Well Balanced Hybrid Schemes for Shallow Water Flows

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Numerical schemes for Shallow Water Flows

Motivation

- Engineering applications of Shallow Water Flows
 - ocean and hydraulic engineering
 - Flows in rivers, reservoirs and coastal areas
- Stationary or nearly stationary solutions need to be accurately computed.
- Numerical treatment of source terms derived from the bottom topography is a very delicate issue
 - Well Balanced Schemes [Greenberg&Leroux,96]
 - Source term Upwinding [Roe-86, Bermudez&Vazquez-Cendon-94]

Shallow Water Flow

Derived from the Navier-Stokes model after:

- depth averaging
- hydrostatic hypothesis
- neglecting viscosity and turbulence
- not considering wind effects nor Coriolis force.



System of conservation laws plus a source term:

$$U_t + F(U)_x + E(U)_y = S$$

$$\begin{pmatrix} h \\ q_1 \\ q_2 \end{pmatrix}_t + \begin{pmatrix} q_1 \\ \frac{q_1^2}{h} + \frac{1}{2}gh^2 \\ \frac{q_1q_2}{h} \end{pmatrix}_x + \begin{pmatrix} q_2 \\ \frac{q_1q_2}{h} \\ \frac{q_2^2}{h} + \frac{1}{2}gh^2 \end{pmatrix}_y = \begin{pmatrix} 0 \\ -ghz_x \\ -ghz_y \end{pmatrix}$$

$$q_1 = hv_x, \quad q_2 = hv_y$$

Numerical Methods - Source term upwinding

$U_t + F(U)_x + E(U)_y = S(U)$

[Roe, Proc. NL. Hyp. Prob., 86]

[LeVeque, JCP-98]

[Bermudez, Vazquez, C&F, 94]

[Greenberg, Le Roux, SINUM 96]

Point-wise evaluation of S(U) does NOT work Fractional splitting is NOT a good idea C-property Well Balanced schemes



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Well Balanced schemes

LeVeque's test, JCP98: Quasi-stationary flow.



Numerical Methods - Source term upwinding

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[Roe, Proc. NL. Hyp. Prob., 86]	Point-wise evaluation of $S(U)$ does NOT work	
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[Greenberg, Le Roux, SINUM 96]	Well Balanced schemes	

The spurious waves generated by schemes that are not well-balanced can distort completely the numerical solution.

General Strategy: Combine

- conservative scheme for homogeneous conservation laws
- adequate, upwind discretization of the source term

Things to address:

- Well-Balanced, High order schemes
- Wet/Dry fronts, existing and forming.

Contents







Numerical Treatment of source terms: Automatic Well Balancing





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2. The scheme of Gascon and Corberan JCP 2001





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- 2. The scheme of Gascon and Corberan JCP 2001
- 3. The 1J-2J scheme for Balance Laws, Caselles, Haro, RD, C&F 2008



Numerical Treatment of source terms: Automatic Well Balancing

- 2. The scheme of Gascon and Corberan JCP 2001
- 3. The 1J-2J scheme for Balance Laws, Caselles, Haro, RD, C&F 2008
- 4. A second order WB hybrid scheme, Martinez-Gavara, RD, (in preparation)



$$\partial_t \vec{U} + \vec{\nabla} \vec{\mathcal{F}}(\vec{U}) = \mathcal{B}(\vec{U}, \nabla \vec{U})$$

 $\vec{U} = \vec{U}(x, y, t)$ $(x, y, t) \in \Omega \times]0, T[$ + appropriate initial and boundary conditions

Discretization on Cartesian meshes.

$$\frac{\vec{U}_{ij}^{n+1} - \vec{U}_{ij}^{n}}{\Delta t} + \mathcal{D}_{ij} = \mathcal{B}_{ij}$$

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D_{*ij*} Numerical Divergence, Shock Capturing capabilities.



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\$\mathcal{D}_{ij}\$ Numerical Divergence, Shock Capturing capabilities.
 \$\vec{U}_{ij}\$ \$\approx \frac{1}{|c_{ij}|}\$ \$\sum_{c_{ij}}\$ \$\vec{U}(x,y,t)dxdy\$ Finite Volume Schemes

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- \mathcal{D}_{ij} Numerical Divergence, Shock Capturing capabilities.
- $\vec{U}_{ij} \approx \frac{1}{|c_{ij}|} \int_{c_{ij}} \vec{U}(x, y, t) dx dy$ Finite Volume Schemes
- $\vec{U}_{ij} \approx \vec{U}(x_i, y_j, t)$. Point-Value Schemes, [Shu & Osher, JCP-86].

Shu-Osher SC schemes [Shu, Osher JCP 86]

Dimension by dimension discretization in Multi-dimensions.

$$\mathcal{D}_{ij} = \frac{\vec{F}_{i+\frac{1}{2},j} - \vec{F}_{i-\frac{1}{2},j}}{\Delta x} + \frac{\vec{G}_{i,j+\frac{1}{2}} - \vec{G}_{i,j-\frac{1}{2}}}{\Delta y}$$

