

# Lenosky's energy and the phonon dispersion of graphene

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# Lenosky Model

- Introduced for negatively curved graphene
- Takes into account orbital overlap due to curvature

$$UL_d = \frac{\epsilon_0}{2} \sum_{\langle ij \rangle} (|\mathbf{r}_{ij}| - |\mathbf{e}_{ij}|)^2$$

Stretching

$$UL_a = \epsilon_1 \sum_i \left( \sum_{\langle j \rangle} \hat{\mathbf{r}}_{ij} \right)^2$$

~ Dangling bond

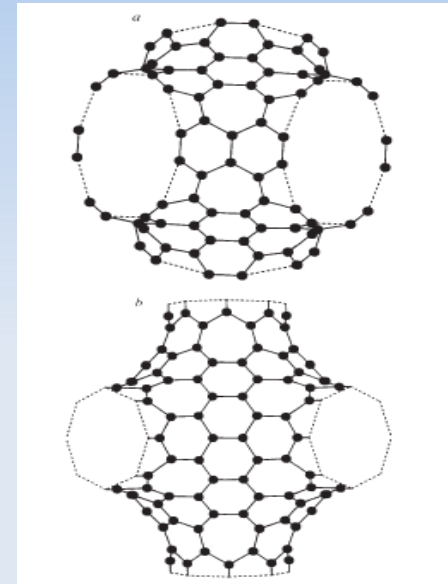
$$UL_{b1} = \epsilon_2 \sum_{\langle ij \rangle} (1 - \hat{\mathbf{n}}_i \cdot \hat{\mathbf{n}}_j)$$

$\pi$ - $\pi$  overlap

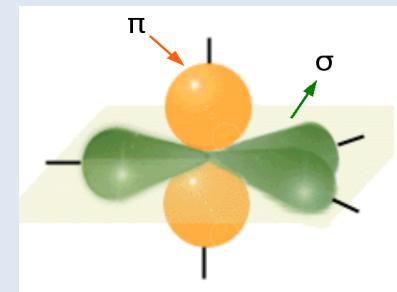
$$UL_{b2} = \epsilon_3 \sum_{\langle ij \rangle} (\hat{\mathbf{n}}_i \cdot \hat{\mathbf{r}}_{ij}) (\hat{\mathbf{n}}_j \cdot \hat{\mathbf{r}}_{ji})$$

$\pi$ - $\sigma$  overlap

Schwarzite



Lenosky *et al.*  
Nature **355** 333 (1992)

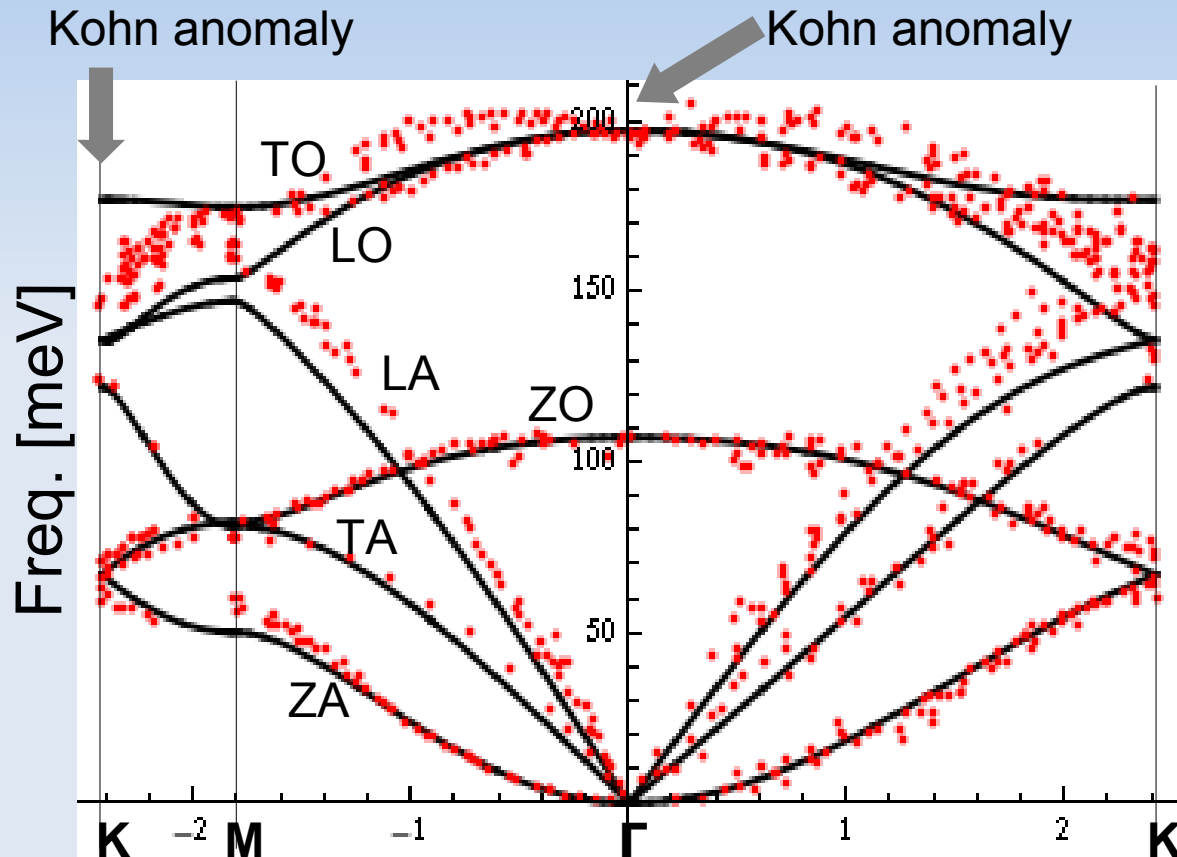


We calculate the phonon dispersion for the Lenosky model



# Hybrid model

We add the bond-bending term to Lenosky's energy:



**Good Fit!**

**Five free parameters:**

$\epsilon_0=30$  eV  $\rightarrow$  stretch

$\epsilon_1=1.34$  eV  $\rightarrow$  dangling

$\epsilon_2=0.23$  eV  $\rightarrow$   $\pi$ - $\pi$

$\epsilon_3=1.16$  eV  $\rightarrow$   $\pi$ - $\sigma$

$\beta=2.40$  eV  $\rightarrow$  bending

$\pi$ - $\sigma$   $\rightarrow$  overlap responsible for ZO flattening  
 $\rightarrow$  dominates over  $\pi$ - $\pi$  overlap