1 - \text{tr} |s_{11} \rho_1^{\otimes \mathcal{F}} - s_{00} \rho_0^{\otimes \mathcal{F}}|)$$

$$\sim e^{-\mathcal{F} \xi_{QCB}}$$

$$\xi_{QCB} = -\ln \min_{0 \leq s \leq 1} \text{tr} \rho_0^s \rho_1^{1-s}$$

QCB: Audenaert et al., PRL 98, 160501 (2007)

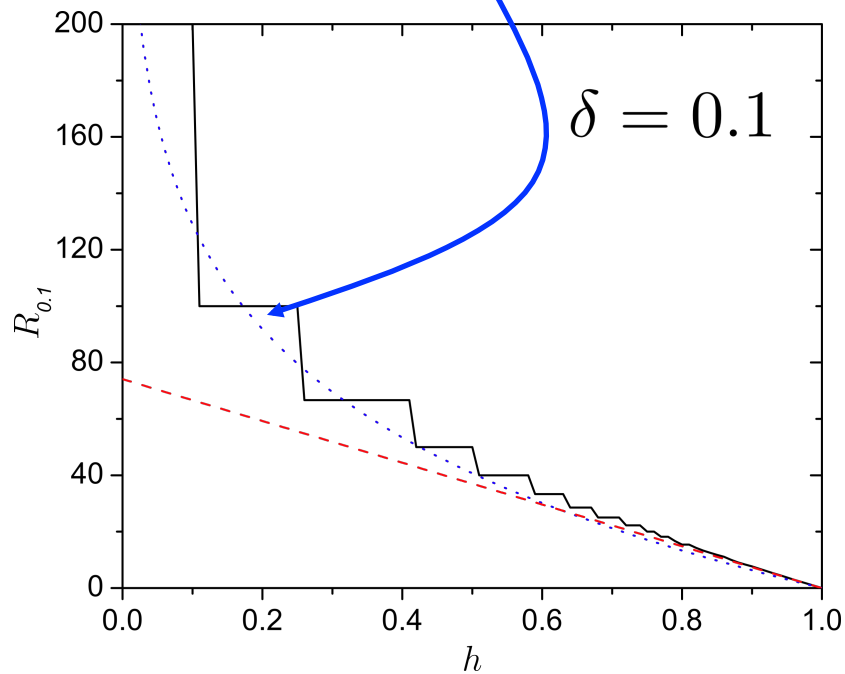
$$\Rightarrow R_\delta = \mathcal{E} \frac{\xi_{QCB}}{\ln(1/\delta)}$$

“Chernoff information”  
gives the Redundant  
Information

# Redundancy, hazy $\mathcal{E}$ , $t = \pi/2$

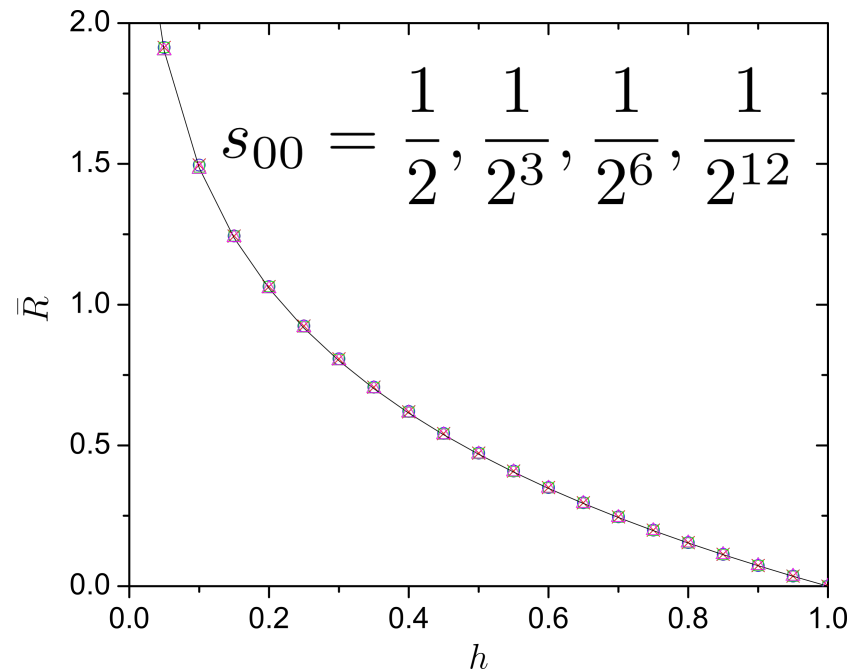


$$R_\delta = \mathcal{E} \frac{\ln 2 \sqrt{\lambda_- \lambda_+}}{\ln(\delta)}$$



## Normalized Redundancy

$$\bar{R} = \lim_{\delta \rightarrow 0} \frac{-R_\delta \ln \delta}{\mathcal{E}} = \xi_{QCB}$$