- Relies on construction of Robust 1-D numerical fluxes.
 - Design first for scalar conservation laws
 - High order nonlinear reconstruction of numerical fluxes for high order accuracy in space (ENO, WENO, PHM ...).
 - Extend to systems via a local characteristic approach: Need to compute a spectral decomposition (eigenvalues, λ^p , and eigenvectors L^p, R^p) of the Jacobian matrices of the flux vectors at each interface.
- Only Uniform grids [Merriman, J. Sci. Comput, 03]
- Robust, but expensive. (Relativistic Flows [RD, Font, Ibañez, Marquina JCP 99, [Marquina, Mulet JCP 2003], [Chiavassa, RD, SISC 01], [Rault, Chiavassa, RD JSC 03], Kinematic flow problems [RD,Mulet Benasque07/JsC 08])

HRSC Numerical Flux Functions

[Shu-Osher JCP-86]

$$F_{i+\frac{1}{2}} = \sum_{p} F_{i+\frac{1}{2}}^{p} R_{i+\frac{1}{2}}^{p}$$

[RD, Marquina JCP-96] Marquina Flux-Splitting = 2-Jacobian Shu-Osher framework [RD, Mulet AMFM-06]

$$F_{i+\frac{1}{2}}^{M} = \sum_{p} F_{i+\frac{1}{2}}^{p,+} R^{p}(U_{L}) + F_{i+\frac{1}{2}}^{p,-} R^{p}(U_{R})$$

Other options

 $F_{i+\frac{1}{2}}^{p}$ characteristic-fluxes (constructed 'as in the scalar-case')

use characteristic information (L^p, R^p, λ^p) obtained from the Jacobian matrix $\frac{\partial F}{\partial U}$.

Numerical treatment of Source terms

Well Balanced Schemes [Greenberg & LeRoux, SINUM-96]: Schemes that preserve steady states at the discrete level.

The C-property [Bermudez & Vazquez-Cendon, C&F-94]: A scheme satisfies the *exact* C-property if it preserves exactly stationary steady states (non-moving water). If it is not exact, but it is accurate of order $O(\triangle x^2)$, it satisfies the approximate C-property.

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Source Term upwinding, Roe-86 $w_t + aw_x = s(x, w) = k(x)w$ (a > 0) Characteristic construction: $\frac{dw}{dt} = k \cdot w$ along $\frac{dx}{dt} = a$

$$w(x,t) = w(x - at, 0) + \int_0^t k(x - a(t - s))w(x - a(t - s), s)ds$$

There is an upwind domain of dependence determined by *a*. 'Upwinding' in the source term discretization is essential for Well Balancing (**WB**)

The scheme of Gascon & Corberan: Automatic WB

Consider the scalar equation

$$w_t + f(w)_x = s(x, w),$$

The scheme of Gascon & Corberan: Automatic WB

Consider the scalar equation $w_t + f(w)_x = s(x, w)$,

When
$$w = w(x) \rightarrow f(w)_x = s(x, w) \rightarrow f(w) = K + \int_{\bar{x}}^x s(y, w(y)) dy$$

Construct schemes that treat the flux and any primitive of the source term in an analogous fashion: Flux upwinding \rightarrow source term upwinding

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Construct schemes that treat the flux and any primitive of the source term in an analogous fashion: Flux upwinding \rightarrow source term upwinding

Define $b(x,w) = -\int_{\bar{x}}^{x} s(y,w(y,t))dy, \rightarrow b_x = -s$

$$w_t + f_x = -b_x \quad \rightarrow \quad w_t + (f+b)_x = 0 \quad \rightarrow \quad w_t + g_x = 0$$

g(x,w) = f(w) + b(x,w)

[G&C JCP-2001] Construction of a second order TVD scheme for $w_t + g_x = 0$,

$$w_j^{n+1} = w_j^n - \lambda \left\{ \bar{g}_{j+1/2}^n - \bar{g}_{j-1/2}^n \right\}$$

G & C JCP-01

Gascon and Corberan start from an explicit three point scheme of the form

$$w_j^{n+1} = w_j^n - \lambda \left(\hat{g}_{j+1/2}^{n+1/2} - \hat{g}_{j-1/2}^{n+1/2} \right)$$

with a Lax-Wendroff type flux

$$\hat{g}_{j+1/2}^{n+1/2} = \frac{1}{2} \left(g_j^n + g_{j+1}^n - \lambda \left. \frac{\partial g}{\partial w} \right|_{j+1/2}^n (g_{j+1}^n - g_j^n) \right)$$

$$g_i^n = f_i^n + b_i^n; \quad b_i^n = \int_{\bar{x}}^{x_i} s(y, w(y, t_n)) dy, \qquad \frac{\partial g}{\partial w} = \frac{\partial f}{\partial w} + \frac{\partial b}{\partial w}$$

$$b_{i+1} = b_i + b_{i,i+1}, \qquad b_{i,i+1} = \int_{x_i}^{x_{i+1}} s(y, w(y, t_n)) dy$$

$$g_{i+1} - g_i = f_{i+1} - f_i + b_{i,i+1}$$

and the scheme can be rewriten so that only $b_{j,j+1}$ need to be computed.

Steady (smooth) flow: $\equiv u_t = 0$, $\Rightarrow f_x = s = -b_x \Leftrightarrow g_x = 0$

$$f_{i+1} - f_i = \int_{x_i}^{x_{i+1}} f_x = \int_{x_i}^{x_{i+1}} s = \int_{x_i}^{x_{i+1}} -b_x = -(b_{i+1} - b_i) = -b_{i,i+1}$$

$$\Rightarrow \quad g_{i+1} - g_i = 0 \quad \equiv \quad f_{i+1} - f_i + b_{i,i+1} = 0$$

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Steady (smooth) flow:
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 $\Rightarrow g_{i+1} - g_i = 0 \equiv f_{i+1} - f_i + b_{i,i+1} = 0$

Let u(x,t) = u(x) be a stationary solution of $u_t + f_x = s$. Let $u_i^0 = u(x_i)$. Then for $n = 0, \forall j$

$$g_{j+1}^{n} = g_{j}^{n}, \quad \forall j, \quad \Rightarrow \quad \hat{g}_{j+1/2}^{n+1/2} = \frac{1}{2} \left(g_{j}^{n} + g_{j+1}^{n} - \lambda \left. \frac{\partial g}{\partial w} \right|_{j+1/2}^{n} \left(g_{j+1}^{n} - g_{j}^{n} \right) \right) = g_{j}^{n}$$

Hence	$g_{j+1}^n = g_j^n,$	$\forall j,\forall n>0,\forall j$	\Rightarrow	$\hat{g}_{j+1/2}^{n+1/2} - \hat{g}_{j-1/2}^{n+1/2} = 0$
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The scheme

- is exactly well balanced (Greenberg& Leroux)
- satisfies the exact C-property (Bermudez& Vazquez-Cendon)

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- is exactly well balanced (Greenberg& Leroux)
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Incorporating the source term into the flux divergence \Rightarrow Flux-Gradient and Source-term Balancing (FSB) \Rightarrow Automatic WB



FSB and Automatic WB

FSB: $f_{i+1} - f_i + b_{i,i+1} = 0$, $\forall i$, $b_{i,i+1} = \int_{x_i}^{x_i+1} s(y, w(y, t)) dy$. Key to WB

OBS! The computation of $b_{i,i+1}$ is carried out by numerical integration,

$$\begin{cases} b_{i,i+1} = \hat{b}_{i,i+1} + O(\triangle x^{p+1}), \\ b_i = \hat{b}_i + O(\triangle x^p) \end{cases} g_i = f(u_i) + b_i \approx \hat{g}_i = f(u_i) + \hat{b}_i, \\ \hat{g}_{j+1} - \hat{g}_j = f_{j+1} - f_j + \hat{b}_{i,i+1} = f_{j+1} - f_j + b_{i,i+1} + O(\triangle x^{p+1}) \\ = g_{j+1} - g_j + O(\triangle x^{p+1}) \end{cases}$$

In steady smooth flow, $\hat{g}_{j+1}^n - \hat{g}_j^n = O(\bigtriangleup x^{p+1})$, hence

$$\hat{g}_{j+1/2}^{n+1/2} - \hat{g}_{j-1/2}^{n+1/2} = O(\triangle x^{p+1}) \quad \to \quad \frac{1}{\triangle x} (\hat{g}_{j+1/2}^{n+1/2} - \hat{g}_{j-1/2}^{n+1/2}) = O(\triangle x^{p})$$

The scheme satisfies the approximate C-property (B&V-C) when $p \ge 2$.

Embid's Problem

Scalar model for gas flow through a duct of variable cross section.

$$\begin{cases} w_t + (w^2/2)_x = (6x - 3)w, & 0 < x < 1\\ w(0, t) = 1; w(1, t) = -.1 \end{cases}$$

Steady solution:
$$w(x) = \begin{cases} 1 + 3x^2 - 3x, & x < x_j \\ -0.1 + 3x^2 - 3x, & x > x_j \end{cases}$$



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Curbing the oscillations: G&C TVD scheme

$$\begin{split} \hat{g}_{j+1/2}^{n+1/2} &= \frac{1}{2} \left(g_j^n + g_{j+1}^n - \lambda \left. \frac{\partial g}{\partial w} \right|_{j+1/2}^n (g_{j+1}^n - g_j^n) \right) \\ & \left. \frac{\partial g}{\partial w} = \frac{\partial f}{\partial w} - \frac{\partial b}{\partial w}, \qquad \alpha \approx \lambda \frac{\partial f}{\partial w}, \qquad \beta \approx \frac{\partial b}{\partial w} \\ \alpha_{j+1/2} &= \lambda \begin{cases} \left. \frac{f_{j+1} - f_j}{w_{j+1} - w_j}, \\ \left. \frac{\partial f}{\partial w} \right|_{j+1/2}, \end{cases} & \beta_{j+1/2} = \lambda \begin{cases} \left. \frac{b_{j+1} - b_j}{w_{j+1} - w_j}, & \text{if } w_{j+1} - w_j \neq 0 \\ 0, & \text{if } w_{j+1} - w_j = 0 \end{cases} \end{split}$$

Curbing the oscillations: G&C TVD scheme

Adjust h(x) so that scheme is TVD \rightarrow CFL-like restrictions on $\alpha_{i+1/2} + \beta_{i+1/2}$ First order scheme!. Upgrade to second order following Harten JCP-97