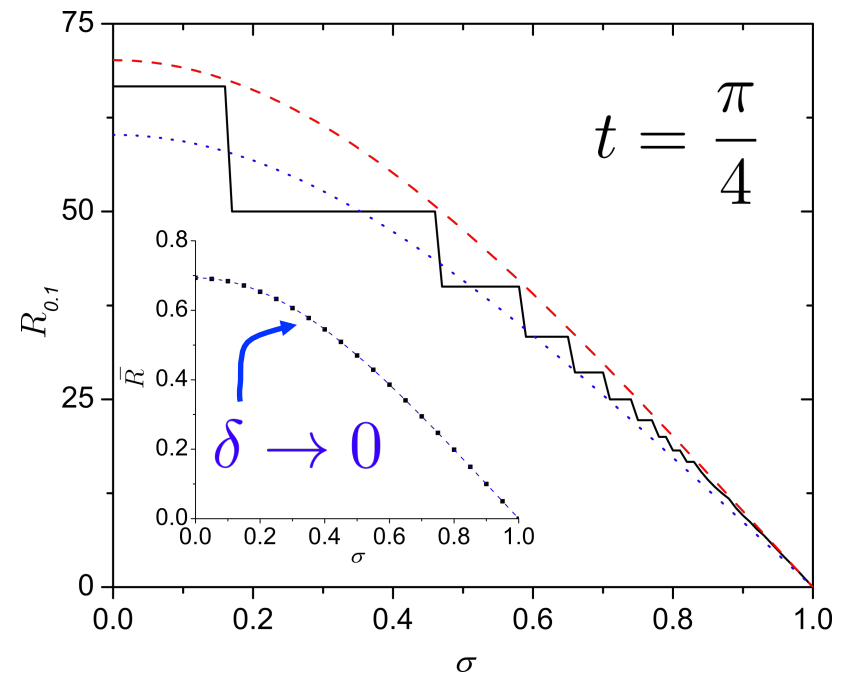
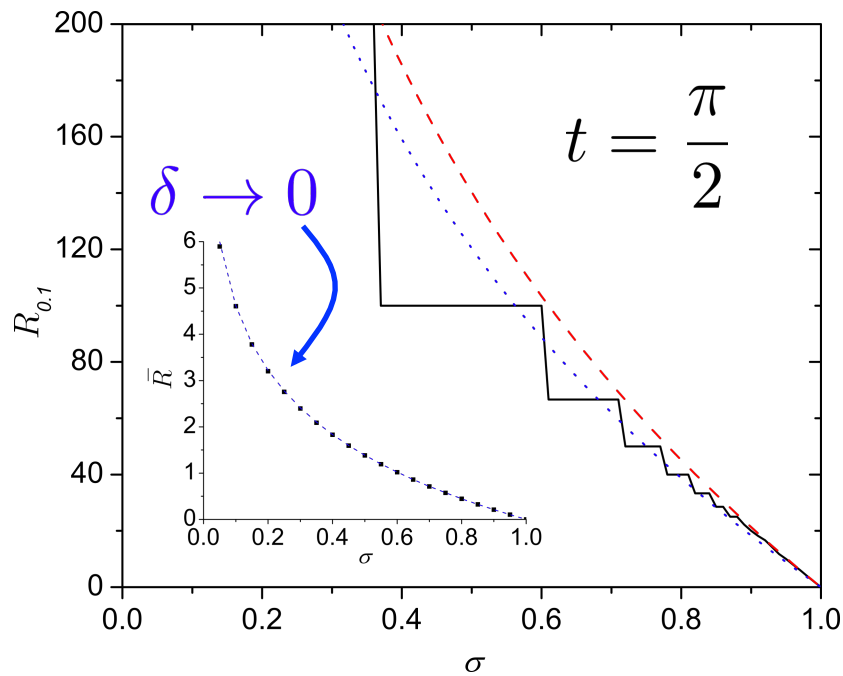


Redundancy collapses  
onto universal curve

# Misaligned, pure $\mathcal{E}$



$$R_\delta \approx \frac{\mathcal{E} \ln |\Lambda_k(t)|^2}{\ln \delta}$$



$$\Lambda_k(t) = \cos(t) - i\sigma \sin(t)$$



# Quantum Darwinism in Practice

- Preferred states – the pointer states - survive and spawn the most information theoretic progeny, which is communicated by  $\mathcal{E}$
- Non-ideal environments:  $I(\mathcal{S} : \mathcal{F}) \approx H_{\mathcal{F}}(t) - H_{\mathcal{F}}(0)$ 
  - Rate of information gain reduced by the initial entropy or misalignment, i.e., noisy channel:  $1-h$
  - Pointer states can still be exhaustively and redundantly determined from the environment
- Redundancy is given by the “Chernoff Information”