Curbing the oscillations: G&C TVD scheme

$$\begin{split} \hat{g}_{j+1/2}^{n+1/2} &= \frac{1}{2} \left(g_j^n + g_{j+1}^n - \lambda \left. \frac{\partial g}{\partial w} \right|_{j+1/2}^n (g_{j+1}^n - g_j^n) \right) \\ & \frac{\partial g}{\partial w} = \frac{\partial f}{\partial w} - \frac{\partial b}{\partial w}, \qquad \alpha \approx \lambda \frac{\partial f}{\partial w}, \qquad \beta \approx \frac{\partial b}{\partial w} \\ \alpha_{j+1/2} &= \lambda \begin{cases} \frac{f_{j+1} - f_j}{w_{j+1} - w_j}, & \beta_{j+1/2} = \lambda \begin{cases} \frac{b_{j+1} - b_j}{w_{j+1} - w_j}, & \text{if } w_{j+1} - w_j \neq 0 \\ 0, & \text{if } w_{j+1} - w_j = 0 \end{cases} \\ \end{split}$$
Search for a TVD scheme of the form
$$\begin{split} w_j^{n+1} &= w_j^n - \lambda (\bar{g}_{j+1/2} - \bar{g}_{j-1/2}) \\ \bar{g}_{j+1/2} &= \frac{1}{2} \left(g_j + g_{j+1} - h(\alpha_{j+1/2} + \beta_{j+1/2})(g_{j+1} - g_j) \right) \end{split}$$

Adjust h(x) so that scheme is TVD \rightarrow CFL-like restrictions on $\alpha_{i+1/2} + \beta_{i+1/2}$

First order scheme!. Upgrade to second order following Harten JCP-97 Applications: Nozze-flows

Revisiting Gascon and Corberan TVD scheme

Anna Martinez Gavara PhD Thesis.

BUT!!

- Solutions to scalar balance laws might not be TVD! TVD restrictions might be too stringent on numerical schemes for general balance laws.
- Solution CFL-like restrictions based on $\alpha_{i+1/2} + \beta_{i+1/2}$ might be too restrictive

FSB for Shallow Water Flows

Incorporate the source term in the flux divergence for shallow water flows

Split source term,
$$S = S_1 + S_2 = \begin{pmatrix} 0 \\ -ghz_x \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ 0 \\ -ghz_y \end{pmatrix}$$

$$B(x, y, t) = -\int_0^x S_1(s, y, t)ds; \qquad \partial_x B = -S_1$$

Define functions:

$$C(x,y,t) = -\int_0^y S_2(x,s,t)ds; \qquad \partial_y C = -S_2$$

Formally Rewrite the original system in 'conservation form' as follows,

 $U_t + (F + B)_x + (E + C)_y = 0$ $U_t + (G)_x + (H)_y = 0$

 \rightarrow combined fluxes: physical fluxes + primitive of source term

Shu-Osher HRSC approach for FSB schemes

Semi-discrete formulation on Cartesian meshes (separate spatial and temporal accuracy).

$$\partial_t \vec{U}_{ij} + \frac{\vec{G}_{i+\frac{1}{2},j} - \vec{G}_{i-\frac{1}{2},j}}{\Delta x} + \frac{\vec{H}_{i,j+\frac{1}{2}} - \vec{H}_{i,j-\frac{1}{2}}}{\Delta y} = 0$$

- Accuracy in time: Runge-Kutta schemes $\vec{U}_{ij}^n \Rightarrow \vec{U}_{ij}^{n+1}$.
- Relies on construction of Robust 1-D numerical fluxes.
 - Design first for scalar conservation laws
 - High order nonlinear reconstruction of numerical fluxes for high order accuracy in space (ENO, WENO, PHM ...).
 - Extend to systems via a local characteristic approach: Need to compute a spectral decomposition (eigenvalues, λ^p , and eigenvectors L^p, R^p) of the Jacobian matrices of the flux vectors at each interface.


[Caselles, Haro, RD C& F 09]

Start with the scalar equation

$$w_t + f(w)_x = s(x, w),$$

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Define $b(x,t) = -\int_{\bar{x}}^{x} s(y,w(y,t))dy, \rightarrow b_x = -s$

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Rewrite $w_t + g_x = 0$, g(x, t) = f(w(x, t)) + b(x, t)

[Caselles, Haro, RD C& F 09]

b(x,t)

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Rewrite $w_t + g_x = 0$, g(x,t) = f(w(x,t)) + b(x,t)

Semi-discrete formulation: $\partial_t w_i + \frac{1}{\Delta x}(g_{i+1/2} - g_{i-1/2}) = 0$

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Semi-discrete formulation: $\partial_t w_i + \frac{1}{\Delta x}(g_{i+1/2} - g_{i-1/2}) = 0$

 $g_{i+1/2}$: Shu-Osher RF algorithm with upwind direction determined by f(w)

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Semi-discrete formulation: $\partial_t w_i + \frac{1}{\Delta x}(g_{i+1/2} - g_{i-1/2}) = 0$

 $g_{i+1/2}$: Shu-Osher RF algorithm with upwind direction determined by f(w)Higher order: Use nonlinear reconstruction process on $g_i = f_i + b_i$

OBS!: $b_i = \int^{x_i} s(w) dx$ has to be computed by numerical integration!

Seek to involve only $b_{l,l+1} = \int_{x_l}^{x_{l+1}} s(w) dx$

Shu-Osher RF algorithm:

 $g_{i+1/2} - g_{i-1/2} = \left[\mathcal{G}_{i+1/2} + \mathsf{HOT}_{i+1/2} \right] - \left[\mathcal{G}_{i-1/2} + \mathsf{HOT}_{i-1/2} \right]$



Shu-Osher RF algorithm:

$$g_{i+1/2} - g_{i-1/2} = [\mathcal{G}_{i+1/2} + \mathsf{HOT}_{i+1/2}] - [\mathcal{G}_{i-1/2} + \mathsf{HOT}_{i-1/2}]$$

 $\mathcal{G}_{i\pm 1/2}$ first order contributions

$$\mathcal{G}_{i+1/2} = \begin{cases} g_i = f_i + b_i & \text{if } f' > 0 \text{ in } [w_i, w_{i+1}] \\ g_{i+1} = f_{i+1} + b_{i+1} & \text{if } f' < 0 \text{ in } [w_i, w_{i+1}] \\ \frac{1}{2}(g_i^+ + g_{i+1}^-) & \text{else} \end{cases}$$

$$\alpha_{i+1/2} = \max\{|f'(w)|, w \in [w_i, w_{i+1}]\}, \qquad \begin{cases} g_i^+ &= (g_i + \alpha_{i+1/2}w_i)\\ g_{i+1}^- &= (g_{i+1} - \alpha_{i+1/2}w_{i+1}) \end{cases}$$

HOT_{$i\pm 1/2$} Higher order terms, based on divided differences of g. Divided differences of g only depend on $b_{i,i+1}$. Shu-Osher RF algorithm:

$$g_{i+1/2} - g_{i-1/2} = [\mathcal{G}_{i+1/2} + \mathsf{HOT}_{i+1/2}] - [\mathcal{G}_{i-1/2} + \mathsf{HOT}_{i-1/2}]$$

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First order scheme, applied to $w_t + aw_x = s(x, w) = aw$ (a > 0) $w_t = \frac{1}{\Delta x} \left(\mathcal{G}_{i+1/2} - \mathcal{G}_{i-1/2} \right) = \frac{1}{\Delta x} \left(f_i - f_{i-1} + \int_{x_{i-1}}^{x_i} s(w(z,t)) dz \right)$ [Roe, 86]

Shu-Osher RF algorithm:

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Notice that $\mathcal{G}_{i+1/2} - \mathcal{G}_{i-1/2}$ can be written in terms of

$$b_{l+1} - b_l = \int_{x_l}^{x_{l+1}} s - \int_{x_l}^{x_l} s = \int_{x_l}^{x_{l+1}} s(x, w(x)) dx = b_{l,l+1}, \quad l = i - 1, i$$

[Caselles, Haro, RD C & F 09] $\mathcal{G}_{i+1/2} - \mathcal{G}_{i-1/2} = \mathcal{G}_{i+1/2}^+ - \mathcal{G}_{i-1/2}^-$

$$\mathcal{G}_{i+1/2} - \mathcal{G}_{i-1/2} = (\mathcal{G}_{i+1/2} - b_i) - (\mathcal{G}_{i-1/2} - b_i) =: \mathcal{G}_{i+1/2}^+ - \mathcal{G}_{i-1/2}^-$$

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$$\mathcal{G}_{i+1/2}^{+} = \begin{cases} f_i & \text{if } f' > 0 \text{ in } [w_i, w_{i+1}] \\ f_{i+1} + b_{i,i+1} & \text{if } f' < 0 \text{ in } [w_i, w_{i+1}] \\ \frac{1}{2}(f_i^+ + f_{i+1}^-) + \frac{1}{2}b_{i,i+1} & \text{else} \end{cases}$$

$$\mathcal{G}_{i-1/2}^{-} = \begin{cases} f_{i-1} - b_{i-1,i} & \text{if } f' > 0 \text{ in } [w_{i-1}, w_i] \\ f_i & \text{if } f' < 0 \text{ in } [w_{i-1}, w_i] \\ \frac{1}{2}(f_i^+ + f_{i-1}^-) - \frac{1}{2}b_{i-1,i} & \text{else} \end{cases}$$

$$\begin{cases} f_i^+ = f_i + \alpha_{i+1/2} w_i, \\ f_i^- = f_i - \alpha_{i-1/2} w_i \end{cases}$$



[Caselles, Haro, RD C & F 09] $\mathcal{G}_{i+1/2} - \mathcal{G}_{i-1/2} = \mathcal{G}_{i+1/2}^+ - \mathcal{G}_{i-1/2}^-$

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$$\mathcal{G}_{i+1/2}^+ - \mathcal{G}_{i+1/2}^- = b_{i,i+1} = \int_{x_i}^{x_{i+1}} -s(w(x), x) dx$$

 $\mathcal{G}_{i+1/2}^{\pm}$ split the source term contribution at the i + 1/2 interface according to the wind, as specified by f'(u).

$$\begin{split} g_{i+1/2} &- g_{i-1/2} &= \left[\mathcal{G}_{i+1/2}^+ + \operatorname{HOT}_{i+1/2} \right] - \left[\mathcal{G}_{i-1/2}^- + \operatorname{HOT}_{i-1/2} \right] \\ &= G_{i+1/2}^+ - G_{i-1/2}^- \\ \mathcal{G}_{i+1/2}^+ &= \begin{cases} f_i & \text{if } f' > 0 \text{ in } [w_i, w_{i+1}] \\ f_{i+1} + b_{i,i+1} & \text{if } f' < 0 \text{ in } [w_i, w_{i+1}] \\ \frac{1}{2}(f_i^+ + f_{i+1}^-) + \frac{1}{2}b_{i,i+1} & \text{else} \end{cases} \\ \mathcal{G}_{i-1/2}^- &= \begin{cases} f_{i-1} - b_{i-1,i} & \text{if } f' > 0 \text{ in } [w_i, w_{i+1}] \\ f_i & \text{if } f' < 0 \text{ in } [w_i, w_{i+1}] \\ \frac{1}{2}(f_i^+ + f_{i-1}^-) - \frac{1}{2}b_{i-1,i} & \text{else} \end{cases} \end{split}$$

 $HOT_{i-1/2}$ Higher order terms, based on divided differences of g.

only $b_{l,l+1} = \int_{x_l}^{x_{l+1}} s(w) dx$ are involved !!

For f(w) nonlinear, $G_{i+1/2}^{\pm}$ lead to an automatic upwind splitting of the source term contribution at the i + 1/2 cell.

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Extension to 1D-systems:

$$U_t + G(U)_x = 0 \quad \Rightarrow \quad \partial_t U_i + \frac{1}{\Delta x} \left(G_{i+1/2}^+ - G_{i-1/2}^- \right) = 0$$

Construction of $G_{i+1/2}^{\pm}$: Use spectral decomposition of $\frac{\partial F}{\partial U}$

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Follow the basic design structure of Marquina Flux-Splitting technique:
The 2J scheme: $G_{i+1/2}^{\pm} = \sum_{p} (G_{i+1/2}^{p,\pm})^{L} R^{p} (U^{L}) + (G_{i+1/2}^{p,\pm})^{R} R^{p} (U^{R})$

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- If U^L = U^R = U^{*} (Plain Extension of Shu-Osher to HCL+source term) The 1J Scheme: G[±]_{i+1/2} = ∑_p G^{p,±}_{i+1/2} R^p(U^{*}) U^{*} ≡ U_{i+1/2} interface state.

 $G^{p,\pm}$ characteristic fluxes, computed 'as in the scalar case' for the *p*th field.

Numerical integration crucial for Well Balancing

$$\left(\begin{array}{c}h\\q\end{array}\right)_t + \left(\begin{array}{c}q\\\frac{q^2}{h} + \frac{1}{2}gh^2\end{array}\right)_x = \left(\begin{array}{c}0\\-ghz_x\end{array}\right)$$

$$B_{i,i+1}^{n} = \begin{pmatrix} 0 \\ \int_{x_{i}}^{x_{i+1}} ghz_{x} dx \end{pmatrix}, \quad \int_{x_{i}}^{x_{i+1}} ghz_{x} dx \approx \frac{g}{2}(z_{i+1}-z_{i})(h_{i}+h_{i+1}) + O(\triangle x^{3})$$

for quiescent steady flows: q = 0, h + z = constant,

$$\int_{x_i}^{x_{i+1}} ghz_x dx = \int_{x_i}^{x_{i+1}} gh(-h)_x dx = -\frac{g}{2}(h_i^2 - h_{i+1}^2) = \frac{g}{2}(z_{i+1} - z_i)(h_i + h_{i+1})$$

- JJ-scheme exactly well balanced for stationary flows.
- **2**J-scheme only approximately well balanced, order ≥ 2 .



LeVeque's test, JCP98: Quasi-stationary flow.





The 1J-2J scheme can manage wet/dry fronts.



Drain on a non-flat bottom

Dry bed generation + non-flat bottom

Slight loss of conservation in the 2J extension to systems.

A. Martinez-Gavara PhD Thesis Start with the scalar equation

$$w_t + f(w)_x = s(w),$$

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Define

$$b(x,t) = -\int_{\bar{x}}^{x} s(y,w(y,t))dy, \quad \rightarrow \quad b_x = -s$$

Rewrite $w_t + g_x = 0$, g(w, x, t) = f(w) + b(x, t)

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Look for a Lax-Wendroff type second order scheme:

$$U_{i}^{n+1} = U_{i}^{n} - \lambda (\tilde{g}_{i+\frac{1}{2}}^{n+\frac{1}{2}} - \tilde{g}_{i-\frac{1}{2}}^{n+\frac{1}{2}}) \qquad \text{with} \qquad \tilde{g}_{i+\frac{1}{2}}^{n+\frac{1}{2}} := \hat{g}_{i+\frac{1}{2}}^{n} + \frac{\Delta t}{2} \frac{\partial \hat{g}}{\partial t} \Big|_{i+\frac{1}{2}}^{n}.$$

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Second order accuracy:
$$\begin{cases} \hat{g}_{i+\frac{1}{2}}^{n} := \frac{1}{2} (g_{i+1}^{n} + g_{i}^{n}), \\\\ \hat{g}_{t} \Big|_{i+\frac{1}{2}}^{n} := -f_{w} \Big|_{i+\frac{1}{2}}^{n} \frac{g_{i+1}^{n} - g_{i}^{n}}{\Delta x} + b_{t} \Big|_{i+\frac{1}{2}}^{n}, \\\\ b_{t} \Big|_{i+\frac{1}{2}}^{n} = \int_{0}^{x_{i+\frac{1}{2}}} \frac{\partial s}{\partial u} (y, u(y, t_{n})) g_{y}(y, t_{n}) dy, \end{cases}$$

Using TRAPEZOIDAL RULE

$$b_t \Big|_{i+\frac{1}{2}}^n - b_t \Big|_{i-\frac{1}{2}}^n = \int_{x_{i-1/2}}^{x_{i+1/2}} s_u(y, u(y, t_n)) g_x(y, t_n) dy$$
$$= \left(s_u \Big|_{i+\frac{1}{2}}^n \frac{g_{i+1}^n - g_i^n}{\Delta x} + s_u \Big|_{i-\frac{1}{2}}^n \frac{g_i^n - g_{i-1}^n}{\Delta x} \right) \frac{\Delta x}{2} + \mathcal{O}(\Delta x^3)$$

$$U_i^{n+1} = U_i^n - \frac{\Delta t}{\Delta x} (\mathcal{G}_{i+\frac{1}{2}}^n - \mathcal{G}_{i-\frac{1}{2}}^n) - \frac{\Delta t}{\Delta x} \mathcal{S}_i^n$$

$$\begin{aligned} \mathcal{G}_{i+\frac{1}{2}}^{n} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{S}_{i}^{n} &= \frac{1}{2} \left(\beta_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) + \beta_{i-\frac{1}{2}}^{n} (g_{i}^{n} - g_{i-1}^{n}) \right) \end{aligned}$$



Automatic Well Balancing

Let u(x) be a stationary solution $\Leftrightarrow u_t = 0 \Leftrightarrow g_x = 0 \Leftrightarrow g_{i+1} = g_i$

$$\begin{aligned} \mathcal{G}_{i+\frac{1}{2}}^{n} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) = g_{i}^{n} \\ \mathcal{S}_{i}^{n} &= \frac{1}{2} \left(\beta_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) + \beta_{i-\frac{1}{2}}^{n} (g_{i}^{n} - g_{i-1}^{n}) \right) = 0 \end{aligned}$$

$$U_i^{n+1} = U_i^n - \frac{\Delta t}{\Delta x} (\mathcal{G}_{i+\frac{1}{2}}^n - \mathcal{G}_{i-\frac{1}{2}}^n) - \frac{\Delta t}{\Delta x} \mathcal{S}_i^n = U_i^n \qquad \text{exactly WB}$$

$$g_{i+1} - g_i = f_{i+1} - f_i + b_{i,i+1}$$
 if $b_{i,i+1} = \int_{x_i}^{x_{i+1}} s = \hat{b}_{i,i+1} + O(\triangle x^{p+1})$

 $\mathcal{G}_{i+\frac{1}{2}}^n - \mathcal{G}_{i-\frac{1}{2}}^n = O(\triangle x^{p+1}) \qquad \mathcal{S}_i^n = O(\triangle x^{p+1}) \quad \Rightarrow \quad U_i^{n+1} = U_i^n + O(\triangle x^{p+1})$

The scheme satisfies the approximate C-property ($p \ge 2$).

Greenberg&Leroux tests [SINUM,96]



Curbing the oscillations: A WB second order hybrid scheme



$$\begin{aligned} \text{Second order method} \qquad U_i^{n+1} &= U_i^n - \frac{\Delta t}{\Delta x} (\mathcal{G}_{i+\frac{1}{2}}^n - \mathcal{G}_{i-\frac{1}{2}}^n) - \frac{\Delta t}{\Delta x} \mathcal{S}_i^n \\ \mathcal{G}_{i+\frac{1}{2}}^n &= \mathcal{G}_{i+\frac{1}{2}}^{HI} = \frac{1}{2} \left(g_{i+1}^n + g_i^n - \alpha_{i+\frac{1}{2}}^n (g_{i+1}^n - g_i^n) \right), \qquad \alpha_{i+\frac{1}{2}}^n = \lambda \frac{\partial f}{\partial u} \Big|_{i+\frac{1}{2}}^n \end{aligned}$$

Second order method
$$U_i^{n+1} = U_i^n - \frac{\Delta t}{\Delta x} (\mathcal{G}_{i+\frac{1}{2}}^n - \mathcal{G}_{i-\frac{1}{2}}^n) - \frac{\Delta t}{\Delta x} \mathcal{S}_i^n$$
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Lax-Wendroff-type scheme: Oscillatory behavior at discontinuous fronts

Apply the Flux-Limiter technology to curb oscillations.

$$\begin{split} \text{Second order method} \qquad U_{i}^{n+1} &= U_{i}^{n} - \frac{\Delta t}{\Delta x} (\mathcal{G}_{i+\frac{1}{2}}^{n} - \mathcal{G}_{i-\frac{1}{2}}^{n}) - \frac{\Delta t}{\Delta x} \mathcal{S}_{i}^{n} \\ \mathcal{G}_{i+\frac{1}{2}}^{n} &= \mathcal{G}_{i+\frac{1}{2}}^{HI} = \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right), \qquad \alpha_{i+\frac{1}{2}}^{n} = \lambda \frac{\partial f}{\partial u} \Big|_{i+\frac{1}{2}}^{n} \\ \mathcal{G}_{i+\frac{1}{2}}^{n} &= \mathcal{G}_{i+\frac{1}{2}}^{LO} + \phi_{i+\frac{1}{2}} (\mathcal{G}_{i+\frac{1}{2}}^{HI} - \mathcal{G}_{i+\frac{1}{2}}^{LO}) \\ \mathcal{G}_{i+\frac{1}{2}}^{n} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \operatorname{sign}(\alpha_{i+\frac{1}{2}}^{n}) (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}}^{n} - g_{i+\frac{1}{2}^{n}} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n} - g_{i+\frac{1}{2}^{n}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n} - g_{i+\frac{1}{2$$

$$\begin{split} \text{Second order method} \qquad U_{i}^{n+1} &= U_{i}^{n} - \frac{\Delta t}{\Delta x} (\mathcal{G}_{i+\frac{1}{2}}^{n} - \mathcal{G}_{i-\frac{1}{2}}^{n}) - \frac{\Delta t}{\Delta x} \mathcal{S}_{i}^{n} \\ \mathcal{G}_{i+\frac{1}{2}}^{n} &= \mathcal{G}_{i+\frac{1}{2}}^{HI} = \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right), \qquad \alpha_{i+\frac{1}{2}}^{n} = \lambda \frac{\partial f}{\partial u} \Big|_{i+\frac{1}{2}}^{n} \\ \mathcal{G}_{i+\frac{1}{2}}^{n} &= \mathcal{G}_{i+\frac{1}{2}}^{LO} + \phi_{i+\frac{1}{2}} (\mathcal{G}_{i+\frac{1}{2}}^{HI} - \mathcal{G}_{i+\frac{1}{2}}^{LO}) \\ \mathcal{G}_{i+\frac{1}{2}}^{n} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \operatorname{sign}(\alpha_{i+\frac{1}{2}}^{n}) (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{G}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{O}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{O}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} + g_{i}^{n} - \alpha_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{O}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} - g_{i}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{O}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+1}^{n} - g_{i}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+1}^{n} - g_{i}^{n}) \right) \\ \mathcal{O}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}}^{n} - g_{i}^{n}) \right) \\ \mathcal{O}_{i+\frac{1}{2}}^{HI} &= \frac{1}{2} \left(g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} - g_{i+\frac{1}{2}}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}}^{n} - g_{i+\frac{1}{2}^{n}}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}}^{n} - g_{i+\frac{1}{2}^{n}^{n} - g_{i+\frac{1}{2}^{n}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n}}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n}}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n} (g_{i+\frac{1}{2}^{n}^{n} - g_{i+\frac{1}{2}^{n}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n}^{n} (g_{i+\frac{1}{2}^{n} - g_{i+\frac{1}{2}^{n}^{n}^{n} (g_{i+\frac{1}{2}^{n}^{n} - g_{i+\frac{1}{2}^{n}^{n}^{n} (g_{i+\frac{1}{2}^{n}^{n} (g_{i+\frac{1}{2}^{n}^{n} - g_{i+\frac{1}{2}^{n}$$

 $g_{i+1} = g_i, \quad \forall i \to U^{n+1} = U^n$ Well Balanced



The WBH2 scheme for 1D Shallow water flows

$$\left(\begin{array}{c}h\\q\end{array}\right)_t + \left(\begin{array}{c}q\\\frac{q^2}{h} + \frac{1}{2}gh^2\end{array}\right)_x = \left(\begin{array}{c}0\\-ghz_x\end{array}\right)$$

$$U_i^{n+1} = U_i^n - \frac{\Delta t}{\Delta x} \left(\mathcal{G}_{i+1/2}^n - \mathcal{G}_{i-1/2}^n \right) - \frac{\Delta t}{\Delta x} \mathcal{S}_i^n$$
$$\mathcal{G}_{i+1/2}^n = \sum_{p=1}^2 \mathcal{G}_{i+1/2}^{n,p} R_{i+1/2}^{n,p}$$

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- Roe's average state at the interface.
- **J** TVDB $\mathcal{G}_{i+1/2}^{n,p}$ *p*th characteristic flux
- $\alpha_{i+\frac{1}{2}}^n = p$ -th eigenvalue of the Jacobian matrix $\frac{\partial F}{\partial U}$.

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$$\mathcal{G}_{i+1/2}^n = \sum_{p=1}^2 \mathcal{G}_{i+1/2}^{n,p} R_{i+1/2}^{n,p}$$

$$\mathcal{S}_{i}^{n} = \frac{1}{2} \left(\beta_{i+\frac{1}{2}}^{n} (F_{i+1}^{n} - F_{i}^{n} + B_{i,i+1}^{n}) + \beta_{i-\frac{1}{2}}^{n} (F_{i}^{n} - F_{i-1}^{n} + B_{i-1,i}^{n}) \right)$$

$$\beta_{i+1/2} = \frac{\partial S}{\partial U}\Big|_{i+\frac{1}{2}}^{n} = \left(\begin{array}{cc} 0 & 0\\ -gz_{x}\Big|_{i+\frac{1}{2}}^{n} & 0 \end{array}\right) \qquad B_{i,i+1}^{n} = \left(\begin{array}{c} 0\\ \int_{x_{i}}^{x_{i+1}} ghz_{x} dx \end{array}\right)$$

$$\int_{x_i}^{x_{i+1}} ghz_x dx \approx \frac{g}{2} (z_{i+1} - z_i)(h_i + h_{i+1}); \qquad z_x \big|_{i+\frac{1}{2}} = \frac{z_{i+1} - z_i}{\Delta x}$$

Dam Break over a discontinuous topography

Surface level h + z. Left: Numerical h + z using 400 grid cells, and bottom topography. Right: Numerical h + z using 400 and 4000 grid cells. Top: t = 15s, Bottom t = 60s


2D-extension: The C-property

Grid size	l_1 -error
128×128	$1.215 \cdot 10^{-15}$
256×256	$2.627 \cdot 10^{-15}$
512×512	$7.690 \cdot 10^{-16}$



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2-D Test: Quasi-stationary flow

Same topography: $z(x, y) = 0.5e^{-50((x-0.5)^2 + (y-0.5)^2)}$. Initial Data

$$h(x,y) = \begin{cases} 1.01 - z(x,y), & 0.1 < x < 0.2; \\ 1 - z(x,y), & otherwise. \end{cases} \begin{pmatrix} q_1(x,y) \\ q_2(x,y) \end{pmatrix} = 0$$





Perspectives

Treatment of Dry areas.

Combine WBH2 with Adaptive Mesh Refinement techniques for complex situations.

Thanks for your attention!